Experimental quantification of the Fe-valence state on amosite-asbestos boundaries using acSTEM Dual-Electron Energy-Loss Spectroscopy

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Abstract

The determination of the oxidation state and the coordination geometry of iron expands the knowledge obtained by standard mineralogical characterization. It provides information that is crucial in assessing the potential of minerals to interact with the surrounding environment and to generate reactive oxygen species, which can disrupt the normal function of living organisms. Aberration-corrected Scanning Transmission Electron Microscopy Dual-Electron Energy-Loss Spectroscopy (acSTEM Dual-EELS) has only rarely been applied in environmental and medical mineralogy, but it can return data that are essential for the description of near-surface and surface
mechanisms involved in many environmental and health-related processes. In this study, we have applied the energy-loss near-edge structure (ELNES) and the L\textsubscript{2,3} white-line intensity-ratio methods using both the universal curve and progressively larger integrating windows to verify their effectiveness in satisfactorily describing the state of iron in amosite grain boundaries, and, at the same time, to estimate thickness in the same region of interest. The average valence state obtained from acSTEM Dual-EELS and from a simplified geometrical model were in good agreement, and within the range defined by the bulk and the measured surface-valence states. In the specific case presented here, the use of the universal curve was most suitable in defining the valence state of iron in amosite grain boundaries. The study of ELNES revealed an excellent correspondence with the valence state determined by the L\textsubscript{2,3} white-line intensity-ratio method through the use of the universal curve, and it seems that the spectra carry some information regarding the coordination geometry of Fe. The combination of visual examination, reconstruction of the grain boundaries through a simple geometrical model, and Dual-EELS investigation is a powerful tool for characterizing the grain boundaries of hazardous minerals and foreseeing their potential activity in an organism, with the possibility to describe toxic mechanisms in a stepwise fashion.

Keywords

Dual Electron Energy-Loss Spectroscopy; Fe-valence state; amphibole; asbestos; surface chemistry; spatially resolved crystal-chemistry
Introduction

The valence state of Fe in members of the amphibole supergroup plays an important role when elucidating mineralogical and geological history since it can be influenced by multiple geochemical reactions in a wide range of conditions (Cavé et al., 2006). It can further contribute to the understanding of the redox conditions in which these minerals crystallized (Garvie and Busek, 1998), and of weathering, dissolution and recrystallization processes to which these minerals were exposed in the natural environment. When in contact with living cells, the mixed-valence state of Fe in amphiboles can determine the mineral’s potential in disrupting Fe homeostasis and in generating reactive oxygen species (ROS) in cells and at the systemic level (Jablonski et al., 2017). The generation of ROS may lead to several respiratory diseases (e.g., asbestosis, pleural abnormalities) and malignancies, such as squamous cell carcinoma, small-cell and large-cell carcinoma, and adenocarcinoma (Pacella et al., 2012), including mesothelioma, which is strictly related to asbestiform minerals. Amphiboles that are not necessarily classified as asbestos or asbestiform were recently also suggested to be related to autoimmune disease (Li et al., 2012; Ferro et al., 2013; Pfau et al., 2014), which makes the description of the toxicity caused by these minerals even more pressing. During interaction with the biological environment, Fe in asbestiform minerals plays different roles depending on its oxidation state, coordination, and surface-site occupancy (Gualtieri et al., 2016). ROS may be generated when Fe is released into cells or biofluids, or during surface-bound Fe-promoted reactions (Schoonen et al., 2006). Therefore, the possibility to determine both valence and bonding of Fe at an atomic- or nano-scale allows for a detailed description of the mechanisms that can potentially lead to generation of ROS, understanding the transformation of these minerals, and foreseeing the sequence of
events causing related pathology, while bearing in mind that Fe is just one of the variables contributing to these mechanisms.

The determination of Fe-valence state, coordination geometry and local morphometry in a specific observable region of interest (ROI) of a (potentially) hazardous mineral can thus provide a very important contribution to the comprehension of the early mechanisms by which diseases may be triggered.

The use of aberration-corrected scanning transmission electron microscopy (acSTEM) and Dual-Electron Energy-Loss Spectroscopy (Dual-EELS) allows for simultaneous collection of two spectra at different energy-loss ranges in a specific ROI on individual particles. This translates to an immediate alignment of the zero-loss peak and, therefore, side-steps the need to collect two consecutive spectra on the same sample area to perform the alignment (Potapov and Schryvers, 2004), resulting in a shorter exposure of the ROI to the electron beam, thus reducing potential electron-beam damages and operational time. Application of Dual-EELS has a great advantage compared to conventional EELS because it eliminates the need of measuring the low-loss and the high-loss spectra consecutively, allowing for the \textit{in-loc\text{\textup{c}}} determination of the ROI thickness simultaneously with its valence state.

The Fe-valence state is commonly determined by techniques with poor spatial resolution, such as Mössbauer spectroscopy and wet chemical analysis, or by indirect stoichiometry-based calculations from electron probe micro-analysis data (e.g., Droop, 1987). Despite the accuracy of redox titration and thermogravimetric approaches in the determination of the oxidation state of an element in a phase, these techniques provide averaged and compound-dependent results (Lee et al., 1980) without satisfactory spatial resolution. Advanced spectroscopic techniques, such as X-Ray Absorption Near-Edge Structure (XANES), Extended X-Ray Absorption Fine Structure...
(EXAFS) or Synchrotron Mössbauer Source (SMS) spectroscopy, can provide Fe-valence state
data that are accurate and comparable to those obtained from Electron Energy-Loss Near-Edge
Structure (ELNES), but again with a far lower spatial resolution. With respect to asbestos-related
disease and disease-mechanism studies, X-Ray Photoelectron Spectroscopy (XPS) has been used
to track the surface evolution of asbestos fibers in contact with media that simulate biofluids
(e.g., Pacella et al., 2015), limiting the investigation to the first ≈10 nm at the surface of the
minerals. Even though EELS has a higher spatial resolution than X-Ray Absorption
Spectroscopy (XAS) (Taftø and Krivanek, 1982) and related techniques, it has been used in a
limited manner in mineralogy, primarily because EELS is commonly intended as a technique
mainly dedicated to light elements and complementary to Energy-Dispersive X-ray (EDX)
spectroscopy (Andreozzi and Pollastri, 2017). However, EELS can provide additional and more
detailed information: acSTEM Dual-EELS is a powerful technique, which allows for
simultaneous acquisition of atomic- and nano-scale images and related data on chemical
composition, coordination and oxidation state for the exact same area. The technique, therefore,
provides invaluable local information on the nature of the observed specimen. Development and
application of this technique, thus, give an opportunity to describe surfaces, mineral interfaces,
and potentially the atomic-scale sequence of processes within the mineral or at the interface
between the mineral and the surrounding chemical environment.

In this study, we systematically apply methods that allow us to determine the Fe-oxidation state
from Dual-EELS spectra recorded in several ROIs of naturally occurring amosite, which has
been characterized previously (Pollastri et al., 2015). By using acSTEM, we were able to select
and orient fiber boundaries to obtain accurate and precise, spatially resolved information. In
addition, we used XPS as an independent technique to assess and discuss the results obtained through acSTEM Dual-EELS on the fibers.

The most common procedure used to determine the valence state of transition metals through EELS is known as the white-line ratio or L_{2,3}-ratio method. In this procedure, the formal valence state is correlated with the L_{3}- and L_{2}-white lines emerging from the spin-orbit splits that correspond to the 2p_{3/2}→3d and 2p_{1/2}→3d transitions (Colliex et al., 1991; Riedl et al., 2007).

The 3d transition metal white-line intensity ratio has been demonstrated to be different from the expected value of 2 (Leapman and Grunes, 1980). This is a consequence of the interaction between the 3d valence electrons, the ejected electrons, and the core hole. The intensity ratio I(L_3)/I(L_2) assumes maximum values in the 3d^5 configuration, decreasing toward both the 3d^0 and 3d^{10} configurations (Sparrow et al., 1984; Thole and Van der Laan, 1988). This simply results in a proportional reduction of the white-line intensity ratio while the Fe-valence state approaches its minimum value. The white-line intensity-ratio procedure has several limitations since the I(L_3)/I(L_2) does not necessarily always have a one-to-one correspondence to the Fe-valence state (Sparrow et al., 1984; Graetz et al., 2004; Tan et al., 2012). One of the reported problems is related to the sample thickness, which is described as \( t/\lambda \), where \( t \) is the average thickness of the investigated ROI and \( \lambda \) is the inelastic mean free path of the electrons passing through the analyzed region. The removal of plural-scattering through deconvolution or the investigation of a specific area with \( t/\lambda < 0.5 \) should be satisfactory in limiting the thickness influence on the results (Schmid and Mader, 2006; Varela et al., 2009). While some authors suggest that the I(L_3)/I(L_2) is relatively insensitive to thickness (Wang et al., 2000) below \( t/\lambda < 0.5 \) (Varela et al., 2009), other authors sustain that spectral deconvolution is necessary to obtain an
accurate evaluation of the valence state (Tan et al., 2012). On the other hand, Loomer et al. (2007) reported that there was no thickness effect for $t/\lambda$ values slightly greater than 1.

In our contribution, the white-line intensity-ratio investigation has been coupled with ELNES in the spectral regions of the O–K and Fe–L$_{2,3}$ edges in an attempt to retrieve information on the Fe-coordination geometry and the contribution of other factors to the relation of spectral shape with the valence state. The applicability of the so-called “universal curve” (Van Aken et al., 1998) was compared to the use of different calibration curves obtained from standards with known valence states. This comparison was made because the investigated ROIs are characterized not only by the complex crystal structure of amphibole but also by the presence of an amorphous shell, which can induce unexpected edge modifications in the ELNES spectra that in turn may affect the extracted valence-state values. In particular, the universal curve uses a narrow window (2eV wide), which can be influenced by both spectral noise and small differences in the ELNES shape (Tan et al., 2012). In view of these possible complications we tested several approaches to assess which L$_{2,3}$ method and integration-window width are most suitable in describing the complexity of the amphibole-fiber boundaries and their modified surfaces.

In addition, we created a simple geometric model to estimate the valence state at particle boundaries starting from Mössbauer spectroscopy and XPS data so that it could be compared with the results obtained through acSTEM Dual-EELS.

The possible influences of thickness and beam damage on the results have been considered.

**Materials and methods**
Samples of amosite (or asbestiform grunerite) were investigated by two different techniques, XPS and acSTEM Dual-EELS. We need to emphasize that the XPS data were obtained by an independent laboratory (ENEA, Casaccia Research Center, Materials and New Technologies Unit, Rome, Italy), and the data were not made accessible to us until after we had completed our acSTEM Dual-EELS investigation. This approach was chosen to guarantee a non-biased, independent evaluation of our acSTEM Dual-EELS data.

The XPS spectra collected on the bulk fibers, i.e. not separated into short and long fibers, were acquired with a V.G. ESCALAB MKII spectrometer and recorded using an Mg anode (Kα radiation line at 1253.6 eV, non-monochromatized) as X-ray source at a voltage of 10 kV and a power of 240 W. Spectra were collected in the constant-analyzer energy mode at a pass energy of 10 eV and a step channel of 0.1 eV. The sample was suspended in analytical-grade acetone and gently ground in an agate mortar. The powdered material was then transferred onto a double-sided sticky tape on a standard sample holder. To determine the oxidation state of iron, the Fe–2p3/2 peak was deconvoluted following the approach described in Fantauzzi et al. (2010).

Two samples of amosite were investigated by acSTEM Dual-EELS, one consisting of short fibers, the other of long fibers. These samples were prepared for this study as described in Pollastri et al. (2014) and from the same batch that was also analyzed by XPS. The two samples were suspended in 2-propanol and then transferred onto separate 300 mesh holey-carbon TEM copper grids (SPI supplies, West Chester, PA, U.S.A.). All experiments were conducted at an operational voltage of 80 kV with an aberration-corrected STEM (JEOL, model ARM 200 F) equipped with a high-brightness Cold-Field Emission Gun (CFEG) and an energy filter (Quantum ER GIF, Gatan, USA) with Dual-EELS capability. The instrument was tuned to a collection semi-angle of 22.65 mrad (8 cm camera length – working distance) to increase the...
intensity of the signal on core-loss edges (Colliex et al., 1991) and a convergence semi-angle of 24.00 mrad, using an aperture of 5 mm and a dispersion of 0.25 eV/channel. EELS was performed in dual mode on 100 separate, randomly selected areas of 625 nm² each. These areas were located along the grain boundaries of the fibers (50 ROIs each in the short- and long-fiber grids). The first spectral region was recorded over a period of $2 \cdot 10^{-4}$ s, whereas the core-loss region was registered for a duration of 10 s, and in both cases, 3 frames were summed up to generate a spectrum. The recording of 3 consecutive frames allowed us to first screen for the possible “morphological” transformation of ELNES due to beam damage resulting from the exposure of a ROI to the electron beam. The first spectral region ranged from -50 to 462 eV, the core-loss region of the spectra from 500 to 1012 eV, which allowed for simultaneous visualization of the O and Fe edges.

**Determination of the ROI thickness**

The low-loss region was used to determine the relative sample thickness ($t/\lambda$). The relative thickness of each individually analyzed ROI was calculated by determining the ratio between the zero-loss electrons and the total transmitted intensity, according to Poisson’s statistics (Malis et al., 1988). All the ROIs with $t/\lambda > 0.5$ were arbitrarily discarded from the study. The real sample thickness was then calculated by estimating the $\lambda$ value for amosite crystals through the Iakoubovskii et al. (2008) formulation, which is appropriate for a collection semi-angle ($\beta$) greater than 20 mrad. The core-loss region was used to study the O- and Fe-ELNES, and to determine the Fe-oxidation state by applying the selected methods. These data, together
with literature data, imaging of fiber boundaries, and morphometry were also used to reconstruct
the geometry and estimate the expected valence state of the studied amphiboles.

$L_{2,3}$-ratio method

As a first step, we used an inverse power law to remove the pre-edge background of the collected
core-loss EEL spectra (Supplementary S.1). The intensity due to transitions into the continuum
was subtracted by applying a double arctangent step function as continuum for the removal of
the post-edge background (Supplementary S.1).

Three progressively larger integration windows (Supplementary S.2) were tested on spectra
obtained on standards using the script “Double Atan EELS Background” created by Mitchell
(2015). In the first case, we used two integrating windows (2 eV wide), one located at the
maximum of the L$_3$ edge, and one located at the L$_2$ edge, as suggested by the modified integral
Fe-L$_{2,3}$ white-line intensity method proposed by Van Aken et al. (1998), analogous to the
universal curve. In the second case, we applied two integrating windows, which were centered at
the peak maxima and which were 4 eV wide (Schmid and Mader, 2006). In the third case, we
used two 8 eV wide windows (similar to Tan et al., 2012) with the lower limit of the windows
located just after the onset of the L$_3$ and L$_2$ edges, respectively (which frequently corresponded
to centering the windows on the peak maxima).

This procedure was used to determine the correspondence between L$_3$/L$_2$ ratios and the valence
state in three standards with the known Fe-oxidation states of 2+, 2.7+, and 3+ (Supplementary
S.3). The L$_3$/L$_2$ ratios versus valence-state data were fitted using an exponential function to
obtain calibration curves. The universal curve (Van Aken et al., 1998) and the equations of the calibration curves were then used to determine the average valence state of ROIs on amphibole fiber boundaries. In this study, the calibration curves were generated using the same standards and the same Atomic Resolution Microscope (ARM) as in Rojac et al. (2017), because the stability of the standards in terms of valence state was already verified on the same instrument.

Results

XPS reference analyses

The XPS analysis, performed on fibers that were not separated into two dimensional categories, revealed a predominance of Fe$^{3+}$ in the forms of Fe$^{3+}$ and of FeOOH (Table 1). From these data, the corresponding average Fe-valence state was calculated at 2.67 with $\sigma_{n-1}=0.04$.

ROI thickness, features and geometrical model for amosite

The fibers exhibit straight boundaries, which are in many cases covered by a discontinuous and partially amorphized “shell” of variable thickness (Fig. 1A). This “shell” is very irregular, presenting a rough surface, and exhibits an average thickness of 3.74 nm ($\sigma_{n-1}=2.30$) for short fibers and 1.99 nm ($\sigma_{n-1}=0.88$) for long fibers. The overall thickness of the investigated ROIs at the fiber boundaries, evaluated from the low-loss spectra, varies between 5.73 and 37.56 nm, with an average value of 20.09 nm ($\sigma_{n-1}=9.68$). The average ROI thickness of the short fibers is 16.25 nm ($\sigma_{n-1}=7.69$), whereas that of the long fibers is 25.30 nm ($\sigma_{n-1}=9.77$). All those areas, where we were not able to orient the crystal or where the thickness $t/\lambda$ was $>$0.5 (roughly
corresponding to \( t > 50 \text{ nm} \), were discarded from the subsequent analyses, effectively decreasing the total number of useable areas from 100 to 66 (38 for short fibers, 28 for long fibers).

A geometric model of the ROIs for short and long fibers was created based on the average ROI thickness, average “shell” thickness, and the characteristic cleavage angles of the amphiboles (Fig. 1B).

In our simple geometrical model, we assumed that the features seen in the section perpendicular to the amphibole’s \( c \)-axis remain the same when moving along \( c \)-axis, namely average oxidation state, “shell” thickness, and morphometry. The section perpendicular to the amphibole’s \( c \)-axis is represented as a trapezoid with acute angles of 62° and obtuse angles of 118°. The ROI width is treated as the height \( (h) \) of the trapezoid, whereas the ROI thickness is equivalent to \( \frac{B + b}{2} \), where \( B \) is the long base and \( b \) the short base of the trapezoid. The bases are parallel to the (010) crystallographic plane.

Furthermore, if we assume – as a simplification – that the external shell is completely oxidized and the bulk material is completely reduced, the calculation of the ratio between the perpendicular-section area of the “shell” \( (A_{\text{shell}}: \text{red in Fig. 1B}) \) and the total perpendicular section area \( (A_{\text{tot}}: \text{red plus green in Fig. 1B}) \), yields a rough estimate of the expected average valence state, i.e., \( \text{valence state} = \frac{A_{\text{shell}}}{A_{\text{tot}}} \), for the given geometry and “shell”-to-bulk ratio (1st approach).

As a second approach, we can calculate the expected average valence state assuming that the “shell” has a Fe-valence state of 2.67+ (XPS data), whereas the bulk valence state of amosite, determined by Mössbauer spectroscopy, is 2.08+ (Pollastri, 2015). This approach should be
valid, because the XPS data are usually representative of the first 0.3 to 3 nm (Vansant, 1995),
and possibly up to 10 nm of the material starting from the surface (Elmi et al., 2016).
Both approaches returned similar results, documenting that the near-surface regions of the short
fibers are more oxidized than those of the long fibers (Tab. 2).

Fe–L$_{2,3}$ white-line intensity ratio: standard calibration curves

The calibration curves for the Fe–L$_{2,3}$ white-line intensity ratio were constructed by fitting an
exponential function to the $I$(L$_3$)/$I$(L$_2$) data points obtained from the following standards (Fig. 2):
FeTiO$_3$ with all Fe as Fe$^{2+}$ (99.8%, Alfa Aesar, product number 1317), Co$_{0.6}$Fe$_{2.4}$O$_4$ with Fe as
Fe$^{2.7+}$ (synthesized according to Gyergiek et al., 2010), and Fe$_2$O$_3$ with all Fe being Fe$^{3+}$
(99.945%, Alfa Aesar, product number 14680).
The average L$_{2,3}$ ratios and $\sigma_{n-1}$ values for each standard were determined using multiple
analytical points (Fig. 2 and Supplementary S.4). The calibration curve for the standard data
points, determined using a 4 eV window (blue dotted line in Fig. 2), shows the smallest $\sigma_{n-1}$ with
respect to the L$_3$/L$_2$ ratio. On the other hand, when the error is calculated for the valence state
(horizontal error bar in Fig. 2), the 8 eV window shows the smallest values (excluding the Fe$^{2+}$
data points for which the $\sigma_{n-1}$ ranges are comparable). The smaller horizontal error bar of the 8
eV calibration curve compared to the 4 eV calibration curve for the Fe$^{2.7+}$ and Fe$^{3+}$ data points is
a consequence of the steeper slope of the first one (0.54) compared to the latter one (0.40). The 2
eV calibration curve has the highest slope value (1.27), which limits the error propagation on the
valence state, but the large error with respect to the L$_3$/L$_2$ ratio generates the largest error on the
calculated valence state (Fig. 2 and supplementary figure S.4.1.). In the next section, the valence
Fe–L₂,₃ white-line intensity ratio: amosite ROIs

The overall average valence states of the investigated ROIs of amosite, obtained by applying the white-line intensity-ratio method to the spectra of short fibers and long fibers, are similar when using the universal curve, the 4 eV window calibration curve or the 8 eV window calibration curve, but the use of a 2 eV window calibration curve returns slightly lower Fe-valence state values (Tab. 3 and supplementary S.5). The σₙ₋₁ values are in each case larger for short amosite ROIs than for long amosite ROIs.

The use of the \( I(L_3)/I(L_2) \) obtained with 4 eV windows to determine valence state of an ROI, showed a good match with the observed shape and intensity of the L₃ amphibole edges. The application of the universal curve yields a similar overall average valence state to the one obtained by using 4 eV windows on analyses performed on short fibers (Table 3). The long-fiber ROIs have a similar average valence state when using the universal curve, the 4 eV window and the 8 eV window curves.

The valence state for the ROIs of the long fibers are distributed in an area of the chart (Fig. 3) that ranges from 2.10+ to 2.27+, whereas the valence state for the ROIs of the short fibers shows a large variation, ranging from 2.09+ to 3.10+ (Fig. 3). The average valence state for the ROIs of
both long and short fibers is within the bulk valence state (from Mössbauer spectroscopy) and
the surface/near-surface valence state (from XPS).

The valence state determined on the amphibole ROIs using the universal curve represents the
best match with regard to the intensity and shape of the L3–edge ELNES and thus, was used for
all further evaluations of ELNES spectra (see next section and Discussion).

O–K edge ELNES of amosite

A representative core-loss spectrum (Fig. 4) of an amosite ROI with an Fe-valence state of
2.97+, as assigned through the use of the universal curve (red spectrum in Fig. 4), is
characterized by a pre-edge peak, labeled as (a), a dominant peak (b), a weaker maximum at
≈545-550 eV (c) followed by a large bump (d) (not shown), consistent with literature data
(Colliex et al., 1991). The pre-edge peak has two main components: one located at 528.25 eV
(σn-1=0.33), the other one at 529.51 eV (σn-1=0.62). The detected pre-edge is not observed in
amosite ROIs with a valence state that is lower than 2.44+, as clearly visible in a representative
spectrum from a ROI with a valence state of 2.12+ (green spectrum in Fig. 4). Spectra with a
mixed valence state ≥2.44 show variable peak intensity in the O-K pre-edge region (e.g., orange
spectrum in Fig. 4).

Fe–L2,3 edge ELNES of amosite

The Fe–L3 and Fe–L2 edges show features that can be attributed primarily to the oxidation state
and, to a lesser extent, to the coordination geometry of Fe in amosite-fiber boundaries (Fig. 5).
The L3 edge is composed of a first peak located at 709.74 (σn-1=0.14) eV, and a second peak
located at 711.45 (σn-1=0.15) eV. These two peaks are usually labeled as Fe\(^{2+}\) peak and Fe\(^{3+}\) peak, respectively since their intensity is characteristic of the related valence state. The L\(_3\) edge of the Fe\(^{3+}\)-dominant amosite ROI (red spectrum in Fig. 5) shows a first, less intense shoulder followed by a peak maximum. In the Fe\(^{2+}\)-dominant L\(_3\) edge (green line in Fig. 5), the shape of the edge is the inverse, with a more intense first peak, followed by a second, less intense shoulder at higher energy-loss values. The mixed-valence state, assigned using the universal curve, of a representative amosite ROI in between these two extremes (≈2.44+, orange line in Fig. 5) shows an L\(_3\) edge composed of two peaks with similar intensity.

The L\(_2\) edge shows a similar trend, but it tends to be sharper at lower valence states compared to that at higher valence state. The first component of the L\(_2\) edge, representing the more intense peak for Fe\(^{2+}\)-dominant ROIs, is located at 722.19 (σn-1=0.27) eV. The second, less intense component is located at 724.37 (σn-1=0.18). The intensity maxima are reversed for highly oxidized ROIs (e.g., red spectrum of Fe\(^{2.97+}\) in Fig. 5).

The average separation of the L\(_3\)- and L\(_2\)-edge maxima is 12.54 eV (σn-1=0.33).

**Thickness effect**

Since the white-line intensity-ratio method may be influenced by the thickness of the studied ROI, we investigated this possibility by plotting the L\(_3\)/L\(_2\)-intensity ratio versus thickness (Fig. 6). The plot of the obtained L\(_2,3\) ratio versus thickness, however, did not reveal any correlation between the two quantities. This diagram further documents a large vertical dispersion of the data points corresponding to a similar thickness.
The acSTEM observations show that the amphiboles are far from exhibiting perfect crystal shapes at their grain boundaries. The crystals are in many cases covered by Fe$^{3+}$-dominated, crystalline or partially amorphized material of variable thickness. The visually observed features are consistent with those observed by ELNES of both the O–K and the Fe–L$\text{2,3}$ edges, as well as the associated peaks, and by the valence state, which was determined applying the universal curve. The O–K pre-edge, observed in representative ROIs (Fe$^{2.44+}$ and Fe$^{2.97+}$), is characteristic of ROIs with a valence state equal or greater than 2.44$^+$. This pre-edge region of the O–K ELNES spectra is worth exploring in more detail: the sharp peak at $\approx 530$ eV is related to a high concentration of high-valence 3d transition metals (de Groot et al., 1989). In our experiment, the pre-edge peak was located at 529.51 eV ($\sigma_{n-1}=0.62$) and is indicative of Fe$^{3+}$-rich ROIs, because this peak is visible only in areas with an average valence state $\geq 2.44^+$. This peak is assumed to originate from the transition of O–1s core states to O–2p states, which are hybridized with the Fe–3d orbitals (de Groot et al., 1989; Kurata and Colliex, 1993; Garvie 2010). The other pre-edge component, located at $\approx 528$ eV (OH peak), is more controversial. The hypothesis that a 528 eV peak can be directly related to the quantity of OH or H$_2$O in minerals was refuted by Garvie (2010). The peak we observed in Fe$^{3+}$-dominant ROIs of amosite at 528.25 eV ($\sigma_{n-1}=0.33$) was stable in terms of intensity and did not show the sequence of “appearance” and “disappearance” described in the literature for the so-called OH peak located at $\approx 528$ eV (Garvie, 2010). The observed stability of this peak may be related to several factors, including (i) the relatively low electron dose we applied, estimated at $\approx 1\times10^3$ electrons/nm$^2$, and (ii) the peculiar and complex structure at the grain boundaries of the studied fibers. Since we never observed this “OH peak” in the O–K pre-edge region of spectra collected.
in “shell”-free ROIs, we conclude that it is characteristic of the oxidized “shell” that covers the mineral. This peak may therefore be interpreted as a qualitative indicator of the presence of OH in the “shell” or at the boundary between the “shell” and the real amphibole surface, or as an indicator of a continuous and small loss of hydrogen because of a low electron fluence in scanning mode, which results in a stable intensity of the OH peak. At this level of investigation and conditions, we do agree with Garvie (2010) that it is not possible to use the observed peak for quantification purpose, but this peak could indicate that the Fe$^{3+}$-rich shell is hydroxylated. On the other hand, we believe that the beam damage, if present, is negligible, since we did not observe a direct increase or decrease in intensity of this peak and the Fe-valence state remained stable. When interrogating the literature for materials that are similar to the investigated shells, a “safe” electron dose that avoids reduction on synthetic 6-line ferrihydrite was determined to be $1 \cdot 10^8$ electrons/nm$^2$ (Pan et al., 2006), whereas mixed-valent Fe (oxyhydr)oxides (green rust) showed a threshold of $4 \cdot 10^3$ electrons/nm$^2$ where oxidation begins (Freeman et al., 2019). We thus conclude that at this level of investigation and with our experimental setup, we can reasonably exclude an influence of beam damage on our results.

The shape of the Fe$^{3+}$–L$^3$ edge is easily recognizable due to the presence of an intense peak at 711.45 (σ$_{n-1}$=0.15) eV and the visibility of a less intense peak (shoulder) at lower energy-loss values (709.74 ± 0.14 eV). On the other hand, the Fe$^{2+}$-dominant L$^3$ edges show an opposite trend, with a more intense first peak and a second, less intense shoulder. In the case in which the valence state approaches 2.50+, the two main components of the L$^3$ edge show similar intensities (Fig. 5). As outlined in the results section, the ELNES shape has an acceptable match with the valence-state values obtained using the calibration curve with a window at 4 eV. Unfortunately, the use of this curve fails when evaluating a valence state that approaches 2.50+. The spectra to
which we assigned an oxidation state of 2.44+ using the universal curve (Fig. 5) has a first intense peak at a lower energy-loss value ($709.74 \pm 0.14$ eV) and a second less intense peak at $711.45$ ($\sigma_{n-1}=0.15$) eV, which is consistent with the assigned valence state and other spectra of similar minerals in the literature (e.g. ac50hd50 and ac40hd60 pyroxene in Fig. 2b, Van Aken and Liebscher, 2002). In the same spectrum, using the calibration curve with a 4 eV window rather than the universal curve would lead to a valence state of 2.53+. The assigned valence state, thus, would have been in disagreement with the observed ELNES shape since the intensity of the two components of the L₃ peak (located at $709.74 \pm 0.14$ eV, and $711.45$ $\sigma_{n-1}=0.15$ eV) indicates a slightly higher amount of Fe²⁺.

The measured peak separation between the L₃ and L₂ maxima on the amosite ROIs agrees with the work of Van Aken et al. (1998). The peak separations observed in ROIs of fiber boundaries are in the same range, but shifted by $\approx 0.25$ eV to lower electron-loss values. This shift is within the $\sigma_{n-1}$ range and thus not significant. This small deviation between literature values and experimental results may also be a consequence of symmetry variations or the presence of extra-fine structures (Colliex et al., 1991).

The presence of more distinguishable features occurring in the Fe²⁺–L₂ edges can be related to six-fold coordination. In more oxidized ROIs, the L₂ edges have poor features, a condition that can be related to a coordination geometry of 6 (structural Fe³⁺) or lower, which may correspond to the presence of amorphous material or surface-bonded Fe³⁺ (Van Aken and Liebscher, 2002).

In addition to valence state and coordination geometry, there is a possibility that the presence of Al³⁺ can shape the ELNES, leading to an increase in the intensity of the right flank of the Fe–L₃ edge located at $711.45$ ($\sigma_{n-1}=0.15$) eV, thus affecting the results of the white-line intensity-ratio calculation (Frost and Langenhorst, 2002; Langenhorst et al., 2013). In our case, we believe that
The influence of Al$_2$O$_3$ is minimal, if not negligible, since the Al$_2$O$_3$ concentration determined by wavelength-dispersive X-ray spectroscopy is only 0.04 wt % (Pollastri et al., 2014).

The use of a different size and position of the integration window when applying the Fe–L$_{2,3}$ white-line intensity-ratio method for the study of different compounds may affect the results depending on small differences in the ELNES shape (Tan et al., 2012). This effect may be larger for a narrow window, and could be influenced by both the ELNES shape and the spectral noise.

The use of wider windows for the definition of the calibration curve through the standards resulted in a reduced $\sigma_{n-1}$ for the windows at 4 eV and 8 eV for each of the valence states (2+, 2.7+, and 3+) of the standards, and thus a better fit of the calibration curve. The data series obtained with the 4 eV window showed slightly lower $\sigma_{n-1}$ values, but a poor qualitative correspondence to the spectral ELNES. The use of the universal curve returned a better correspondence among the measured valence state and the ELNES shape. Despite the similarity of average results obtained using the universal curve and the calibration curves derived from our standards (with the 4 eV and 8 eV windows), in this specific case the use of a larger window (8 eV) probably yields less good results because of the possible presence of additional structures in the low- and high-end tails of the considered edges generated by the influence of multiplet splitting and crystal-field effects (Van Aken et al., 1998).

Our experiment did not provide any evidence for a correlation between L$_{2,3}$ ratios and the sample thickness (Fig. 6), consistent with the results of other studies (Wang et al., 2000; Schmid and Mader, 2006; Loomer et al., 2007; Varela et al., 2009). However, we think that moving to higher magnification and approaching an atomic scale investigation, the thickness may considerably influence the results when applying the white-line intensity ratio, as noted by Tan et al. (2012).

The experimentally calculated average valence state on the surface of amosite fibers using XPS
Fe$^{2.67+}$ is considerably higher than the one reported for the bulk of the same sample (Fe$^{2.08+}$) (Pollastri et al., 2015). This discrepancy between the bulk and the surface, however, is not unusual, since surfaces are typically more oxidized than the bulk (Fantauzzi et al., 2010). Moreover, it reflects the presence of a hydrated Fe$^{3+}$-dominant “shell”, which we observed in acSTEM images.

Our Dual-EELS results revealed an average valence state of 2.44+ ($\sigma_{n-1}=0.31$) and 2.19+ ($\sigma_{n-1}=0.05$) for the studied ROIs in the short and long fibers, respectively, when using the universal curve. These results (Fig. 3) are within the range defined by the bulk amosite analysis (Fe$^{2.08+}$; Pollastri et al., 2015) and the data obtained by XPS for the surface and near-surface layers of amosite (Fe$^{2.67+}$, Table 1). Short fibers, characterized by a larger surface area ($\approx9.5$ m$^2$/g; Pollastri et al., 2014), are more likely to be heterogeneously covered by an oxidized “shell”, consistent with our observations and the larger range of valence states that we have recorded (Fig. 3 and supplementary S.5.1.). Long fibers, however, have a smaller surface area ($\approx3.9$ m$^2$/g; Pollastri et al., 2014), and a thinner or less extended “shell”, resulting in a lower average valence state and $\sigma_{n-1}$ (Fig. 3 and supplementary S.5.2.), which is a consequence of the major contribution of the nearly completely reduced bulk amosite. The Dual-EELS investigation allows a researcher to effectively relate visual observations to the valence state of a certain ROI, thus enabling a better description of the crystallochemical state and transformation of a fiber boundary compared to that derived from XPS data.

The use of Mössbauer spectroscopy and XPS is fundamental and faster when the valence state of an entire fiber population needs to be estimated. On the other hand, Dual-EELS provides unique information on the high heterogeneity of the valence state, geometry and amorphous shell...
coverage of the ROIs, and further allows for visualization of the crystal structure and boundaries of the analyzed crystals. We have shown a good correspondence between the valence state estimated through our geometrical model and the valence state obtained through Dual-EELS and the associated statistical errors ($\pm\sigma_{n-1}$): for short fibers, the valence state calculated from geometric considerations (1st approach) was $2.49+ (\sigma_{n-1}=0.44)$, i.e. similar to the real valence state of 2.44 ($\sigma_{n-1}=0.31$), as determined by Dual-EELS (universal curve). For long fibers, the valence state calculated from geometric considerations was $2.21+ (\sigma_{n-1}=0.10)$, again very similar to the one determined experimentally (universal curve), which was equal to $2.19+ (\sigma_{n-1}=0.05)$. As observed for the experimental values, the standard deviation calculated from the theoretical model for the ROIs of short amosite fibers is larger than the one for the ROIs of long amosite fibers (Tab. 2 and 3).

The 2nd geometric approach resulted in a valence state of $2.35+$ for short fibers and $2.15+$ for long fibers, both lower than the values obtained by the 1st approach and experimentally. Whereas the long-fiber valence state is similar for the experimental results and the geometry-based calculation (Difference of 0.04), there is a difference of 0.09 between the experimental valence state and the calculated valence state in short fibers when using the 2nd geometric approach. This discrepancy may be a consequence of the unknown depth of analysis of the XPS starting from the surface, causing the impossibility to know the “shell” versus bulk ratio using this approach.

Implications
The use of acSTEM Dual-EELS on amosite asbestos fibers represents a powerful tool, which allows for relating the morphometry, crystallochemistry, and valence state of a specific investigated area. The use of different integration-window widths and positioning may be dependent on the analyzed material, standards, and instrument, and thus it is important to test the universal curve versus different calibration curves and verify their correspondence to the observed ELNES shape. In this comparison, the coordination geometry and elements that may affect the ELNES shape (e.g., Al$^{3+}$) should be considered. We demonstrated that the Dual-EELS results for amosite boundaries are in good agreement with the valence-state estimates obtained using a simplified geometrical model, which was derived by combining the visual information acquired by acSTEM and the determination of the bulk (Mössbauer spectroscopy) and surface/near-surface (XPS) valence-state.

The determination of the Fe-valence state at the surface of mineral fibers is of pivotal importance since Fe may act as electron donor in ROS generation within biological systems. Surface-bonded Fe$^{2+}$ has been shown to react faster with molecular oxygen than dissolved Fe (Schoonen et al., 2006). This higher reaction speed can play a major role in disrupting the redox state of biological materials (e.g., cells, tissues, and biofluids). This mechanism, therefore, needs to be explored in greater detail, leading to a step-by-step description at the higher available spatial resolution, since surfaces are highly heterogeneous with respect to their topography, physical state (e.g., zeta potential) and chemical composition. This surface heterogeneity is poorly depicted using methods, such as XPS, but acSTEM Dual-EELS can provide fundamental information in this regard, which will help in understanding the overall processes that take place at the boundary between a mineral and its surrounding environment.
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Figure captions

Fig. 1

a) Micrograph collected in acSTEM bright-field mode showing an oriented amphibole-fiber boundary covered by an irregular amorphous “shell”. The ROI width is parallel to [010], whereas the ROI length runs parallel to the c* axis of the fiber. b) Geometric reconstruction of a fiber boundary oriented with respect to the electron beam. The green area is the amphibole bulk (Fe$^{2+}$ dominant), whereas the red area is the “shell” (Fe$^{3+}$ dominant). Both ROI length and ROI width are 25 nm long. The image proportions correspond to a long fiber.

Fig. 2

Calibration curves defined by the three standards using different integration-window widths. The universal curve (from Van Aken et al., 1998) is shown for reference (black continuous line). The red dashed line was determined by using a 2 eV integration window (fit to the data points shown as red dots), the blue dotted line was determined by using a 4 eV window (fit to the data shown as blue squares), and the green dash-dotted line was determined by using an 8 eV window (fit to the data shown as green triangles). Error bars are $\sigma_{n-1}$.

Fig. 3

The green dashed vertical line (M.s.) represents the bulk valence state (Mössbauer spectroscopy – Pollastri, 2015). The red dashed vertical line represents the near-surface valence state (XPS). The universal curve (dashed grey curve) together with the $\Delta$(Fe$^{3+}$/ΣFe) error range equal to ± 0.03 (grey dotted curves)(Van Aken, 2002) is shown for reference. The black crosses represent $L_3/L_2$ ratio and valence state for ROIs of the short fibers, whereas the yellow crosses represent the same data for ROIs of the long fibers. The blue squares and error bars ($\sigma_{n-1}$) show the average $L_3/L_2$ ratio and valence state for the ROIs of short fibers and long fibers. The yellow dashed vertical line represents the valence state obtained by the 1st geometrical model for long fibers (G.l.a.), whereas the orange dashed vertical line shows the valence state of the short fibers (G.s.a.).
Selected O–K edges corresponding to representative amosite ROIs for Fe$^{2.12+}$ (green), Fe$^{2.44+}$ (orange), and Fe$^{2.97+}$ (red). Valence states were assigned by using the universal curve. For clarity, feature labels are shown only in the Fe$^{2.97+}$ spectrum.

Selected Fe–L$_{2,3}$ edges (before arctangent background removal) corresponding to representative ROIs in amosite with valence states of 2.12+ (green), 2.44+ (orange), and 2.97+ (red), respectively. Valence states were assigned by using the universal curve. For clarity, L$_{2,3}$ edge labels are shown only in the Fe$^{2.97+}$ spectrum.

Relationship between the white-line intensity ratio and the thickness in different amosite ROIs ($n = 66$). Each ROI is represented by a square, the dotted line is the linear fit to the data points.
### XPS analysis

<table>
<thead>
<tr>
<th>Sample</th>
<th>Fe(II)/Fe$_{tot}$</th>
<th>Fe(III)/Fe$_{tot}$</th>
<th>FeOOH/Fe$_{tot}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Amosite</td>
<td>33 (3)</td>
<td>26 (4)</td>
<td>41 (4)</td>
</tr>
</tbody>
</table>

**Table 1:** Concentration of Fe$^{2+}$, Fe$^{3+}$ and FeOOH, as determined by XPS, compared to the total iron present in the analyzed region. The standard deviation ($\sigma_{n-1}$) is given in parentheses.
Geometrical model parameters and estimated valence state.

<table>
<thead>
<tr>
<th>Width of ROI ((h_{\text{trapezoid}}))</th>
<th>Thickness of ROI ((\frac{R+b}{2}))</th>
<th>Average “Shell” thickness</th>
<th>Fe-valence state 1st approach</th>
<th>Fe-valence state 2nd approach</th>
</tr>
</thead>
<tbody>
<tr>
<td>Short fibers</td>
<td>16.25 (7.69) nm</td>
<td>3.74 (2.30) nm</td>
<td>2.49 (0.44) nm</td>
<td>2.35 (0.48) nm</td>
</tr>
<tr>
<td></td>
<td>25 nm</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Long fibers</td>
<td>25.30 (9.77) nm</td>
<td>1.99 (0.88) nm</td>
<td>2.21 (0.10) nm</td>
<td>2.15 (0.14) nm</td>
</tr>
</tbody>
</table>

Table 2
Summary of all geometric parameters used to estimate the valence state of amosite ROIs using two different approaches. All values in parentheses are \(\sigma_{n-1}\).
### ROI characteristics summary

<table>
<thead>
<tr>
<th>Valence state</th>
<th>ROI thickness (nm)</th>
<th>White-line intensity ratio (Windows)</th>
<th>Geometrical model (Pollastri, 2015)</th>
<th>XPS</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>Universal curve (2eV)</td>
<td>1st</td>
<td>2nd</td>
</tr>
<tr>
<td>Short amosite</td>
<td>Mean</td>
<td>16.25</td>
<td>2.44</td>
<td>2.33</td>
</tr>
<tr>
<td></td>
<td>$\sigma_{n-1}$</td>
<td>7.69</td>
<td>0.31</td>
<td>0.38</td>
</tr>
<tr>
<td>Long amosite</td>
<td>Mean</td>
<td>25.30</td>
<td>2.19</td>
<td>2.01</td>
</tr>
<tr>
<td></td>
<td>$\sigma_{n-1}$</td>
<td>9.77</td>
<td>0.05</td>
<td>0.07</td>
</tr>
</tbody>
</table>

Tab. 3

Valence states obtained using progressively larger integrating windows for short and long amosite fibers. For the XPS measurement, the standard deviation ($\sigma_{n-1}$) is given in parenthesis.
Figure 1
Figure 2

Calibration curves and universal curve

![Figure 2](image_url)
Figure 4

O–K edge ELNES

Normalized intensity

Energy Loss (eV)

Fe$^{2.97+}$
Fe$^{2.44+}$
Fe$^{2.12+}$
Figure 5

Fe–L$_{2,3}$ edge ELNES

Normalized intensity

Fe$^{2.97+}$
Fe$^{2.44+}$
Fe$^{2.12+}$

Energy Loss (eV)
Figure 6

White-line ratio versus thickness (universal curve)

Slope = -0.01
R² = 0.07