1	Revision 1
2 3	Calorimetric study of skutterudite (CoAs $_{2.92}$ ) and heazlewoodite (Ni $_3S_2$ )
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7 9 10 11 12 13 14 15 16 17 18 19	<ul> <li><sup>1</sup> Institute of Geosciences, Friedrich-Schiller University, Burgweg 11, 07749 Jena, Germany *corresponding author, email: Juraj.Majzlan@uni-jena.de; ORCID: 0000-0002-9151-4394</li> <li><sup>2</sup> School of Molecular Sciences and Center for Materials of the Universe, Arizona State University, Tempe, Arizona 85287, USA</li> <li><sup>3</sup> Czech Geological Survey, Geologická 6, 152 00 Prague 5, Czech Republic</li> <li><sup>4</sup> Institute of Geochemistry, Mineralogy and Mineral Resources, Faculty of Science, Charles University, Albertov 6, 128 00 Prague</li> <li><sup>5</sup> Institute of Experimental Mineralogy (IEM RAS), 142432 Chernogolovka, Moscow Region, Russia</li> <li><sup>6</sup> Ural Federal University, Ekaterinburg 620002, Russia</li> <li><sup>7</sup> Dubna State University, Dubna 141982 Russia</li> <li><sup>8</sup> Department of Chemistry and Physics of Materials, University of Salzburg, Jakob-Haringer-Strasse 2a, 5020 Salzburg, Austria</li> </ul>
20	ABSTRACT – Nickel and cobalt arsenides, sulfarsenides, and sulfides occur in many
21	hydrothermal ore deposits but the thermodynamic properties of these phases are not well known,
22	in some cases not known at all. In this work, we determined a full set of thermodynamic
23	properties for heazlewoodite and skutterudite. Both phases were synthesized in evacuated silica
24	tubes at elevated temperatures and electron microprobe analyses gave their composition as $Ni_3S_2$
25	and CoAs2.92, respectively. Enthalpies of formation were measured by high-temperature oxide-
26	melt solution calorimetry. The reference phases were pure elements, thus eliminating any
27	systematic errors related to such phases. The enthalpies of formation at $T = 298.15$ K and $P = 10^5$
28	Pa are $-216.0\pm8.4(2\sigma)$ and $-88.2\pm6.1$ kJ·mol <sup>-1</sup> for Ni <sub>3</sub> S <sub>2</sub> and CoAs <sub>2.92</sub> , respectively. Entropies
29	were calculated from low-temperature heat capacity ( $C_p$ ) data from relaxation (PPMS)
30	calorimetry and are 133.8 $\pm$ 1.6 and 106.4 $\pm$ 1.3 J·mol <sup>-1</sup> ·K <sup>-1</sup> , respectively. The calculated Gibbs
31	free energies of formation are $-210.0\pm8.4$ and $-79.9\pm6.2$ kJ·mol <sup>-1</sup> for Ni <sub>3</sub> S <sub>2</sub> and CoAs <sub>2.92</sub> ,
32	respectively. The PPMS $C_p$ data, together with a set of differential scanning calorimetry

- measurements, were used to derive  $C_p$  polynomials up to 700 K with the Kieffer model based on previously published frequencies of acoustic and optic modes. Equilibrium constants for selected
- reactions with an aqueous phase were calculated up to 700 K. Geochemical modeling in these
- 36 systems, however, should await until more reliable data for other phases from the system Co-Ni-
- 37 As-S are available.
- **Keywords**: heazlewoodite; skutterudite; enthalpy; entropy; geochemical modeling
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- 40

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#### INTRODUCTION

42 Nickel and cobalt arsenides, sulfarsenides, and sulfides occur in many hydrothermal ore deposits (Dolansky 2007, Belkin et al. 2008, Ahmed et al. 2009, Gervilla et al. 2012, Kreissl et al. 2018, 43 Scharrer et al. 2019, Tourneur et al. 2021, Horn et al. 2021). They are found especially in 44 orthomagmatic Cu-Ni-PGE ores (Naldrett 2004), hydrothermal ores, commonly of the "five-45 element association" (Burisch et al. 2017), or in stratiform Cu-Co ores (Dewaele et al. 2006). 46 47 They represent a large and complex group of minerals, with sulfides, disulfides, mono-, di- and triarsenides, and sulfarsenides (a list of the minerals with further references can be found in 48 Table 1 in Hem 2006). In addition to Co, Ni, and As, many of these minerals show also the As-49 50 Sb substitution, for example in the gersdorffite-ullmanite (NiAsS-NiSbS) series (Stevko a Sejkora 2020). Occasionally, minerals like gersdorffite may contain economically interesting 51 52 concentrations of gold or platinum-group elements (e.g., Pašava et al. 2013, Cabri et al. 2017). 53 The presence of reduced arsenic is likely to drive gold incorporation in the sulfarsenide minerals (Pokrovski et al. 2021). Arsenides also commonly form in serpentinization processes at high T-P 54 in subduction zones and may be important carriers of As and PGE in such settings (Gonzáles-55 Jiménez et al. 2021). Some of the compositions show polymorphism, for example, NiAs<sub>2</sub> is 56 57 known as the minerals rammelsbergite, pararammelsbergite, and krutovite. Structural variations are encountered even within one mineral species; gersdorffite, for example, can have an ordered 58 structure (space group  $P2_{13}$ ) or two disordered structures (space groups  $Pa_{3}$  and  $P_{1}$ ) (Bayliss 59 and Stephenson 1967, 1968, Steger et al. 1974, Bayliss, 1982). 60 61 Experimental work on dry systems (Klemm, 1965; Maurel and Picot, 1974, Hem and Makovicky 2004a, b, see also Table 2 in Hem 2006) showed that complete solid solutions can be 62

expected at temperatures above 500 °C but such high temperatures were rarely encountered 63 during the formation of the ores with Co-Ni arsenides and sulfarsenides (Scharrer et al. 2019). 64 Thermodynamic properties of the Co-Ni-As-S phases are mostly not well constrained or 65 unknown. For the obvious reason of their abundance, more attention was paid to the Fe-dominant 66 67 members of arsenides and sulfides. For example, thermodynamic properties of löllingite (FeAs<sub>2</sub>) were determined early by Barton (1969) from high-temperature equilibria and extrapolation to 68 room temperature. A calorimetric measurement of heat capacity ( $C_{\rm p}$ ) of a synthetic löllingite 69 70 sample and the resulting entropy by Panishkin et al. (1991) deviated much from the datum in Barton (1969). In their compilation, Robie and Hemingway (1995) adopted the S° from 71 Panishkin et al. (1991). They also selected  $\Delta_f G^o$  which is numerically equal to  $\Delta_f H^o$  in Barton 72 (1969) and adjusted their  $\Delta_{\rm f} {\rm H}^{\rm o}$  accordingly. The reasons behind this adjustment are not clear. 73 Bessinger and Apps (2005) modified the  $\Delta_f G^\circ$  from Barton (1969) slightly and took S<sup>o</sup> from 74 Barton (1969), even though they referred to Pokrovski et al. (1996) as their data source. Perfetti 75 et al. (2008) accepted S<sup>o</sup> from Panishkin et al. (1991) and recalculated  $\Delta_f H^o$  and  $\Delta_f G^o$ , based on 76 their new solubility data and combustion calorimetry data of Stolyarova (1977). When reviewing 77 the available source of thermodynamic data for löllingite, one notes differences of 40 kJ·mol<sup>-1</sup> in 78  $\Delta_{\rm f}$ H° and 40 J·mol<sup>-1</sup>·K<sup>-1</sup> in S° (see also Perfetti et al. 2008). Such differences are larger than 79 80 expected experimental uncertainties in thermodynamic measurements. This example shows the difficulties and gaps in our understanding of these minerals, even though the Fe-As-S system 81 received much more attention than the systems Co-As-S or Ni-As-S. 82 83 In this work, we have derived thermodynamic properties of skutterudite (nominally CoAs<sub>3</sub>) and heazlewoodite (nominally Ni<sub>3</sub>S<sub>2</sub>). Enthalpies of formation were measured by high-84 temperature oxide-melt solution calorimetry. The reference phases were pure elements, thus 85

eliminating potential systematic errors from reference chemical compounds. This work was also 86 meant as another test of this type of calorimetry on samples with highly volatile components (As, 87 S). Flushing with oxygen maintains oxidative conditions and all elements are converted rapidly 88 89 and reproducibly to their highest oxidation state in the solvent (Majzlan 2017). The entropies were calculated from low-temperature heat capacities measured by relaxation calorimetry. 90 Differential scanning calorimetry was used to measure  $C_{\rm p}$  above the ambient temperature and the 91  $C_{\rm p}$  data were extrapolated to higher temperatures using the appropriate treatment (Kieffer 1985). 92 The data were then used to derive log K's for selected chemical reactions. These equilibrium 93 constants can be used for geochemical modeling involving systems with Co-Ni sulfides and 94 95 arsenides.

96

### MATERIALS

97 The samples used in this study were prepared in evacuated silica tubes. The protocols
98 described below are the results of many trials and errors, optimized to yield the desired products.
99 Many other syntheses were carried out but did not run to completion, were contaminated by
100 other minor sulfides/arsenides or oxygen.

A synthetic analog of heazlewoodite (Ni<sub>3</sub>S<sub>2</sub>) was prepared from a stoichiometric (3:2) mixture of pure elements (Ni 99.999 %, S 99.999 %) that were placed into a silica-glass tube. The tube with the charge was evacuated and sealed. The ampoule was placed into a preheated horizontal furnace, and the sample was tempered (200 °C for 32 days; 300 °C for 19 days). Afterwards, the sample was quenched in a cold-water bath, homogenized in an agate mortar, evacuated, sealed again, and heated (300 °C for two months). Finally, the sample was quenched, and the content purity and composition were checked using powder X-ray diffraction.

108	The synthetic analog of skutterudite was synthesized in evacuated silica-glass ampoules from
109	pure elements (Co 99.9%, As 99.5%). The elements were weighed in an atomic ratio of 1:4 and
110	placed in a silica-glass ampoule. The ampoule was evacuated, sealed and placed in a horizontal
111	tube furnace. The temperature of that part of the ampoule, in which the charge was located, was
112	500 °C. The temperature of the opposite part was 450 °C. The synthesis lasted 50 days. The
113	excess arsenic condensed in the colder part of the ampoule and was easily separated. After that
114	the reacted mixture was removed, ground in an agate mortar under acetone, placed in another
115	silica-glass ampoule, and heated at the same temperature for 30 days.
116	METHODS
117	Powder X-ray diffraction (PXRD) patterns of all compounds used in this work were collected
118	using a Bruker D8 Advance DaVinci diffractometer (Institute of Geoscience of the University of
119	Jena, Germany), employing Cu K $\alpha$ radiation ( $\lambda = 1.54058$ Å). The patterns were collected at
120	room temperature between 5 and 90 °2 $\theta$ , with a step size of 0.02 °2 $\theta$ and a time per step of 1 s.
121	Lattice parameters and quantitative fractions of the studied phases were refined by a full-profile
122	fit using the software suite JANA2006 (Petříček et al., 2014).
123	The powdery synthetic phases were embedded in a resin and polished. They were briefly
124	examined for impurity phases or inhomogeneities by reflected-light microscopy. The phases
125	were investigated by energy- and wavelength-dispersive (WDX) electron microprobe (EMP)
126	analyses with the electron microprobe JEOL JXA-8230 Superprobe (Institute of Geoscience of
127	the University of Jena, Germany). All samples were carbon-coated and measured with an
128	accelerating voltage of 20 kV and 20 nA. The spot size was 1 $\mu$ m and the measurement time on
129	peak and on background 40 s. Peak overlap correction was used to avoid peak interference
130	between the lines of S and Co. The elements were analyzed on the following emission lines,

standards, with the estimated detection limits (in wt.%): Sb – L $\alpha$ , InSb, 0.07; S – K $\alpha$ , pyrite, 131 0.04; Fe – K $\alpha$ , pyrite, 0.04; As – K $\alpha$ , arsenopyrite, 0.07; Cu – K $\alpha$ , chalcopyrite, 0.05; Co – K $\alpha$ , 132 133 Co metal, 0.05; Ni – K $\alpha$ , Ni metal, 0.04. High-temperature oxide melt solution calorimetry has been described in detail by Navrotsky 134 (1997, 2014). The experiments were performed at the Arizona State University, USA. The 135 136 samples were dropped into a molten sodium molybdate  $(3Na_2O \cdot 4MoO_3)$  solvent, maintained at 1073 K. Oxygen gas was flushed over the solvent at 90 mL min<sup>-1</sup> and bubbled through it at 5 137  $mL \cdot min^{-1}$ . The choice of the calorimetric solvents is limited to a few compositions, namely 138 sodium molybdate, lead borate, or alkali borate (see Navrotsky 1997). In the recent work of 139 140 Abramchuk et al. (2020), sodium molybdate at 1073 K was used to dissolve a suite of elements, 141 thus enabling to close the thermochemical cycle without the need for reference phases. We used 142 this advantage that also dictated the choice of the solvent and operation temperature. 143 Low temperature heat capacity  $(C_p)$  was measured by relaxation calorimetry using a commercial Physical Properties Measurement System (PPMS, from Quantum Design, San 144 Diego, California) at the University of Salzburg, Austria. With due care, accuracy can be within 145 1 % at 5 K to 300 K, and 5 % at 0.7 K to 5 K (Dachs and Bertoldi, 2005; Kennedy et al., 2007). 146 The powdered samples were wrapped in a thin Al foil and compressed to produce a  $\approx 0.5$  mm 147 148 thick pellet, which was then placed onto the sample platform of the calorimeter for measurement. Differential scanning calorimetry (DSC) was used to measure heat capacities near and above 149 room temperature using a Perkin Elmer Diamond DSC. Samples were measured in a flow of dry 150 151 argon. Further details of the method are described by Benisek et al. (2012). The entropy was calculated by integration of the  $C_{\rm p}/T$  function in the interval from 0 K to 298.15 K. 152

153

## RESULTS

#### **Characterization: Powder X-ray diffraction, chemical analysis** 154

155	All samples used in this study were crystalline and consisted of a single phase. No indications
156	of an amorphous impurity were found. The peaks were sharp, with their width (measured as
157	FWHM) comparable to that of the LaB6 standard. Refined lattice parameters of all phases are
158	given in Table 1. The PXRD patterns are shown in Figs. S1 and S2.
159	Examination with the electron microprobe confirmed that the samples consist of the elements
160	of interest, either Ni-S or Co-As, with trace amount of other elements (Table 1). Images in back-
161	scattered electrons (BSE) showed that the heazlewoodite and skutterudite samples were
162	homogeneous. They were made of porous, angular aggregates, up to several hundreds of
163	micrometers large, and their shape reflects the shape of the starting material (metal filing).
164	Chemical composition of the synthetic heazlewoodite sample corresponds to the nominal
165	formula $Ni_3S_2$ . For skutterudite, a slight deviation from the nominal formula was found. The
166	composition, determined by EMP analyses, is CoAs2.92. All thermodynamic data reported in this
167	work are normalized to the molecular masses of $Ni_3S_2$ (240.2102 g·mol <sup>-1</sup> ) and CoAs <sub>2.92</sub>
168	$(277.7043 \text{ g} \cdot \text{mol}^{-1}).$
169	Determination of enthalpies of formation

Determination of enthalples of formation 168

170 Enthalpies of formation were determined *via* appropriate thermochemical cycles (Tables 2, 3) 171 from the experimentally measured drop solution enthalpies. No difficulties or anomalies were encountered during calorimetric experiments. The samples dissolved rapidly and reproducibly. 172 173 The large drop solution enthalpies (Table 4) relate to the fact that the samples are oxidized and dissolved in the oxygen atmosphere in the active zone of the calorimeter. The large signals cause 174 also larger uncertainties attached to the data; these uncertainties are expressed in calorimetry as 175

two standard deviations of the mean. Relative to the magnitude of the signal, the uncertainties
are smaller than usually encountered. They are only 0.09 % and 0.24 % for the drop solution
enthalpies of Ni<sub>3</sub>S<sub>2</sub> and CoAs<sub>2.92</sub>, respectively, which is smaller than the usual precision of 1-2
%. The resulting enthalpies of formation are given in Table 5.

# 180 **Determination of low-temperature** $C_p$ **and entropies**

181 Heat capacity at temperatures between 2 and 303 K was measured for the synthetic

heazlewoodite and skutterudite samples with the commercial PPMS system. The  $C_p$  data for both

synthetic phases show no anomalies. Attempts to fit the low-temperature portion of the data (T  $\leq$ 

184 20 K) with a Debye function ( $C_p = A_3T^3$ ) or an extended Debye function ( $C_p = A_3T^3 + A_5T^5 + A_5T^5$ )

185  $A_7T^7$ ) failed. Satisfactory fits required the inclusion of a linear (A<sub>1</sub>T) or quadratic (A<sub>2</sub>T<sup>2</sup>) term,

probably reflecting the semi-conducting nature of the phases studied (see Gopal 1966). The data

between 20 and 300 K were fitted by two orthogonal polynomials. The adjustable parameters of

these functions, together with the raw data and the results of integration, are available in the

supporting electronic information. The calculated standard entropies (S<sup>o</sup> at T = 298.15 K) are

190 listed in Table 5.

## 191 **Determination of high-temperature** $C_p$

192 Heat capacity above 280 K was measured by differential scanning calorimetry (DSC). In the

region of overlap of the DSC and PPMS data, the agreement between the data sets was good.

194 The deviation was 1.3 % for heazlewoodite and 2.2 % for skutterudite. Upon increasing

- temperature, the data became unusually scattered. At  $T \approx 366$  K (skutterudite) or  $\approx 466$  K
- 196 (heazlewoodite), the data were so erratic that the measurements were stopped. Even though the
- samples were measured in an inert atmosphere of purified argon, it has to be assumed that traces

198	of oxygen caused partial oxidation of the samples. Of particular problem is sulfur or arsenic that
199	co-exists in the gas phase with the studied phases (see Leegaard and Rosenqvist 1964). With
200	increasing temperature, more sulfur or arsenic is transferred into the gas phase and can rapidly
201	react with small amounts of oxygen.
202	Extrapolation of $C_p$ up to 700 K
203	In order to be able to use the presented data for geochemical models of hydrothermal systems,
204	heat capacity must be determined up to higher temperatures than those that were reached
205	experimentally in this work. For this purpose, the available experimental $C_p$ data were
206	augmented by heat capacities at constant volume ( $C_v$ ) computed using a model proposed by
207	Kieffer (1985, her equation 60)
208	$C_{\rm v} \propto S({\rm x}_{\rm i}) + K({\rm x}_{\rm u},{\rm x}_{\rm i}) + E({\rm x}_{\rm E}) \tag{7}$
203	
209	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function,
210 211	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the
<ul><li>210</li><li>210</li><li>211</li><li>212</li></ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail
<ul> <li>210</li> <li>210</li> <li>211</li> <li>212</li> <li>213</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here.
<ul> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally
<ul> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> <li>215</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally or by a lattice-dynamic calculation. We emphasize that the function (7) is not a fit to data but is
<ul> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> <li>215</li> <li>216</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally or by a lattice-dynamic calculation. We emphasize that the function (7) is not a fit to data but is based purely on experimental or calculated frequencies of low- to high-frequency modes in a
<ul> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> <li>215</li> <li>216</li> <li>217</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally or by a lattice-dynamic calculation. We emphasize that the function (7) is not a fit to data but is based purely on experimental or calculated frequencies of low- to high-frequency modes in a structure. The function (7), as used in this work, contains no adjustable parameters.
<ul> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> <li>215</li> <li>216</li> <li>217</li> <li>218</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally or by a lattice-dynamic calculation. We emphasize that the function (7) is not a fit to data but is based purely on experimental or calculated frequencies of low- to high-frequency modes in a structure. The function (7), as used in this work, contains no adjustable parameters. For the calculation for skutterudite, frequencies of the Raman bands were taken from https://
<ol> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> <li>215</li> <li>216</li> <li>217</li> <li>218</li> <li>219</li> </ol>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally or by a lattice-dynamic calculation. We emphasize that the function (7) is not a fit to data but is based purely on experimental or calculated frequencies of low- to high-frequency modes in a structure. The function (7), as used in this work, contains no adjustable parameters. For the calculation for skutterudite, frequencies of the Raman bands were taken from https:// rruff.info/Skutterudite/R050593 (accessed on September 10, 2021). The frequencies for the
<ul> <li>210</li> <li>211</li> <li>212</li> <li>213</li> <li>214</li> <li>215</li> <li>216</li> <li>217</li> <li>218</li> <li>219</li> <li>220</li> </ul>	where $S(x_i)$ is the heat capacity function of a monatomic solid with a sine dispersion function, $K(x_u,x_i)$ refers to heat capacity of a continuum between frequencies $x_u$ and $x_i$ , and $E(x_E)$ is the heat capacity function of an Einstein oscillator. These functions and their prefactors are in detail defined by the equations 57-60 in Kieffer (1985) and will not be repeated here. The frequencies necessary to calculate $C_v$ according to (7) can be determined experimentally or by a lattice-dynamic calculation. We emphasize that the function (7) is not a fit to data but is based purely on experimental or calculated frequencies of low- to high-frequency modes in a structure. The function (7), as used in this work, contains no adjustable parameters. For the calculation for skutterudite, frequencies of the Raman bands were taken from https:// rruff.info/Skutterudite/R050593 (accessed on September 10, 2021). The frequencies for the acoustic branches were taken from Li and Mingo (2014) who performed first-principles

frequencies were taken from https://rruff.info/Heazlewoodite (accessed on September 10, 2021).

223 The acoustic modes were characterized by first-principles calculations by Tian et al. (2021).

The isochoric heat capacity  $C_v$  can be used to calculate the isobaric heat capacity  $C_p$  as

225 
$$C_{\rm p} = C_{\rm v} + {\rm VT}\alpha^2 {\rm B}\kappa \qquad (8)$$

where V is the molar volume, T is the thermodynamic temperature,  $\alpha$  is thermal expansion

227 coefficient, and B is the bulk modulus. The variable  $\kappa$  is the only adjustable parameter that enters

the calculations. It is necessary, because the calculated  $C_p$  is slightly lower than the experimental

229  $C_p$  data. Such situation can be encountered in these calculations and is corrected by  $\kappa$  that can be

regarded as an anharmonicity term in the equation (8).

For skutterudite,  $\alpha$  was taken from Rogl et al. (2010) for CoSb<sub>3</sub> skutterudite, and B from

Zouablia et al. (2020) for CoP<sub>3</sub> skutterudite. The value of  $\kappa$  was set to 2.4. For heazlewoodite,  $\alpha$ 

was taken from Tian et al. (2021), B from Zhang et al. (2020), and  $\kappa$  was set to 2.4.

The fits are compared to the experimental data and the extrapolations are shown in Fig. 1. The

results in the range of 280-700 K were re-fitted with polynomials accepted by the program

236 SUPCRTBL (Johnson et al. 1992, Zimmer et al. 2016):

237 
$$C_{p} = a + bT + cT^{-2} + dT^{-0.5}$$
 (9)

The coefficients *a*-*d* are listed in Table 6 in the format required by SUPCRTBL. Note that the values for the coefficients *a*, *c*, and *d* are to be used in SUPCRTBL as listed; the value of *b* should be multiplied by  $10^{-5}$  before inserting into the equation (9). The resulting  $C_p$  is in kJ·mol<sup>-</sup>  $^{1}\cdot$ K<sup>-1</sup>.

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#### DISCUSSION

#### 243 Comparison of the data for heazlewoodite with earlier data

The enthalpies of formation of the two phases studied here can be compared to some previous 244 results. The  $\Delta_{\rm f} {\rm H}^{\rm o}$  value for heazlewoodite (Table 5) is in excellent agreement with the value 245 selected by Gamsjäger et al. (2005), based on the calorimetric work of Cemic and Kleppa (1986). 246 They used reaction calorimetry with Ni+NiS mixtures sealed in evacuated silica tubes. The 247 entropy of heazlewoodite, derived in this work, is also in excellent agreement with the selection 248 of Gamsjäger et al. (2005). Our value is essentially identical with that from Weller and Kelley 249 (1964) and only slightly higher than that from Stølen et al. (1991) (Table 5). 250 251 The perfect agreement for heazlewoodite confirms the accuracy of the data for the reference 252 compounds and the applicability of high-temperature oxide-melt solution calorimetry. It lends 253 support for the data for skutterudite where earlier data are missing.

## 254 Equilibria between nickel sulfides

The accuracy of the data for the nickel sulfides can be, to some extent, tested by comparison to known phase diagrams, as already presented in Gamsjäger et al. (2005). Sulfur pressure over nickel sulfide and the T-P(S<sub>2</sub>,g) data were collected by Rosenqvist (1954) and Leegaard and Rosenqvist (1964). The pair heazlewoodite-millerite ( $\beta$ -NiS) is related by a simple reaction

259 
$$2Ni_3S_2 + S_2(g) = 6\beta - NiS$$
 (10)

The data for millerite were taken from Gamsjäger et al. (2005). The log  $K_{10}$  values calculated with SUPCRTBL and the resulting sulfur gas fugacities are compared to the measured data in Fig. 2. The temperature ranges of the calculations and measurements do not overlap but linear extrapolation is possible. The difference between the log  $fS_2$  values predicted with our data and

264	those from Rosenqvist (1954) (Fig. 2) at $T = 900$ K is $\approx 1.5$ , corresponding to $RT\Delta \ln fS_2$ of $\approx 26$
265	$kJ \cdot mol^{-1}$ . Optimization of the data to fit the high-temperature measurements of Rosenqvist
266	(1954) would be misleading because the nickel sulfides tend to non-stoichiometry at high
267	temperatures (as $Ni_3S_{2\pm x}$ , $NiS_{1+x}$ , see Leegaard and Rosenqvist 1964). The deviation from
268	stoichiometry is perhaps responsible for the mismatch seen in Fig. 2. Additionally, these phases
269	have high-temperature polymorphs with slightly different thermodynamic properties (see
270	Gamsjäger et al. 2005). All in all, the result of this test is encouraging for the accuracy of the
271	data for heazlewoodite and millerite. More detailed calculation of the sulfur pressures in nickel-
272	sulfide equilibria is provided by Gamsjäger et al. (2005) (their figure V-32) at high temperatures.
273	No attempt was made to incorporate the thermodynamic data for vaesite (NiS2) into these
274	calculations, even though Chase (1998) and Gamsjäger et al. (2005) list data for this phase.
275	These data, however, are associated with unusually large uncertainties. They indicate another,
276	underlying problem with these data. As a side note, we attempted to synthesize vaesite also for
277	this study. The sample turned out to be a mixture of $NiS_2$ and $Ni_3S_4$ . These two phases should
278	have pyrite-like and spinel-like structures, respectively. Powder XRD data for our sample show
279	only pyrite-like structure. This puzzling structural relationship will be more closely investigated
280	but could be potentially a reason (or one of the reasons) for the problems related to the
281	measurements on vaesite. Finally, the database Thermoddem (Blanc et al. 2012) contains
282	temperature-dependent log $K$ values for a dissolution reaction involving values. Yet, a closer
283	look at the values in this database (https://thermoddem.brgm.fr/species/nis2, accessed on
284	September 9, 2021) shows that $C_p$ of vaesite was taken as a constant over 0-300 °C. Such
285	treatment of thermodynamic data is highly misleading and should be avoided.

# 286 Equilibrium constants for dissolution reactions

Geochemical software used for modeling of aqueous solutions, either around room temperature or at elevated temperatures, commonly requires tabulation of equilibrium constants for selected dissolution reactions involving the phase of interest. For this purpose, all data presented above were inserted into SUPCRTBL (Zimmer et al. 2016) and the equilibrium constants were calculated.

The calculations are slightly complicated by the fact that the nominal valences, calculated from the formulae, differ from the valences usually assigned to the participating elements (such as  $Ni^{2+}$ ,  $S^{2-}$ , and so on). For heazlewoodite, the reactions to consider are

295 
$$Ni_3S_2(cr) + 6H^+(aq) + 0.5O_2(aq) \rightarrow 3Ni^{2+}(aq) + 2H_2S^0(aq) + H_2O(l)$$
 (11)

296 
$$Ni_3S_2(cr) + 4H^+(aq) + 0.5O_2(aq) \rightarrow 3Ni^{2+}(aq) + 2HS^-(aq) + H_2O(l)$$
 (12)

Under most hydrothermal conditions, the aqueous species  $H_2S^0$  will predominate over  $HS^-$ . The aggregate state of O<sub>2</sub> and H<sub>2</sub>O is important inasmuch the software packages usually do not explicitly specify it for the master O<sub>2</sub> and H<sub>2</sub>O species. The values of log  $K_{11}$  and log  $K_{12}$  in the temperature range 0-350 °C are listed in Table S1.

## 301 For skutterudite, the oxidation reactions are

302 
$$CoAs_{2.92}(cr) + 3.38H_2O(1) + 2.69O_2(aq) + 2H^+(aq) \rightarrow Co^{2+}(aq) + 2.92As(OH)_3^{0}(aq)$$
 (13)

303 
$$CoAs_{2.92}(cr) + 3.38H_2O(l) + 2.69O_2(aq) \rightarrow Co^{2+}(aq) + 2.92H_2AsO_3(aq) + 0.92H^+(aq)$$
 (14)

For the arsenical species,  $As(OH)_3^0$  is prevalent in hydrothermal fluids, and  $AsO(OH)_2^-$  (this

species is listed as H<sub>2</sub>AsO<sub>3</sub><sup>-</sup> in GWB or AsO<sub>2</sub><sup>-</sup> in SUPCRTBL) is subordinate, predominant only

at pH > 8 at T = 300 °C (Testemale et al. 2011). The values of log  $K_{13}$  and log  $K_{14}$  in the

temperature range 0-350 °C are listed in Table S1.

308 Some geochemical software accepts a function that expresses log *K* in terms of absolute

- temperature T in Kelvin instead of log *K* values at a pre-selected temperature grid. PHREEQC,
- for example, uses the function in its **-analytic** option

311 
$$\log K = A_1 + A_2T + A_3/T + A_4 \log T + A_5/T^2$$
 (15)

GWB expresses the temperature dependence of  $\log K$  as

313 
$$\log K = a + b(T-T_r) + c(T^2-T_r^2) + d(1/T-1/T_r) + e(1/T^2-1/T_r^2) + f \ln(T/T_r)$$
(16)

where T is absolute temperature in Kelvin and  $T_r = 298.15$  K.

The fit parameters  $A_1$ - $A_5$  for equation (15) and the reactions (11) through (14) are listed in

Table S2. The fit parameters a-f for equation (16) and the reactions (11) through (14) are listed in

317 Table S3.

## 318 Relationship of skutterudite to other cobalt arsenides

The  $\Delta_f H^o$  value for skutterudite (Table 5) can be only compared to an estimate which is

significantly different. We are not aware of thermodynamic measurements on skutterudite.

321 The GWB and LLNL databases contain log *K* values for reactions involving safflorite

322 (CoAs<sub>2</sub>) and modderite (CoAs), even though no references are given. The log  $K_{T=298.15}$  values

match those given in Charykova et al. (2011) who also gave no references. Wagman et al. (1982)

listed only  $\Delta_{\rm f} {\rm H}^{\rm o}$  for these phases, also without references. For example, they gave  $-61.5 \, {\rm kJ \cdot mol^{-1}}$ 

for CoAs<sub>2</sub>, but Naumov et al. (1974) selected –83 kJ·mol<sup>-1</sup>. These data, even though of uncertain

- 326 origin, allow the calculation of phase diagrams whose informative value should not be overrated.
- 327 Scharrer et al. (2019) constructed elaborate models with these phases and generated estimates for
- the missing data. Ironically, they themselves do not recommend the use of some of the data they

329	published (see their Table 3). We refrain from such practice and hold the opinion that
330	geochemical modeling should be based on sound data. As pointed out by Henke (2009, appendix
331	C) with respect to thermodynamics of arsenical phases, "once high-quality laboratory data are
332	available, they may be used in geochemical computer models".

333

## IMPLICATIONS

334 This work enhances the capabilities of geochemical modeling of hydrothermal ores, especially those of the so-called 'five-element association' (e.g., Kreissl et al. 2018, Scharrer et al. 2019, 335 336 Horn et al. 2021, Tourner et al. 2021). Skutterudite is one of the major sources of cobalt in these 337 ores, being used for the construction of lithium-ion batteries (Li et al. 2018). Therefore, understanding of these ores, the ore-forming processes, and identification of new resources is an 338 important future task for the geosciences. This work also documented the capability of 339 340 calorimetric techniques to obtain highly accurate data for chalcogenides. Combination of the experimental data with the Kieffer model generated accurate representation of high-temperature 341 heat capacities and can be used to circumvent experimental difficulties associated with the 342 measurements at high temperatures. This approach can be used to resolve some long-standing 343 controversies in the thermodynamic data sets for common and less common chalcogenide 344 minerals. 345

346

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538 Table 1. Space groups, lattice parameters (a, c, unit-cell volume V), and chemical composition (wt.%) of 539 the synthetic phases used in this work. Chemical composition is recalculated to the molar ratio in the ant digits.

540	bottom part of the table.	The numbers in	parentheses are t	the uncertainties on	the last significa
541	nominal				

542	composition	Ni <sub>3</sub> S <sub>2</sub>	CoAs <sub>3</sub>		
543 E 4 4	actual	NI; C	Cala		
544	mineral name	INI302 heazlewoodite	COAS <sub>2.92</sub> skutterudite		
545					
540	space group $\tilde{a}(\lambda)$	K32 5 7476(1)	Im-3 8 2050(2)		
547 E10	$a(\mathbf{A})$	3.7470(1) 7.1207(1)	8.2030(2)		
540	$\mathcal{C}(\mathbf{A})$	12	11		
549	n Sh	13 < DI	 < DI		
	50 S	< DL	< DL 0.06(1)		
221	5	20.83(28)	0.00(1)		
552	AS Ca	0.04(0)	/8.21(8/)		
555		< DL	21.0/(27)		
554		< DL 72 (5(40)	< DL		
555	N1	/2.65(40)	< DL		
556	Fe	< DL	0.04(2)		
557	total	99.54(38)	99.41(102)		
558		normalized to 1 N	or Co atom		
559	S	0.676			
560	As		2.920		
561	Co		1		
562	Ni	1			
563	n = number of anal	yses averaged			
564	< DL = below dete	ction limit			
565					
566	Table 2 Thermod	lynamic cycle used t	o calculate the enthalpies	of formation for hear	zlewoodite
567			•••••••		210 11 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
568	Ni S (s 298 K) -	$+(2y+r/2)\Omega_{2}(\sigma -107)$	$(3 \text{ K}) \rightarrow$		
560	$rNi^{2+}$ (sol 1073 k	$(2y + x/2)O_2(g, 107)$ $(x) + xO^{2-}(g, 1073)$	$K + v S O c^{2-}$ (sol 1073 K)	<b>\</b>	ΔН.
570	$N_{i}$ (s01, 10751 $N_{i}$ (s 208 K) + 0	$5 \Omega_{\alpha} (\sigma \ 1073 \text{ K}) \rightarrow 1073 \text{ K}$	$N_{i}^{2+}$ (sol 1073 K) + $\Omega^{2-}$ (	( sol 1073 K)	
570	NI(3, 290  K) + 0 S(a, 209  K) + 20	$(\sigma 1072 \text{ K}) \rightarrow 0$	$^{2-}(aol 1072 K)$	501, 1075  K	
571	S(S, 290  K) + 20	$V_2(\mathbf{g}, 1073 \text{ K}) \rightarrow 30$	4 (501, 1075  K)	$\Lambda U^0 - \Lambda U \perp r\Lambda$	$\Delta I I_3$
572	$\lambda NI (S, 298 K) + 2$	$y_{\mathcal{S}}(s, 298 \text{ K}) \rightarrow \text{NI}_x$	Dy (8, 290 K)	$\Delta_{\rm f} \Pi = -\Delta \Pi_1 + \chi \Delta$	$n_2 + y \Delta H_3$
573					
574	Table 3. Thermoo	dynamic cycle used t	o calculate the enthalpy o	of formation for skutte	erudite.
575					

576	$\overline{\text{CoAs}_z(s, 298 \text{ K}) + (2z+1/2)O_2(g, 1073 \text{ K})} \rightarrow$		
577	$Co^{2+}$ (sol, 1073 K) + $O^{2-}$ (sol, 1073 K) + $zAsO_4^{3-}$ (sol, 1073 K)	$\Delta H_4$	
578	Co (s, 298 K) + 0.5O <sub>2</sub> (g, 1073 K) → Co <sup>2+</sup> (sol, 1073 K) + O <sup>2−</sup> (sol	, 1073 K) ΔH <sub>5</sub>	
579	As $(s, 298 \text{ K}) + 2O_2 (g, 1073 \text{ K}) \rightarrow AsO_4^{3-} (sol, 1073 \text{ K})$	$\Delta H_6$	
580	$Co (s, 298 \text{ K}) + zAs (s, 298 \text{ K}) \rightarrow CoAs_z (s, 298 \text{ K})$	$\Delta_{\rm f} {\rm H}^{\rm o} = -\Delta {\rm H}_4 + \Delta {\rm H}_5 + z \Delta {\rm H}_6$	
581			

582

1	solution enthalple	s of the sample	es and refere	ence phases. All v	values in kJ·mol <sup>-1</sup> .
samples	1(05,073+1,5)	(10)6			
$N_{13}S_2$	$-1685.9/^{\circ} \pm 1.56^{\circ}$	$(10)^{\circ}$			
$COAS_{2.92}$	-1424.46±3.40 (	10)			
reference phas	ses				
Ni	-214.52±2.06 (H	Iayun et al. 20	20)		
Со	-227.75±2.18 (H	Iayun et al. 20	20)		
5	-629.18±2.73 (A	bramchuk et a	al. 2020)		
As	-440.03±1.57 (A	bramchuk et a	al. 2020)		
mean	1				
' two standard	deviations of the i	nean			
number of m	easurements				
Fable 5. Summer	nary of thermodyn	amic data for	heazlewoodi	ite and skutterudi	ite, obtained in this wo
Calculated ent	halpies of formation	on are compar	ed to the pre	viously publishe	d literature values.
Coefficients fo	or the C. nolynomi	als are listed i	n Table 6	,	
	or the ep porynomia	als are listed i	II Tuble 0.		
	٨cHo	So		$\Delta \epsilon G^0$	Vº
	/				•
	$kJ \cdot mol^{-1}$	J·mol <sup>-1</sup> ·K	-1	$kJ \cdot mol^{-1}$	J·bar <sup>-1</sup>
Ni <sub>2</sub> S <sub>2</sub>	$kJ \cdot mol^{-1}$ -216.0±8.4	$\frac{J \cdot \text{mol}^{-1} \cdot \text{K}}{133.8 \pm 1.6}$	5	$kJ \cdot mol^{-1}$ -210.0±8.4	$\frac{J \cdot bar^{-1}}{4.095}$
Ni <sub>3</sub> S <sub>2</sub>	$kJ \cdot mol^{-1}$ -216.0±8.4 -217.2±1.6 <sup>a</sup>		57 <sup>a</sup>	$\frac{100}{\text{kJ} \cdot \text{mol}^{-1}}$ -210.0±8.4 -211.2±1.6 <sup>a</sup>	J·bar <sup>-1</sup> 4.095
Ni <sub>3</sub> S <sub>2</sub>	$\frac{kJ \cdot mol^{-1}}{-216.0 \pm 8.4}$ -217.2 \pm 1.6 <sup>a</sup>		5 7 <sup>a</sup>	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ -211.2±1.6 <sup>a</sup>	J·bar <sup>-1</sup> 4.095
Ni <sub>3</sub> S <sub>2</sub>	$\frac{kJ \cdot mol^{-1}}{-216.0 \pm 8.4}$ -217.2 \pm 1.6 <sup>a</sup>	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2°	5 7 <sup>a</sup> .84 <sup>b</sup>	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ -211.2 \pm 1.6 <sup>a</sup>	J·bar <sup>-1</sup> 4.095
Ni <sub>3</sub> S <sub>2</sub> CoAs2 92	$\frac{\text{kJ} \cdot \text{mol}^{-1}}{-216.0 \pm 8.4}$ $-217.2 \pm 1.6^{\text{a}}$ $-88.2 \pm 6.1$	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	7 <sup>a</sup> .84 <sup>b</sup>	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{a}$ $-79.9 \pm 6.2$	J·bar <sup>-1</sup> 4.095 4.159
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub>	$\frac{kJ \cdot mol^{-1}}{-216.0 \pm 8.4}$ $-217.2 \pm 1.6^{a}$ $-88.2 \pm 6.1$ $-104.6^{d}$	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	5 7 <sup>a</sup> .84 <sup>b</sup> 3	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{a}$ $-79.9 \pm 6.2$	J·bar <sup>-1</sup> 4.095 4.159
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub>	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ $t al. (2005)$	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	5 7 <sup>a</sup> .84 <sup>b</sup> 3	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{a}$ $-79.9 \pm 6.2$	J·bar <sup>-1</sup> 4.095 4.159
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> <sup>1</sup> Gamsjäger et <sup>2</sup> Weller and K	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ t al. (2005) Celley (1964)	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	5 7 <sup>a</sup> .84 <sup>b</sup> 3	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{a}$ $-79.9 \pm 6.2$	J·bar <sup>-1</sup> 4.095 4.159
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ $t al. (2005)$ $Celley (1964)$ $(1991)$	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	1 5 7 <sup>a</sup> .84 <sup>b</sup> 3	$\frac{kJ \cdot mol^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{a}$ $-79.9 \pm 6.2$	J·bar <sup>-1</sup> 4.095 4.159
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically	$\frac{\text{kJ} \cdot \text{mol}^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{\text{a}}$ $-88.2\pm 6.1$ $-104.6^{\text{d}}$ $\text{t al. (2005)}$ $\text{Kelley (1964)}$ $(1991)$ $\text{calculated for Code}$	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	1 5 7 <sup>a</sup> .84 <sup>b</sup> 3	$kJ \cdot mol^{-1}$ -210.0±8.4 -211.2±1.6 <sup>a</sup> -79.9±6.2	J·bar <sup>-1</sup> 4.095 4.159
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically accessed on	$kJ \cdot mol^{-1}$ -216.0±8.4 -217.2±1.6 <sup>a</sup> -88.2±6.1 -104.6 <sup>d</sup> t al. (2005) Celley (1964) (1991) calculated for CoA	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	1 5 7 <sup>a</sup> .84 <sup>b</sup> 3	kJ·mol <sup>-1</sup> -210.0±8.4 -211.2±1.6 <sup>a</sup> -79.9±6.2	J·bar <sup>-1</sup> 4.095 4.159 ol/JVASP-20306.html
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically accessed on	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ t al. (2005) Celley (1964) (1991) calculated for CoA n July 21, 2021	J·mol <sup>-1</sup> ·K 133.8±1.6 133.5±0.7 133.89±0 133.2° 106.4±1.3	1 5 7 <sup>a</sup> .84 <sup>b</sup> 3	kJ·mol <sup>-1</sup> -210.0±8.4 -211.2±1.6 <sup>a</sup> -79.9±6.2	J·bar <sup>-1</sup> 4.095 4.159 ol/JVASP-20306.html
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically accessed on Fable 6 Adius	$\frac{\text{kJ} \cdot \text{mol}^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{\text{a}}$ $-88.2\pm 6.1$ $-104.6^{\text{d}}$ $\text{t al. (2005)}$ $\text{Celley (1964)}$ $(1991)$ $\text{calculated for CoA}$ $\text{n July 21, 2021}$ $\text{stable coefficients}$	$\frac{J \cdot mol^{-1} \cdot K}{133.8 \pm 1.6}$ $133.8 \pm 0.7$ $133.89 \pm 0$ $133.2^{\circ}$ $106.4 \pm 1.3$ As <sub>3</sub> , <u>https://www</u> for the C <sub>2</sub> poly	1 5 7 <sup>a</sup> .84 <sup>b</sup> 3 <u>vw.ctcms.nis</u>	$\frac{kJ \cdot mol^{-1}}{-210.0\pm 8.4}$ $-211.2\pm 1.6^{a}$ $-79.9\pm 6.2$ $\frac{kt.gov}{-knc6/jsma}$	$\frac{J \cdot bar^{-1}}{4.095}$ 4.159 <b>ol/JVASP-20306.html</b> + $dT^{-0.5}$ , equation 9) for
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Veller and K Stølen et al. ( theoretically accessed of <b>Table 6</b> . Adjus	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ t al. (2005) Celley (1964) (1991) calculated for CoA n July 21, 2021 stable coefficients ad heazlewoodite	$\frac{J \cdot mol^{-1} \cdot K}{133.8 \pm 1.6}$ $133.8 \pm 1.6$ $133.5 \pm 0.7$ $133.89 \pm 0$ $133.2^{c}$ $106.4 \pm 1.3$ As <sub>3</sub> , <u>https://www</u> for the <i>C</i> <sub>p</sub> poly	$c^{-1}$ $c^{7a}$ $c^{7$	$\frac{\text{kJ} \cdot \text{mol}^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{\text{a}}$ $-79.9 \pm 6.2$ $\frac{\text{kt.gov}}{-79.9 \pm 6.2}$ $= a + b\text{T} + c\text{T}^{-2} - \frac{1}{2}$ $= c + b\text{T} + c\text{T}^{-2} - \frac{1}{2}$ $= c + b\text{T} + c\text{T}^{-2} - \frac{1}{2}$	$\frac{J \cdot bar^{-1}}{4.095}$ $4.159$ $dT^{-0.5}, equation 9) for the superstant of the $
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically accessed or rable 6. Adjus skutterudite ar letails	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ $t al. (2005)$ $(1991)$ $calculated for CoA$ $h July 21, 2021$ $stable coefficients$ $h d heazlewoodite.$	$\frac{J \cdot mol^{-1} \cdot K}{133.8 \pm 1.6}$ $133.8 \pm 1.6$ $133.5 \pm 0.7$ $133.89 \pm 0$ $133.2^{c}$ $106.4 \pm 1.3$ As <sub>3</sub> , <u>https://www</u> for the <i>C</i> <sub>p</sub> poly The numerical	$\frac{c^{-1}}{c^{7^{a}}}$ .84 <sup>b</sup> .84 .84 <sup>b</sup> .84 .84 <sup>b</sup> .84 .84 <sup>b</sup> .84 .84 <sup>b</sup> .84 <sup>b</sup>	$\frac{kJ \cdot mol^{-1}}{-210.0\pm 8.4}$ $-211.2\pm 1.6^{a}$ $-79.9\pm 6.2$ $= a + bT + cT^{-2} - b$ scaled as required	$\frac{J \cdot bar^{-1}}{4.095}$ 4.159 61/JVASP-20306.html + $d T^{-0.5}$ , equation 9) for the superstant set of the
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically accessed or Table 6. Adjus skutterudite ar details.	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ t al. (2005) Celley (1964) (1991) calculated for CoA n July 21, 2021 stable coefficients ad heazlewoodite.	$\frac{J \cdot mol^{-1} \cdot K}{133.8 \pm 1.6}$ $133.8 \pm 1.6$ $133.5 \pm 0.7$ $133.89 \pm 0$ $133.2^{c}$ $106.4 \pm 1.3$ As <sub>3</sub> , <u>https://www</u> for the <i>C</i> <sub>p</sub> poly The numerical b	$\frac{c^{-1}}{c^{7^{a}}}$ $\frac{c^{-1}}{c^{7^{a}}}$ $\frac{c^{-1}}{c}$ $\frac{c^{-1}}{c}$ $\frac{c^{-1}}{c}$	$\frac{\text{kJ} \cdot \text{mol}^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{\text{a}}$ $-79.9 \pm 6.2$ $\frac{\text{kt.gov/~knc6/jsma}}{\text{scaled as required}}$	$\frac{J \cdot bar^{-1}}{4.095}$ 4.159 <b>ol/JVASP-20306.html</b> + $d T^{-0.5}$ , equation 9) for by SUPCRTBL. See
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Weller and K Stølen et al. ( theoretically accessed on Fable 6. Adjus skutterudite ar letails.	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ tal. (2005) Celley (1964) (1991) calculated for CoA n July 21, 2021 stable coefficients nd heazlewoodite. a 0.1096	$\frac{J \cdot mol^{-1} \cdot K}{133.8 \pm 1.6}$ $133.8 \pm 0.7$ $133.89 \pm 0$ $133.2^{c}$ $106.4 \pm 1.3$ As <sub>3</sub> , <u>https://www</u> for the <i>C</i> <sub>p</sub> poly The numerical b 0.8507	$\frac{c^{-1}}{7^{a}}$ $\frac{c^{-1}}{7^{a}}$ $\frac{c^{-1}}{7^{a}}$ $\frac{c^{-1}}{7^{a}}$	$\frac{kJ \cdot mol^{-1}}{-210.0\pm 8.4}$ $-211.2\pm 1.6^{a}$ $-79.9\pm 6.2$ $= a + bT + cT^{-2} - c$ scaled as required $\frac{d}{-0.2695}$	$\frac{J \cdot bar^{-1}}{4.095}$ 4.159 <b>ol/JVASP-20306.html</b> + $d T^{-0.5}$ , equation 9) for 1 by SUPCRTBL. See
Ni <sub>3</sub> S <sub>2</sub> CoAs <sub>2.92</sub> Gamsjäger et Veller and K Stølen et al. ( theoretically accessed of Table 6. Adjus skutterudite ar letails. bhase skutterudite neazlewoodite	$\frac{kJ \cdot mol^{-1}}{-216.0\pm 8.4}$ $-217.2\pm 1.6^{a}$ $-88.2\pm 6.1$ $-104.6^{d}$ t al. (2005) Celley (1964) (1991) calculated for CoA n July 21, 2021 stable coefficients ad heazlewoodite. a 0.1096 0.1361	$\frac{J \cdot mol^{-1} \cdot K}{133.8 \pm 1.6}$ $133.8 \pm 1.6$ $133.5 \pm 0.7$ $133.89 \pm 0$ $133.2^{c}$ $106.4 \pm 1.3$ As <sub>3</sub> , <u>https://www</u> for the <i>C</i> <sub>p</sub> poly The numerical <i>b</i> 0.8507 2.300	$\frac{c^{-1}}{c^{a}}$ $\frac{c^{-1}}{c^{a}}$ $\frac{c^{-1}}{c^{a}}$ $\frac{c^{-1}}{c^{a}}$	$\frac{\text{kJ} \cdot \text{mol}^{-1}}{-210.0 \pm 8.4}$ $-211.2 \pm 1.6^{\text{a}}$ $-79.9 \pm 6.2$ $= a + b\text{T} + c\text{T}^{-2}$ $= a + b\text{T} + c\text{T}^{-2}$ scaled as required $\frac{d}{-0.2695}$ $-0.2490$	$\frac{J \cdot bar^{-1}}{4.095}$ 4.159 61/JVASP-20306.html + $d T^{-0.5}$ , equation 9) for the superstanding superstance of the superstanding superst

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## FIGURE CAPTIONS

Fig. 1. Experimental heat capacities from PPMS (circles) and DSC (squares) with the Kieffer
model used to extrapolate C<sub>p</sub> to higher temperatures. Note that this is not the fit used to
determine entropy at T = 298.15 K.
Fig. 2. Equilibrium sulfur pressure for the pair heazlewoodite-millerite (reaction 10) calculated
from the data presented and discussed in this work (diamonds) and experimentally measured
by Rosenqvist (1954, circles). The linear fit is extrapolated to higher temperatures to allow for

a visual judgment of the agreement between the two data sets.



