1 Revision 1

Macroscopic electrostatic effects in ATR-FTIR spectra of modern and 2 archeological bones 3 Julie Aufort^{1,*}, Matthieu Lebon², Xavier Gallet², Loïc Ségalen³, Christel Gervais⁴, Christian 4 Brouder¹, Etienne Balan¹ 5 ¹ IMPMC, Sorbonne Université, UMR CNRS 7590, UMR IRD 206, MNHN, 4 place Jussieu, 75252 Paris, cedex 6 05 France 7 ² HNHP, Sorbonne Université, UMR CNRS 7194, MNHN – Univ. Perpignan Via Domitia, 17 Place du 8 Trocadéro, 75116 Paris - France 9 10 ³ ISTEP, Sorbonne Université, UMR CNRS 7193, 4 place Jussieu, 75252 Paris, cedex 05 France 11 ⁴ LCMCP, Sorbonne Université, UMR CNRS 7574, 4 place Jussieu, 75252 Paris cedex 05, France 12 13 ABSTRACT 14 Bones mostly consist of composite materials based on almost equivalent volume fractions of mineral (apatite) and organic (collagen) components. Accordingly, their infrared 15 spectroscopic properties should reflect this composite nature. In this letter, we show by theory 16 17 and experiment that the variability of the strong phosphate bands in the ATR-FTIR spectra of 18 a series of modern and archeological bone samples can be related to electrostatic interactions affecting apatite particles and depending on the bone collagen content. Key parameters 19 20 controlling the shape of these bands are the mineral volume fraction and the dielectric 21 constant of the embedding matrix. The magnitude of these effects is larger than the one related to microscopic changes of the apatite structure. Consequently, the interplay of 22 23 microscopic and macroscopic parameters should be considered when using FTIR spectroscopy to monitor the preservation state of bioapatite during diagenetic and fossilization 24 25 processes, especially during the degradation of the organic fraction of bone.

26 Keywords: ATR-FTIR spectroscopy, bone, fossilization, electrostatic properties

27 Corresponding author: Julie Aufort, IMPMC, case 115, 4 place Jussieu, 75252 Paris, cedex 05 France

28 E-mail : Julie.Aufort@impmc.upmc.fr Phone: 00.33.1.44.27.42.27 Fax: 00.33.1.44.27.37.85

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INTRODUCTION

30 Vertebrate skeletons mostly consist of composite materials containing a mineral 31 component and a fraction of organic and water molecules (Pasteris et al. 2008). The inorganic 32 component is usually described as nanocrystalline carbonate-bearing hydroxylapatite 33 $(Ca_5(PO_4)_3OH)$, although more complex structures and chemical compositions have been 34 observed (e.g. Pasteris et al. 2008; Li and Pasteris 2014). Bones consist of about 50 vol% of 35 \sim 2 nm thick, 10-20 nm wide and 20-50 nm long apatitic nanocrystallites and a near equivalent 36 volume of collagen fibrils (Glimcher 2006; Pasteris et al. 2008). Compared to bone, tooth 37 enamel lacks collagen, contains ~1-2 vol% of organic compounds and about 2 vol% of water, 38 and displays larger apatite crystallites making up for ~90 vol% of the material (LeGeros and 39 LeGeros 1984, Elliott 2002; Skinner 2005).

40 Key information on the atomic-scale organization of such complex materials is 41 provided by spectroscopic methods (e.g. Pasteris et al. 2008, Yi et al. 2014). As it is a 42 versatile technique requiring only small sample amounts, Fourier-transform infrared (FTIR) 43 spectroscopy has been widely applied to the study of bones and teeth. Although most of these 44 studies have biomedical goals, FTIR spectroscopy also allows efficient screening of the 45 preservation state of skeleton remains prior to isotopic and chemical analysis for paleoenvironmental reconstructions or radiocarbon dating (e.g.; Trueman et al. 2008; Roche et al. 46 47 2010; Lebon et al. 2014; Snoeck and Pelligrini 2015). Diagenetic transformation of bones and teeth following burial often leads to an increase in crystallinity as well as a decrease in 48 49 organic fraction and structural carbonate contents (e.g. Hedges 2002; Trueman et al. 2008).

In practice, FTIR spectra of bioapatite are increasingly recorded in the attenuated total
reflection (ATR) mode (e.g. Stathopoulou et al. 2008; Thompson et al. 2009; Lebon et al.
2014; Beasley et al. 2014; Surmik et al. 2016). Unlike conventional transmission geometry,
ATR spectroscopy does not require the multistep preparation of dry and homogeneous pellets

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containing a percent-level fraction of solid particles dispersed in a KBr matrix. The 54 55 absorption-like ATR spectrum is recorded on the packed sample powder by measuring its 56 reflection coefficient at the interface with a highly refractive material. However, the different 57 geometry and nature of the sample may impede a straightforward comparison of ATR with 58 transmission spectra (Aufort et al. 2016). Differences in the crystallinity degree and 59 carbonate-to-phosphate ratio of samples measured in transmission, ATR and diffuse reflectance geometries have been reported (Beasley et al. 2014). The position and intensity of 60 61 the ATR signal also depend on the refractive index of the ATR crystal (Boulet-Audet et al. 62 2010), typically diamond (n=2.4) or germanium (n=4). Use of Ge usually requires larger amounts of sample and leads to a comparatively smaller band intensity, thus limiting the 63 detection of relatively weak bands, such as those related to carbonate impurities and organic 64 65 matter. In addition, the ATR spectra depend on characteristics of the sample microstructure, 66 such as porosity, nature of the embedding medium, aggregation state, and shape of particles, 67 as recently shown for apatite (Aufort et al. 2016). This dependence arises from the long-range electrostatic interactions within and between particles, interactions which also affect the 68 transmission spectra (e.g. Iglesias et al. 1990; Balan et al. 2011, Kendrick and Burnett 2016). 69

Based on these observations, it can be anticipated that the FTIR spectra of composite biomaterials such as bone should depend on the relative fractions of mineral and organic constituents through macroscopic electrostatic interactions. In this contribution, we show that the line-shape variability of the ATR-FTIR spectrum of modern and archaeological bones can be accounted for by combining the description of the intrinsic vibrational properties of the mineral component with an electrostatic model taking into consideration their variable collagen fraction. 77

MATERIALS AND METHODS

Three archaeological animal bone and a modern ox bone samples (Table 1) previously 78 79 investigated using ATR-FTIR spectroscopy by Lebon et al. (2014) were chosen for their various states of collagen preservation as well as a modern tooth enamel sample previously 80 investigated by Yi et al. (2014). For each sample, about 50 mg of powder were finely ground 81 82 with agate mortar and pestle. The ATR-FTIR spectra were recorded using a Quest ATR 83 device (Specac) and a Nicolet 6700 FTIR spectrometer. Pure powder samples were packed at the surface of the ATR crystal and spectra recorded between 400 and 4000 cm⁻¹ by averaging 84 100 scans with a resolution of 1 cm⁻¹. For each sample, two different spectra were recorded on 85 the same Quest ATR device using interchangeable diamond and germanium ATR crystals. 86 The raw ATR spectra are reported in absorbance units, without any correction treatment or 87 baseline subtraction. 88

89 The modeling strategy has been extensively detailed in Aufort et al. (2016). Briefly, the theoretical ATR spectrum is obtained by computing the frequency-dependent reflection 90 coefficient with a 45° incidence angle at the interface between the ATR crystal, defined as a 91 92 homogeneous isotropic medium with constant refractive index (2.4 for diamond and 4 for 93 Ge), and the sample. The powder sample is characterized by a dielectric function computed 94 for a mixture of spherical apatite particles and embedding matrix, using the Bruggeman 95 effective medium model (Bruggeman 1935). In this model, both the particles and matrix 96 domains are considered as being inserted in an effective medium, leading to a self-consistent 97 relation between the effective medium properties and those of individual constituents. The 98 uniaxial dielectric tensor of bulk apatite is obtained from its microscopic vibrational and 99 dielectric properties using a Drude-Lorentz expression (Table 2). All microscopic sources of 100 broadening are assumed to be effectively accounted for by the damping parameter of the 101 Drude-Lorentz model. The embedding matrix is described using a constant refractive index

varying between 1 for air in a porous mineral-air sample to ~1.6 for pure collagen (Wang et 102 al. 1996). No attempt was made to consider the specific vibrational properties of collagen. 103 104 The varying parameters used to calculate the theoretical spectra are thus restricted to the 105 refractive index of the ATR crystal, the mineral fraction and the dielectric constant of the 106 embedding matrix (ε_h) (Table 1). Theoretical spectra computed for systematic variations of 107 these parameters are reported as Supplementary Information. The effective medium 108 approximation considers the sample homogeneous at the length scale of the experiment, 109 typically requiring particle sizes smaller than a few µm. To compare calculation with experiment, the theoretical intensities were multiplied by a scale factor, to account for 110 inhomogeneities of the samples at larger length scales, e.g. related to the imperfect grinding 111 112 of organic-mineral mixtures and/or imperfect contact between the powder sample and the 113 ATR crystal (Table 1).

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RESULTS AND DISCUSSION

116 Experimental diamond and Ge ATR spectra of modern ox bone (Figure 1) display the 117 usual absorption bands ascribed to the internal vibrations of apatitic orthophosphate groups, to 118 structural carbonate impurities located at the A and B sites (hydroxyl and phosphate sites, 119 respectively) of apatite structure, as well as to the amide bands related to the collagen matrix 120 (e.g. Elliott 2002). Its collagen content is estimated to be ~ 23 wt% as assessed from nitrogen concentration and intensity of amide bands (Lebon et al. 2014). A weak and broad feature at 121 ~ 700 cm⁻¹ is related to collagen absorption (Suppl. Info.). Similar spectral features for the 122 123 phosphate, carbonate and amide bands are observed in archaeological bones HAR and BZ-124 PJ1-09 which still display a significant collagen content (16 and 11 wt%, respectively). The 125 spectrum of altered archaeological bone (EH2-20) and modern tooth enamel display similar

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126 absorption bands except that the collagen amide bands are absent. The ATR spectra of these 127 two samples are about four times more intense than those recorded on the collagen-rich bone 128 samples. As previously observed by Aufort et al. (2016), significant differences are observed 129 between the spectra recorded using a Ge or a diamond crystal for the whole series of samples. 130 Beside a stronger intensity, the intense v_3 stretching and v_4 bending PO₄ bands of the 131 diamond-ATR spectra are broadened on their low-frequency side, and the v₃ band overlaps with the weaker $v_1 PO_4$ and $v_2 CO_3$ bands. They also display a greater variability among the 132 133 samples, their low-frequency asymmetry being more pronounced for the collagen-rich 134 samples (Figure 1). In contrast, the variability of the Ge spectra is weaker and mostly related 135 to variations in the relative intensities of the amide and structural carbonate bands. Thus, the 136 presence of the collagen matrix affects the intense phosphate bands of the diamond ATR 137 spectra of bone samples but the similarities of Ge-ATR spectra suggest that this effect is not 138 associated with a strong modification of the microscopic structure of apatitic crystallites.

139 Based on a systematic analysis of spectral changes related to variations of the 140 macroscopic parameters describing the composite nature of the samples (Suppl. Info.), further 141 insight in the observed variations can be gained from the comparison of experimental and 142 modeled spectra. We stress that our modeling strategy involves a compromise between a 143 minimal variation of model parameters (Table 1) and a reproduction of relative variations 144 observed among the spectra. The bulk apatite properties (Table 2) are modeled using 145 frequencies and force oscillator strengths close to those previously used to model fluorapatite 146 spectra (Aufort et al. 2016). However, significantly higher broadening parameters must be 147 used to account for the lower degree of crystalline order and higher strain of biogenic apatite 148 particles. The v_2 CO₃ bands which can overlap with phosphate-related bands of apatite have 149 also been introduced in the model. The B-type resonance corresponds to substitution of CO_3 for PO₄, presumably occupying a triangular face of the tetrahedral phosphate sites, and has 150

151 been assumed to be isotropic. Its broadening parameter is similar to that used for v_4 and v_1 152 PO_4 bands. The A-type bands correspond to carbonate groups occupying channel sites with an 153 orientation parallel to the c-axis. A comparatively smaller broadening parameter was found to 154 account for the shape of the A-type band, suggesting that the local environment of channel 155 carbonates is less sensitive to variations in the apatite structural order. The stretching CO₃ and 156 amide bands observed at higher frequencies (Figure 1) have not been included in the models 157 as they do not overlap with the apatite-related phosphate bands. A single set of parameters 158 (Table 2) describing the microscopic properties of bulk apatite is used to model the ATR 159 spectra of the whole sample series. Each sample is thus considered as a binary composite 160 defined by three parameters (Table 1): the mineral fraction, the dielectric constant of the other 161 component, and a scale factor. The Ge and diamond ATR spectra are then computed by only 162 varying the refractive index of the ATR crystal.

163 Collagen-poor enamel and EH2-20 bone samples are well-described as an apatite/air 164 composite, with a larger mineral fraction for enamel than for the heavily transformed 165 archeological bone (55 % vs. 45 %). These samples require almost no scaling of theoretical 166 intensities to match the experimental ones (Table 1). When comparing the diamond and Ge 167 spectra, the model reproduces the stronger spectral intensity and the downshift of v_4 and v_3 PO_4 bands observed in the diamond spectrum. Although the maximum of the $v_3 PO_4$ bands 168 was observed at 1026 cm⁻¹ in the Ge spectrum, close to the corresponding transverse optical 169 frequencies (1027 & 1032 cm⁻¹; Table 2), it is shifted to 1014 and 1006 cm⁻¹ in the diamond-170 ATR spectrum of the EH2-20 and modern enamel samples. Similar but smaller ($\sim 6 \text{ cm}^{-1}$) 171 shifts are observed and modeled for the v_4 PO₄ band at ~560 cm⁻¹. 172

In contrast, the spectroscopic properties of collagen-rich modern bone samples are better reproduced considering a binary composite containing 40 vol. % of apatite particles and a second component with a dielectric constant $\varepsilon_h = 2.5$, consistent with the dielectric

properties of a collagen matrix (Wang et al. 1996). Compared with an apatite-air composite 176 177 and for intermediate volume fractions, the increase in dielectric constant of the matrix leads to a prominent asymmetry of the v_4 and v_3 PO₄ bands in the diamond-ATR spectrum but only 178 179 weakly affects the line shape of the Ge-ATR spectrum (Suppl. Info.), which is consistent with 180 the experimental observations. This pertains to the stronger effects due to anomalous 181 dispersion in the case of a reflection on an ATR crystal with lower refractive index such as 182 diamond (Boulet-Audet et al. 2010). The low-frequency asymmetry of the v_4 PO₄ band in the diamond ATR spectrum is particularly sensitive to the dielectric properties of the embedding 183 184 matrix (Figure 1). Note that their overlap with the broad and intense $v_3 PO_4$ band results in an asymmetric line shape of the weak $v_2 CO_3$ and $v_1 PO_4$ bands in the diamond-ATR spectrum, a 185 186 feature reproduced in the modeled spectrum. The spectra of HAR and BZ-PJ1-09 samples 187 with intermediate amounts of collagen are close to those of modern ox bone. The properties of 188 sample HAR, the estimated collagen content of which is ~16 wt. %, are well reproduced using the same mineral fraction and matrix dielectric constant (ε_h) as for the modern bone. 189 For the more transformed sample BZ-PJ1-09, the estimated collagen content (~11 wt. %) of 190 191 which is about half that of a modern bone, a better agreement with the experimental spectra is 192 achieved for a lower value of the dielectric constant of the host matrix (ε_{h} = 1.9), keeping the 193 mineral fraction constant (40 vol. %). The decrease in collagen content is correlated to a less 194 pronounced broadening of the v_3 PO₄ and v_4 PO₄ ATR bands toward lower frequencies, reproduced in the model by lowering the dielectric constant of the host matrix. 195

In agreement with experimental observations, variations of ε_h result in little changes in the line shape of modeled Ge ATR spectra (Suppl. Info.). However, the theoretical spectra of collagen-rich samples overestimate the intensity of their experimental counterparts, which are about four times less intense than those of collagen-poor samples (e.g. for modern ox bone the theoretical intensities require a 0.175 scale factor to match the experimental ones). As discussed in Aufort et al. (2016), this theoretical overestimation is most likely related to sample heterogeneities but further grinding of the collagen-rich samples did not allow recovering the theoretical intensities. Note that a decrease in the mineral fraction modifies the line shape (Suppl. Info.) and cannot account for a simple decrease of the whole spectrum intensity.

206 Summarizing the present results, an effective medium modeling of the powder sample's 207 dielectric properties based on the Bruggeman scheme and using a minimum number of free 208 parameters accounts for the variability of major phosphate bands in the ATR spectrum of 209 bone and tooth enamel samples with various collagen contents. Importantly, the microscopic 210 properties of the apatite are not used as free parameters in the modeling of the different 211 spectra, underlining the key influence of effective medium parameters. This further suggests 212 that the strong ATR-FTIR bands related to phosphate vibrations are significantly affected by 213 electrostatic effects occurring at a mesoscopic scale in the composite samples and much less 214 by molecular-level interactions between the organic and inorganic components. Owing to 215 these electrostatic effects, the experimental ATR-FTIR spectrum of a bone cannot be 216 considered as a simple superposition of the ATR-FTIR spectra of apatite and collagen.

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IMPLICATIONS

The study of bone material using ATR-FTIR spectroscopy is carried out in a broad range of fields, from paleontology, archaeology and forensics to biomedical applications. Our results highlight yet another consequence of the complex composite nature of bone and of the intricate relationship between the closely intertwined organic and inorganic components at the different scales of bone's hierarchical structure. From the preservation assessment of archaeological finds to the conception of hydroxylapatite-based biomaterials for biomedical

225	applications, the interpretation of ATR-FTIR spectra of bone materials should carefully
226	consider mesoscopic scale effects in addition to the atomic scale organization of the sample.
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301 Table 1: Sample source and description, N wt%, and parameters used to define the binary

302 composite for the corresponding modeled ATR spectra shown in Fig. 1.

Sample	Description - origin	Age	Ν	Min. fraction	ϵ_{host}	scale factor
			wt. % ^a	vol. %		
Ox	Ox bone	Modern	3.96 ^a	40	2.5	0.175
HAR	Harsova Tell, Romania	Chalcolithic	2.75 ^a	40	2.5	0.175
BZ-PJ1-09	Bize-Tournal Cave, France	Prehistoric	1.80 ^a	40	1.9	0.33
EH2-20	El Harroura Cave, Morocco	Neolithic	0.16 ^a	45	1.0	1.0
Enamel	Camelidae tooth enamel ^b	Modern	n.d.	55	1.0	0.8

303 ^aLebon et al. (2014), ^bYi et al. 2014

Table 2: Parameters used to model the dielectric tensor of carbonated hydroxylapatite. The Cartesian components of the unixial tensor are defined by a Drude-Lorentz expression:

$$\varepsilon_{\alpha}(\omega) = \varepsilon_{\infty} + \sum_{i} \frac{A_{i,\alpha}^{2}}{(\omega_{i}^{2} - \omega^{2} - i\omega\Gamma_{i})}$$

307 where α refers to the Cartesian axis, $\varepsilon_{\infty} = 2.65$ is the high-frequency dielectric contribution 308 considered here to be real, constant and isotropic, ω_i is the vibrational transverse optical 309 frequency, $A_{i,\alpha}^2$ the effective mode oscillator strength, and Γ_i is the damping parameter of 310 vibrational mode i.

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Symmetry	Mode	$\omega_i (cm^{-1})$	$A_{i,a}(cm^{-1})$	$\Gamma_{i}(cm^{-1})$
A _u	$v_3(PO_4)$	1027	665	40
	$v_2(CO_3)$	874	80	16
	$v_4(PO_4)$	560	300	15
	$v_2(PO_4)$	473	45	15
E_{1u}	$v_3(PO_4)$	1090	270	40
	$v_3(PO_4)$	1032	630	40
	$v_1(PO_4)$	958	110	18
	$v_2(CO_3)$	880	30	5
	$v_2(CO_3)$	874	80	16
	$v_4(PO_4)$	600	215	15
	$v_4(PO_4)$	580	120	15

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Figure 1: Experimental and modelled ATR-FTIR spectra of modern ox bones (A,B), archaeological animal bones with intermediate levels of collagen preservation (C,D); low collagen containing archaeological bone and tooth enamel (E,F). Spectra A, C and E have been recorded using a Ge ATR crystal. Spectra B, D and F have been recorded using a diamond ATR crystal. Experimental ATR-FTIR spectra have been shifted vertically for comparison.

