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### **Evaluating the physicochemical conditions for gold occurrences in pyrite**

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### **Abstract**

While noble metals often occur as minor components in host minerals in various ore deposits, little theoretical assessment exists to predict the occurrence of these metals. Here, we probe the fundamental controls responsible for the occurrence of trace elements in host minerals through first-principles calculations. We apply the theoretical model to understanding the debated issues concerning the occurrence of gold (Au) in pyrite, in which the valence of Au is ascribed to either

21 positive or negative values. Our results indicate that (1) both positive and negative valent Au may  
22 occur in pyrite and (2) higher sulfur fugacity and lower temperature lead to more Au<sup>+</sup> occupying Fe sites  
23 in pyrite. These findings suggest that chemical states and speciation of the Au in host pyrite are  
24 ultimately controlled by temperature and sulfur fugacity, providing insight into the formation conditions  
25 of ore deposits and facilitating strategy design for beneficiation.

26 **Keywords:** gold-bearing pyrite, physicochemical condition dependence, gold occurrence,  
27 density functional theory, thermodynamics, sulfur fugacity

## 28 **Introduction**

29 Gold (Au) not only is of great importance for financial reserves but also, like other trace  
30 elements including indium, germanium, and tellurium, etc., is required in unprecedented quantities in  
31 industrial applications (Gulley et al. 2018). Due to low abundance (parts to sub-parts per million), Au,  
32 as well as other trace elements, frequently disperse in host minerals as solute constituents of solid  
33 solutions (Zhai et al. 2019). The physicochemical state and speciation that these metals assume in the  
34 corresponding hosts often carry critical information indicative of the ore formation processes (Xing et al.  
35 2019; Du et al. 2020), in addition to being a key gauge for downstream metal recovery (Wang et al.  
36 2008). Despite such general understanding, strategies rationalizing the occurrence of noble metals in  
37 host minerals remain largely empirical. The lack of a theoretical basis to investigate the occurrence of  
38 valuable metals impedes the interpretation of genetic information of metallogenic processes and  
39 high-efficient recovery of these metals.

40 Pyrite (FeS<sub>2</sub>) occurs in almost all metalliferous deposits and is a predominant gold-bearing

41 mineral in many giant gold deposits such as the Carlin-type (Reich et al. 2005). Although the  
42 occurrence of Au in pyrite has been of wide concern, the chemical state and site occupation of Au are  
43 still in debate. Previous studies, based on spectroscopy (Li et al. 1995; Simon et al. 1999; Cabri et al.  
44 2000; Trigub et al. 2017; Merkulova et al. 2019; Pokrovski et al. 2019) or charge balance consideration  
45 (Chouinard et al. 2005), focused primarily on samples from specific deposits while ignoring the  
46 substitutional physics of pyrite and the physicochemical conditions under which it formed. This results  
47 in a historical paradox of the oxidation state of Au in pyrite, i.e., either positive or negative oxidation  
48 state. The positive +1 and +3 valences are interpreted from X-ray absorption near edge structure  
49 (XANES) spectroscopy and charge balance consideration (Simon et al. 1999; Cabri et al. 2000; Trigub et  
50 al. 2017; Merkulova et al. 2019; Pokrovski et al. 2019), while the negative oxidation state of gold is  
51 confirmed by X-ray photoelectron spectroscopy (XPS) and Mössbauer spectroscopy (Li et al. 1995,  
52 2002). Therefore, we hypothesize that the inconsistency of the interpretation of gold occurrence in  
53 pyrite may be caused by variations of site occupations and formation conditions.

54       Herein, we propose a theoretical framework, which links the occurrence of Au in pyrite structure  
55 to its metallogenic environments, to test the hypothesis. For a host mineral, the incorporated elements  
56 are *de facto* impurities whose site occupancy in host minerals can be treated as a distribution of the  
57 impurity elements in a crystal field, which can be described by the Maxwell-Boltzmann distribution.  
58 In such a system, the occupation probability of the impurity elements at different lattice sites of host  
59 minerals hinges upon both the incorporation-energy-level and the temperature. The incorporation  
60 energy of impurity elements into a host mineral is commonly controlled by the chemical potential of the

61 elements in the host lattice, which in turn is related to the environments where the host crystals grew.  
62 Hence, the occupation probability and electron structure of metals in their host minerals can be precisely  
63 assessed, with the employment of modern first-principles techniques, under various growth  
64 environments.

65 We quantify the probability of Au occurrence in pyrite by evaluating the Boltzmann factor  
66 governing the atomic distribution in the lattice of the host mineral. The Boltzmann factor ( $w$ ) is  
67 defined as:

$$68 \quad w = \frac{p_i}{p_0} = e^{\frac{-\varepsilon_i}{kT}}, \quad (1)$$

69 where  $p_0$  and  $p_i$  denote the probabilities of Au occurrence locating at the first (i.e., the ground state) and  
70  $i^{\text{th}}$  incorporation-energy-level ( $\varepsilon_i$ ), respectively;  $k$  and  $T$  are the Boltzmann constant and temperature,  
71 respectively. The sum of the occurrence probabilities,  $p_i$ , equals to 1, i.e.,

$$72 \quad \Sigma p_i = 1. \quad (2)$$

73 By relating the occurrence probability to the corresponding electronic structure, we find that the stability  
74 field can be associated with either positive or negative valence state in a wide range of physicochemical  
75 environments, suggesting the site occupation and valence state of Au in pyrite can serve as a genetic  
76 proxy for the formation of pyritic Au deposits.

## 77 **Methods**

### 78 **Density functional theory calculations**

79 The density functional theory (DFT) calculations were performed using the Vienna Ab initio  
80 Simulation Package (VASP.5.4.4) (Kresse and Furthmuller 1996) along with the projector augmented

81 wave (PAW) method (Kresse and Joubert 1999). Standard PAW pseudopotentials (Fe:  $d^7s^1$ , S:  $s^2p^4$ , Au:  
82  $s^1d^{10}$ , and Cs:  $5s^5p^6s$ ) were used. The generalized gradient approximation (GGA) formulation by  
83 Perdew, Burke, and Enzerhoff (PBE) functional (Perdew et al. 1996) was used to describe the  
84 exchange-correlation interaction among electrons, and the energy of 400 eV was used as a cutoff value  
85 for the plane wave basis. Total energies were converged at  $10^{-6}$  eV. The simplified Hubbard  $U$   
86 correction (GGA+ $U$ ) method (Dudarev et al. 1998) was employed to describe the 3d electrons of Fe  
87 element. The parameter  $U-J= 1.6$  eV was employed, and the feasibility of such a parameter was  
88 confirmed by previous studies on pyritic systems (Zhang et al. 2012; Xian et al. 2019b, 2019a, 2021).  
89 The calculated lattice constant of 5.423 Å is in good agreement with both the experimental data and  
90 previous calculations (Table 1). All the structures employed in this study were relaxed until forces  
91 became smaller than 0.02 eV/Å. The ab initio molecular dynamics (AIMD) calculations were  
92 performed in NVT ensemble with Nosé-Hoover chain thermostats (Martyna et al. 1992) under 298 K.  
93 The time step of the AIMD runs was 1 fs. The visualization and analyses of the AIMD trajectories  
94 were performed by the package of VASPKIT (Wang et al. 2021) and VMD (Humphrey et al. 1996).  
95 The theoretical valence state was treated by the Bader charge-division scheme based on the electronic  
96 charge density (Bader 1990), which was embedded in the code developed by the Henkelman Group  
97 (Henkelman et al. 2006; Tang et al. 2009). The qualitative oxidation state of Au can be identified from  
98 the Bader charge, although the Bader charge cannot directly represent the chemical oxidation state.

99 The pyrites with Au located at various sites were modeled by  $2 \times 2 \times 2$  supercells with 96 (for  
100 the substitution cases) or 97 (for the interstitial incorporation model) atoms. For the interstitial model,

101 the site was randomly selected from an interstitial site in the pyrite structure, as all the interstitial sites  
102 are symmetrically equivalent. The configuration files of reference materials (including AuS, CsAu,  
103 and elemental Au, Fe, and S) were downloaded from the Materials Project website (Jain et al. 2013).  
104 Structure relaxation of these configurations was modeled by the unit cells. The calculated lattice  
105 parameters of the reference materials are in good agreement with experimental or previous calculated  
106 values (Table 1).

### 107 **Calculations of the incorporation-energy-level and the occurrence probability**

108 The Au incorporation energies ( $E_i$ ) in these pyritic models were calculated through the equation:

$$109 \quad E_i = E_{gpy} - E_g - E_{py} + \mu_{atom}^{sub}, \quad (3)$$

110 where  $E_{gpy}$ ,  $E_g$ , and  $E_{py}$  denote the DFT total energies of the gold-bearing pyrite, one Au atom in gold  
111 metal, and the pure pyrite, respectively.  $\mu_{atom}^{sub}$  denotes the chemical potential of specific substitutional  
112 atom (i.e., the chemical potential of Fe ( $\mu_{Fe}^{py}$ ) and S ( $\mu_S^{py}$ ) of pyrite in this study). Mathematically,  $\mu_{Fe}^{py}$   
113 and  $\mu_S^{py}$  obey the constraint

$$114 \quad \mu_{Fe}^{py} + 2\mu_S^{py} = \mu_{FeS_2}^{py}, \quad (4)$$

115 where  $\mu_{FeS_2}^{py}$  denotes the chemical potential of one FeS<sub>2</sub> formula unit in pyrite. For simulations  
116 relevant to various sulfur environments,  $\mu_S^{py}$  refers to the alpha S elementary crystal phase ( $\alpha$ -S).  
117 Thus, the sulfur environments were characterized by the chemical potential difference,  $\Delta\mu_S$ .

$$118 \quad \Delta\mu_S = \mu_S^{py} - \mu_S^{\alpha-S}. \quad (5)$$

119 The relative sulfur chemical potential  $\Delta\mu_S$  is linked to sulfur fugacity through the equation:

$$120 \quad \Delta\mu_S = RT \ln f_{S_2} \quad (6)$$

121 where the  $R$  and  $fS_2$  denote gas constant and  $S_2$  vapor fugacity, respectively.

122 The incorporation-energy-levels,  $\varepsilon_i$ , were calculated from the difference between the  
123 incorporation energies of any incorporation type and that of the ground state ( $E_0$ , the lowest  
124 incorporation energy).

$$125 \quad \varepsilon_i = E_i - E_0. \quad (7)$$

126 If ground state is estimated to be the substitution of Au for the Fe site, the  
127 incorporation-energy-levels of the Fe, S, and interstitial sites are  $\varepsilon_0 = \varepsilon_{Fe} = 0$  eV,  $\varepsilon_S$ , and  $\varepsilon_{int.}$ , respectively.

128 Combining Eqs. (1) and (2), one can get an equation set:

$$129 \quad \begin{cases} p_S = P_{Fe} e^{-\frac{\varepsilon_S}{kT}} \\ p_{int.} = P_{Fe} e^{-\frac{\varepsilon_{int.}}{kT}} \\ p_S + P_{Fe} + p_{int.} = 1 \end{cases}, \quad (8)$$

130 from which the occurrence probability of the substitution of Au for Fe ( $p_{Fe}$ ), S ( $p_S$ ), and the interstitial  
131 sites ( $p_{int.}$ ) can be derived. In a given condition of  $\Delta\mu_S$  and  $T$ , the occurrence probabilities,  $p_i$ , can be  
132 calculated by resolving the equation set (8).

### 133 **Simulations of HAADF-STEM images and EELS spectra**

134 The fast simulation of high-angle annular dark field transmission electron microscope  
135 (HAADF-STEM) images was performed on a program based on kinematic-scattering method (He et al.  
136 2015) using cluster models containing more than 1000 atoms along the [001] direction. We employed  
137 a Gaussian width (i.e., probe size of a real transmission electron microscope (TEM)) of 0.7 Å for  
138 simulations, because such a probe size can be easily reached in a modern TEM with spherical aberration  
139 corrector. The simulations of electron energy loss spectroscopy (EELS) were performed using the

140 Wien2K code (Blaha et al. 2020) with operation voltage at 200 kV. The input configurations for  
141 Wien2K calculations were relaxed from the VASP code. In Wien2K calculations, the PBE functional  
142 and a default  $R_{MT}$  and  $RK_{max}$  were employed, which were sufficient to reach the same convergence. A  
143 0.5 eV of broaden width was employed for all the EELS spectra.

## 144 **Results**

### 145 **Site occupation and valence of gold in pyrite**

146 The structure of pyrite consists of  $[FeS_6]$  octahedra connected by S dimers through corner  
147 sharing (Fig. 1a). There are three possible pyrite sites for Au to reside in, i.e., the Fe, the S, and the  
148 interstitial sites (Figs. 1b-1d). The first two are sites where one Au can substitute for one Fe or S atom,  
149 while the third one requires extra space for Au to be incorporated into the host pyrite lattice. The Fe  
150 site has the first coordination shell composed of six S ions (Fig. 1b) while the coordination shell of S  
151 consists of three Fe and one S ions (Fig. 1c). In the cases that the interstitial sites are occupied by Au,  
152 the interstitial Au atoms are usually coordinated with four Fe and six S ions (Fig. 1d).

153 We constructed separate configurations for the above-mentioned triple-occupation models. The  
154 optimized structural properties show that incorporation of Au can increase the lattice size of pyrite from  
155 159.31 to 159.93, 160.54, and 163.49  $\text{\AA}^3$  per unit cell with Au incorporation in the Fe, S and interstitial  
156 sites, respectively. The increase in cell volume may lead to lattice strain, and thereby decrease the  
157 stability of the system (Han et al. 2019). Thus, pyrite with Au at the interstitial sites could be the most  
158 unstable of the three configurations, consistent with the estimated occurrence probability discussed  
159 below.



160 We analyzed the valence state of Au in pyrite using the Bader charge (Bader 1990). The results  
161 show (Figs. 1b-1d) that gold can be negative (-0.09), positive (+0.30), and zero valence when it occupies  
162 the S, Fe, and interstitial sites, respectively. These results are consistent with reference materials with  
163 known valence state of Au, including CsAu (Au valence negative, Bader charge = -0.76), AuS (Au  
164 valence positive, Bader charge = +0.24), and metallic gold (Au valence 0, Bader charge = 0).

165 The structural stability of the Au-bearing pyrite was further evaluated via AIMD at room  
166 temperature (298 K). The analyses of AIMD trajectories show that the substitution models of  
167 Au-bearing pyrite are stable while the one with interstitial sites occupied by Au is unstable. The Au  
168 atoms are localized at either Fe or S sites during the 10 ps of AIMD. Meanwhile, the interstitial Au  
169 atom diffuses in the lattice and thereby exchanges with one Fe atom in pyrite in 4 ps of AIMD.  
170 Ultimately, the Au atom occupies the Fe site while the exchanged Fe atom relocates to the interstitial site  
171 ([see details in supplementary video1](#)).

172 We calculated the radial distribution functions (RDF) and coordination numbers (CN), using  
173 AIMD, to characterize local structures of Au-bearing pyrites. The results (Figs. 1e-1g) shows that Au  
174 incorporation can only influence the bond length as the CN of Au at the Fe and S sites remain the same  
175 with the substituted Fe and S ions. Since the interstitial model is unstable and readily evolves to the Fe  
176 substitution configuration, we only computed the RDF and CN for the substitution models after reaching  
177 equilibrium of the AIMD runs. When Au occupies the Fe or S sites in pyrite, the pair separation  
178 distances of the first coordination shell (i.e., the bond length) of Au-S are 2.46 and 2.26 Å, respectively,  
179 longer than those of Fe-S (2.26 Å) and S-S (2.16 Å) in pure pyrite. The increase in the bond length

180 (from Fe-S and S-S to Au-S) presumably is responsible for the increase in the lattice parameter as  
181 described above.

182 We also simulated measurable structural properties of gold-bearing pyrite. The occurrence of  
183 gold in the pyrite lattice sites (Figs. 2a-2c) can be readily visualized by simulated HAADF-STEM  
184 images. The intensity distribution of the HAADF-STEM image of pure pyrite (Fig. 2a) possesses a  
185 uniform nature with small variation induced by thermal fluctuation. The intensity distributions of  
186 HAADF-STEM images of gold-bearing pyrite (Figs. 2b-2c) show sharp peaks where a gold atom occurs.  
187 The peaks are symmetrical and asymmetrical when Au atoms locate at Fe and S sites, respectively. The  
188 intensity distributions of HAADF-STEM (profiles in Figs. 2a-2c) are proportional to  $Z^n$  which describes  
189 the departure of scattering cross-section to the ideal Rutherford scattering, where  $Z$  is the atomic number  
190 of atoms involved and  $n=2$  is the power exponent (He et al. 2015).

191 Comparable to the three reference materials having various valence states of Au (Fig. 2e), the  
192 simulated Au  $M_{4,5}$  EELS spectra of gold-bearing pyrites with Au located at various lattice sites (Fig. 2d)  
193 are distinguishable from each other. When Au substitutes for S in pyrite, three notable characteristic  
194 peaks occur in the Au  $M_4$  EELS spectrum, similar to those of CsAu with gold in the -1 valence. The  
195 interstitial gold in pyrite possesses similar features of gold metal (one sharp with two shoulder peaks at  
196 the low energy side), suggesting the zero valence of the Au atom. However, the Au  $M_4$  EELS spectrum  
197 of the gold-bearing pyrite with gold located at Fe sites is slightly different from that of AuS although the  
198 Bader charges of both materials are positive.

199 The presence of Au in the Fe and S sites in pyrite can also shift the vibrational property of the

200 host mineral. The calculated vibrational spectra (Fig. 3), based on the AIMD trajectories, show that the  
201 spectra of gold substitutional pyrites shifted to lower wavenumber compared to that of pure pyrite,  
202 indicating that the incorporation of gold weakened the vibrations. When Au substitute for either Fe or  
203 S in pyrite, the vibrations at the wavenumber greater than  $380\text{ cm}^{-1}$  and smaller than  $320\text{ cm}^{-1}$  weakened,  
204 while those ranged from  $320$  to  $380\text{ cm}^{-1}$  strengthened, indicating that the substitution of Au into pyrite  
205 can vastly shift the lattice dynamics.

206

### 207 **Physicochemically controlled occurrence of gold in pyrite**

208 The calculated incorporation energy (Fig. 4) based on Eq. (3) shows that the ground state (i.e.,  $\epsilon_0$ ,  
209 the lowest energy) of the incorporation-energy-levels varies at different ranges of sulfur chemical  
210 potential. The ground state was estimated to be the S sites when the sulfur chemical potential is less  
211 than  $-0.566\text{ eV}$  relative to the  $\alpha$ -S crystal. In contrast, it changes into the Fe sites when the system  
212 possesses higher sulfur chemical potential ( $> -0.566\text{ eV}$ ). Combining Eqs. (1)–(8), the occurrence  
213 probability of Au in pyrite is linked to sulfur fugacity. The results (Fig. 5) show that gold atoms mainly  
214 occupy either Fe or S sites under various physicochemical conditions. The sum of the occurrence  
215 probabilities of Au sitting at Fe and S sites under certain  $fS_2$  and  $T$  is very close to 1, while the highest  
216 occurrence probability of Au located at the interstitial sites is smaller than  $10^{-49}$ . Gold atoms prefer to  
217 occupy Fe sites under relatively high  $fS_2$  and low  $T$  while the preferred occupying shifts to S sites under  
218 inverse conditions.

219

## **Discussion**

## 220 **Chemical speciation of “invisible” gold in pyrite.**

221 The results presented above may offer important clues through which the inconsistency in the  
222 existing interpretations of Au speciation in pyritic structures can be readily understood. The most  
223 widely accepted oxidation state of Au in pyritic structures is +1 from XANES results (Simon et al. 1999;  
224 Merkulova et al. 2019). This may be caused by that most pyrite formed under sulfur-rich conditions,  
225 and Au prefers to occupy the Fe sites in pyrite (Fig. 5). Recent results from HR-XANES show that Au  
226 locates at Fe sites and the Au coordinates with  $5.7 \pm 0.7$  S atoms at 2.41 Å (Merkulova et al. 2019),  
227 consistent with our theoretical model for Au substitution of Fe in this study (Fig. 1f). The negative  
228 valence of Au in both pyrite and arsenopyrite was first confirmed by XPS (Li et al. 1995) and further  
229 supported by  $^{197}\text{Au}$  Mössbauer spectra (Li et al. 2002). This interpretation is consistent with our  
230 theoretical model for Au substitution of S in this study (Fig. 1g), which confirms that the Au substituting  
231 S in pyrite possesses negative valence.

## 232 **Implications for gold recovery and interpretation of formation conditions of ore deposits**

233 The relation between  $f\text{S}_2$  and the occurrence probability of Au in the pyrite lattice (Fig. 5),  
234 presented in this study, is consistent with previous experimental data showing that Au occupies the Fe  
235 sites when the  $\log f\text{S}_2$  varies from -3 to -1 bar under 450°C and 1 kbar (Trigub et al. 2017). This is  
236 critical for gold recovery, as the recovery of refractory gold ore is mainly restrained by the occurrence of  
237 gold in pyrite. Most of Au in refractory ore occurs as nanoparticles or lattice-bound form. Since the  
238 lattice-bound form is stable under wide physicochemical conditions with active interstitial form (Fig. 5),  
239 the interstitial Au could be considered as the initial form of Au nanoparticles formation when the gold

240 concentration surpasses its solubility in pyrite (Reich et al. 2005). Therefore, the recovery of the gold  
241 nanoparticles from pyrite may be implemented without destroying pyrite structure, while the recovery of  
242 lattice-bound gold involves the ruin (e.g., oxidation) of pyrite lattice.

243 The occurrence of Au in the pyrite lattice, experimentally measurable through either  
244 spectroscopy (such as XPS, XANES, in-situ EELS, and Mössbauer) or microscopy such as  
245 HAADF-STEM (Fig. 2), can be used to directly interpret the physicochemical conditions of ore deposits  
246 formation based on the relation shown in Fig. 5. Pyrite with dominant substitution of Au for Fe should  
247 be formed under lower  $T$  and higher  $fS_2$  in comparison to those with Au mainly occupying S sites. For  
248 instance, the refractory gold ore (in which gold occurs as +1 valence and occupies Fe sites) collected  
249 from a sandy mine tailing in Colombia, South America (Merkulova et al. 2019) may form under either  
250 relative lower  $T$  or higher  $fS_2$  than those (in which gold occurs as negative valence and occupies S sites)  
251 from Jinya gold deposits, Guangxi province, China (Li et al. 1995, 2002). Further, the correlation  
252 between Fe/S and Au concentrations in pyrite can also be used to indirectly substantiate the formation  
253 conditions of ore deposits as Au can occupy either Fe or S sites. From a Spearman rank correlation  
254 analysis (Xie et al. 2018), Au is found to negatively correlate with Fe, while no obvious correlation with  
255 S for Carlin-type ore pyrites from both Nevada, USA and Guizhou, China, suggesting that Au mainly  
256 occupies the Fe sites in pyrite. We thus infer that the Carlin-type gold deposits formed under relative  
257 low temperature and high  $fS_2$ . The stronger correlation of Au to Fe (-0.544) in Nevada deposits relative  
258 to that (-0.466) in Guizhou implies that the latter formed at higher  $T$ . This deduction is supported by  
259 previous experimental observations (S.CLINÉ et al. 2013; Xie et al. 2018) that show the trapping

260 temperature of ore fluid in Guizhou is about 220° to 345°C, ~40° to 100°C higher than that in Nevada.

261

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370

371 **Figure captions**

372 **Figure 1.** Structural models and partial radial distribution function (RDF,  $g(r)$ ) of gold-bearing pyrite.

373 A-D: Structural models of bulk pyrite (**a**), Au substitution of Fe (**b**) and S (**c**) sites, and Au occupation at  
374 interstitial site (**d**) of pyrite; the yellow, brown, and orange spheres denote sulfur, iron, and gold atoms,  
375 respectively. (**e**)–(**g**): Partial RDF of pure pyrite (**e**) and substitutional (**f** and **g**) Au-bearing pyrite.

376 **Figure 2.** Predicted structural properties of gold-bearing pyrite. (**a**)–(**c**) denote the atomic images of  
377 HAADF-STEM for pure pyrite and pyrite with Au occupation at Fe and S sites, respectively. The scale  
378 bars in (**a**)–(**c**) represent 5 nm. (**d**) and (**e**) denote EELS spectra of Au  $M_{4,5}$  edge in gold-bearing pyrite  
379 and reference materials, respectively.

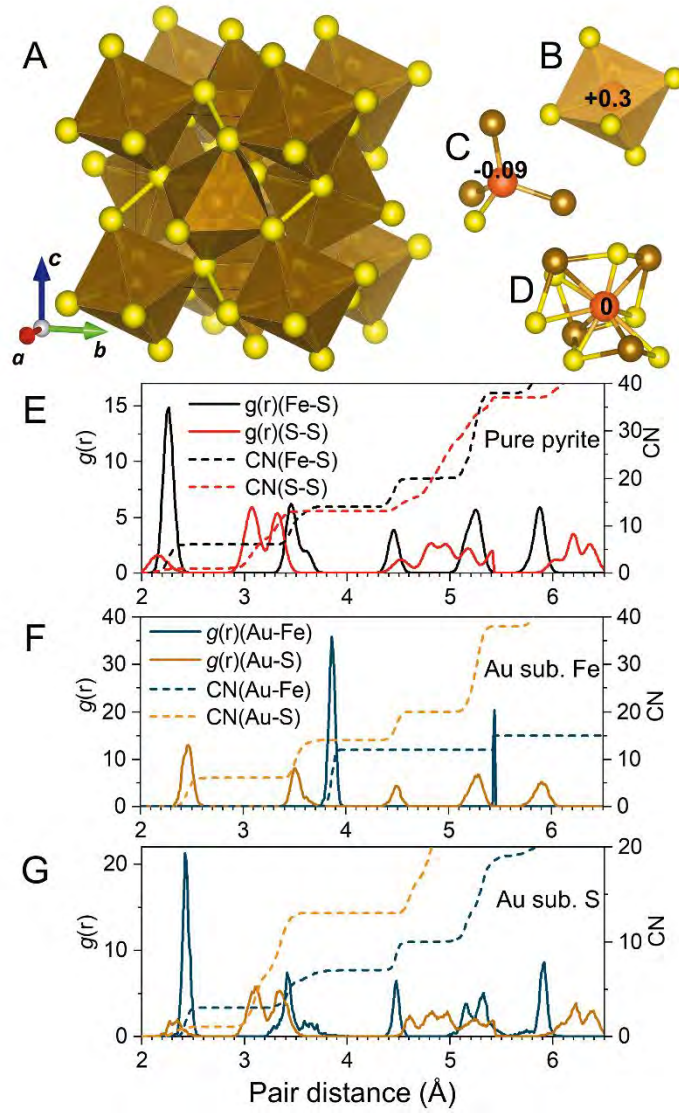
380 **Figure 3.** Vibrational density of state of substitutional gold-bearing pyrite and pure pyrite.

381 **Figure 4.** Incorporation energy variation versus sulfur chemical potential change ( $\Delta\mu_S$ ) of gold-bearing  
382 pyrite.

383 **Figure 5.** Sulfur fugacity and temperature dependent occurrence probability of gold in pyrite. (**a**)–(**c**)  
384 denote the Au occupation at Fe, S and interstitial sites in pyrite, respectively. The dash lines denote the  
385 contour line of sulfur fugacity ( $fS_2$ ) and the numbers on the dash lines are the corresponding  $\log fS_2$  in  
386 bar. The color bars denote the occurrence probability of gold in the corresponding site.

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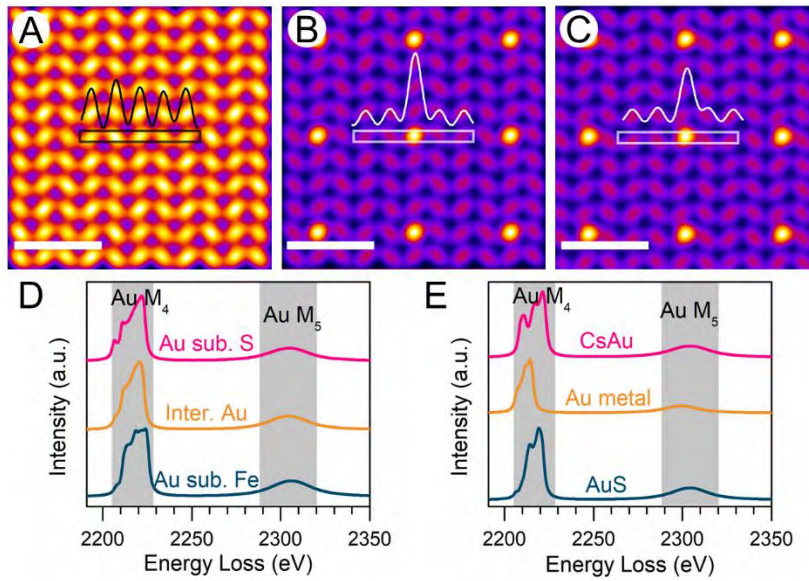
388 Figure 1



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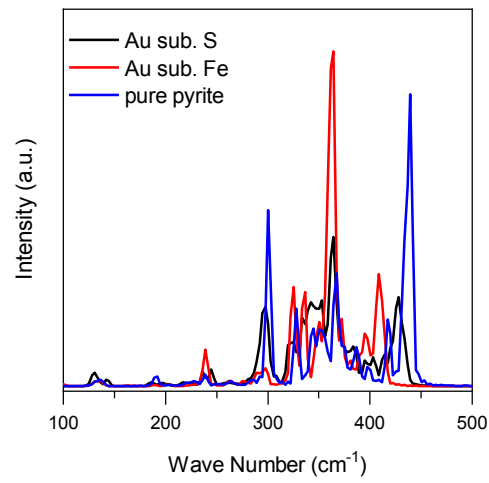
391 Figure 2



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394 Figure 3

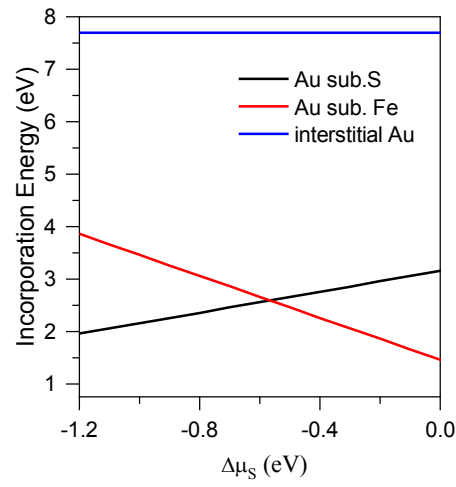


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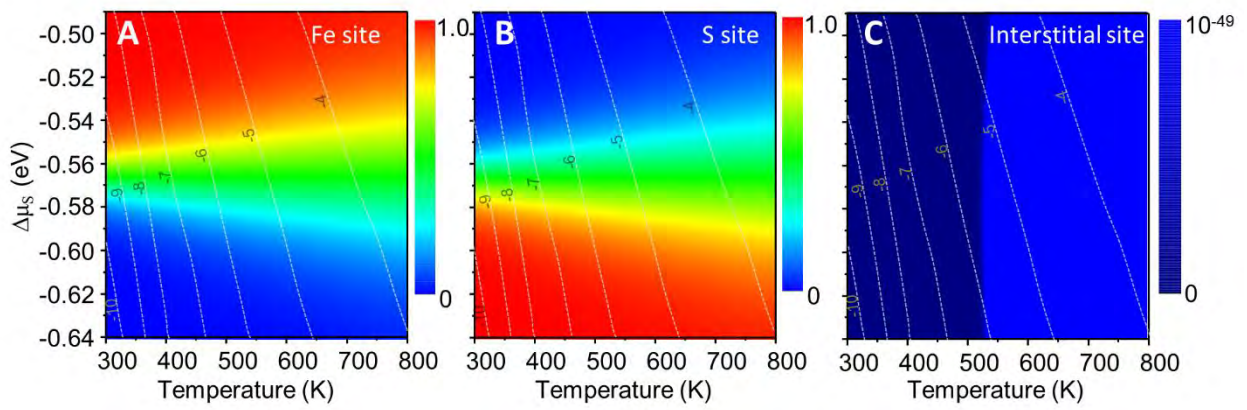
398 Figure 4



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401 Figure 5



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404 **Table 1.** Comparison of calculated and experimental lattice parameters for minerals and structures  
 405 employed in this study.

Mineral and structures	Lattice parameters		Data source
	$a, b, c / \text{\AA}$	$\alpha, \beta, \gamma / ^\circ$	
Experimental pyrite	5.418, 5.418, 5.418	90, 90, 90	Benbattouche et al. (1989)
Theoretical pyrite	5.423, 5.423, 5.423	90, 90, 90	This study
Theoretical pyrite with Au sub. Fe	10.8880, 10.8880, 10.8880	90, 90, 90	This study
Theoretical pyrite with Au sub. S	10.8701, 10.8701, 10.8701	90, 90, 90	This study
Theoretical pyrite with interstitial Au	10.9012, 10.9326, 10.9732	90.1, 90.3, 90.1	This study
Experimental metallic gold	4.072, 4.072, 4.072	90, 90, 90	Suh et al. (1988)
Theoretical metallic gold	4.0783, 4.0783, 4.0783	90, 90, 90	This study
Experimental metallic iron	2.866, 2.866, 2.866	90, 90, 90	Wilburn & Bassett (1978)
Theoretical metallic iron	2.8664, 2.8664, 2.8664	90, 90, 90	This study
Experimental alpha-S8	10.4646, 12.8660, 24.4860	90, 90, 90	Rettig & Trotter (1987)
Theoretical alpha-S8	10.4646, 12.8660, 24.4860	90, 90, 90	This study
Referred theoretical AuS	3.9331, 3.9331, 6.8182	90, 90, 120	Material Project-984598_AuS
Theoretical AuS	3.8972, 3.8972, 6.8257	90, 90, 120	This study
Referred theoretical CsAu	4.3811, 4.3811, 4.3811	90, 90, 90	Material Project-2667_CsAu
Theoretical CsAu	4.3811, 4.3811, 4.3811	90, 90, 90	This study