

# Molecular dynamics simulation of phase transitions and melting in MgSiO<sub>3</sub> with the perovskite structure—Comment

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Chaplot et al. (1998) purport to simulate MgSiO<sub>3</sub>-perovskite melting. Instead, the paper contains simulations of a thermal instability.

One can define thermal instability of a crystal lattice as a mechanical instability induced by temperature. The temperature of thermal instability can be found by gradual increase of temperature until the sudden change of properties appears. This is exactly what was performed by Chaplot et al. The temperatures of thermal instability of the ideal lattice were high compare to experimental. Recognizing the possibility of overheating, Chaplot et al. carried out simulations to determine temperatures of thermal instability of the lattice containing 1% vacancies. The lattice with defects is less stable than the ideal lattice, therefore, the temperature of thermal instability for defect structure was lower than for the ideal lattice. Still, neither the simulation of ideal nor defect structures answered the question: What is the melting temperature of the MgSiO<sub>3</sub>-perovskite with the accepted model of interatomic interaction?

Melting and thermal instability are different phenomena. Temperature of melting ( $T_{\text{melt}}$ ) is defined as the temperature where Gibbs free energies of crystal and liquid phases are equal. Morris et al. (1994) showed that calculating  $T_{\text{melt}}$  by calculating Gibbs energies is equivalent to finding  $T_{\text{melt}}$  from two-phase simulations. The question is how much melting temperatures could differ from temperatures of thermal instability calculated by Chaplot et al. for ideal lattice? An exact answer could be provided by the direct comparison of two methods, as was done by Belonoshko (1998) for melting of corundum. In absence of this comparison, we can refer to other studies to obtain an approximate values of overheating in Chaplot et al. calculations.

Matsui and Price (1991) and Belonoshko (1994) have studied the melting of perovskite. Matsui and Price (1991) has calculated  $T_{\text{inst}}$ , whereas Belonoshko (1994) applied the two-phase method to calculate  $T_{\text{melt}}$ . Interesting, an application of correct method allows to reach good agreement with experiment (Fig. 1). The same model of interatomic interaction was applied in both studies. From the comparison it follows that  $T_{\text{inst}}$  is higher than  $T_{\text{melt}}$  by about 1500 K at the pressure of 25 GPa. Because overheating increases with pressure (Belonoshko 1998), the value of 1500 K is likely to be a lower bound for overheating.

The potential applied by Chaplot et al. is different from that applied by Matsui and Price (1991) and Belonoshko (1994). However, the common thing about both potentials is that they do reproduce experimental properties in good agreement with experiment (Belonoshko and Dubrovinsky 1996c). Belonoshko and Dubrovinsky (1996b) has demonstrated that a different potential may produce about the same melting curve if experimental data is reproduced. Therefore, we can rather safely accept the estimate of 2000 K for overheating in calculations by Chaplot et al. when the ideal structure was studied. As for temperatures of thermal instability of defect lattice, it is not clear how much they differ from true melting temperatures (true with respect to correctly calculated for the accepted model). In fact,  $T_{\text{inst}}$  of defect structure can be higher or lower than  $T_{\text{melt}}$ . It is possible to determine how many atoms should be taken out from an ideal structure to make it collapse at, e.g., room tem-

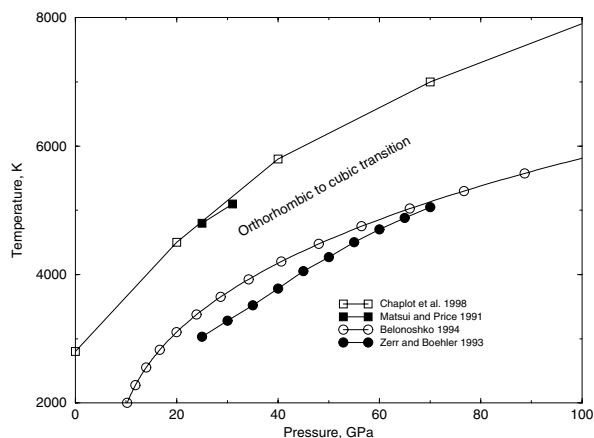


FIGURE 1. Pressure dependence of melting temperatures of MgSiO<sub>3</sub>-perovskite as suggested in the studies indicated on the legend. Matsui and Price (1991) and Chaplot et al. (1998) calculated curves provide temperatures of thermal instability, not melting temperatures, as purported. The 2-phase simulations by Belonoshko (1994) for the Matsui and Price (1991) model produces melting temperatures in close agreement with experimental data. The curve shown is a digitization of the curve used by Zerr and Boehler (1993) to represent their experimental data. The text “orthorhombic to cubic transition” is placed as close as possible to the range where it was calculated by Matsui and Price (1991) and Chaplot et al. (1998).

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perature.

From the Figure 1 the pressure- $T_{\text{inst}}$  curve is higher than experimental melting curve by about 2000 K, depending on pressure. It is likely, that if Chaplot and co-authors would use two-phase simulation method (e.g., Belonoshko and Dubrovinsky 1996a) then their melting curve would match experimental curve closely, as is the case with the simulation by Belonoshko (1994).

The orthorhombic to cubic perovskite transition found by Chaplot et al. occurs at the temperatures close to  $T_{\text{inst}}$ . It is likely that the transition is above the melting temperature and should be regarded as an artifact of erroneous simulation of melting. The same transition was observed by Matsui and Price (1991). Again, above the melting temperature (Belonoshko 1994). Chaplot et al. observed that the orthorhombic to cubic perovskite transition is accompanied by melting of the oxygen sublattice. A similar feature was observed by Belonoshko (1998). However, because Belonoshko (1994), in contrast to Chaplot et al., has calculated melting temperature, he was able to distinguish that phenomena as an artifact. Similar high ionic diffusion was found by Zhou et al. (1997) when conducting MD simulation of neighborite in the close vicinity of “melting” (as we now understand, possibly above  $T_{\text{melt}}$  but below  $T_{\text{inst}}$ ). Zhou et al. (1997) wrote that the phenomena was similar to what was observed by Matsui and Price (1991)—and now we know that that was an artifact.

The major point of this comment is that one has to model melting correctly. Otherwise it is not clear if results of simulation are meaningful or rather concerned with phenomena in the metastably overheated range. Several studies calculated melting in the style of Chaplot et al., known also under the name HUIY (heat-until-it-yields). Using HUIY method, MD simulated “melting” was reported in a number of studies (e.g., Boehler et al. 1997; Cavazzoni et al. 1999; Matsui and Price 1991; Vocadlo and Price 1996; Zhou et al. 1997).

A huge variation in melting temperatures exists in part because some studies claim to calculate melting (e.g., Chaplot et al. 1998), but instead provide calculations of instability whereas others indeed simulate melting (e.g., Belonoshko 1994).

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## REFERENCES CITED

- Belonoshko, A.B. (1994) Molecular dynamics of MgSiO<sub>3</sub> perovskite at high pressures: Equation of state, structure, and melting transition. *Geochimica et Cosmochimica Acta*, 58, 4039–4047.
- (1998) Melting of corundum using conventional and two-phase molecular dynamic simulation method. *Physics and Chemistry of Minerals*, 25, 138–141.
- Belonoshko, A.B. and Dubrovinsky, L.S. (1996a) Molecular dynamics of NaCl (B1 and B2) and MgO (B1) melting: Two-phase simulation. *American Mineralogist*, 81, 303–316.
- (1996b) Molecular and lattice dynamics study of the MgO-SiO<sub>2</sub> system using a transferable interatomic potential. *Geochimica et Cosmochimica Acta*, 60, 1645–1656.
- (1996c) Equations of state of MgSiO<sub>3</sub>-perovskite and MgO (periclase) from computer simulations. *Physics of the Earth and Planetary Interiors*, 98, 47–54 (1996).
- Boehler, R., Ross, M., and Boecker, D.B. (1997) Melting of LiF and NaCl to 1 Mbar: Systematics of ionic solids at extreme conditions. *Physical Review Letters*, 78, 4589–4592.
- Cavazzoni, C., Chiarotti, G.L., Scandolo, S., Tosatti, E., Bernasconi, M., and Parrinello, M. (1999) Superionic and metallic states of water and ammonia at giant planet conditions. *Science*, 283, 44–46.
- Chaplot, S.L., Choudhury, N., and Rao, K.R. (1998) Molecular dynamics simulation of phase transitions and melting in MgSiO<sub>3</sub> with the perovskite structure. *American Mineralogist*, 83, 937–941.
- Matsui, M., and Price, G.D. (1991) Simulation of the pre-melting behaviour of MgSiO<sub>3</sub> perovskite at high pressures and temperatures. *Nature*, 351, 735–737.
- Morris, J.R., Wang, C.Z., Ho, M.K., and Chan, C.T. (1994) Melting line of aluminum from simulations of coexisting phases. *Physical Review B*, 49, 3109–3113.
- Vocadlo, L., and Price, G. D. (1996) The melting of MgO—computer simulations via molecular dynamics. *Physics and Chemistry of Minerals*, 23, 42–49.
- Zerr, A. and Boehler, R. (1993) Melting of (Mg,Fe)SiO<sub>3</sub> perovskite to 625 kilobars: Indication of a high melting temperature in the lower mantle. *Science*, 262, 553–555.
- Zhou, L.X., Hardy, J.R., and Cao, H.Z. (1997) Molecular dynamics simulation of superionicity in neighborite, NaMgF<sub>3</sub>. *Geophysical Research Letters*, 24, 747–750.

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