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# **Revision 3**

# Water quantification in olivine and wadsleyite by Raman spectroscopy and study of errors and uncertainties

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#### 6 **ABSTRACT**

7 The study of nominally anhydrous minerals with vibrational spectroscopy, despite its 8 sensitivity, tends to produce large uncertainties (in absorbance or intensity) if the observed 9 dispersion of the values arising from the anisotropy of interaction with light in non-cubic 10 minerals is not assessed. In this study, we focused on Raman spectroscopy, which allows the 11 measurement of crystals down to few micrometers in size in back-scattered geometry, and with 12 any water content, down to 200 ppm by weight of water. Using synthetic hydrous single-crystals 13 of olivine and wadsleyite, we demonstrate that under ideal conditions of measurement and 14 sampling, the data dispersion reaches  $\pm 30\%$  of the average (at  $1\sigma$ ) for olivine, and  $\pm 32\%$  for 15 wadsleyite, mostly because of their natural anisotropy. As this anisotropy is linked to physical 16 properties of the mineral, it should not be completely considered as error without treatment. By 17 simulating a large number of measurements with a 3D model of the OH/Si spectral intensity ratio 18 for olivine and wadsleyite as a function of orientation, we observe that although dispersion 19 increases when increasing the number of measured points in the sample, analytical error 20 decreases, and the contribution of anisotropy to this error decreases. With a sufficient number of

21 points (five to ten, depending on the measurement method), the greatest contribution to the error 22 on the measured intensities is related to the instrument's biases, and reaches 12 to 15% in ideal 23 cases, indicating that laser and power drift corrections have to be carefully performed. We finally 24 applied this knowledge on error sources (to translate data dispersion into analytical error) on 25 olivine and wadsleyite standards with known water contents to build calibration lines for each 26 mineral in order to convert the intensity ratio of the water bands over the structural bands (OH/Si) 27 to water content. The conversion factor from OH/Si to ppm by weight of water (H<sub>2</sub>O) is 28 93108±24005 for olivine, 250868±45591 for iron-bearing wadslevite, and 57546±13916 for iron-29 free wadsleyite, showing the strong effect of iron on the spectral intensities.

30 Keywords: wadsleyite, olivine, nominally anhydrous minerals, Raman spectroscopy,
31 water quantification

### 32 **INTRODUCTION**

33 Even though the major mineral phases of the Earth's mantle are nominally anhydrous, 34 many of them are known to contain water as OH point defects in their structure (Bell and 35 Rossman, 1992; Smyth and Keppler, 2006; Peslier 2010; Demouchy & Bolfan-Casanova, 2016). 36 The water storage capacity of the most abundant upper mantle mineral, namely olivine, has been 37 the subject of many studies, however most reports concern simple systems or single crystal 38 olivine (e.g. Kohlstedt et al. 1996; Withers & Hirschmann 2007 and 2008; Bali et al. 2008; 39 Kovács et al. 2010; Férot & Bolfan-Casanova 2012; Litasov et al. 2014, Yang et al. 2016). 40 Single-phase experiments allow the growth of single crystals large enough to be suitable for any 41 analytical method, especially absorption infrared spectroscopy using Fourier transform infrared (FTIR), secondary ion mass spectrometry (SIMS), or elastic recoil detection analysis (ERDA). 42

43 In experiments with natural mantle compositions on the other hand, olivine, 44 orthopyroxene, clinopyroxene and garnet can coexist, strongly limiting crystal growth. These 45 samples often display crystal sizes from 20 to 50 µm, limiting the use of the conventional 46 methods cited above on the fine-grained samples. Investigation of the storage capacity in more 47 complex systems such as peridotite has been carried out and the small grain size of the grains 48 required the use of SIMS (few tens of cubic micrometers analyzed) or even nano-SIMS (one or two order of magnitude less) (see Ardia et al., 2012; Tenner et al., 2012; Novella et al., 2014). 49 50 This technique is complicated to use because of high background levels of H, complex 51 preparation to avoid H contamination and the need for well-characterized standards (Koga et al. 52 2003, Mosenfelder et al. 2011).

53 FTIR is frequently used sensitive technique to quantify hydroxyl content in nominally 54 anhydrous minerals, which also gives structural information about H point defects. Reliable 55 methods exist to quantify water in anisotropic minerals using FTIR, and a consequent literature 56 exists on the subject (e.g., Libowitzsky and Rossman 1996; Asimow et al. 2006; Kovács et al. 57 2008; Withers et al. 2012, Withers 2013; Qiu et al. 2018). However, in the case of very water-58 rich samples as, for example, wadsleyite, a high-pressure polymorph of olivine, that can contain 59 up to 3.2 weight percent of water (Inoue et al., 1995), infrared spectroscopy requires an important 60 thinning of the samples (which may cause their loss) in order to avoid the entire absorption of the 61 infrared signal.

In this study, we used confocal polarized Raman spectroscopy (the laser source being polarized, the incident beam is therefore also polarized). This technique offers several advantages for water quantification in synthetic minerals with a small grain size and low to very high water contents. Regarding sample preparation, OH quantification using back-scattered Raman

spectroscopy requires only one side of the sample to be polished, while double polishing is 66 67 required in the measurement of absorbance using FTIR. Unfortunately, very thin polishing of the sample can irremediably damage them, especially those sintered under conditions where fluids 68 69 are highly wetting, which is the case of the conditions of the deep upper mantle (Yoshino et al. 70 2007). The detection limit of water quantification using Raman spectroscopy (around 71 50-100 ppm wt) may be a problem with samples synthesized under conditions of the uppermost 72 mantle, where the water solubility is the lowest for many mineral species (Férot & Bolfan-73 Casanova 2012, Yang 2016), but any higher concentration can be measured (Bolfan-Casanova et 74 al. 2014). The spot size of confocal Raman spectroscopy (3-5 µm) allows measurements on very 75 fine-grained samples (down to crystal sizes around 5 um) without difficulties. The time required 76 for reasonably precise measurements (some minutes) is low enough to multiply the number of 77 measured points on different crystals throughout the sample and have a statistically correct 78 coverage of the whole sample. One major drawback of Raman spectroscopy is that the intensity 79 absorption depends on many factors in addition to the concentration, such as the intensity of the 80 incident laser, confocality and lens magnification, which control the volume of sample that is 81 excited by the laser beam. Moreover, the chosen gratings and the optics of the spectrometer 82 (instrument-dependent parameters), as well as the optical properties of the sample or its surface 83 state will affect the efficiency of the measurement (see e.g. Mercier et al. 2009, Schiavi et al. 84 2018; Zarei et al. 2018). In contrast, the absorbance measured using FTIR depends solely on the 85 thickness, concentration and absorptivity of the sample itself, following a relationship known as 86 the Beer-Lambert law. In addition, in absorbance spectroscopy the intensity transmitted by the 87 sample is always normalized to that of the incident beam, which tends to eliminate instrumental 88 biases on the intensity of the signal from the sample.

89 Raman spectroscopy is widely used to quantify H<sub>2</sub>O concentrations in glasses or melt 90 inclusions (e.g., Thomas et al. 2008; Mercier et al. 2009; Schiavi et al., 2018). In these studies, 91 quantification of the water concentration relies on the comparison of the OH/Si of the unknown 92 to that of well-characterized standards measured under identical conditions, as the OH intensity 93 has been shown to increase linearly with water content. Here, OH/Si is defined by the integrated 94 intensity of the water band normalized to that of the silicate vibrations (see the following section 95 for details). As numerous factors may affect the measurement efficiency (such as variation of 96 focusing depth or surface quality), using the ratio of the OH band area over that of silica bands 97 area reduces the data dispersion or scatter, as both regions are equally affected by focusing or 98 surface variations. Calibrations of the method have been proposed by Thomas et al. (2008 and 99 2015) for garnet, Bolfan-Casanova et al. (2014) for olivine, Thomas et al. (2015) for ringwoodite, 100 and Weis et al. (2018) for orthopyroxene (this last study having used also forward-scattering).

101 Previous water content quantification in olivine conducted with Raman spectroscopy 102 often display large error bars on their results (Bolfan-Casanova et al. 2014). In this study, we 103 demonstrate that those error bars are significantly related to the dispersion (used in the statistical 104 meaning of the term throughout this work) of the relative intensities of the different vibrational 105 modes caused by the anisotropy of the minerals. As this anisotropy is a natural consequence of 106 the structure and symmetry of the crystal, it is not directly related to error, and can even be used 107 to get information on orientation (see for example Ishibashi et al. 2008). Firstly, we studied the 108 relative effect of the different error sources and of anisotropy on the statistical dispersion of 109 measurements in hydrated olivine and wadsleyite single crystals and used it to propose a method 110 to estimate the analytical error from the data dispersion of the measurements using polarized

unanalyzed Raman spectroscopy. We then propose a calibration for water quantification inolivine and wadsleyite based on standards characterized by FTIR or ERDA methods.

# 113 METHODS

114 Synthesis of olivine and wadsleyite

115 Olivine and wadsleyite single crystals were synthesized in the multi-anvil press at 12 and 116 15 GPa and at 1200 and 1350 °C from San Carlos olivine as starting material. Olivine single 117 crystals and powder were placed in a folded rhenium foil capsule, placed itself in a welded gold-118 palladium capsule containing brucite powder. The experimental assembly consisted of an MgO 119 octahedron containing  $Cr_2O_3$ , a zirconia thermal insulator, a LaCrO<sub>3</sub> heater with molybdenum 120 electrodes in contact with the anvils and an MgO central part containing the capsule. Temperature 121 was controlled using a W-Re thermocouple (5% and 26% Re). Experiments were conducted at 122 the Laboratoire Magmas et Volcans (LMV, Clermont-Ferrand, France) on a Voggenreiter Mavo-123 press LP 1500 tons multi-anvil press equipped with a Kawai-Endo apparatus. Heating was 124 performed and controlled by a Pacific 140-AMX AC power source. Secondary anvils were 32 125 mm tungsten carbide cubes with 8 and 6 mm truncations. Olivine synthesis duration was 4.5 126 hours at 1200 °C, and wadsleyite synthesis lasted 2 hours at 1350 °C. For both syntheses, 127 temperature was gradually decreased after the experiment (around 50 °C per minute) instead of 128 quenching to prevent crystal fracturation. Recovered crystals were oriented using polarized light 129 microscopy for olivine and X-ray diffraction for wadsleyite, and then cut and mirror polished 130 perpendicular to crystallographic axes (see Figure 1). Olivine and wadsleyite are both 131 orthorhombic so display an anisotropic behavior with respect to light absorption. Wadsleyite has 132 been reported to become monoclinic above 0.5 wt % H<sub>2</sub>O using X-ray diffraction (Jacobsen et al.

133 2005), however we did not detect any noticeable spectral difference that could be due to this134 small change in the β-angle.

# 135 FTIR spectroscopy

Most of the samples used in order to calibrate the Raman quantification method were olivines characterized using FTIR by Férot & Bolfan-Casanova (2012). These water content values were updated using the most recent calibration of the infrared extinction coefficient ( $\epsilon$ ) determined using ERDA by Withers et al. (2012) of 45200 1 mol<sub>H2O</sub><sup>-1</sup> cm<sup>-2</sup> instead of the 28450 1 mol<sub>H2O</sub><sup>-1</sup> cm<sup>-2</sup> value given by Bell et al. (2003). This decreases all the water contents of olivine reported in Férot and Bolfan-Casanova (2012) by a factor of 1.589.

142 Polarized spectra acquired for this study were collected on a Vertex70 Brucker 143 spectrometer coupled to a Hyperion microscope with a 15× lens, condenser and knife-edge 144 apertures creating a rectangular target area of 40-50 µm. Samples were analyzed on a CaF<sub>2</sub> plate at a resolution of 2 cm<sup>-1</sup>, with 200 to 300 accumulations. After background subtraction and 145 atmospheric correction, a cubic baseline correction has been applied on the 1500 to 4000 cm<sup>-1</sup> 146 area, and integration calculated in the 2700 to 3720 cm<sup>-1</sup> area. The absorbance is then used to 147 148 calculate the H<sub>2</sub>O content based on Beer Lambert's law (see equation 1) with the integrated molar 149 absorption coefficient (also called extinction coefficient) cited above, and a density factor X of 5521 l·mol<sub>H20</sub><sup>-1</sup> (used to convert mol<sub>H20</sub>·l<sup>-1</sup> to ppm by weight of H<sub>2</sub>O, equal to  $18.02 \times 10^{6}/\rho$ ). The 150 151 density p has been calculated from Fischer & Medaris (1969) accounting for the olivine 152 composition. The equation used is the following, with  $\varepsilon$  being the infrared extinction coefficient, 153  $A_{tot}$  the total absorbance, and d the thickness of the sample in centimeters:

154 Equation 1

$$C_{H_2O} = X \times \frac{A_{tot}}{\varepsilon \times d}$$

#### 155 Raman spectroscopy

156 Raman spectra were collected at LMV using an InVia confocal Raman micro-157 spectrometer manufactured by Renishaw and equipped with a 532 nm diode laser (200 mW 158 output power), a Peltier-cooled CCD detector of 1024×256 pixels, a motorized XYZ stage and a 159 Leica DM 2500M optical microscope. Scattered light was collected in a back-scattered geometry. 160 An edge filter effectively reduced both Rayleigh scattered photons and photons from the exciting laser source at 0 cm<sup>-1</sup> that had been reflected by the sample surface. A 2400 grooves/mm grating 161 was used for the analyses, which resulted in a spectral resolution between 1.3 cm<sup>-1</sup> (around 100 162 cm<sup>-1</sup>) and 0.72 cm<sup>-1</sup> (around 3700 cm<sup>-1</sup>). All spectra were acquired with polarized light without 163 164 analyzer, and sample was rotated under the beam to change the polarization angle. A ×100 165 microscope lens (numerical aperture 0.9) was used and the slit aperture was set to 20 µm (high 166 confocality setting). Daily calibration of the spectrometer was performed based on a Si 520.5 cm<sup>-1</sup> peak. The effective laser power (changed by filters) used was 65 to 71 mW for olivine, and 167 168 8 to 9 mW for wadsleyite and glass (low enough power to prevent destabilization and 169 dehydration) and was measured to normalize spectra to 1 mW. A power of 125 to 140 mW can 170 cause some damages on water-rich olivine, and a power of 15 mW caused a slight OH intensity 171 reduction in wadsleyite with very high water content (sample 2054, 3.4% in weight), but no 172 dehydration or destabilization effect has been observed with the lower laser power values finally 173 chosen. A hydrous basaltic glass (see Schiavi et al. 2018) is used to check and be able to correct 174 for power drifts and efficiency drifts along the different area of measurement.

175 As already stated above, the quantification of water concentration is based on the measurement of the OH band integrated intensity normalized to the silicate band integrated 176 177 intensity. Under static conditions of measurement, the window analyzed using a grating of 2400 groves/mm decreases from approximately 1200 cm<sup>-1</sup> at low wavenumbers to 800 cm<sup>-1</sup> at high 178 wavenumbers. For all samples and standards, Si area refers to wavenumbers from 61 to 1318 179 cm<sup>-1</sup> centered on 720 cm<sup>-1</sup>, and OH area refers to wavenumbers from 2978 to 3784 cm<sup>-1</sup> centered 180 on 3400 cm<sup>-1</sup>. Both areas were chosen to contain all the needed peaks and bands. Acquisition 181 182 times are short for the Si area because of its high intensity, to prevent saturation of the CCD 183 detector, and longer for the OH area to increase the signal to noise ratio, the OH bands being 184 much weaker. Thus, acquisition times were  $2 \times 5$  s for the Si area, and  $5 \times 60$  s for the OH area of 185 olivine and wadsleyite. For the glass, the acquisition times were  $4 \times 10$  s for the Si area, and 186  $5 \times 120$  s for the OH area. The daily variations of the spectrometer were corrected by normalizing 187 to the average OH/Si intensity ratio of the hydrated glass measured in each measurement session, 188 each of those consisting on two random points on the glass sample, providing a more 189 reproducible OH/Si intensity ratio here called (OH/Si)<sub>Smp Norm</sub>. The normalization operation is 190 shown in equation 2, where (OH/Si)<sub>Smp Meas</sub> is the average OH/Si of the measured sample, 191 (OH/Si)<sub>Glass Meas</sub> is the standard glass OH/Si measured in the same session as the sample, and (OH/Si)<sub>Glass Std</sub> is the measurement of the OH/Si of the glass measured in the same session as the 192 193 standards. The result is the normalized OH/Si of the sample (OH/Si)<sub>Smp Norm</sub>.

194 Equation 2

$$\frac{OH}{Si}_{Smp Norm} = \frac{OH}{Si}_{Smp Meas} \times \frac{\frac{OH}{Si}_{Glass Std}}{\frac{OH}{Si}_{Glass Meas}}$$

. . .

195 The procedure of baseline correction is essential in obtaining reproducible values of 196 OH/Si for each phase (see Figure 2 and supplementary materials 1). The baseline shape is linear 197 for the olivine Si area (rarely cubic), cubic for its OH area, cubic or polygonal for the wadsleyite 198 Si area and cubic for its OH area. The anchor points used to define the baseline are shown in 199 Figure 2. The spectra used afterwards for integrations are the raw spectra subtracted of the 200 baseline, and divided by the total acquisition time and the laser power, to normalize each 201 spectrum to a power of 1 mW and an acquisition time of 1 s. It has to be noted that this is not the 202 unit of the spectra, as the intensity measured on each spectrum depends on various parameters of the instrument such as the grating and the wave number, but can be labelled as  $counts \cdot s^{-1} \cdot mW^{-1}$ . 203 204 Integration consists in the area under the curve of the baseline-corrected and normalized 205 spectrum: [200, 1100] for the Si area and [3050, 3700] for the OH area of olivine, and [300, 206 1100] for the Si area and [3150, 3700] for the OH area of wadsleyite (see Figure 2). The result of 207 this integration defines thereafter the intensity of the spectra. In the case of wadslevite, ironbearing samples display a very intense band at 171 cm<sup>-1</sup> that merges with the white line (at low 208 209 wavenumbers) rendering its baseline correction unreliable due to the difficult anchoring of the 210 baseline at low wavenumbers (see supplementary materials 1). This is why the integration of the Si area of wadsleyite starts at a higher wavenumber of 300 cm<sup>-1</sup>. 211

For all the spectra of glass, the single-crystal spectra and some of the wadsleyite and olivine samples, one measurement point consists in the OH intensity divided by the Si intensity (OH/Si ratio) with the same polarization angle, and is referred as "single points". On most of the olivine and wadsleyite samples, one measurement point will be the average of two OH/Si ratio obtained from two measurements on the same crystal with orthogonal polarization angles, and is referred as "orthogonal couples".

The averaged OH/Si values of several measurement points on a sample is the value used to compare the different samples. Eight olivine and seven wadsleyite standard samples of various known water contents have been used to build the calibration lines allowing the conversion from OH/Si ratio to water content. Water contents of olivine standards have been determined by FTIR, using Withers et al. (2012) extinction coefficient, and those of wadsleyite standards were measured by ERDA (Bolfan-Casanova et al. 2018).

# 224 **Results**

# 225 Quantification of the OH/Si variation related to mineral orientation

In order to assess the effect of orientation alone on the dispersion of the average OH/Si value of each phase, polarized Raman spectra have been acquired on each face of the prepared single crystals of olivine and wadsleyite. We acquired 19 spectra, from 0 to 180°, on each face for both OH and Si area, and values for 190 to 350° were calculated by symmetry. Normalized OH/Si values were then fitted using equation 3 for olivine and equation 4 for wadsleyite (see Figure 3).

232 Equation 3

$$\frac{OH}{Si}(\theta) = H + a_1 \times \cos^2(\theta + d_1)$$

233 Equation 4

$$\frac{OH}{Si}(\theta) = H + a_1 \times \cos^2(\theta + d_1) + a_2 \times \cos^2(2 * (\theta + d_2))$$

Here, *H* is a constant depicting the minimum value for a given face adjusting the vertical position of the curve;  $a_1$  is the amplitude (the difference between the maximum and minimum

236 intensities of the given face for equation 3) of the cosine describing the features with a period of 237 180° (2-fold symmetry) shifted of an angle  $d_1$ ;  $a_2$  is the amplitude of the cosine describing the 238 features with a period of 90° (4-fold symmetry) shifted of an angle  $d_2$ . Each parameter has been 239 optimized to minimize the differences between the measurement data and the fitted curve. It has 240 to be noted that the necessity for introducing parameters d arise from the fact that crystals are not 241 exactly polished following the crystallographic axes. These fits are thus empirical and only serve 242 the purpose of creating a model of the analyzed crystals. For olivine, this misalignment is around 243 10° (as visible in Figure 3), and has later been neglected (the maxima of the faces (100) and (010) 244 have been aligned along crystallographic axes in the 3D model). On the other hand, the 245 misalignment of wadsleyite is high, and orientation has been lost, rendering necessary the use of 246 the shift parameters d. The observed component with a 90° period is necessary in the case of 247 wadsleyite to fit correctly the measured intensity, and can mostly be noticed in the fits of the 248 intensity variations on the faces arbitrarily called F1, F2 and F3 (see M F1 and M F2 fit curves 249 around 150 and 330°).

250 Considering that the single-crystal analyzed is homogeneous in water content, the only 251 source of deviation between the modeled curve and the measured data is the combination of data 252 treatment error and measurement error. Data treatment error has been estimated from the 253 differences between repeated baseline correction and treatments of Raman spectra, and will be 254 referred as  $E_{Tr}$  in following equations, tables and figures. The estimated values for  $E_{Tr}$  on OH/Si 255 are 3.6% for olivine and 4.0% for wadsleyite. This error tends to remain small when the signal to 256 noise ratio is high enough. The second error source taken into account is the one caused by all 257 measurement uncertainties, such as surface irregularities or focus offset, and will be referred as 258  $E_{Meas}$ . Supposing that  $E_{Meas}$  and  $E_{Tr}$  are independent sources of error, the deviation of

259 measurement points from the model is the square root of the sum of the two error sources 260 squared, and is provided by the fit shown in Figure 3. From this calculation, we obtain a 261 measurement related uncertainty  $E_{Meas}$  reaching 4.4% for olivine, and 4.1% for wadsleyite. The 262 last source of error arises from the daily variation of the spectrometer efficiency and the method 263 used to correct it. Hence, the treatment and measurement errors related to this correction are the 264 cause of an additional error of 11.9%, and will be referred as  $E_{Cor}$ . This value has been estimated 265 accounting for the measurement and data treatment errors related to the standard glass 266 measurement. The first two error,  $E_{Meas}$  and  $E_{Tr}$ , sources apply on each point separately, while the 267 correction error  $E_{Cor}$  applies to all the points of a measurement session. The values of the 268 estimated errors are reported in table 1.

269 The intensity versus angle model curves (as shown in Figure 3) were then used to build a 270 3D model of OH/Si variation as a function of crystallographic orientation. In Figure 3 it is visible 271 that the fitted curves do not coincide with the corresponding minima and maxima from the other 272 faces (for example, the maxima of the "a" and "b" curves for olivine are different, but both 273 describe the OH/Si intensity parallel to the "c" axis and should be equal). These differences are 274 very likely to be caused by misalignment of the polished faces with respect to the 275 crystallographic axes. To obtain a continuous set of three curves in 3D (as depicted by the bold 276 lines in Figure 4), the amplitudes of the curves have been increased or decreased so that they 277 cross each other along the crystallographic axes, keeping at the same time the average value of 278 each curve constant. Intermediate orientation values were completed by combining those 279 acquired along crystallographic axes, each point of any orientation being obtained with equation 280 5:

281 Equation 5

$$\left(\frac{OH}{Si}\right)_{\alpha} = \sum \left(\frac{OH}{Si}\right)_{i} \times \cos^{2}(\alpha_{i})$$

282 In this equation, any orientation  $\alpha$  is defined by a vector with a norm equal to the 283 corresponding (OH/Si)<sub>a</sub> value. This vector has one projection on each of the "a", "b" and "c" 284 planes corresponding to values (OH/Si)<sub>i</sub> fitted with equation 3 or 4 for a given angle (see Figure 285 3).  $(OH/Si)_{\alpha}$  is then calculated as being the sum of these three  $(OH/Si)_{i}$ , each multiplied by the 286 squared cosine of the angle between the defined vector and its projection on the corresponding 287 plane. This operation results in a peanut-shaped three-dimensional OH/Si model for each phase when displayed in a 3D polar graph (see Figure 4, depicting OH/Si variation in olivine and 288 289 supplementary materials 2 for wadsleyite). The OH/Si is depicted here as the distance between 290 the peanut-shaped surface and the intersection of the a, b, and c axes. In both olivine and 291 wadsleyite cases, it can be observed that high OH/Si values are located along one axis, identified 292 as the c-axis for olivine, and that variations seem approximately axisymmetric for both of them. 293 This modelling process can be applied to other orthorhombic minerals (or quadratic), each model 294 being based on its own measured OH/Si values (as shown in Figure 3). Minerals with other 295 crystal system could still be modelled, but they might require a different procedure of polishing, 296 measurement and 3D model extrapolation.

# 297 Distinguishing OH/Si natural dispersion and errors

The 3D OH/Si distribution models for olivine and wadsleyite have then been used as basis to simulate measurements, in order to untangle the effects of natural dispersion (i.e. due to crystallographic orientation) and analytical uncertainties on the final error for any type of measurement (single points or orthogonal couples) and as a function of the number of

302 measurement points. To achieve this, the first step is to generate a set of random uniformly 303 distributed crystal orientations that simulates a polycrystalline sample. A random "longitude" 304 angle between 0 and 360° is first generated (uniform distribution), and a second "latitude" angle 305 is generated with a semicircle distribution with a maximal probability for  $0^{\circ}$  and a probability of 306 reaching zero when approaching 90 or -90°. To each orientation corresponds an OH/Si value 307 determined as described for the construction of the 3D model. On each point chosen among those 308 of this surface, the estimated measurement and treatment errors are applied, using equation 6. 309 Here,  $X_{Meas}$  and  $X_{Tr}$  represent normally distributed random values with an average of zero, and a 310 standard deviation equal to the desired error component  $E_{Meas}$  or  $E_{Tr}$ . Each "modeled", or 311 simulated point (OH/Si)<sub>M</sub> has its own  $X_{Meas}$  and  $X_{Tr}$  values, because those errors do apply on 312 every point separately and independently, finally yielding the "real" value  $(OH/Si)_{R}$ .

#### 313 Equation 6

$$\frac{OH}{Si_R} = \frac{OH}{Si_M} \times (1 + X_{Meas}) \times (1 + X_{Tr})$$

For single point measurements, the simulated sample value will be the average of all the (OH/Si)<sub>R</sub>, multiplied in the same way by  $1 + X_{Cor}$  to simulate the errors arising from the laser power drift correction. In the case of orthogonal couples measurements, the two resulting (OH/Si)<sub>R</sub> values of each measurement point are averaged, then all points are again averaged before applying the  $1 + X_{Cor}$  factor.

All the error and dispersion values used in the following section are expressed as relative errors, in percent of deviation from the average. Hereafter, the goal is to get the standard deviation and the average of each simulated measurement session for any number of measurement points and repeat the operation a large number of times to get a result as

323 statistically significant as possible. In this study, final values have been obtained for sets of 2 to 324 30 points in both single point and orthogonal couples, with one million draws for each case. For 325 each number of points, the average of all the standard deviations of all sets of values has been 326 interpreted as the expected measured dispersion  $D_{Th}$ , and the standard deviation of all the 327 averages of all the set of points has been interpreted as the error expected for the measurement 328  $(E_{Th})$ . Both  $D_{Th}$  and  $E_{Th}$  depend on the number of measurement points and on the technique used 329 for the measurement (single point or orthogonal couples). Dispersion increases when adding 330 points, reaching an approximate plateau past 10 points, and error decreases when adding points, 331 but stabilizes much slower past 30 points, limited by the contribution of E<sub>Cor</sub> to the final error (see Figure 5). 332

333 The dispersion and error values obtained by the above simulation were then fitted to 334 obtain error propagation equations as a function of the number of points, the measurement 335 method, and the different errors sources. The natural dispersion caused by anisotropy can be 336 estimated from equation 7 below. Here  $D_{Nat}$  is the natural dispersion, which is the expected 337 relative standard deviation obtained on a random set of measurements without any kind of error added as a function of the number of points (N).  $D_{An}$  represents the total anisotropy, and is the 338 339 same standard deviation as  $D_{Nat}$ , but for an infinite set of points, and finally "a" is a fitting 340 parameter (all the values of the different parameters are shown in Table 1).

341 Equation 7

$$D_{Nat} = D_{An} - \frac{a}{N-1}$$

Equations 8 and 9 also express the dispersion of values and in addition takes into account the errors arising from measurement and data treatment. As the normalization related errors (E<sub>Cor</sub>)

apply on all the points in the same way, it does not have any effect on the relative dispersion. Equation 8 gives the expected dispersion for single points measurements, and equation 9 for orthogonal couples. Here,  $D_{An}$ , "*a*" and *N* are the same as in equation 7,  $E_{Meas}$  is the measurement error, and  $E_{Tr}$  is the data treatment error.

348 Equation 8

$$D_{Th} = \sqrt{D_{An}^2 + E_{Meas}^2 + E_{Tr}^2} - \frac{a}{N-1}$$

#### 349 Equation 9

$$D_{Th} = \sqrt{D_{An}^2 + \frac{E_{Meas}^2 + E_{Tr}^2}{2} - \frac{a}{N-1}}$$

The theoretical uncertainties follow a simple error propagation equation, which is expressed as shown in equation 10 for single points measurements, and equation 11 for orthogonal couples. Here,  $E_{Cor}$  plays a role on the final relative error and tends to account for most of the uncertainty when the number of points increases (see Figure 5).

#### 354 Equation 10

$$E_{Th} = \sqrt{\frac{D_{An}^2 + E_{Meas}^2 + E_{Tr}^2}{N} + E_{Cor}^2}$$

#### 355 Equation 11

$$E_{Th} = \sqrt{\frac{D_{An}^2 + \frac{E_{Meas}^2 + E_{Tr}^2}{2}}{N} + E_{Cor}^2}$$

Each of the above equations can be applied for both olivine and wadsleyite. All the parameters are displayed in Table 1.

358 It can be noted that all the dispersions and errors in this simulation neglect the initial 359 asymmetry of the distribution of the OH/Si values (see distribution histograms in supplementary 360 materials 3 and 4) which, although abnormal and strongly asymmetric (towards high values for 361 single points measurements and low values for the orthogonal couples), rapidly tends toward a 362 Gaussian-shaped distribution above 10 to 15 points. Moreover, the asymmetry of the distribution 363 might be very difficult to observe on a real set of measurement because of the limited set of 364 points in which the distribution of values may not have any statistical significance. Hence, the 365 deviation caused by this hypothesis tends to be negligible above 10 to 15 measurement points, 366 but very low number of points should still follow a very asymmetric distribution, increasing 367 error.

368 It can be noted that the measurement and treatment errors have a very small effect on both 369 the dispersion and the error (see Figure 5). On the other hand, power drift correction related 370 errors (E<sub>Cor</sub>) have a very strong effect on the final error, demonstrating that this step must be 371 carefully performed in order to get correct values. An important result is that the use of averaged 372 orthogonally polarized measurements roughly halves the theoretical dispersion of the values, thus 373 decreasing the associated error. The total expected dispersion directly related to anisotropy 374 (neglecting all the other error sources) for olivine is around  $\pm 29.4\%$  for single points and  $\pm 14.8\%$ 375 for orthogonal couples. Wadsleyite displays slightly higher values, with 32.5% and 16.8% (all 376 these values are shown in Table 1). The same effect could be reached on single points measurements with approximately four times more points. Another advantage of using the 377

orthogonal couples is that it forces to measure more diverse grain orientations in the sample,hence limiting the effect of a potential preferential crystallographic orientation within the sample.

380 To sum up, assuming a chemically homogeneous sample (in terms of water content), if 381 the standard deviation is close to the expected dispersion (calculated with equations 8 or 9), the 382 real analytical error is given by equations 10 or 11. It follows that, in the ideal case of randomly 383 oriented crystals, the error on OH/Si for 5 to 10 measurement points is between 15 and 18% for 384 single points, and 12 and 15% for orthogonal couples. A deviation inferior to the expected 385 deviation (still assuming that the sample is homogeneous) would point out an insufficient 386 sampling of the natural dispersion. A rough estimation of the real error in this case can be obtained with equation 12. Here, the difference between the expected deviation  $D_{Th}$  and the 387 388 measured value  $D_{Smp}$  causes the final error  $(E_{Smp})$  to increase from the ideal value  $E_{Th}$  to the total 389 dispersion  $D_{Th}$ , equal to the error expected on a measurement based on one single measurement 390 point.

In contrast, if the standard deviation is higher than  $D_{Th}$ , whether because the sample is heterogeneous or the measurements are too defective, the error "reduction" from  $D_{Smp}$  to  $E_{Smp}$ (from equations 8 or 9 to equations 10 or 11) may not be reasonable to apply directly. Considering that the sampling is sufficient in this case, an error approximation can be obtained from equation 10 (variables are the same as in equation 13).

396 Equation 12

$$E_{Smp} = D_{Th} - (D_{Th} - E_{Th}) \times \frac{D_{Smp}}{D_{Th}}$$

397 Equation 13

$$E_{Smp} = \sqrt{\left(D_{Smp} \times \frac{E_{Th}}{D_{Th}}\right)^2 + \left(D_{Smp} - D_{Th}\right)^2}$$

#### **398** Olivine and wadsleyite calibration lines

399 The water content of the samples used to build the calibration lines are shown in Table 2, 400 along with the values of OH/Si measured with Raman spectroscopy. The calibration lines (see 401 equation 14 and Figure 6) are built with least-square minimization, always pass by the origin, and 402 take into account the effect of the water content uncertainties arising from FTIR and ERDA 403 measurements as well as the Raman measurement. All the error bars on the OH/Si have been 404 estimated with the method described above. Some samples thus display a relatively large error 405 bar (such as 949 or M380b, see Figure 6 and Table 2) because the measured dispersion was 406 significantly greater than the expected dispersion (obtained from equations 8 to 11). On the other 407 hand, some samples, such as M382, displayed a very low dispersion that may be related to 408 insufficient sampling, once again resulting in a larger error bar. The factor to convert from OH/Si 409 to ppm in weight of water (F in equation 11) is  $93108\pm24005$  for olivine,  $250868\pm45591$  for 410 wadsleyite (with iron), and 57546±13916 for wadsleyite (without iron). This result shows the 411 tremendous effect of iron on the spectrum's shape and intensity, implying that chemical composition has to be verified in order to quantify the water content of wadsleyite. There are two 412 413 identified causes to this difference, the first being that for two samples with comparable water 414 contents, the Si area of an iron-free wadsleyite displays sharper peaks, even if their heights are 415 comparable, implying a lower Si intensity. However, at the same time, its OH area is more 416 intense. This causes the OH/Si of an iron-free wadsleyite to be much higher than the one of an

- 417 iron-bearing wadsleyite with a comparable water content (see supplementary materials 5). Figure
- 418 6 displays the calibration lines for olivine and wadsleyite.
- 419 Equation 14

$$C_{H_2O}(ppm wt) = F \times \frac{OH}{Si}_{Smp Norm}$$

The uncertainty on the slope obtained for each calibration has to be added to the final error  $E_{Smp}$  described above, the total uncertainty on the water content being given by equation 15, where  $E_{H_{2O}}$  is the total error,  $E_{Smp}$  is the total relative error on the OH/Si of the sample, and  $E_{Cal}$  is the calibration-related relative error. The same can be applied to dispersions.

424

#### 425 Equation 15

$$E_{H_2O} = \sqrt{E_{Smp}^2 + E_{Cal}^2}$$

The uncertainty on the slope of the calibration line (as shown in Figure 6) arising from the errors in water quantification of the standards and from the error on their OH/Si *often* becomes the major component of the total uncertainty on the water content compared to the error on the measurement of OH/Si with Raman spectroscopy.

Regarding olivine, the 26% uncertainty on the calibration line is much larger than the 12% error on the OH/Si measurement attained for a large number of points (in the case of orthogonal couples, see Figure 5). Following equation 15, this 12% error ( $E_{Smp}$ ) increases to 29% only because of the uncertainty on the calibration ( $E_{Cal}$ ). The same problem applies to wadsleyite with the 18% error on the calibration line for iron-bearing wadsleyite and the 24% error for iron-

free wadsleyite. The greater is the error on the OH/Si of the sample (because of insufficient sampling or sample heterogeneities), the smaller the contribution of calibration uncertainties to the total error will be (see Figure 7).

# 438 **DISCUSSION**

439 The quantification of as much error and uncertainty sources as possible provides a better 440 understanding of the data dispersion inherent to the use of Raman spectroscopy on anisotropic 441 crystals. The relative impact of each source of error highlights what affects the more the final 442 error, showing where further progress and developments could be made. The first issue is the 443 laser-drift during and between different sessions. The method used to overcome this problem 444 brings an additional source of error (referred as E<sub>Cor</sub> throughout this study), as it implies two 445 OH/Si ratio measurements on a standard glass, thus adding twice data treatment and measurement error to the result (divided by  $\sqrt{2}$ , as the value measured is supposed to be the same 446 447 for the two measurements). Moreover, as this correction has to be applied on all the points of a 448 sample, its contribution to the total error does not decrease with an increasing number of points. 449 Consequently, its relative contribution increases, and becomes the major part of the error on the 450 OH/Si ratio above around 20 to 25 single points, or 5 orthogonal couples (see Figure 5). This 451 power-drift correction error is also responsible for a significant part of the calibration error, thus 452 acting in two different ways in the final error. On the other hand, insufficient sampling of the 453 natural dispersion rapidly causes the final error to be larger than the measured dispersion (see 454 equation 11). Orthogonally polarized averaged measurements, improving by force the sampling 455 of various orientations (and thus, anisotropy) help to prevent from orientation biases.

456 Secondly, the effect of the calibration uncertainty (caused by the errors on each standard 457 point) on the final error acts similarly as the power-drift correction uncertainty (both are

458 insensitive to the number of points), rapidly becoming the major part of the error, particularly for 459 olivine. Focusing on the ideal case where the dispersion observed on the sample  $(D_{Smp})$  equals the 460 theoretical dispersion  $(D_{Th})$ , implying that  $E_{Smp}$  and  $E_{Th}$  are equal, the conversion factors  $D_{H2O}$ 461  $E_{H2O}$  (following equation 15, and applying this to the dispersion by replacing  $E_{Smp}$  with  $D_{Smp}$  to obtain  $D_{H2O}$ ) and  $D_{Smp}/E_{Smp}$  of wadsleyite are superior to those of olivine in all cases (except for 462 463 two points in the orthogonal couples case, see Figure 7). This is caused by the higher effect of the 464 error sources on  $D_{Th}$  than on  $E_{Th}$  in the case of wadsleyite, where  $E_{Cor}$  mitigates their effect, hence 465 increasing  $D_{Smp}/E_{Smp}$ . Adding  $E_{Cal}$  does not change this observation, except that it mitigates these 466 effects, lowering the expected "gain" on errors, with higher attenuation for higher  $E_{Cal}$  (see Figure 467 7). In the case of a poor sampling, accounting for  $E_{Cal}$  tends to mitigate the effect of the error 468 augmentation (see equation 9) by increasing  $E_{Smp}$  to a much higher  $E_{H2O}$ , lowering the effect of 469 equation 15.

470 Water content heterogeneities may cause the measured dispersion to be larger than the 471 expected one. In such cases, when the ratio of observed over expected dispersion is close to one, 472 even if the gain may seem high, it is important to consider that chemical heterogeneities should 473 remain as low as possible (see the discussion around equation 13). The present water 474 quantification method by Raman spectroscopy relies on the hypothesis that only the orientation 475 varies from one crystal to another, and that these orientations are distributed randomly and 476 uniformly. This means that chemically heterogeneous samples (should it be a major element 477 content variation, or in water content variation) are unsuitable for precise water quantification 478 with this method. In the case of wadsleyite, we observed that the iron content has a tremendous 479 effect on the calibration line slope, implying that composition has to be measured to verify if all 480 the analyzed crystals have a similar iron content. The problem of the random distribution of

481 orientations in samples also has to be considered, as preferential orientation may occur in some 482 samples. However, if a sample is chemically homogeneous, a sample with a strong disparity of 483 orientations could be identified if it displays values dispersion much smaller than the expected 484 dispersion (see discussion around equation 12).

485 **CONCLUSION** 

486 Raman spectroscopy allows the water content measurement of samples of any size above 487 few micrometers, and with water contents down to 150-200 ppm in weight (see Bolfan-Casanova 488 et al. 2014), with no observed upper limit, where many other methods may require large samples, 489 or are unsuitable for very high water contents (FTIR). Even if vibrational spectroscopy is very 490 useful to discriminate OH point defects over contamination, the sensitivity to crystallographic 491 orientation in anisotropic samples often leads to high standard deviations in the measurements of 492 water content (Férot and Bolfan-Casanova 2012). Here we show that a large part of this 493 dispersion may be related to the natural anisotropy of the measured mineral, and is thus normal. 494 The main objective of this work was to be able to calculate the error associated to the 495 measurement of water content in olivine and wadsleyite, using Raman spectroscopy, knowing the 496 anisotropy of the OH/Si ratio for each phase. We identified throughout this study various sources 497 of error. The major parameters (beside anisotropy) affecting errors on the final values are the 498 uncertainties on the calibration line, insufficient sampling of anisotropy, and the laser-drift 499 correction errors. Nevertheless, most of the errors obtained with this procedure fall in a range of 500 20-25% of relative error for wadsleyite, and 25-30% for olivine, making Raman spectroscopy 501 still suitable for quantification water content in olivine and wadsleyite. The detailed study of the error sources has provided a greater understanding of the method, and has allowed an 502

503 improvement of the treatment methods, reducing significantly the uncertainties arising from this 504 part of the procedure.

# 505 **IMPLICATIONS**

The high spatial resolution of Raman spectroscopy allows the study of water distribution among different phases of fine polymineralic samples of complex (natural) composition, with a wide range of measurable water contents. Although relatively high, the uncertainties on water concentration are sufficiently low to infer the effect of thermodynamic intensive parameters on water incorporation. The method proposed here can also be applied to all orthorhombic minerals.

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633	Table 1: Fitting parameters $(D_{An} \text{ and } a)$ and estimated error sources values $(E)$ for
634	dispersion and error equations (equations 6 to 13). $D_{An}$ is the calculated anisotropy contribution to
635	the observed values deviation; $E_{Meas}$ is the measurement related error, $E_{Tr}$ the treatment related
636	error, and $E_{Cor}$ the laser-drift correction related error.

	0	livine	Wadsleyite			
	Single points	Orthogonal couples	Single points	Orthogonal couples		
D <sub>An</sub>	28.1	14.1	32.5	16.8		
а	6	3	6.74	3.37		
<b>E</b> <sub>Meas</sub>	4.44	4.44	4.25	4.25		
E <sub>Tr</sub>	3.58	3.58	3.82	3.82		
E <sub>Cor</sub>	11.9	11.9	11.9	11.9		

Table 2: Standards used for the olivine and wadsleyite calibrations of OH/Si versus  $H_2O$ concentration in ppm by weight of water. Wadsleyite\* standards are iron-free. Raman measurement type "1" refers to the case of single point measurement and "2" to the case of orthogonal couples.

		Source measurement			Raman measurement			
Sample	Phase	Method	ppm wt H₂O	error	Туре	Points	OH/Si	error
949	Olivine	FTIR	1304	290	1	8	0.01534	0.00371
M497	Olivine	ERDA	750	38	1	10	0.00405	0.00081
M589	Olivine	FTIR	1244	596	1	57	0.01127	0.00141
1033	Olivine	FTIR	522	176	2	11	0.00806	0.00118
895b	Olivine	FTIR	180	21	2	12	0.00382	0.00156
1044b	Olivine	FTIR	1766	674	2	10	0.01924	0.00250
M817	Olivine	FTIR	640	215	2	10	0.00350	0.00103
M818	Olivine	FTIR	291	72	2	10	0.00269	0.00104
M230	Wadsleyite	ERDA	1045	52	1	10	0.01591	0.00270
M226A	Wadsleyite	ERDA	4209	210	2	12	0.03863	0.00642
M380b	Wadsleyite	ERDA	12206	610	1	10	0.04219	0.01833
M382	Wadsleyite	ERDA	27271	1364	2	13	0.10122	0.01864
M226B	Wadsleyite*	ERDA	4000	200	2	10	0.10574	0.01457
2053	Wadsleyite*	ERDA	21600	1080	1	10	0.46945	0.10382
2054	Wadsleyite*	ERDA	34000	1700	1	10	0.49357	0.07957

# 643 **FIGURE CAPTIONS**

644 Figure 1 : Optical image of the olivine and wadsleyite samples used for the anisotropy 645 quantification.

646

Figure 2 : Raman spectra of olivine (grey) and wadsleyite (black) in the silicate region and the OH region. The wide bars underneath each graph (light bar for olivine, dark bar for wadsleyite) depict the integration window used for water quantification. The darker short lines depict the anchor point area used for baseline correction. The baseline shape is linear for olivine Si area, and cubic for OH area, and is polylinear for wadsleyite Si area, and cubic for OH area.

652

Figure 3 : OH/Si values of olivine (top) and wadsleyite (bottom) as a function of the orientation of the crystal relative to the beam, for three perpendicular faces. Error bars are the estimated uncertainties linked to measurement errors and data treatment (see text). Horizontal error bars are an estimation of the angle's error. Solid and dashed lines (M) represent the fitted curves obtained through equation 3 (or 4 for wadsleyite). Olivine crystallographic faces were identified, whereas for wadsleyite, faces are named arbitrarily (F1, F2 and F3).

659

Figure 4 : Three-dimensional plot of OH/Si values of olivine (as fitted in Figure 3) as a function of crystallographic orientation. Values obtained with incident beam parallel to the *a* axis are displayed in the (100) (bc) plane, and so on for the other faces. Grey points (which due to

their density may be displayed as light grey lines here) are the extrapolations of the solid lines(derived from Figure 3) for any given orientation (see equation 5).

665

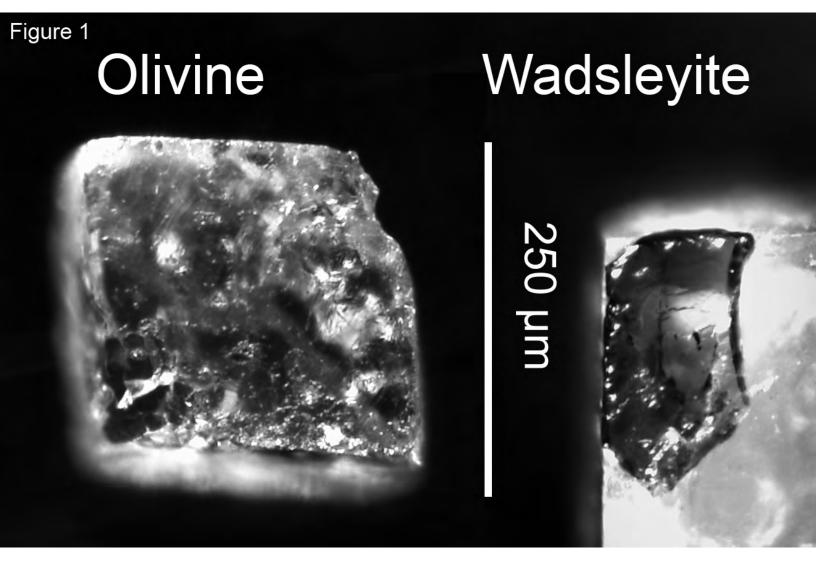
Figure 5 : Simulated relative dispersions (empty circles) and errors (solid circles) for olivine (top) and wadsleyite (bottom), for the case of single point measurements (left) and orthogonal couples (right), expressed in percent deviation from the average. Light grey symbols stand for the contribution of anisotropy (*An*) alone. Adding the measurement (*Meas*) uncertainties yields the intermediate grey symbols (mostly hidden by the dark grey ones). Adding treatment (*Tr*) uncertainties gives the dark grey symbols, and finally adding the uncertainty arising from the power correction (*Cor*) yields the black circles.

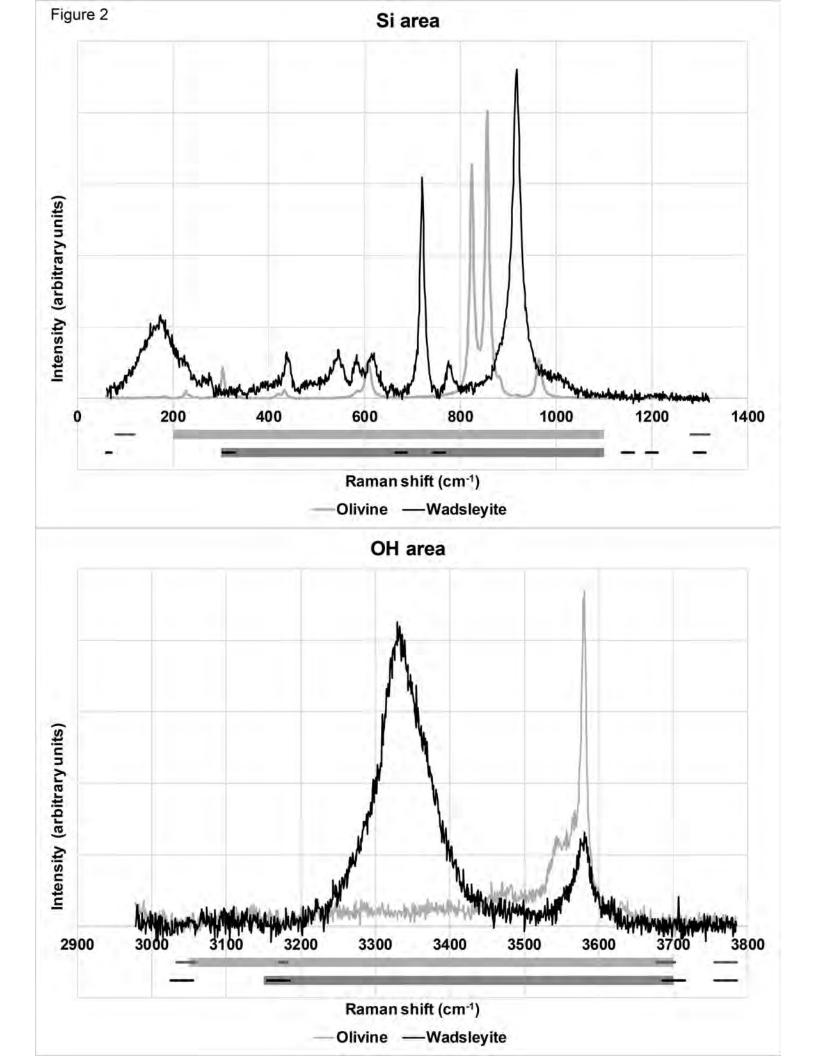
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Figure 6 : Calibration lines for olivine (top), iron-bearing wadsleyite (center) and iron-free wadsleyite (bottom). The straight line corresponds to the fitted value. The dashed and dotted lines show the  $1\sigma$  and  $2\sigma$  uncertainties respectively.

677

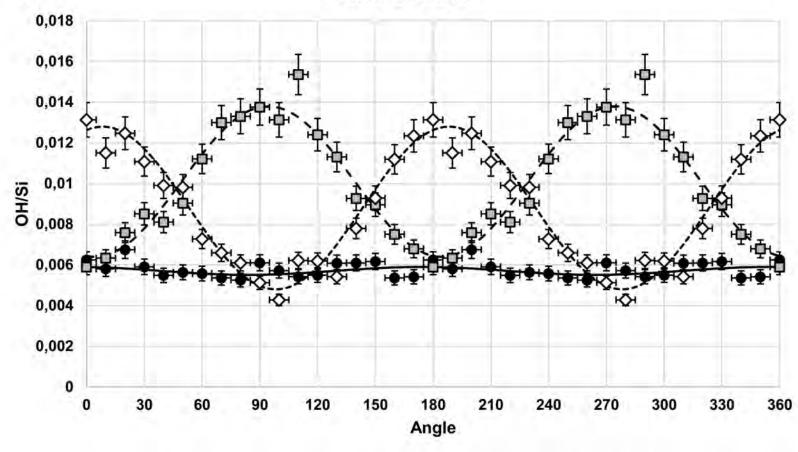
Figure 7 : Dispersion over error ratio (D/E) as a function of the number of points, for analytical uncertainties only (grey circles and discs, representing  $D_{Smp}$  and  $E_{Smp}$ ), and for total uncertainty, including calibration related uncertainty (black circles and discs, for  $D_{H2O}$  and  $E_{H2O}$ ). Full circles represent the ideal case where  $D_{Smp} = D_{Th}$  (ideal sampling), thin circles represent the insufficient sampling case (here,  $D_{Smp}$  is 75% of  $D_{Th}$ ), and thick circles represent the case of a dispersion  $D_{Smp}$  50% superior to  $D_{Th}$ .





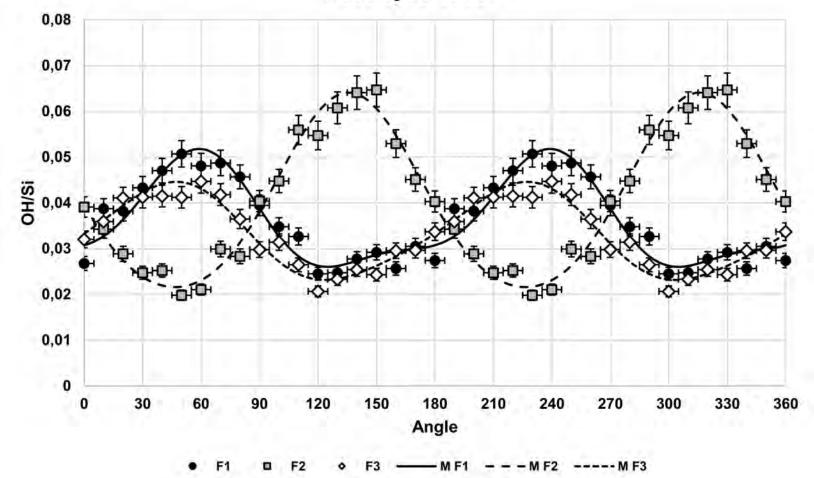


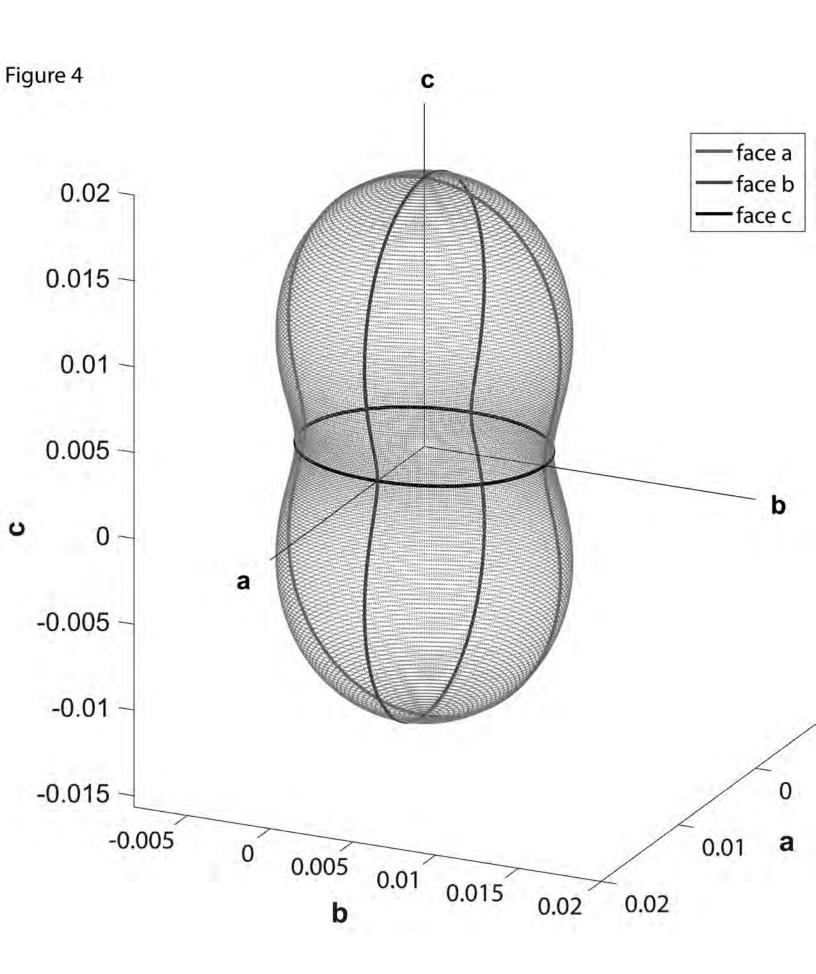
**Olivine OH/Si** 

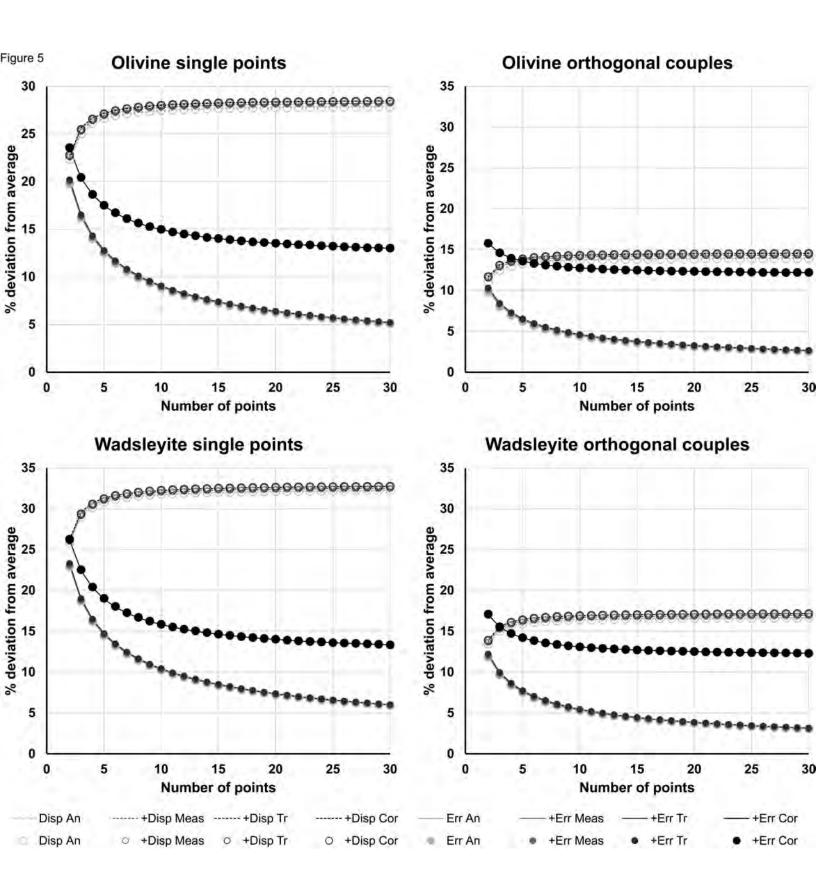


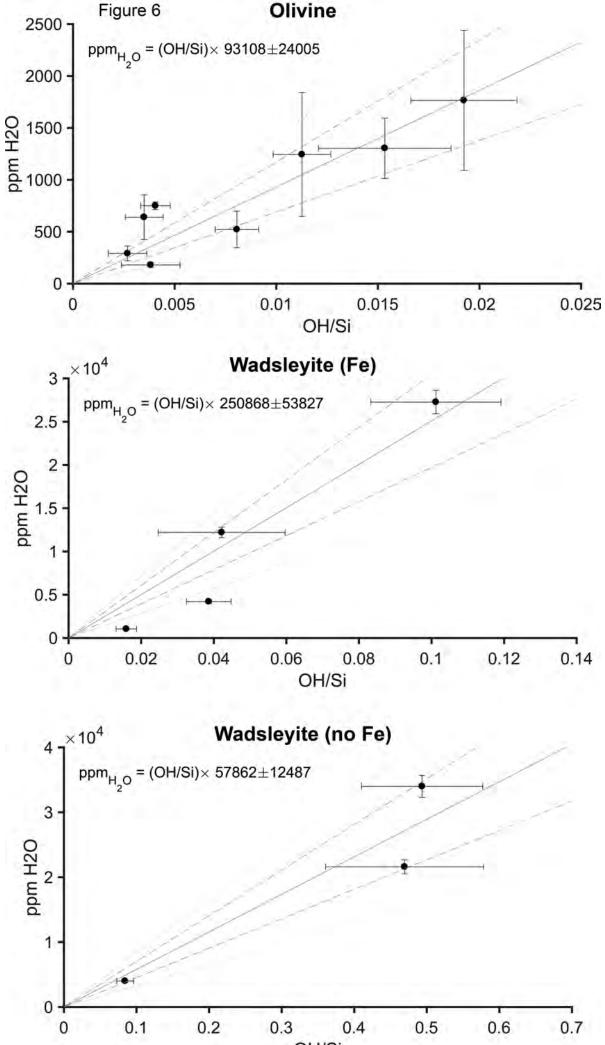
● c □ a ◆ b -----Ma -----Mb

Wadsleyite OH/Si









OH/Si

