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3	Ultrasonic Studies of alkali-rich hydrous silicate glasses: elasticity, density, and
4	implications for water dissolution mechanisms
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Abstract

20 The acoustic velocities of three series of alkali-rich hydrous silicate glasses were 21 determined at ambient conditions using ultrasonic interferometry. The sound velocities and calculated elastic properties are nearly linear functions of dissolved water content. 22 The water content derivatives of both compressional $(\partial(V_P)/\partial(X_{H_2O}))$ and shear wave 23 velocity $(\partial(V_S)/\partial(X_{H_2O}))$ decrease with increasing Na₂O content, which suggests that 24 increasing sodium content might weaken the effect of water on acoustic velocities. For 25 26 each glass series, the shear modulus decreases with increasing water content, whereas the 27 adiabatic bulk modulus (K_s) varies little with water content. By comparing our results of K_S to previously published data on hydrous felsic alkalic (rhyolite, phonolite, and trachyte) 28 29 glasses, we speculate that H₂O may have a negligible effect on the compressibility of felsic alkalic glasses/melts. The measured densities of each series of glasses were 30 converted to molar volumes and extrapolated to the hypothetical pure water to derive the 31 partial molar volume of water (\bar{V}_{H_2O}) . The resulting \bar{V}_{H_2O} was found to be independent 32 of composition, with a value of 10.4 ± 0.5 cm³/mol at room temperature. Whereas the 33 partial molar bulk modulus (K_S) of water is composition-dependent. Furthermore, the 34 Poisson's ratio (v) increases with increasing water content for all three composition series. 35 Therefore, the similar behavior of Al-free and Al-bearing glasses implies that water 36 37 depolymerizes melts in both systems. The decrease of water content derivative of the Poisson's ratio $(\partial(v)/\partial(X_{H_2O}))$ with increasing Al/(Al + Si) (Q⁴(4Al) accordingly) 38

39	indicates that water interaction with Al-O bonds to produce Al-OH depolymerizes
40	peralkaline silicate melts less effectively than with Si-O bonds to form Si-OH. Hence, we
41	speculate that water prefers to interact with Si-O bonds rather than Al-O bonds in
42	peralkaline silicate systems. The dissolution mechanisms of water between peraluminous
43	and peralkaline melts might be different, which merits further research.
44	Keywords: ultrasonic interferometry, hydrous silicate glasses, alkali-rich, the partial
45	molar volume of water, elastic properties, water speciation; dissolution mechanisms

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Introduction

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49 The presence of water in magmas has long been considered an important control of igneous activities. Water is the most dominant volatile constituent of natural magmas and 50 plays a dramatic role in the physicochemical properties of silicate melts and glasses. For 51 instance, the density and the partial molar volume of water (\bar{V}_{H_20}) of hydrous silicate 52 melts have been studied *in-situ* using various methods, including the P-V-T cell in two 53 54 interconnected internally heated pressure vessels (Burnham and Davis 1971), X-ray 55 absorption (Malfait et al. 2014a, 2014b; Sakamaki et al. 2009; Seifert et al. 2013), and the sink-float method (Agee 2008; Jing and Karato, 2012; Matsukage et al. 2005; Sakamaki 56 57 et al. 2006). The investigations of hydrous supercooled liquids have provided further 58 constraints on the volumetric characteristics of water dissolved in silicate melts (Ochs and Lange 1997, 1999; Bouhifd et al. 2001, 2015). Furthermore, a review of available 59 density data of hydrous silicate glasses shows that \overline{V}_{H_2O} is independent of composition 60 61 by the intercept method (Richet et al. 2000). It was further investigated for a range of glasses under ambient conditions using the same method (Ardia et al. 2014; Balzer et al. 62 63 2020; Wu et al. 2017).

Due to the technical difficulties in measurements of melts at crustal or mantle conditions, several ambient condition investigations have focused on silicate glasses to deduce the structure and behavior of silicate melts. The pronounced effects of H₂O on

67 silicate glasses could provide some insights into the hydrous melts. Recently, Le Losq et 68 al. (2017) reported that strong hydrogen bonds (H-bonds) are present in Na-bearing 69 glasses below glass transition temperature (T_g) in an *in-situ* Raman spectroscopic study, which indicates glasses are inadequate to describe the structure of the corresponding 70 71 melts. However, the missing of strong H-bonds in the Ca-bearing glasses indicates that the formation of H-bonds in glasses depends on glass composition. Therefore, further 72 studies on the structure and properties of glasses are needed to understand the similarities 73 74 and differences between glasses and melts.

75 Understanding the high compressibility for hydrous melts can aid in predicting density increase with depth. Sound velocities are the prime data to provide direct 76 information on the elastic moduli and compressibility. Nevertheless, the effect of 77 dissolved water on the acoustic velocities and elastic behaviors of silicate glasses are still 78 poorly constrained. A few experiments on geological relevant compositions have been 79 conducted by Brillouin spectroscopy at ambient conditions. According to Richet and 80 81 Polian (1998), the partial molar bulk modulus (K_s) of water in andesite glasses (18 ± 3) GPa) is similar to that of Ice VII (23.7 ± 0.9 GPa) (Hemley et al. 1987), suggesting that 82 water behaves as the densest form of ice. The partial molar compressibility of water is 83 shown to be strongly dependent on composition over a range of hydrous silicate glasses 84 85 (Malfait et al. 2011; Whittington et al. 2012), implying \bar{V}_{H_2O} may be affected by composition at elevated pressures. Wu et al. (2017) showed that the effect of dissolved 86

water on the sound velocities of basalt glasses has tremendous discrepancies with previous studies. Recently, Gu et al. (2021) found that reductions in velocities due to hydration are larger in shear velocity than those in compressional velocity.

To better constrain the role of water on acoustic velocity and elasticity, the range of 90 glass compositions needs to be expanded. Many models of water dissolution in silicate 91 92 systems assume that water ruptures T-O-T links and behaves as a network modifier, much like alkali oxides such as Na₂O (Whittington et al. 2001). In addition, it was 93 94 demonstrated that alkali oxides and water have similar effects on volume (Richet and 95 Polian 1998; Richet et al. 2000) and viscosity of silicate glasses (Dingwell et al. 1996). However, relationships between the effects of water and alkali oxides have not been 96 97 revealed by elasticity measurements. Systematic investigations of the elastic behavior are 98 required to clarify the effect of alkali oxides and water on silicate glasses.

99 The main purpose of this study is to constrain the effect of water content on the sound velocities and elastic behaviors of three series of alkali-rich silicate glasses with up 100 101 to 5.25 wt% H₂O by ultrasonic interferometry at ambient conditions. The feasibility of using this method for the first successful determinations of acoustic velocities on hydrous 102 103 silicate glasses is presented. Another goal of this study is to investigate whether \overline{V}_{H_2O} in silicate glasses depends on composition at room temperature. The dependence of the 104 105 partial molar K_S and volume of water on compositions in silicate glasses and melts are compared and the causes for their contrasting behaviors are explained. Moreover, two 106

107	possible interpretations are proposed for linear functions of physical properties with
108	water content based on water speciation. Finally, the Poisson's ratio and its water content
109	derivative $(\partial(v)/\partial(X_{H_2O}))$ are used to infer the dissolution mechanisms of water.
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111	Experimental methods
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113 Synthesis of hydrous glasses

114 Three anhydrous alkali-rich multicomponent silicate glasses (Nelson9, Nelson15, 115 and NAS15) were selected as the starting materials, which were originally synthesized by Nelson and Carmichael (1979) and Riebling (1966) (please refer to the two papers for 116 117 details). We only give a brief description of synthesis procedures here. Reagent grade oxide and carbonate powders were weighed in the proportions required and then fused in 118 119 Pt crucibles at 1500-1700°C in a box furnace. After a few hours, the melts were poured on graphite plates and quenched to glasses. Repeated cycles of grinding and fusion were 120 implemented to ensure homogeneity of the glasses. The nominal compositions of three 121 anhydrous glasses are listed in Table 1. All the starting materials used in this work were 122 remelted to drive off the adsorbed water and ground into powder. 123

For each starting material, 3-5 hydrous specimens containing different amounts (0-5.45 wt%) of water were synthesized (See Behrens et al. 1996 for a detailed description). Briefly, the samples were hydrated by sealing powdered glasses and

distilled water in platinum capsules, which were welded shut and isostatically hot pressed under hyperliquidus conditions at 5 kbar and 1507-1673 K for a few hours in internally-heated gas pressure vessels. The power supply was turned off at the end of the experiments, and the glasses were quenched isobarically to prevent the formation of bubbles and water exsolution. The recovered run products used in this study were quenched to homogeneous glasses, free of crystals and bubbles, at least to the resolution of optical microscopy.

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135 Characterization of hydrous glasses

The water contents were determined by Karl Fischer Titration (KFT) analysis (Table 136 137 1), with the apparatus and procedures described by Behrens (1995) and Holtz et al. 138 (1992). The hydrous glasses were rapidly heated up in a Pt sample holder using a 139 high-frequency generator. The released H_2O was transported with a pre-dried argon 140 stream to the titration cell, where the water content was quantified based on the reaction 141 with coulometrically generated iodine and chemically bound sulfur dioxide already 142 present. To check the homogeneity of dissolved water distribution within silicate glasses, thin slices cut from both ends of several recovered sample cylinders were analyzed. The 143 144 uncertainty on the reported water content is typically smaller than 0.1 wt%. Although the 145 amounts of dissolved water were not measured after sound velocity experiments, there is 146 no reason to believe that water was lost during the measurements at ambient conditions.

The chemical compositions of samples were determined by electron microprobe 147 analyses (EMPA), and the reported values are averages of at least 10-point analysis 148 149 (Table 1). Morgan and London (1996, 2005) proposed that sodium (Na) migration is more evident in hydrous alkali aluminosilicate glasses than in anhydrous ones. To 150 151 minimize alkali mobilization, the EMPA was performed on a Jeol JXA-8230 with 15 kV 152 acceleration voltage, 5 nA beam current, and 30 µm diameter defocused beam. The densities (ρ) of samples were further characterized by the Archimedean method 153 154 at room temperature (298 K). Each glass was weighed 5-15 times with an electronic 155 analytical balance (METTLER TOLEDO XS204) in air and while immersed in toluene. The density is expressed as $\rho = \frac{m_1}{m_1 - m_2} \rho_0$, where ρ_0 is the known density of toluene; 156 157 m_1 and m_2 are the measured weights of investigated silicate glasses in air and toluene.

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159 Ultrasonic interferometry: transfer function method

The ultrasonic experiments were conducted using the ultrasonic interferometry system at Peking University. Ultrasonic interferometry combined with the transfer function method was used to determine the round-trip travel times of ultrasonic waves through the specimens. A detailed description of the transfer function technique can be found elsewhere (Li et al. 2002, 2004). The pure-mode longitudinal (Z-cut) and shear (41° X-cut) LiNbO₃ piezoelectric transducers were used to measure travel times ($t_{(P,S)}$), respectively. The longitudinal (V_P) and transverse sound velocity (V_S) can be calculated

with measured travel time $(t_{(P,S)})$ and specimen lengths (L) at ambient conditions: $V_{(P,S)} = 2L/t_{(P,S)}$.

The experimental setup for ultrasonic interferometry travel time measurements is 169 shown in Fig. 1. The piezoelectric transducer transmits and receives a broadband 170 171 radio-frequency (RF) pulse. The propagation time within the specimen is determined by the time delay between two consecutive echoes reflected from the front and back surfaces 172 173 of the specimen (Fig. 1). With sandpaper and diamond paste, the two surfaces along the 174 acoustic path of each glass were polished to be optically flat and parallel. The transfer function method of ultrasonic interferometry allows rapid collection of real-time acoustic 175 data (Fig. 2a) and subsequent analysis offline (Li and Liebermann 2007). After that, sine 176 177 wave tone bursts (30 MHz, 10 cycles for longitudinal wave and 20 MHz, 5 cycles for transverse wave) are used to convolute with the transfer function to simulate pulse-echo 178 overlap (PEO) measurement. Finally, the travel time is determined by shifting the second 179 180 copy relative to the initial echo pattern until the specimen echo and buffer rod echo overlap (Fig. 2b). 181

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183 Experimental precision and accuracy

Measurements were first carried out on a working standard (fused silica), whose sound velocity and density have been determined previously by other techniques, to check the accuracy of all facilities before measuring samples (Table 2). Our results are

187	consistent with available data (Andreatch and McSkimin 1976; Meister et al. 1980;
188	Manghnani et al. 2020) and the average error in the velocity and density measurements
189	are about 0.34% and 0.07%, respectively. In addition, the measurements of water-free
190	andesite glass (chemical composition presented in Table 1) compare well with data from
191	the literature with similar compositions (Table 2). The excellent agreement enables us to
192	provide a well-defined constraint on the role of H2O on the elastic properties of
193	alkali-rich silicate glasses. The reported precisions of measured values are obtained as the
194	mean standards deviation from repetitive measurements, while errors of derived values
195	are calculated by propagating experimental uncertainties (Table 3).
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197	Results
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199	The measured densities, lengths, travel times of alkali-rich hydrous silicate glasses
200	are reported in Table 3. Within experimental uncertainty, the density of each glass series
201	shows a linear dependence on the water content (Fig. 3). The decrease in glass density
202	with the addition of H_2O can be attributed to the low atomic mass of hydrogen, which
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	causes the H ₂ O component dissolved in glass to be much less dense than other oxide
204	causes the H_2O component dissolved in glass to be much less dense than other oxide components. Density is reduced by 0.3%, 0.43%, and 0.53% for Nelson9, Nelson15, and

NAS15 respectively when H₂O is increased by 1 wt%. In addition, the density decreases

with increasing SiO₂ content over the composition range investigated ($\rho_{Nelson9} < \rho_{Nelson15} < \rho_{NAS15}$, Fig. 3).

The acoustic velocities are plotted against water content in Fig. 4. For Nelson9 glasses, V_P and V_S decrease at the rate of 8.2 and 12.3 m/s per mol% dissolved H₂O, respectively. For Nelson15 glasses (0-9.89 mol% H₂O), V_P decreases from 5813 to 5781 m/s and V_S decreases from 3484 to 3423 m/s. For NAS15 glasses, the data suggest that V_S decreases slightly with increasing dissolved H₂O, whereas V_P is virtually independent of dissolved H₂O. An important feature of Fig. 4 is the lack of a systematic behavior of the variation of V_P and V_S with water content in various compositions.

The adiabatic bulk modulus (K_S) is derived from the measured density and sound 215 velocities: $K_s = \rho \cdot (V_P^2 - 4/3 V_S^2)$. According to the equation: $G = \rho \cdot V_S^2$, the shear 216 modulus (G) magnifies the variations of the shear velocity. As shown in Fig. 5, the shear 217 moduli for all three series decrease with increasing dissolved H_2O . However, dissolved 218 water has no discernible effect on the bulk modulus of alkali-rich silicate glasses. The 219 effect of chemical composition on K_S is more obvious than on G. For instance, the bulk 220 and shear modulus increase from 39 to 48 (23.1%) and 26 to 30 (15.4%) GPa for 221 222 anhydrous glasses, respectively.

The molar volume (*V*) was calculated from the gram formula weight (*g.f.w.*) and density: $V = g.f.w./\rho$ (Table 3). The dependence of the molar volume on dissolved H₂O is found to be linear with a low degree of scattering (Fig. 6a).

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Discussion

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229 The effect of water on wave velocities and elastic moduli

The different trend in wave velocities is likely attributed to the different Na₂O 230 231 content. For instance, it appears that the water content derivatives of the shear $(\partial(V_S)/\partial(X_{H_2O}))$ and compressional wave velocity $(\partial(V_P)/\partial(X_{H_2O}))$ decrease with 232 233 increasing Na₂O content (Fig. 4, Table 4). This implies that the increase in sodium 234 content might diminish the influence of water on acoustic velocities. For each composition, the effect of increasing water content on acoustic velocities was weaker for 235 236 V_P than for V_S (Fig. 4, Table 4). The larger decreases in V_S can be ascribed to depolymerization of the silicate network by hydroxyl groups, which reduces its shear 237 resistance (Deschamps et al. 2011). V_P is not affected as much because the 238 compressibility of the glass structure is not influenced as much by H₂O_{mol} in interstitial 239 240 voids or depolymerization by hydroxyl groups (Gu et al. 2021). The wave velocities of basalt glasses, on the other hand, suggest that dissolved water has a smaller effect on V_{s} 241 than on V_P (Malfait et al. 2011; Whittington et al. 2012; Wu et al. 2017). This contrasting 242 behavior might be due to water speciation in glasses, which is affected by chemical 243 composition (the degree of polymerization), cooling history (T_g) , and temperature 244 (Behrens 2020; Behrens and Yamashita 2008; Le Losq et al. 2015; Nowak and Behrens 245

246 1995; Romano et al. 1995).

247 The elastic moduli of basalt and andesite glasses systematically decrease as the 248 water content increases (Malfait et al. 2011; Richet and Polian 1998; Whittington et al. 2012; Wu et al. 2017). The adiabatic compressibility (β_s) was determined from the 249 relation: $\beta_s = 1/K_s$. Therefore, the addition of water increases compressibility. This 250 suggested that H₂O reacts with the bridging oxygen of tetrahedra in basalt and andesite 251 glasses, resulting in more compressible silicate structures. However, for each series of 252 253 glasses in this study, the bulk modulus is nearly constant with the addition of H₂O, which 254 is broadly consistent with the elastic behaviors of rhyolite (Malfait et al. 2011), phonolite, 255 and trachyte (Whittington et al. 2012) glasses (Fig. 5a). Likewise, the compressibility is 256 also nearly constant, coinciding with the viewpoint that the H₂O component becomes less compressible with increasing alkalinity (Whittington et al. 2012). The minimal effect of 257 dissolved H₂O on K_S observed in rhyolite glasses was interpreted that V_S decreases more 258 strongly than V_P , which counterbalances the concomitant decrease in density upon 259 260 hydration (Malfait et al. 2011).

According to a prior study, both K_s and G of silicate glasses slightly decrease with increasing SiO₂ concentrations (Malfait et al. 2011). Our observations report the decrease in K_s as silica content increases (Fig. 5a). Nonetheless, the fitted line from the shear moduli of the NAS15 glass series, which have the lowest SiO₂ content, is located in the middle (instead of the upper) range of the three series of glasses (Fig. 5b). This abnormal

behavior is most likely attributed to the high Na₂O content (up to 29 wt%) in NAS15 glasses. Malfait and Sanchez-valle (2013) demonstrated that bulk modulus is almost independent of Na₂O content, while shear modulus decreases systematically with increasing Na₂O content. Thus, the high Na₂O content in NAS15 glasses might cause the anomaly in shear moduli, but not in bulk moduli.

The number of non-bridging oxygens per tetrahedrally coordinated cation (NBO/T) 271 provides a measurement of the degree of polymerization, which is one of the most critical 272 273 characteristics for controlling magmatic processes in the Earth's interior. Whittington et 274 al. (2012) proposed that acoustic velocity and elastic modulus increase with decreasing 275 degrees of polymerization. This trend obtained from aluminosilicate glasses with a broad 276 range of polymerization does not apply to our samples. For instance, Fig. 4b and Fig. 5b 277 display the opposite trend that shear wave velocity and shear modulus decrease with decreasing polymerization (NBO/T value for anhydrous samples are shown in Table 4). 278 The reason for this may be due in part to the compositions of our samples falling outside 279 280 the composition range of previous studies, for example, Nelson15 and NAS15 are alkaline-earth-elements-free and Nelson9 is alumina-free. These behaviors appear to be 281 more complex than the previous conclusion. Thus, the trends of sound velocity and 282 elastic modulus on silicate glasses are not independently determined by the SiO₂ content 283 or the degree of polymerization, both or even more factors should be considered. For 284 instance, it is demonstrated that the temperature at which the glass transition takes place 285

 (T_g) affects the densities (hence the elastic properties) of glasses (Richet et al. 2000).

287	Notably, the elastic moduli of the felsic alkalic (phonolite, trachyte, and rhyolite)
288	glasses derived from Brillouin scattering spectroscopy (Malfait et al. 2011; Whittington et
289	al. 2012) show a similar trend to our results (Fig. 5), with all glasses sharing the common
290	feature of being alkali-rich. These felsic alkalic rocks are a relatively small component of
291	ocean island volcanic suites (Thompson et al. 2001). The role of water in the
292	physicochemical properties of those melts is critical for modeling the magmatic activities
293	at crustal depths (Ardia et al. 2014). If the investigated multicomponent compositions are
294	extrapolated to natural compositions, the observations of this study suggest that water has
295	a negligible effect on bulk modulus and compressibility in felsic alkalic glasses/melts.

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297 The room-temperature partial molar volume of water (\overline{V}_{H_20}) in glasses

By linearly extrapolating molar volume to the hypothetical pure water component, the room-temperature \overline{V}_{H_20} in glasses can be derived to be 10.0 ± 0.5 , 10.6 ± 0.3 , and $10.6 \pm 0.4 \text{ cm}^3/\text{mol}$ for Nelson9, Nelson15, and NAS15 glasses, respectively (Table 5). This extrapolation method to obtain \overline{V}_{H_20} is based on the assumption of ideal mixing between water and the silicate components (Richet et al. 2000). As shown in Fig. 6b, the linear dependence of molar volume on the dissolved water demonstrates that \overline{V}_{H_20} is independent of composition, which is consistent with the ideal mixing model.

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Table 5 displays \overline{V}_{H_2O} and synthesis P-T conditions in various studies. It appears

that \bar{V}_{H_2O} varies with the synthesis pressure. For example, the \bar{V}_{H_2O} values obtained at 2 306 GPa (~9 cm³/mol) are remarkedly lower than those obtained from 1 GPa runs (~10-11 307 cm³/mol) (Wu et al. 2017). In contrast, Ardia et al. (2014) showed that the \overline{V}_{H_2O} 308 increases to ~ 14 cm³/mol at the highest synthesis pressure (2.5 GPa). For the three 309 composition series in this study, \bar{V}_{H_2O} is almost constant at 10.4 ± 0.5 cm³/mol, which 310 311 might be attributed to the same synthesis pressure, eliminating the potential effect of 312 synthesis pressure. Due to the large extrapolation to 100 mol% H₂O, the data obtained in this study are considered to be broadly consistent with the constant value of 12.0 ± 0.5 313 cm^3/mol by Richet et al. (2000). 314

315 From density and thermal expansion measurements on hydrous silicate glasses, it has been reported that \bar{V}_{H_2O} is nearly independent of composition in supercooled liquids 316 near T_g (Ochs and Lange 1997, 1999; Bouhifd et al. 2001, 2015). Malfait et al. (2014a) 317 have combined literature data to conclude that \overline{V}_{H_2O} is independent of composition for 318 silicate melts at simultaneous high P-T conditions, at least within experimental 319 uncertainties. The constant \bar{V}_{H_2O} in glasses, supercooled liquids, and melts will simplify 320 321 the construction of a volumetric model for hydrous magmas under geologically relevant conditions. 322

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324 The partial molar bulk modulus of water in glasses and melts

325 The partial molar volume of water appears to be insensitive to the anhydrous

326 composition for melts as well as glasses, but this is not the case for the partial molar 327 elastic modulus of water. The partial molar bulk modulus of water for each glass series 328 derived by extrapolation does depend on the anhydrous composition (Table 4), which is 329 also reported by previous studies (Malfait et al. 2011; Whingttington et al. 2012). As a 330 density models that assume ideal mixing of the water component result. (composition-independent \bar{V}_{H_2O}) may result in larger uncertainties for the density of 331 hydrous melts at higher pressures. Under the assumption of an ideal mixing model, it is 332 333 now well established that the partial molar volumes of major oxide components are 334 independent of composition in anhydrous silicate melts (Lange and Carmchael 1987). 335 Schilling et al. (2001) suggested that the ideal mixing model is also applicable to elastic 336 properties and acoustic velocities of anhydrous glasses in the anorthite-diopside-forsterite 337 (An-Di-Fo) system. However, they cautioned that extrapolating molar elastic moduli to drastically different compositions can lead to significant errors due to the different 338 339 structures. The reasoning behind this was further illustrated by the fact that the measured elastic moduli for silica glass ($K_S = 36.5$ GPa, Bass 1995) are markedly different from the 340 inferred molar K_S (47.4 GPa) of a virtual SiO₂ component in the An-Di-Fo glasses 341 (Schilling et al. 2001). We speculate that the ideal mixing model may not apply to the 342 bulk modulus of hydrous glasses. Therefore, it seems that the partial molar K_S of water 343 depends on the anhydrous composition of the glasses, which is in contrast to \bar{V}_{H_2O} . 344

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Nevertheless, the partial molar K_S of water for andesite and rhyolite melts (2.5 and

346	2.7 GPa, respectively) are independent of silicate composition at 1273 K (Malfait et al.
347	2014). This suggests that the behavior of water in melts and glasses may differ
348	significantly. This discrepancy is likely related to configurational contributions (unique to
349	melts) to the physical properties of melts (Richet and Neuville, 1992). In particular,
350	Stebbins et al. (1984) indicated that the ideal mixing model can be used to reproduce the
351	heat capacities of silicate glasses within error, but is less appropriate for that of silicate
352	melts due to the development of new structural configurations at temperatures above T_g .
353	Moreover, the effect of water speciation with temperature cannot be excluded. For
354	instance, the concentration of OH groups in the melt varies strongly with temperature and
355	is far greater than predicted by spectroscopic studies on glasses at ambient conditions
356	(Nowak and Behrens 1995). Consequently, there is a fundamental difference in the effect
357	of water on silicate melts versus glasses, systematic investigations are necessary to
358	understand the volumetric and elastic properties both in glasses and melts in the future.

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360 Water speciation

The role of water in silicate glasses is directly linked to its speciation. It has long been recognized that water dissolves in silicate glasses as hydroxyl groups (OH) and molecular water (H_2O_{mol}) (Newman et al. 1986; Stolper 1982). As a traditional view, hydroxyl groups predominate in silicate glasses at lower water content, while the fraction of H_2O_{mol} increases with water content (Schmidt et al. 2001; Stolper 1982). The strongly

nonlinear role of water on the melt viscosity has been interpreted as the structural 366 367 consequences of the decreasing proportion of water dissolved as OH groups within the 368 melt with increasing water content (Stolper 1982). It appears that H₂O_{mol} does not affect melt viscosity, and the amount of OH groups in the melt structure solely determines the 369 370 viscosity (Dingwell et al. 1996). Additionally, Richet and Polian (1998) found that the 371 variations of the shear modulus of hydrous andesite glasses are discontinuous with increasing water content and ascribed it to rapidly increasing H₂O_{mol} concentration when 372 373 the total water content is above 1 wt%. 374 However, the molar volume, sound velocities, and elastic properties of the hydrous 375 silicate glasses from this study are all linear functions of water content. The same trends 376 were observed in a variety of glasses in previous studies (Malfait et al. 2011; Whittington 377 et al. 2012; Wu et al. 2017). Because we did not carry out spectroscopic measurements, two possible interpretations can be proposed for the observations: (1) Water affects these 378 379 properties in the same manner regardless of its speciation because the linear trend was 380 observed in ranges where either OH groups or H_2O_{mol} predominate. The hydroxyl groups dominate at low water contents before being overtaken by H₂O_{mol}, which occurs in 381 quenched glasses at water contents of 3-4 wt% (Stolper 1982). Therefore, the speciation 382 change of water has little effect on the volume or K_S , implying that the volume change of 383 the reaction between bridging oxygen and H₂O_{mol} to form OH is small in all glass series. 384 It further indicates that the changes in water speciation are driven by the entropy rather 385

than the volume of the system (Richet and Polian 1998). Moreover, Whittington et al. 386 387 (2012) proposed that the effects of water on compression mechanisms are determined by 388 water-silicate interactions, but not the interconversion of water speciation. (2) Like 389 viscosity, the molar volume and elastic properties are sensitive to water speciation. Recently, Cody et al. (2020) used ¹H MAS nuclear magnetic resonance (NMR) 390 391 spectroscopy to measure water speciation in aluminosilicate glasses. They suggested that H₂O_{mol} species are predominant over OH initially, while OH becomes dominant at very 392 393 high total water content (>20 mol%). This recent finding is dramatically different from 394 the traditional view of the water speciation trend. Thus, based on the new speciation trend, 395 the linear correlation of the molar volume, wave velocities, and elastic properties with water content can be attributed to the relatively constant ratio of OH to H₂O_{mol} with no 396 397 interconversion between the two water species in the range of water content (up to 5.45 398 wt% H₂O) in our glasses. 399 400

Implications for the dissolution mechanisms of water

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The OH and H_2O_{mol} species behave in distinct ways within the silicate network of glasses and melts (Behrens et al. 2007). Consequently, not only the quantification of water speciation but also the understanding of the water dissolution mechanisms in the

silicate structure is important. The composition of the silicate system affects the water 406 dissolution mechanism. Water prefers to be present as H2Omol that interacts with the 407 408 bridging oxygens to produce OH linked to tetrahedral cation (T-OH) in highly polymerized melts and free hydroxyl groups (M-OH, OH species bonded only to 409 410 network-modifying cations) may also be present in strongly depolymerized melts (Cody et al. 2005; Le Losq et al. 2015; Moretti et al. 2014; Mysen and Cody 2005; Xue and 411 Kanzaki 2004). In addition, NMR studies for silicate glasses demonstrated the dissolution 412 413 mechanisms of water in Al-bearing silicate networks are essentially similar to those in Al-free systems, and multiple dissolution mechanisms for water are present 414 415 simultaneously (Malfait and Xue 2010; Xue and Kanzaki 2008).

416 Indirect evidence, such as the reduction in viscosity (Friedman et al. 1963; Giordano et al. 2004; Hess and Dingwell 1996; Richet et al. 1996) and T_g (Giordano et al. 2005) as 417 water is incorporated into the melts, may reflect the structural change in the silicate 418 network. In particular, water is much more effective at reducing the viscosity of 419 polymerized melts than depolymerized ones (Whittington et al. 2000), which indicates 420 that the water dissolution mechanism may depend on composition. The Poisson's ratio (ν) 421 is a direct proxy for melt structure that correlates with network connectivity and is 422 expressed as $v = 1/2 \cdot [(V_P/V_S)^2 - 2]/[(V_P/V_S)^2 - 1]$. According to experimental 423 424 evidence, ν increases with rising atomic packing density (C_{σ}) and decreases with increasing average coordination number $\langle n \rangle$ or the number of bridging oxygen $n_{\rm BO}$ 425

426	(Greaves et al. 2011; Rouxel 2007). Hence, the increase in Poisson's ratio could be used
427	to qualitatively predict the network depolymerization. As might be expected, within
428	experimental errors, the Poisson's ratio for each glass series is a linear function of water
429	content (Fig. 7a). The increase of ν with the addition of H ₂ O together with the NMR
430	spectroscopic evidence for the formation of T-OH ($T = Si$, Al) through the rupture of
431	oxygen bridges: T-O-T + H_2O = 2T-OH (Farnan et al. 1987; Malfait and Xue 2010; Xue
432	2009; Xue and Kanzaki 2006), supports the hypothesis that dissolved water
433	depolymerizes melts in both Al-free (Nelson9) and Al-bearing (Nelson15 and NAS15)
434	silicate melts.

The Al/(Al + Si) ratio is one of the most important compositional variables of 435 natural magmatic systems. It has been proposed that the distribution of Al³⁺ among the 436 structural unit $(Q^n, n = number of bridging oxygens to other tetrahedral)$ depends on the 437 Al/(Al + Si) of hydrous aluminosilicate systems (Mysen 2007). The Al³⁺ preferentially 438 occupies Q⁴-species and the abundance of Q⁴(4Al, i.e. the number of next-nearest Al to Si 439 is 4) increases rapidly with Al/(Al + Si), which plays a dominant role in raising the 440 fraction of Q^4 -species with increasing bulk Al/(Al + Si) in peralkaline glasses and melts 441 (Mysen et al. 2003). In addition, dissolved water appears to exhibit preferences for 442 interaction with bridging oxygen bonds in Q⁴-species (Mysen and Cody 2005; Zotov and 443 Keppler 1998). 444

445

The Poisson's ratio increases with H₂O at a higher rate in the Al-free Nelson9 series

than in the other two Al-bearing systems (Fig. 7a, Table 4). Based on this behavior, we 446 speculate that the water concentration derivative of Poisson's ratio $(\partial(v)/\partial(X_{H_2O}))$ are 447 possibly sensitive to the Al/(Al + Si) ratio. Fig. 7b shows that the $\partial(v)/\partial(X_{H_2O})$ is 448 negatively correlated with Al/(Al + Si) when our data are plotted with previous data for 449 rhyolite, phonolite, and trachyte glasses. The observation that $\partial(v)/\partial(X_{H_2O})$ decreases 450 with increasing Al/(Al + Si) ratio may be explained by assuming that water prefers to 451 interact with Si-O bonds to form Si-OH rather than Al-O bonds to produce Al-OH in 452 peralkaline systems. This is a reasonable assumption because the rate of v-increase with 453 water content $(\partial(v)/\partial(X_{H_2O}))$, which represents the rate of depolymerization, decreases 454 with increasing Al/(Al + Si) (therefore $Q^4(4Al)$ abundance). This implies that water 455 interaction with Al-O bonds depolymerizes peralkaline silicate melts less effectively than 456 with Si-O bonds. A recent NMR investigation revealed that the Al-OH/(Al-OH + Si-OH) 457 proportion is smaller than the Al/(Al + Si) ratio in rhyolite and phonolite melts at lower 458 459 temperatures (Malfait and Xue 2014). This also supports the idea that water prefers to form Si-OH rather than Al-OH in peralkaline silicate melts. Nevertheless, from the 460 increased water solubility with Al/(Al + Si) observed by Holtz et al. (2000), Mysen and 461 Richet (2019) hypothesized that water prefers to generate Al-OH rather than Si-OH for 462 alkali meta-aluminosilicates. Moreover, the dependence of water solubility on 463 composition has also been observed, particularly for peralkaline and peraluminous 464 silicate composition (Behrens and Jantos 2001; Mysen and Armstrong 2002). Hence, the 465

dissolution mechanisms of water between peraluminous and peralkaline magma systems

467 might be different.

In summary, the decrease of water content derivative of the Poisson's ratio 468 $(\partial(v)/\partial(X_{H_2O}))$ with increasing Al/(Al + Si) (Q⁴(4Al) accordingly) indicates that water 469 interaction with Al-O bonds depolymerizes peralkaline silicate melts less effectively than 470 471 with Si-O bonds. Hence, we speculate that water prefers to interact with Si-O bonds rather than Al-O bonds in peralkaline silicate systems. Further spectroscopic analyses are 472 473 required to obtain more detailed information on the speciation and dissolution 474 mechanisms of water. This study on the density, elastic properties of alkali-rich hydrous silicate glasses could be the beginning of more comprehensive investigations on other 475 476 physical properties. Additional constraints at elevated P-T conditions are needed to 477 explore the effect of pressure and temperature on the structure and properties of hydrous silicate melts/glasses, as well as to better understand the dynamic behavior of natural 478 479 magmas at depth.

480

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Figure captions:

731	Figure 1. A schematic diagram shows the principle of travel time measurements
732	using the ultrasonic interferometric technique (modified after Li and Liebermann 2014).
733	Figure 2. (a) Received signals for andesite glass on the buffer rod; (b) Simulation of
734	the pulse-echo overlap (PEO) method.
735	Figure 3. The density of the three series of hydrous silicate glasses against water
736	content. The dashed lines are linear fits to our data.
737	Figure 4. The variation of (a) compressional and (b) shear wave velocities of the
738	three series of hydrous silicate glasses against water content. The dashed lines are linear
739	fits to our data.
740	Figure 5. (a) Adiabatic bulk and (b) shear modulus of silicate glasses against water
741	content. The data from this study are shown as solid symbols, the data of rhyolite glass
742	from Malfait et al. (2011) are shown as open circles, the data of phonolite and trachyte
743	from Whittington et al. (2012) are shown as open squares and triangles, respectively. The
744	dashed lines are linear fits to our data.
745	Figure 6. (a) Room-temperature molar volume of hydrous silicate glasses against
746	water content. The dashed lines are linear fits to our data. (b) The fitted dashed lines yield
747	the partial molar volume of water (\bar{V}_{H_2O}) as they intercept with the y-axis at 100% H ₂ O.

- The constant \bar{V}_{H_2O} (12 \pm 0.5 cm³/mol, Richet et al. 2000) is shown for a comparison.
- 749 Errors of measurements are smaller than the symbol size.

750	Figure 7. (a) Variation of the Poisson's ratio with water content; (b) The relationship
751	between the water content derivative of Poisson's ratio $(\partial(v)/\partial(X_{H_2O}))$ and Al/(Al + Si)
752	ratio; data taken from M11 (Malfait et al. 2011) and W12 (Whittington et al. 2012) are
753	also plotted as open symbols. The dashed line(s) in (a) are linear fits to our data and in (b)
754	is a linear fit to all the data.
755	

Sample	SiO ₂	Al_2O_3	MgO	CaO	Na ₂ O	K ₂ O	H_2O	Total	<i>g.f.w.</i> [‡]
Nelson9*	72.74		8.2			18.42		99.36	61.730
Nelson9-1	73.72		8.09			18.33	0.00^{\dagger}	100.14	61.731
Nelson9-2	71.24		7.92			18.26	2.02 (4)	99.44	58.904
Nelson9-3	71.15		7.69			18.18	2.92 (5)	99.94	57.767
Nelson9-4	70.53		7.53			17.72	4.05 (4)	99.83	56.287
Nelson9-6	69.30		7.79			17.35	5.45 (6)	99.89	54.445
Nelson15 [*]	60.36	21.47			11.25	6.44		99.52	67.936
Nelson15-1	61.58	21.21			10.17	6.70	0.00^{\dagger}	99.66	67.882
Nelson15-2	60.98	21.06			10.14	6.36	1.12 (4)	99.65	65.800
Nelson15-3	60.53	20.84			9.94	6.21	2.22 (4)	99.74	63.880
Nelson15-4	59.83	20.52			9.75	6.73	2.82 (8)	99.66	62.978
NAS15 [*]	46.89	23.06			29.88			99.83	67.060
NAS15-1	46.22	24.01			28.81		0.00^{\dagger}	99.05	67.394
NAS15-2	47.55	23.58			26.09		2.57 (6)	99.79	62.908
NAS15-3	46.34	23.21			26.44		3.39 (10)	99.38	61.586
Andesite	59.47	20.91	5.24	9.69	2.92	1.72	0.00^{\dagger}	99.96	63.954

Table 1. Compositions (wt %) of the silicate glasses

⁷⁵⁷ † Nominal water content of dry samples, other H₂O contents were determined by KFT

758 analysis.

of each oxide component, and M_i is the molecular weight of each oxide component

761 * Nominal composition of the anhydrous starting materials

762

ρ (g/cm ³)	V_P (m/s)	V_{S} (m/s)	$V_P \operatorname{Diff.}^{\$}$	V_S Diff.% [§]	Reference
2.203 (5)	5955 (13)	3746 (5)			This study *
2.202	5933	3764	-0.37	0.47	Andreatch and McSkimin (1976) ^{\$}
2.204	5955 (5)	3763 (5)	0.00	0.44	Meister et al. $(1980)^{\dagger}$
2.2007	5947 (6)	3769 (4)	-0.14	0.60	Manghnani et al. (2020) ^{\$}
			Andesite	e	
2.589 (13)	6344 (4)	3696 (5)			This study *
2.571	6152 (5)	3616 (5)	-3.03	-2.16	Meister et al. (1980) [†]
2.513	6240	3700	-1.64	0.11	Richet and Polian (1998) [‡]
2.536	6213 (10)	3689 (12)	-2.07	-0.19	Malfait et al. $(2011)^{\ddagger}$

Table 2. Densities and sound velocities of fused silica and anhydrous andesite glasses

764 Sound velocities were measured using different methods: * ultrasonic interferometry;

^{\$} pulse superposition; [†] ultrasonic phase comparison; [‡] Brillouin scattering.

766 [§] $V_{(P, S)}$ Diff.% corresponds to $(V_{(P, S)} \text{ (literature data)} - V_{(P, S)} \text{ (this study)}) \times 100/ V_{(P, S)}$

767 (this study).

Sample	ρ (g/cm ³)	H ₂ O (mol	molar V (cm ³ /mo	L (mm)	<i>t</i> _P (μs)	t _s (µs)	<i>V_P</i> (m/s)	<i>V</i> _S (m/s)	K _S (GPa)	$\frac{\beta_S}{(10^{-2}\text{GPa}^{-1})}$	G (GPa)	v
Nelson9-1	2.417 (13)	0.00	25.54 (4)	3.791 (8)	1.3694 (3)	2.2940 (7)	5537 (12)	3305 (7)	38.9 (2)	2.57 (2)	26.4 (2)	0.223 (4)
Nelson9-2	2.409 (10)	6.64	24.45 (5)	6.622 (4)	2.4042 (5)	4.0788 (12)	5509 (4)	3247 (2)	39.2 (2)	2.55 (1)	25.4 (1)	0.234 (1)
Nelson9-3	2.396 (9)	9.37	24.11 (3)	4.018 (6)	1.4684 (10)	2.5114 (8)	5473 (9)	3200 (5)	39.1 (2)	2.56(1)	24.5 (1)	0.240 (3)
Nelson9-4	2.380 (3)	12.68	23.65 (1)	5.128 (8)	1.8814 (5)	3.2348 (13)	5451 (9)	3171 (5)	38.8 (1)	2.58 (1)	23.9 (1)	0.244 (3)
Nelson9-6	2.377 (7)	16.58	22.9 (2)	2.219 (4)	0.8218 (10)	1.4310 (11)	5400 (12)	3101 (6)	38.8 (2)	2.58 (1)	22.9 (1)	0.254 (4)
Nelson15-1	2.459 (4)	0.00	27.61 (1)	3.33 (2)	1.1457 (9)	1.9117 (12)	5813 (6)	3484 (3)	43.3 (1)	2.31 (1)	29.8 (1)	0.220 (2)
Nelson15-2	2.448 (7)	4.09	26.88 (2)	1.086 (2)	0.3758 (11)	0.6302 (15)	5806 (20)	3462 (10)	43.4 (3)	2.30 (1)	29.4 (2)	0.224 (7)
Nelson15-3	2.436 (6)	7.89	26.22 (2)	3.515 (6)	1.2120 (7)	2.0417 (7)	5800 (10)	3443 (6)	43.4 (2)	2.30 (1)	28.9 (1)	0.228 (4)
Nelson15-4	2.429 (12)	9.89	25.92 (4)	3.245 (2)	1.1226 (10)	1.8958 (7)	5781 (6)	3423 (2)	43.2 (2)	2.31 (1)	28.5 (2)	0.230 (2)
NAS15-1	2.546 (12)	0.00	26.47 (4)	2.505 (12)	0.8644 (30)	1.5002 (28)	5796 (34)	3340 (17)	47.7 (4)	2.10 (2)	28.4 (3)	0.252 (12)
NAS15-2	2.516 (7)	8.99	25.00 (2)	3.996 (4)	1.375 (15)	2.4078 (15)	5812 (9)	3319 (4)	48.0 (1)	2.08 (1)	27.7 (1)	0.258 (3)
NAS15-3	2.500 (12)	11.66	24.63 (4)	1.784 (2)	0.6178 (18)	1.0912 (1)	5775 (18)	3270 (4)	47.8 (2)	2.10(1)	26.7 (1)	0.264 (6)

768 **Table 3.** Densities, acoustic velocities, and elastic properties of the investigated silicate glasses

769	Table 4.	The water	content	derivatives	of wave	velocities	and Poiss	on's ratio,	the partial
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Sample	NBO/T*	$\partial(V_P)/\partial(X_{H_20})$	$\partial(V_S)/\partial(X_{H_2O})$	$\overline{K}_{S_{H_2O}}$	$\partial(v)/\partial(X_{H_2O}) \times 10^{-3}$
Nelson9	0.64	-8.2 (11)	-12.2 (10)	38.9 (3)	1.83 (1)
Nelson15	0.04	-2.8 (10)	-5.8 (6)	43.3 (1)	1.03 (9)
NAS15	0.37	-0.9 (29)	-5.0 (30)	47.7 (15)	0.99 (29)

770 molar bulk modulus of water, and NBO/T of the investigated glass series

* Calculated by NBO/T = (Na + K + 2Mg - Al)/(Si + Al) for the anhydrous compositions.

Table 5. The partial molar volume of water and synthesis conditions of hydrous silicate

774 glasses

Sample	H ₂ O (wt%)	Pressure (GPa)	Temperature (°C)	$\overline{V}_{H_2O}(\mathrm{cm}^3/\mathrm{mol})$ #
This study				
Nelson9	0-5.45	0.5	1234-1240	10.0 ± 0.5
Nelson15	0-2.82	0.5	1240-1250	10.6 ± 0.3
NAS15	0-3.39	0.5	1250	10.6 ± 0.4
Ardia et al. (2014)				
Haplogranite	0.15-5.24	0.5	1000	11.23 ± 1.00
Haplogranite	0.15-5.24	1.0	1000	12.17 ± 1.00
Haplogranite	0.15-5.24	1.5	1000	12.82 ± 1.00
Haplogranite	0.15-5.24	2.5	1000	13.74 ± 1.00
Balzer et al. (2020)				
NAS	0.02-7.90	0.5	1250-1600	11.8 ± 0.1
NKAS(2:1)	0.02-8.00	0.5	1250-1600	12.0 ± 0.1
NKAS(1:1)	0.01-7.96	0.5	1250-1600	12.1 ± 0.2
Bouhifd et al. (2001, 2015)				
Albite [#]	0-5.2	0.2	1200	$10.2 \pm 0.5^{*}$
Tephrite [#]	0-3	0.2-0.3	1300	$10.4 \pm 0.5^{*}$
Trachyte [#]	0-5	0.2-0.3	1300	$11.2 \pm 0.8^{*}$
Foidite [#]	0.7-2.3	0.3	1300	$10.3 \pm 1.0^{*}$
Phonolite	0-4.72	0.2-0.3	1200-1300	$11.0 \pm 0.5^{*}$
Malfait et al. (2011)				
Andesite	0-3.01	0.7	1400	13.9 ± 0.4
Basalt	0.05-4.40	0-0.20	1170-1290	12.0 ± 0.3
Wu et al. (2017)				
Basalt (FX) ^{\$}	0-4.01	1.0	1400	$11.4 \pm 0.5 \ (11.6 \pm 0.8^*)$
Basalt (FW) ^{\$}	0-5.23	1.0	1400	$11.2 \pm 0.4 \ (10.9 \pm 0.6^*)$
Basalt (SHT) ^{\$}	0-3.45	1.0	1400	$10.3 \pm 0.6 (11.5 \pm 0.5^*)$
Basalt (FX)	0-5.48	2.0	1400	9.2 ± 1.2
Basalt (FW)	0-6.82	2.0	1400	9.3 ± 1.0
Basalt (SHT)	0-4.48	2.0	1400	8.9 ± 1.5

775 # \overline{V}_{H_2O} is derived by intercept method described in the text. This table only summarize the data on

hydrous silicate glasses collected after Richet et al. (2000).

777 ^{*} \bar{V}_{H_2O} of relaxed glasses after relaxation during dilatometry measurements up to T_g

- ^{\$} FX(Feixian), FW(Fuxin), and SHT(Sihetun) are three kinds of basaltic glasses from the North China
- 779 Craton (NCC).
- 780 **Figure 1.**
- 781



784 Figure 2.



788 **Figure 3.**



790 **Figure 4.**



792 Figure 5.



793



795 **Figure 6.**



797 Figure 7.



799

798