1 Revision 1

2 Precise determination of the effect of temperature on the density of solid and liquid

3 iron, nickel, and tin

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9 Abstract

Density and thermal expansion coefficient of metals are one of the most 10 fundamental characteristics to describe the equation of state. Especially for liquid metals, 11 12the reported data for density and thermal expansion coefficient vary in the literature, even at ambient pressure. To determine density of solid and liquid metals precisely at 13high temperatures and ambient pressure, we newly developed a high-temperature 14furnace. The densities of solid Sn, Ni, and Fe were determined from sample image with 15uncertainty of 0.11%-0.7% in the temperature range of 285-1803 K with increments of 16171-10 K. The density of solid Sn decreased linearly with increasing temperature up to 493 K, and the decrease became drastic above that to the melting temperature (T_m) . By 18

19	contrast, for solid Ni and Fe, the densities decreased linearly with increasing					
20	temperature up to the T_m (1728 and 1813 K) without any drastic density drop near T_m .					
21	This suggests that Ni and Fe do not exhibit the "premelting effect".					
22	For liquid, the density of liquid Fe was determined with uncertainty of 0.4%-					
23	0.7% in the range of 1818–1998 K with temperature increments of 5 K. The obtained					
24	thermal expansion coefficient (α) of liquid Fe was well approximated as either a					
25	constant value of $\alpha = 2.42(1) \times 10^{-4} \text{ K}^{-1}$ or a linear function of temperature (T); $\alpha =$					
26	$1.37(10) \times 10^{-3} - [6.0(6) \times 10^{-7}]T [\text{K}^{-1}]$ up to at least 2000 K.					
27						
28	Keywords: Density, Thermal expansion, Metal, Liquid					
29						
30	INTRODUCTION					
31	Density is one of the most fundamental physical properties for understanding					
32	planetary interiors. To estimate the structure and composition of a planetary core, the					
33	equation of state (EOS) of Fe alloys is an indispensable information, along with					
34	planetary exploration data. The density of Fe alloys at ambient pressure is one of the					
35	most fundamental characteristics used to describe the EOS, as well as its elastic					
36	properties (bulk modulus, K, its pressure derivative $[dK_T/dP]$, and the thermal expansion					

37	coefficient, α). Although the density and thermal expansion of the core constituting
38	metals (Fe and Ni) have been reported, they have not been determined precisely near the
39	melting temperature (T _m). An unusual drastic decrease in the elastic wave velocities
40	near T _m , the so-called "premelting effect," has been reported for some solid metals (Sn
41	and Fe) based both on experiments and theoretical calculations. It is reported to occur at
42	$T/T_m > 0.99$ for Sn at ambient pressure (Nadal and Le Poac, 2003) and at $T/T_m > 0.96$
43	for Fe at 360 GPa (Martorell et al., 2013a). If the premelting effect occurs in the elastic
44	wave velocity, it could also occur in density, which is closely related to the elastic wave
45	velocity, and thermal expansion. This could significantly influence the EOS and elastic
46	properties of the core. However, the density behaviors of Sn and Fe near $T_{\rm m}$ have not
47	been precisely known.
48	For liquid density, the density (ρ) of liquid Fe at ambient pressure has been
49	reported by many authors (e.g., Luca, 1960; 1972; Kirshenbaum and Cahil, 1962;
50	Frohberg and Weber, 1964; Brillo and Egry, 2004; Assael et al., 2006; Kobatake and
51	Brillo, 2013; Le Maux et al., 2019). However, the reported ρ values have some

significant deviation ($\Delta \rho \sim 0.2 \ g/cm^3 \sim 2.8\%$) and the thermal expansion coefficient has not been precisely constrained yet.

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In this study, density measurements of solid Sn, Ni, Fe, and liquid Fe were

55	performed using a newly developed high-temperature furnace with observation
56	windows. The effect of temperature on the density of solid Sn, Ni, Fe, and liquid Fe was
57	precisely investigated at ambient pressure.
58	
59	METHODS
60	Experimental methods
61	The samples were metal rods of Sn (99.9% in purity, Nilaco Corp.), Ni (99.5%
62	in purity, Nilaco Corp.), and Fe (99.9% in purity, Nilaco Corp.). The diameter and
63	height of the sample were 2.99-3.00 and 2.09-3.00 mm, respectively. The top and
64	bottom interfaces of the samples were mirror-polished.
65	The experiments were conducted using a high-temperature gas furnace
66	(Nagano Co Ltd., Tokyo) installed at Osaka University (Fig. 1a). Heating was
67	performed using a composite resistance heater composed of carbon
68	fiber-reinforced carbon composite (C/C). The C/C heater is superior to generate high
69	temperatures (~2173 K) stably under vacuum or inert gas conditions, and has a small
70	thermal expansion coefficient, which enables rapid heating (~1500 K/min). The C/C
71	heater plate is located around the sample stage and has a slit for sample observation.
72	The sample is set on a sintered alumina plate located at the center of the chamber (Fig.

73	1b). During the experiment, the temperature was monitored using a K-type
74	thermocouple (TC). The TC junction was located at same height as the sample and 18.6
75	mm away from the sample center horizontally. The $T_{\rm m}$ values of Ni and Fe were
76	determined to be 1728 and 1812 K (TC temperature), respectively, based on the shape
77	change of the sample. These results are in good agreement with the literature. Therefore,
78	the temperature at the TC junction was identical to that at the sample center.
79	Temperature was controlled using a proportional integral differential controller. The
80	heating rates were in the range of 3-25 K/min. The measurement was performed under
81	an argon atmosphere at 0.05 MPa (flow rate: 0.8–1.0 l/min).
82	The chamber had two optical windows for observation of the sample in the
83	horizontal direction. For the density measurements, a charge-coupled device (CCD)
84	(1296(H)x966(V), 12 bit, acA1300-30um; Basler AG) or complementary metal-oxide
85	semiconductor (CMOS) camera (2048(H)x1536(V), 12 bit, BU302MG; Toshiba Teli
86	corp.) was set on one side for optical observation, and a backlight (Fiber optic
87	illuminator) was set on the other (Fig. 1a). The image magnification was optimized
88	using a telecentric zoom lens (x0.5-1.3) (VS Technology Corp.). The pixel size of the
89	image was 3.9–4.5 μ m/pixel depending on the magnification. Details of the pixel

exposure time of each image was in the range of 1–80 ms. Images of the sample were
taken with 1–10 K steps and analyzed using the method described in next section.

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94 Analytical methods

The density of the sample was obtained from the sample volume at each temperature and initial mass. The sample volume was measured from the 2-dimensional (2D) sample image assuming rotational (cylindrical) symmetry. Image analysis was carried out using ImageJ software (https://imagej.nih.gov/ij/). The sample image was binarized, and brightness profiles of the image were obtained to identify the boundary between the sample and surrounding part.

101 For solid samples, the diameter and height were measured from the sample 102 brightness profiles (Figs. 2a and 2b). To obtain a clear image of sample bottom boundary which is required for precise determination of the sample height, a small 103 104 alumina disk is placed underneath the sample (Fig. 2a). In the run AK12 that the small alumina disk was not used, height error of AK12 is much (one-order) larger than that of 105 other runs in which the small disk was used (Supplemental¹ Table S4). The validity of 106107 the cylindrical symmetry of the sample was evaluated from images observed from different directions in separate measurements. The sample was put on the rotation stage 108

at room temperature and the image was observed at rotation angles of 0°, 45°, and 90°. 109 For 0–90° rotation, the variations were $(1.0-2.0)x10^{-4}$ cm in diameter and $(1.0-3.0)x10^{-4}$ 110 cm in height, which corresponds to a volume variation of $\Delta V = (1.4-4.6) \times 10^{-5} \text{ cm}^3$ and a 111 density variation of $\Delta \rho = (8.4-20.5) \times 10^{-3} \text{ g/cm}^3$ (0.11~0.28%). Therefore, these density 112variations derived from rotation were small, and the cylindrical symmetry used in this 113 study could be valid for obtaining the sample volume. We included this density 114variation in the density error (density errors are discussed in next paragraph). The 115116 obtained volumes and calculated densities of solid metals (Sn, Ni, and Fe) are summarized in Supplemental¹ Table S1–S3. In some runs, the density of solid sample 117was measured at room temperature from sample image (see Table S1-S3). The 118measured density (Sn: 7.278(15), 7.289(13) g/cm³, Fe: 7.868(9) g/cm³) agree well with 119 the density database at room temperature (Sn: 7.265 g/cm³, Fe: 7.874 g/cm³, 120 Chronological Scientific Tables, 2020). Therefore, it is confirmed that the density 121122obtained from measurement used in this study is consistent with the value of density 123database.

Density error of solid samples contains following contributions. (1) *Errors of sample lengths*. The sample diameter and height were obtained from horizontal and vertical brightness profiles of the sample image, respectively (Fig. 2b). The length

errors were taken from difference between maximum (minimum) and average measured 127 lengths. The volume (density) error was derived from the propagation of these length 128errors. The density error derived from length errors are summarized in Supplemental¹ 129130 Table S4. (2) Rotation error. The sample volume variation derived from sample rotation of $0-90^{\circ}$, as stated in previous paragraph, was included in volume (density) error as 131 summarized in Table S4. (3) Error of pixel calibration. As mentioned in supplement (I), 132133length error derived from pixel calibration are 0.004–0.01 µm. The density error derived from this length error is calculated to $3x10^{-4}$ -8x10⁻⁴% based on error propagation, 134suggesting that the density error from pixel calibration error is negligibly small. Thus, 135the density error of solid sample in this study correspond to a sum of contribution of 136137 length errors and that of rotation error. All these errors are summarized in Table S4. In 138run AK19, for example, total density errors are 0.137-0.162%, which is obtained from a sum of density error from length errors (0.004–0.027%) and that from rotation error 139140 (0.132-0.135%).

For liquid samples, the volume of a sample droplet was obtained from a 2D image using Pappus's centroid theorem, which describes the volume of rotating bodies as:

$$V = 2\pi RS \tag{1}$$

where S and R indicate a half area of the sample cross section divided by the rotating 144 axis and the distance between the rotation axis and center of gravity, respectively. The 145value of S was determined from the 2D image, as shown by the shaded area in Fig. 3a. 146147The sample boundary was identified from the vertical scan of horizontal brightness profiles. The initial image was cut at the rotation axis. Next, the image was binarized, 148and the area (S) and coordinates of the center of gravity were obtained. More detailed 149analysis procedure of liquid sample is provided in Supplement¹ (II). The obtained V and 150 ρ of liquid Fe are summarized in Supplemental¹ Table S5. The error in volume was 151derived from the following errors: (1) error in sample length and rotational symmetry, 152and (2) error in boundary threshold. (1) We calculated the sample volume using the 153images of the left and of right halves with equation (1) to verify the assumption of 154rotational symmetry. The volume (density) difference from left and right half images 155corresponds to density error derived from rotation. Sample length errors are included in 156the calculated volume. The results are listed in Supplemental¹ Table S6. In run AK20, 157for example, the density error from rotation is 0.01–0.038%. This small rotation error 158indicates that the liquid sample was symmetric to the rotation axis. (2) The error in the 159160 boundary threshold is derived mainly from gray pixels (6 to 14 pixels) at the edge boundary of the droplet as shown in the inset of Fig. 3b. We estimated the volume error 161

162	by changing the threshold value in the range of the gray pixels. The error in boundary
163	threshold was improved by taking a shadow picture of the sample using a backlight
164	following the method of Le Maux et al. (2019). The backlight has been used in previous
165	studies (e.g., Brillo and Egry 2003; Langstaff et al. 2013; Le Maux et al. 2019). The
166	estimated density error from boundary threshold are listed in Table S6. Thus, the density
167	errors of liquid sample correspond to a sum of these errors (i.e., length, rotation
168	symmetry, and boundary threshold) and summarized in Table S6. For example, in AK20
169	total density errors are 0.400-0.670%, which is obtained from a sum of density error
170	from rotation and length errors (0.001–0.038%) and that from boundary threshold error
171	(0.027–0.666%).

The recovered samples were mounted in epoxy and polished for textural 172observation and chemical analysis, which were carried out using a scanning electron 173microscope (JSM-6010LA; JEOL) equipped with an energy-dispersive spectrometer 174175(NORAN NSS EDX; Thermo Fischer Scientific) with an accelerating voltage of 20 kV. Chemical composition of starting and recovered samples are listed in Supplemental¹ 176Table S7. Supplemental¹ Fig. S3 shows example of backscattered electron image of the 177178Fe samples quenched from liquid (AK20). From the sample textures of backscattered electron images and chemical composition of the recovered samples, no contamination 179

180 was found in the recovered Sn, Ni, and Fe samples (Supplemental¹ Table S7).

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RESULTS AND DISCUSSION

183 Density of tin

The measured density of Sn (body-centered tetragonal [bct] structure) is plotted 184as a function of temperature in Fig. 4 (Data is found in Supplemental¹ Table S1). The 185density values from two measurements (AK11 and AK14) were reproducible. The 186density decreased almost linearly with increasing temperature up to 493 K, which 187 corresponds to $T/T_m < 0.98$. This trend is consistent with a previous study (Touloukian 188et al., 1975). However, above 493 K, a drastic decrease in density was observed (see Fig. 1894). The decrease in density was approximately 0.7%, which is much larger than the 190 191density error (0.18–0.22%). This phenomenon was possibly derived from the premelting effect. In a previous study, the premelting effect observed in V_P and V_s at T/T_m > 0.99 192was considered to occur as a result of decrease in shear modulus (Nadal and Le Poac, 193 2003). It was found that Sn has a premelting effect near the T_m for both the density and 194 elastic wave velocities. 195196 The thermal expansion coefficient (α) was determined by fitting the density

197 data up to 493 K with equation (2):

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$$\rho = \frac{\rho_0}{\exp(\int_{T_0}^T \alpha dT)}$$
(2),

where ρ_0 denotes the density at the reference temperature T₀. The α of this study and a previous study are listed in Table 1. In this study, ρ_0 was 7.278(15) at T₀ = 300 K, and the obtained α in the temperature range of 285–490 K was 6.22(3)x10⁻⁵ K⁻¹. This result was slightly smaller than that in a previous study (6.438x10⁻⁵ K⁻¹; Touloukian et al. 1975).

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205 Density of nickel

The measured density of Ni is plotted as a function of temperature in Fig. 5 206(Data is found in Supplemental¹ Table S2). The densities of two measurements (AK12 207and AK13) are reproducible. Previous results of Ni density are also plotted in Fig. 5. 208 The density in this study was consistent with that of previous studies (Kohlhaas et al., 209 1967; Suh et al., 1988; Abdullaev et al., 2015) with an uncertainty of 0.2-0.5%, except 210for a result of Newkirk et al. (1958). All of these results indicate that there is no drastic 211decrease in density near T_m, suggesting that face-centered cubic (fcc) Ni does not have 212the premelting effect on density. This result is different from the density behavior of bct 213214Sn.

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The density decreases linearly with increasing temperature up to the $T_m(1728)$

K). The α in this study was determined using eq. (2) to $6.18(6) \times 10^{-5} \text{ K}^{-1}$ in the temperature range of 1083–1713 K. The α of the present and previous studies are listed in Table 1. The result of this study is consistent with the previous result (α = $6.07(8) \times 10^{-5} \text{ K}^{-1}$; Suh et al., 1988) within the error.

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221 Density of solid iron

Densities of solid Fe are plotted as a function of temperature in Fig. 6a together 222with the result of liquid Fe. The present data is found in Supplemental¹ Table S3. The 223reproducibility of the measured density was confirmed from four measurements (AK15, 224AK16, AK18, and AK19) (see Fig. 6a). The two discontinuous density jumps observed 225at 1162(1) and 1666.0(5) K were derived from phase transitions of the α phase 226227 (body-centered cubic (bcc))/ γ phase (fcc) and γ phase (fcc)/ δ phase (bcc), respectively. Previously reported densities of solid Fe are also plotted in Fig. 6a. Our results are in 228good agreement with previous studies of α and γ phases (α phase: Straumanis and Kim, 2291969; Kohlhaas et al., 1967; y phase: Onink et al., 1993; Basinski et al., 1955; Kohlhaas 230et al., 1967) within the error. In the stability field of the δ -phase (1668–1803 K), the 231232density decreases almost linearly with increasing temperature up to near T_m (1813 K) (see Fig. 6b). Thus, the density of solid Fe does not show a drastic decrease near T_m , i.e., 233

the premelting effect. This suggests that Fe (bcc) shows a similar behavior with Ni (fcc),

while it differs from the behavior of Sn (bct).

The thermal expansion coefficient (α) was determined in each phase of Fe by 236fitting with equation (2) (Table 1). The α of the α -phase was 4.49(3) x10⁻⁵ K⁻¹ in the 237range of 823-1163 K, which is consistent with previous result of 4.54(2) x10⁻⁵ K⁻¹ 238(Straumanis and Kim, 1969). The α of the γ -phase was 6.37(5) x10⁻⁵ K⁻¹ in the range of 2391173–1673 K, which is consistent with $\alpha = 6.40 \times 10^{-5} \text{ K}^{-1}$ (Komabayashi and Fei, 2010). 240The α of the δ -phase was 7.02(9) x10⁻⁵ K⁻¹ in the range of 1675–1793 K, which is 35% 241larger than the result of Kohlhaas et al. (1967) (5.20(21) $\times 10^{-5}$ K⁻¹). Because the stability 242field of the δ -phase is a narrow temperature region at high temperature, the number of 243measured data points in previous studies was limited. By contrast, results of this study 244245supplied more data points and a precise temperature dependence of the density because it was measured with much smaller temperature increments (1–10 K) (Fig. 6b). 246

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248 **Density of liquid iron**

The density of liquid Fe is plotted as a function of temperature in Fig. 7 together with the results of previous studies. The data of this study is found in Supplemental¹ Table S5. The reproducibility of the measured density was confirmed

from two measurements (AK18 and AK20). The present results are in good agreement with the results of Kirshenbaum and Cahill (1962), Lucas (1960), Brillo and Egry (2004), and Le Maux et al. (2019) within the error. Compared with the previous results, the present results give the density with a smaller error and were measured in the smallest temperature increments of 5 K (Fig. 7). As shown in Fig. 7, the temperature dependence of the density of liquid Fe

258shows a concave shape. Thus, the α can be well approximated as a linear function of temperature (T); $\alpha = 1.37(10) \times 10^{-3} - (6.0(6) \times 10^{-7})T [K^{-1}]$. If the α is assumed to be a 259constant, the average α of AK18 and AK20 is 2.42(1) × 10⁻⁴ K⁻¹. The α of the present 260and previous studies are listed in Table 1. The temperature dependence of the α of liquid 261metals was reported in previous studies (e.g., Kumai and Dass, 1993). The linear 262263temperature dependence of α shows a better fit to the measured density variation (coefficient of determination, R^2 , is 0.996) than does the constant value of α (R^2 is 264 0.988). The regression equation of α as a function of temperature is valid, at least, in the 265temperature range of this study (1818–1998 K). 266

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The amount of light elements in the Earth's outer core has been estimated

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IMPLICATIONS

based on a comparison between the density of the outer core from the seismic model and the EOS of liquid Fe alloys. A constant value of the α of liquid Fe has been widely used in the EOS. However, in this study, the linear temperature dependence of α could be more appropriate for the density variation of liquid Fe. The temperature dependence of α is an important issue to describe the EOS of liquid Fe, and it also needs to be studied at high pressures.

276In terms of the premelting effect, it is reported that the premelting effect is connected with rapid increase of defectively coordinated atoms or lattice defects based 277on numerical calculations (Delogu, 2006a; 2006b; Manai and Delogu, 2007; Martorell 278279et al. 2013a). To evaluate the deviation from solid density, $\rho/\rho_{solid}(T)$ is plotted as a function of temperature in Supplemental¹ Fig. S4, where ρ and $\rho_{solid}(T)$ are measured 280281density and solid density, respectively. The extrapolated value from the linear trend before the premelting is used for the value for $\rho_{solid}(T)$. The $\rho/\rho_{solid}(T)$ decreases 282suddenly at 485–493 K before melting temperature of 505 K. This may indicate that the 283fraction of defect (melt) in solid Sn increases by the premelting effect. Alternate 284expression of melt fraction in solid Sn is given in Supplement¹ (III). Occurrence of the 285286premelting effect in this study is found to be different depending on crystal structures. This is possibly due that the fraction of defect (melt) is characteristic of lattice geometry 287

288	(Delogu, 2006b). Further study is necessary to investigate occurrence of the premelting
289	effect for various crystal structures of Fe-alloys at high pressures.

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376 Endnote:

377 Deposit item ---, Supplemental Material. Deposit items are free to all readers and found

378 on the MSA website, via the specific issue's Table of Contents.

379

380 Figure captions

- Fig. 1 (a) High-temperature furnace and measurement set-up. (b) Heater and sample
 view from the top inside the chamber. The C/C heater has a slit for sample
 observation. The sample rod is 3 mm in diameter. Cross denotes the position
 of thermocouple junction.
- Fig. 2 (a) Binarized image of solid Fe at 1473 K (AK16). Small alumina disk was
 inserted between the sample and large alumina plate. The sample bottom
 interface without alumina support was kept in flat even high temperatures
 (Fig. 2a). (b) Horizontal brightness profile of the area shown by dashed
 square in (a). Sample boundary was determined from the brightness profile.
 Diameter error corresponds difference between max. and min. diameter
 values.
- Fig. 3 (a) Schematic image of liquid sample droplet for application to Pappus's centroid theorem. Shaded area (S) denotes a half area of the sample cross section divided by the rotating axis (y). R denotes the distance between the rotation axis (y) and center of gravity (G). (b) Raw image of liquid Fe at 1818 K (AK20). (Inset) Gray pixels (1–6 pixels) at the top boundary of the sample.
 Fig. 4 The effect of temperature on the density of solid Sn. Open and solid circles

398		respectively denote the results of AK11 and AK14 of this study. The results of
399		two measurements are reproducible. Dashed line represents the trend reported
400		by Touloukian et al. (1975). Vertical dotted lines denote T/T _m =0.98 (493 K)
401		and 1.0 (505 K).
402	Fig. 5	The effect of temperature on the density of solid Ni. Bold dotted and thin
403		dashed lines denote linear fit to the density data of this study and that of
404		Abdullaev et al. (2015), respectively. Vertical dotted line represents the $T_{\rm m}$
405		(1728 K). Since small alumina disk was not put beneath the sample in AK12,
406		height error of this run was relatively large, causing the larger density error
407		(see Supplemental ¹ Table S4).
408	Fig. 6	(a) The effect of temperature on the density of solid and liquid Fe. Vertical
409		lines at 1162, 1666, and 1813 K denote the phase transitions of $\boldsymbol{\alpha}$ phase
410		(bcc)/ γ phase (fcc), γ phase (fcc)/ δ phase (bcc), and δ phase (bcc)/liquid,
411		respectively. (b) A zoom-in plot of high temperature part of γ and δ phases
412		(1600–1820 K).
413	Fig. 7	The effect of temperature on the density of liquid Fe. Bold dotted curve
414		represents fit to the density data using a linear temperature dependence of α
415		$(\alpha = 1.37(10) \times 10^{-3} - (6.0(6) \times 10^{-7})T [K^{-1}]).$





Figure 1



Figure 2



Figure 3



Figure 4







Figure 6



Figure 7

Sample	phase	$\rho_0 \left[g/cm^3\right]$	$T_0[K]$	T range [K] a [*]	^{**} [10 ⁻⁵ /K] ł	p^{**} [10 ⁻⁷ /K] Method	Reference
Sn	bct	7.278	300	285-490	6.22 (3)	Image	This study (AK11,14)
Ni	fcc	8.584	1083	1083-1713	6.18 (6)	Image	This study (AK12,13)
Fe	a-bcc	7.868	295	295-1153	4.49 (3)	Image	This study (AK15-18)
Fe	γ-fcc	7.653	1173	1163-1665.5	6.37 (5)	Image	This study (AK15-18)
Fe	δ-bcc	7.365	1668	1668-1803	7.02 (9)	Image	This study (AK16-19)
Fe	liquid	7.081	1818	1818-1998	24.2 (1)	Image	This study (AK18,20)
Fe	liquid	7.081	1818	1818-1998	137 (10)	-6.0(6) Image	This study (AK18,20)
Sn	bct	7.290	293		6.44		Touloukian et al. (1975)
Ni	fcc	8.626	976	976-1678	6.07 (8)	X-ray diff.	Suh et al. (1988)
Ni	fcc	8.497	1200	1200-1728	6.40 (6)	Gamma metho	d Abdullaev et al. (2015)
Fe	a-bcc	7.840	427	427-1023	4.54 (2)	X-ray diff.	Straumanis and Kim (1969)
Fe	γ-fcc	7.533	1400	1400-1600	6.4(1)	X-ray diff.	Komabayashi and Fei (2010)
Fe	δ-bcc	7.362	1663	1663-1781	5.20 (21)	X-ray diff.	Kohlhaas et al. (1967)
Fe	liquid	7.164	1822	1822-1972	23.7 (9)	Archimedean	Lucas (1972)
Fe	liquid	6.95	1860	1860-2477	14.3 (5)	Archimedean	Lucas (1960)
Fe	liquid	7.005	1845	1845-2448	16.0 (2)	Archimedean	Kirshenbaum and Cahill (1962)
Fe	liquid	6.99	1808	1808-1961	7.9 (7)	Image	Kobatake and Brillo (2013)

Table 1. Thermal expansion coefficient (α) of solid and liquid metals in this study and previous studies.^{*}

*Numbers in parenthesis indicate an error in last digit(s).

** α is expressed using a linear function of T as $\alpha = a+bT$