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REVISION 1

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**A Mössbauer-based XANES calibration for hydrous basalt glasses reveals
radiation-induced oxidation of Fe**

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ABSTRACT

Oxygen fugacity (fO_2) exerts first-order control on the geochemical evolution of planetary interiors, crusts, and hydrospheres, and the $Fe^{3+}/\Sigma Fe$ ratios of silicate glasses provide a useful proxy for fO_2 . Fe K-edge micro X-ray Absorption Near-Edge Structure (XANES) spectroscopy allows researchers to micro-analytically determine the $Fe^{3+}/\Sigma Fe$

25 ratios of silicate glasses with high precision. In this study we characterize hydrous and
26 anhydrous basalt glass standards with Mössbauer and XANES spectroscopy and show
27 that synchrotron radiation causes progressive changes to the XANES spectra of hydrous
28 glasses as a function of radiation dose (here defined as total photons delivered per square
29 micron), water concentration, and initial $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio.

30

31 We report experiments from 8 different radiation dose conditions and show that Fe in
32 hydrous silicate glasses can undergo rapid oxidation upon exposure to radiation. The rate
33 and degree of oxidation correlates with radiation dose and the product of water
34 concentration and ferrous/ferric iron oxide ratio on a molar basis
35 ($\Phi = \text{XHO}_{1.5} * \text{XFeO}/\text{XFeO}_{1.5}$). For example, a basalt glass with 4.9 wt.% dissolved H_2O
36 and $\text{Fe}^{3+}/\Sigma\text{Fe} = 0.19$ from its Mössbauer spectrum may appear to have a $\text{Fe}^{3+}/\Sigma\text{Fe} \geq 0.35$
37 when analysed over several minutes at a nominal flux density of $\sim 2 \times 10^9$ photons/sec/ μm^2 .
38 This radiation-induced increase in $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio would lead to an overestimation of $f\text{O}_2$
39 by ~ 2 orders of magnitude, with dramatic consequences for the interpretation of
40 geological processes.

41

42 The sample area exposed to radiation shows measureable hydrogen loss, consistent with
43 radiation-induced breaking of O—H bonds, associated H migration and loss, and
44 oxidation of Fe^{2+} . This mechanism is consistent with the observation that anhydrous
45 glasses show no damage under any beam conditions. Cryogenic cooling does not
46 mitigate, but rather accelerates, iron oxidation. The effects of beam damage appear to
47 persist indefinitely.

48

49 We detect beam damage at the lowest photon flux densities tested (3×10^6
50 photons/sec/ μm^2); however, at flux densities $\leq 6 \times 10^7$ photons/sec/ μm^2 , the hydrous glass
51 calibration curve defined by the centroid (derived from XANES spectra) and $\text{Fe}^{3+}/\Sigma\text{Fe}$
52 ratios (derived from Mössbauer spectra), is indistinguishable from the anhydrous
53 calibration curve within the accuracy achievable with Mössbauer spectroscopy. Thus,
54 published $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios from hydrous glasses measured at low photon flux densities are
55 likely to be accurate within measurement uncertainty with respect to what would have
56 been measured by Mössbauer spectroscopy.

57

58 These new results demonstrate that to obtain accurate $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios from hydrous,
59 mafic, silicate glasses, it is first necessary to carefully monitor changes in the XANES
60 spectra as a function of incident dose (e.g., fixed-energy scan). Defocusing and
61 attenuating the beam may prevent significant oxidation of Fe in mafic water-bearing
62 glasses.

63

64 **INTRODUCTION**

65 The $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio of silicate glass may serve as a direct proxy for the oxygen fugacity
66 ($f\text{O}_2$) of an igneous system and exerts control on phase equilibria, physical properties, and
67 the valence state of other heterovalent elements in solution (e.g. Carmichael, 1991; Kress
68 and Carmichael, 1991; Sack et al., 1980). The use of Fe-K-edge XANES to obtain
69 $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios of hydrous silicate glasses has proliferated in recent years owing to the
70 outstanding precision achievable and micron-scale spatial resolution afforded (e.g.

71 Botcharnikov et al., 2005; Brounce et al., 2014a; Cottrell and Kelley, 2011; Gaetani et al.,
72 2012; Helz et al., 2017; Kelley and Cottrell, 2009; Kelley and Cottrell, 2012; Le Voyer et
73 al., 2015; Matjuschkin et al., 2015; Moor et al., 2013; Moussallam et al., 2014; Shorttle et
74 al., 2015; Stamper et al., 2014; Wilke et al., 2006). Most Fe-XANES calibration studies
75 have focused on anhydrous glasses, in large part because fO_2 can be easily controlled in
76 1-atmosphere gas-mixing furnaces (Berry et al., 2003; Cottrell et al., 2009; Dyar et al.,
77 2016; Fiege et al., 2017; Wilke et al., 2001; Wilke et al., 2005; Zhang et al., submitted;
78 Zhang et al., 2016). Given the recent surge in Fe-XANES-based research, however, it is
79 important to understand and quantify how dissolved water and radiation-induced
80 chemical reactions may modify the spectral features used to determine $Fe^{3+}/\Sigma Fe$ ratios.
81
82 Fe-XANES pre-edge features are sensitive to the coordination environment and oxidation
83 state of Fe (e.g. Henderson et al., 2014 and references therein). The potential effects of
84 water dissolved in silicate glass on the structural environment of Fe are therefore of great
85 interest to practitioners of XANES. In their Mössbauer studies on glasses, Wilke et al.
86 (2002, 2006) observed differences in the hyperfine fields and coordination environments
87 between hydrous and anhydrous glasses, but postulated that the differences emerged
88 during quench. Wilke and co-workers concluded that there is no detectable influence of
89 water on the local structural environment of Fe in melts. The effect of dissolved water on
90 the oxidation state of Fe in melts has been the subject of more study and debate, but
91 XANES and wet-chemical studies of mafic glasses synthesized at fixed fO_2 have shown
92 water to have no effect on $Fe^{3+}/\Sigma Fe$ ratios (Botcharnikov et al., 2005; Gaillard et al.,
93 2001; Moore et al., 1995). Given that Fe coordination affects pre-edge intensities, but not

94 the pre-edge centroids most commonly employed to extract $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios (Bajt et al.,
95 1994; Cottrell et al., 2009; Wilke et al., 2005), and the difficulty of independently varying
96 $f\text{O}_2$ and $f\text{H}_2\text{O}$ in experimental systems, relatively little effort has been applied to
97 calibrating hydrous glasses for XANES determinations of $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios.

98

99 Electron and X-ray beam interactions with materials can shift the oxidation state of
100 various elements in biological, amorphous, and crystalline materials (e.g. Corbett et al.,
101 2007; Eeckhout et al., 2005; Shimizugawa et al., 2001; Gonçalves Ferreira et al., 2013).
102 For example, sulfur dissolved in silicate glass may become oxidized or reduced upon
103 exposure to electron and X-ray beams (Klimm et al., 2012; Rowe et al., 2007; Wilke et
104 al., 2008). Relatively little work has been done to investigate the potential effects of
105 radiation on Fe speciation in silicate glasses. Wilke et al. (2005) and Cottrell et al. (2009)
106 assessed, but saw no evidence for, radiation-induced beam damage in XANES spectra of
107 anhydrous basalt glasses collected at bending magnet sources. Shorttle et al. (2015)
108 assessed beam damage in glasses with <1 wt.% water, and observed an increase in the
109 intensity of the second pre-edge multiplet with time (radiation dose) using an
110 unattenuated undulator-source that was “barely above instrument stability” (i.e. was
111 within the noise). The study of Moussallam et al. (2016) was not able to detect beam
112 damage by analyzing successive XANES spectra collected at a single spot on glasses
113 with <1 wt.% water at an attenuated undulator-source. However, as new radiation sources
114 come online that can deliver higher flux densities to a sample, and studies of hydrous
115 silicate glass proliferate, we find it prudent to examine the potential for radiation to
116 induce beam damage.

117

118 **SAMPLES AND METHODS**

119 *Samples.* We examine 27 glasses of basaltic composition in this study: 13 anhydrous
120 basalt glasses from the study of (Cottrell et al., 2009), which are available to the
121 community upon request to Smithsonian's National Rock Collection, and 14 hydrous
122 glasses (0.72 to 4.82 wt.% dissolved total H₂O) from the study of (Botcharnikov et al.,
123 2005). For this study, all analyses were performed on double-side polished glass wafers
124 of approximately 50-100µm in thickness. Details of the experimental methods and
125 complete sample characterizations can be found in Cottrell et al. (2009), Cottrell and
126 Kelley (2011) (including wet-chemical analyses), and Botcharnikov et al. (2005).

127

128 *Analytical Methods.*

129 *Mössbauer Spectroscopy.* Cottrell et al. (2009) and Botcharnikov et al. (2005) report
130 Fe³⁺/ΣFe ratios for each glass based on room temperature Mössbauer spectroscopy,
131 without correction for recoilless fraction. We re-report their results in Table 1, along with
132 new fits to the hydrous B glasses discussed below. Method details for previously reported
133 fits can be found in Cottrell et al. (2009) and Botcharnikov et al. (2005). Cottrell et al.
134 (2009) applied the distribution fitting method of Alberto et al. (1996), which adopts a
135 two-dimensional (2D) Gaussian probability distribution function for the isomer shift and
136 quadrupole splitting. This routine minimizes the error (reduced χ^2 values) between the
137 data and the absorption envelope that arises from the contribution of paramagnetic
138 quadrupole doublets. The area of the fitted ferric doublet relative to the entire absorption
139 envelope provides the Fe³⁺/ΣFe ratio. Using the software package RECOIL (Lagarec and

140 Rancourt, 1998), Botcharnikov et al. (2005) applied an extended Voigt-based fitting
141 (xVBF) method that also assumes a two-dimensional Gaussian distribution of hyperfine
142 parameters for ferric and ferrous iron (Lagarec and Rancourt, 1997). The relative area
143 ratios provide the $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios, and the uncertainties estimated by Botcharnikov et al.
144 (2005) that we re-report in Table 1 reflect counting statistics, variation in the results from
145 the application of multiple fitting models, and the possible effect of neglecting the
146 recoilless fraction.

147

148 To isolate the effects of beam damage, deconvolution of the Mössbauer spectra for
149 hydrous and anhydrous glasses must be consistent. Both xVBF and Alberto et al. (1996)
150 methods fit multiple distributions (based on a Gaussian distribution of Lorentzian line
151 shapes). The site distributions derive from maxima in the probability functions that
152 describe the ferrous and ferric hyperfine field distributions (Mysen et al., 2006, and
153 references therein). Zhang et al. (submitted) show that both approaches yield consistent
154 $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios for the AII and LW glasses. To test for systematic bias between the
155 Cottrell et al. (2009) and Botcharnikov et al. (2005) approaches, we obtained the raw
156 spectra of Botcharnikov et al. (2005), collected at Bayerisches Geoinstitute, University of
157 Bayreuth, and refit 14 of them using the same methods used by Cottrell et al. (2009). This
158 ensures that our analysis of oxidative beam damage from synchrotron radiation will not
159 be affected by inter-laboratory bias in the Mössbauer fitting protocols.

160

161 *X-ray Absorption Near Edge Structure (XANES) Spectroscopy*. Cottrell et al. (2009) and
162 Botcharnikov et al. (2005) report on Fe K-edge XANES spectra collected previously for

163 the glasses used in this study. The former collected spectra at the bending magnet station
164 X26A at the National Synchrotron Light Source (NSLS), Brookhaven National
165 Laboratory, New York, USA, and the latter at the Hamburger
166 Synchrotronstrahlungslabor (HASYLAB), Hamburg, GE. Details about those facilities,
167 collection procedures, and fitting routines can be found in Cottrell et al. (2009) and
168 Botcharnikov et al. (2005).

169

170 For this study, we collected new Fe K-edge XANES spectra on the AII, LW, and B series
171 glasses in fluorescence mode in six synchrotron beam sessions: two on bending magnet
172 station X26A at NSLS in 2015, and four at station 13-ID-E (GSECARS) at the Advanced
173 Photon Source (APS), Argonne National Laboratory, IL, USA, in 2014 - 2016. We
174 applied 8 radiation dose conditions that we detail in Table 2. We cast our results in terms
175 of radiation dose so that the results can be easily generalized to other facilities and
176 beamline configurations. For these experiments it is more convenient and descriptive to
177 represent the mean energy imparted to the sample in terms of flux density, rather than in
178 the more standard SI units of Gray. However, the absorbed dose in Gray (absorption of
179 one joule of radiation energy per kilogram of sample) can be calculated by assuming an
180 average basaltic glass composition and density and the depth at which the incident
181 radiation falls to 1/e within the glass (Henke et al. 1993). For basaltic glass, the flux
182 density is equivalent to ~75-80 times the dose in Grays/sec (Table 2). As we describe
183 below, in all sessions, and at all facilities, we set the area-weighted average energy of the
184 pre-edge peaks (the centroid) equal to 7112.3 for anhydrous glass standard LW_0
185 (Cottrell et al., 2009). LW_0 is analysed frequently throughout all analytical sessions.

186 Normalization of all centroids to the energy drift-corrected centroid of LW_0 accounts
187 for within-session instrumental drift in the monochromator energy calibration and for
188 differences in monochromator calibration between facilities and between sessions (Bajt et
189 al., 1994; Cottrell et al., 2009). The empirical drift-corrected precision (1σ) of the area-
190 weighted centroid positions collected over a ten-year period on the AII and LW glasses is
191 0.004 ± 0.009 eV (n for each glass = 30-50) (Zhang et al., submitted).

192

193 We did not correct for over-absorption effects (Manceau et al., 2002). In fluorescence
194 mode XANES, over-absorption can result in a decrease in the amplitude of the absorption
195 peaks in samples where the absorber mass is high. This can result in a relative increase in
196 the pre-edge peak intensities relative to the main absorption edge, but it does not impact
197 the relative intensities of the two pre-edge peak multiplets or their energy positions
198 (Botcharnikov et al., 2005; Cottrell et al., 2009; Zhang et al., 2016). Moreover, while
199 over-absorption can affect pre-edge peak intensities, all glass compositions examined
200 here have sufficiently high total FeO contents (7-12 wt.%) and thicknesses ($> 50\mu\text{m}$) to
201 be in the infinite thickness regime.

202

203 *X26A, NSLS, Brookhaven National Laboratory, USA*. NSLS (decommissioned in
204 October, 2014) operated at 2.8 GeV and 300 mA. We collected fluorescence spectra on
205 the B series glasses at bending magnet beamline X26A using summed energy dispersive
206 spectra from a 4-element Vortex ME-4 silicon drift diode detector and two single element
207 Vortex-EX detectors (Hitachi) with pulse processing provided by an XMap digital
208 spectrometer system (XIA). The first derivative of Fe K-edge spectra (E_0) collected on

209 zero-valent Fe E_0 is 7112.0 eV. All other methods details are as reported by Cottrell et al.
210 (2009). The incident beam flux at X26A was consistently $\sim 3 \times 10^9$ photons/sec and the
211 spot size was 9 x 5 microns such that the photon flux density was $\sim 6.7 \times 10^7$
212 photons/sec/ μm^2 (Table 2); after a 15.5 minute collection the absorbed dose per spectrum
213 was $\sim 6 \times 10^{10}$ photons/ μm^2 .

214

215 *13-ID-E, GSECARS, Advanced Photon Source, Argonne National Laboratory, USA.*

216 Details regarding the beamline configuration for 13-ID-E can be found in Sutton et al.
217 (2017). This undulator-based beam line at APS operates at 7 GeV and 85-101 mA. We
218 collected fluorescence spectra with a Vortex ME4 silicon-drift diode detector array
219 coupled to a high-speed digital spectrometer system (Quantum Xpress3). A cryogenically
220 cooled double-crystal Si (311) monochromator provided the monochromatic radiation.
221 The first derivative of Fe K-edge spectra (E_0) collected on zero-valent Fe E_0 is 7110.75
222 eV. We collected spectra from 7012 to 7102 eV in 2.5 eV steps with 1 sec dwell; from
223 7102 to 7132 eV (pre-edge region) in 0.1 eV steps with 1 sec dwell and from 7132 to
224 7492 eV (XAFS region) in 2.3 angstrom steps with 1 sec dwell. We edge-step normalized
225 the spectra collected at 13-ID-E by the average intensity between 7315 and 7320 eV
226 where the spectra were flat.

227

228 The 13-ID-E beamline is capable of delivering and detecting in excess of 1×10^{12}
229 photons/sec within a focused spot size of 2 x 2 microns at the Fe K-edge, so that focused
230 flux densities of $\sim 5 \times 10^{11}$ photons/sec/ μm^2 are possible – four orders of magnitude higher
231 than what was possible at X26A. We collected the experimental data presented here over

232 a $\sim 10^4$ range in flux density to precisely assess the spectral response to absorbed radiation
233 dose. At 13-ID-E we collected spectra using 6 protocols where incident flux density was
234 adjusted, through upstream aperturing of the X-ray beam; through filtering using varying
235 thicknesses of aluminium foil; and through defocusing of the incident beam on the
236 sample. These conditions are reported in Table 2.

237

238 We also conducted one set of analyses at 13-ID-E under cryogenic conditions, using a
239 customized cryogenic stage described in Tappero et al. (2016, 2017), to investigate the
240 degree to which cryogenic cooling may mitigate the observed iron oxidation. We
241 collected a time series at an energy of 7113.3 eV at both ambient and cryogenic
242 conditions reported in Table 2.

243

244 *Fourier Transform Infrared (FTIR) Spectroscopy*

245 Glass chip B10 (Table 1) was analyzed by FTIR for dissolved total H₂O using a Thermo
246 Nicolet iS50 FTIR bench spectrometer coupled with a Continuum IR microscope with
247 automated mapping capability at the Graduate School of Oceanography, University of
248 Rhode Island. We mapped an area on a chip of B10 that had been purposely exposed to
249 high photon flux at beamline 13-ID-E. The map was collected using transmission FTIR
250 with a 20 x 20 μm square aperture and no overlap between pixels. Each spectrum was
251 collected over 1200-6000 cm^{-1} wavenumbers using 256 scans at 1 cm^{-1} resolution.
252 Individual transmission spectra were reduced for H₂O^{Total} after fitting the spectral
253 background with a spline function and calculating the net intensity of the absorbance
254 band at 3500 cm^{-1} . The H₂O^{Total} concentration was then calculated using the Beer-

255 Lambert law, a calculated density of 2.90 g/cm^3 , thickness of $49 \pm 1 \text{ }\mu\text{m}$, and the molar
256 absorption coefficient $\epsilon^{3530} = 63$, which is assumed constant within the compositional
257 range of basalts (Dixon et al., 1995).

258

259 **RESULTS**

260 *Mössbauer Spectra*

261 Mössbauer-derived $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios are not very sensitive to fitting philosophy or method,
262 but in this study we re-fit the B glass spectra acquired by Botcharnikov et al. (2005) to
263 eliminate any potential bias between the disparate fitting routines of Cottrell et al. (2009)
264 and Botcharnikov et al. (2005). In addition, our new fits allow us to report hyperfine
265 parameters for these glasses that may aid in understanding the structural environment of
266 Fe in hydrous and anhydrous glasses. We provide the Mössbauer spectra from
267 Botcharnikov et al. (2005) and our fits from this study (both raw data and plots) in the
268 Supplementary Material.

269

270 Two quadrupole doublets, one ferric with a low IS of 0.4 to 0.6 mm/s (relative to Fe
271 metal) and low QS of 0.8 to 1.1 mm/s, and one ferrous with a high isomer shift of 0.8 to
272 1.0 mm/s and quadrupole splitting of 1.9 to 2.2 mm/s, describe the Mössbauer spectra of
273 the B glasses (Table 1, and Figures S1, S2). The Botcharnikov et al. (2005) study did not
274 report hyperfine parameters, but the hyperfine parameters derived from the glass spectra
275 in this study are similar to those reported previously in the literature for silicate glasses
276 (Alberto et al., 1996; Jayasuriya et al., 2004; Mysen et al., 1985a; Zhang et al., 2016).
277 Ferrous iron IS and QS in the hydrous B glasses are within error of those for the

278 anhydrous glasses for all but the most highly oxidized glasses, for which the quadrupole
279 splitting in the hydrous glasses is higher. The QS for ferric iron is smaller in hydrous
280 glasses than in anhydrous glasses by approximately 0.25 to 0.45 mm/s. The IS of ferric
281 iron is slightly higher in the hydrous glasses than in the anhydrous glasses, and decreases
282 with increasing $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio, as expected (Mysen and Richet, 2005; Mysen and Virgo,
283 1989; Mysen et al., 2006). While many factors influence the coordination of ferric iron in
284 silicate glasses (Mysen and Richet, 2005), the greater ferric IS observed in the hydrous
285 glasses suggests that the O— Fe^{3+} bond length is longer, and weaker (e.g. Johnson et al.,
286 1999), and that ferric iron is in higher average coordination (greater domination of $^{\text{VI}}\text{Fe}^{3+}$)
287 in the hydrous B glasses than in the anhydrous AII and LW glasses. Higher coordination
288 of ferric iron in the hydrous glasses is consistent with results from (Wilke et al., 2002;
289 Wilke et al., 2006), and is corroborated by XANES spectra, as we discuss below.

290

291 The $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios in the B glasses, as determined by the new fits reported here, range
292 from 0.12 to 0.56, compared to 0.10 to 0.60 for those same glasses reported in
293 Botcharnikov et al. (2005). The $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios of all B glasses with nominal $\text{Fe}^{3+}/\Sigma\text{Fe}$
294 <0.50 agreed within 1σ (2.7% absolute) with the values reported by Botcharnikov et al.
295 (2005). The three glasses with $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios >0.50 agreed within 2σ of the
296 uncertainties on the fits (4.4% absolute). These fits are within the commonly accepted
297 uncertainty of the technique, which ranges from ± 1 to 6% in absolute $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio
298 (Dyar et al., 1987; Mysen et al., 1985a; Mysen et al., 1985b). $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios from the
299 Botcharnikov et al. (2005) study correlate very well with ratios obtained in this study (R^2
300 >0.99 , $m = 0.9$, $b = 1.87$). Consistent with the results of Cottrell et al. (2009), Zhang et

301 al. (2015), and Zhang et al. (submitted), inclusion of a second doublet in the present study
302 did not reduce the χ^2 of the fits.

303

304 *XANES*

305 *Bending Magnet Beamlines, X26A NSLS.*

306 Full XANES spectra of hydrous B and anhydrous AII and LW glasses collected at
307 bending magnet beamline X26A are shown in Figure 1. We provide spectra for AII and
308 LW anhydrous glasses for a 2014 NSLS session for reference (there is no substantive
309 change since their initial publication in Cottrell et al., 2009). As demonstrated by
310 Botcharnikov et al. (2005) for previous XANES analyses of the B glasses, the energy of
311 the main absorption edge, the ratio of pre-edge peak intensities, and the area-weighted
312 average energy of the pre-edge peaks (the centroid) in the XANES spectra shift smoothly
313 as a function of $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio (e.g., Berry et al., 2003; Botcharnikov et al., 2005; Cottrell
314 et al., 2009; Wilke et al., 2005). The shift in centroid energy provides a precise proxy for
315 $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio in basaltic glasses, especially for $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios below 0.50 (Henderson et
316 al., 2014). We report the centroids and pre-edge peak intensities for all glasses in Table 3.
317 We provide all fitting parameters in Table S2. Consistent with the results of Cottrell et al.
318 (2009), centroids obtained from spatially independent spectra are precise to $\leq \pm 0.01$ eV.

319

320 For anhydrous LW and AII glasses and hydrous B glasses, we plot the $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios
321 derived from Mössbauer fits determined by Botcharnikov et al. (2005) and this study
322 against the $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios derived from XANES at low flux using the anhydrous
323 calibration curve of Cottrell et al. (2009). Regardless of which set of Mössbauer values

324 are used, the $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios of the hydrous glasses fall within 1σ of a 1:1 relationship
325 (Figure 2a).

326

327 *Undulator-based Beamlines, 13-ID-E, APS.*

328 Full XANES spectra of AII, LW, and B glasses collected at insertion device beamline 13-
329 ID-E (GSECARS) are shown in Figure 1 and we provide fitting parameters in Table S2.

330

331 *Condition XANES₁, Anhydrous Glasses.* We report centroids (area weighted average
332 energy of the two pre-edge multiplets) under “nominal” operating conditions at 13-ID-E,
333 our condition “XANES₁,” for each glass in Table 3.

334

335 Under normal high flux operating conditions at station 13-1D-E, we recover the $\text{Fe}^{3+}/\Sigma\text{Fe}$
336 ratios of anhydrous glasses when we use the calibration of Cottrell et al. (2009), which
337 was developed at low flux. Anhydrous glass $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios derived from XANES at
338 high flux fall within 1σ of a 1:1 relationship with the Mossbauer-derived values (Figure
339 2b). Supplemental Figure S3a and Table 3 show that high flux centroids are on average
340 within $\pm 0.008\text{eV}$ (within 1σ) of low flux centroids. One might conclude from this that the
341 centroid, as a spectral feature, is a robust proxy for assessment of the $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio for
342 anhydrous glasses. We emphasize again, however, the importance of normalizing to a
343 standard. While the centroids at the two facilities are reproducible, the individual spectral
344 features are not. Differences in modelled spectral features (Figures S3b,c,d), and the raw
345 spectra (Figure S4), reflect the smaller vertical emittance of the 13-ID-E undulator
346 electron beam compared to the X26A bending magnet source, and the optical

347 configuration at each beamline. That we reproduce the centroids of anhydrous glasses
348 across facilities to approximately ± 0.01 eV is a non-trivial result. Fiege et al. (2017) state
349 that the “overall uncertainty” in the centroid is an order of magnitude higher (± 0.1 eV).
350 We attribute the high precision we achieve primarily to ensuring that we account for
351 instrumental drift within and between analytical sessions by shifting the energy of all
352 centroids such that the centroid of LW_0 $\equiv 7112.3$ eV. The key result is that the
353 anhydrous glass centroids remain stable within uncertainty at all beam conditions and
354 across all facilities.

355

356 *Condition XANES₁, Hydrous Glasses.* At a given $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio, the centroids of hydrous
357 B glasses are variably shifted to higher values at 13-ID-E under nominal operating
358 conditions (XANES₁), in stark contrast to the highly reproducible anhydrous glasses
359 (Table 3). For some hydrous glasses (e.g., B20), we derive the same centroid within
360 0.003eV (within 1σ) between facilities, but for others (e.g., B11) we derive a centroid up
361 to 0.25eV higher (25σ higher) under high flux conditions than low flux conditions
362 (Figure S4). This manifests as XANES-derived $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios that are up to 1.7 times
363 higher than those derived from Mössbauer spectroscopy (Figure 2b).

364

365 *Intensity of 2nd Multiplet with Time.* To better evaluate how the Fe XANES spectra of
366 these glasses evolve with accumulated radiation dose, we conducted experiments to
367 monitor changes in the intensity of the second pre-edge multiplet at 7113.3 eV (attributed
368 to Fe^{3+}) as a function of time and flux density (i.e. dose). Glass B11 was used for this
369 experiment; it is the most hydrous in the suite under study, with 4.82 wt.% total dissolved

370 water. This glass is relatively reduced relative to other hydrous glasses tested, with a
371 $\text{Fe}^{3+}/\Sigma\text{Fe} = 0.19$.

372

373 Figure 3 plots the peak amplitude of the second pre-edge multiplet peak of hydrous glass
374 B11 with time, measured at high flux density and low flux density under both ambient
375 and cryogenically cooled conditions. We first consider data collected at ambient
376 temperature. Peak amplitude increases with time under high and low flux conditions, with
377 a rapid increase observed within the first 100 seconds of high flux irradiation. The higher
378 photon flux density generated by the focused beam causes a greater increase in the
379 amplitude of the second pre-edge multiplet at all time points. However, when the change
380 in amplitude is measured relative to the dose delivered, the amplitude of the second
381 multiplet (extent of damage) increases at approximately the same rate in high and low
382 flux density conditions. This evidences the dose-dependence of the spectral change.
383 While low flux beam damage does not affect the centroid of these hydrous glasses within
384 the measurement error of our calibration (Figures 2, S5), these data show that damage
385 occurs in hydrous glasses even at low flux density.

386

387 We then explored the degree to which cryogenic cooling impacts spectral evolution with
388 time. Studies over a broad range of incident beam energies show that cryo-cooling of
389 biological samples significantly reduces damage due to mass (water) loss during
390 exposure to ionizing radiation (e.g. Dubochet et al., 1988; Steinbrecht and Zierold, 2012).
391 Cryo-cooling may also potentially reduce the rate of beam-induced changes in metal
392 valence state in some materials (Corbett et al., 2007; Grabolle et al., 2006), although the

393 degree to which this mitigates the effects is debated (Beetz and Jacobsen, 2003; Meents
394 et al., 2010). We repeated the time series experiments described previously on sample
395 B11, with the glass held at -195°C. The measurements show the surprising result that at
396 cryogenic temperatures, under both high and low flux density conditions, the relative rate
397 of Fe oxidation increases. Beam damage is worse at cryogenic temperatures.

398

399 It is also instructive to see that spectral changes plateau under all conditions studied
400 within 3-16 minutes. It takes anywhere from approximately 4 to 30 minutes to collect an
401 entire quality XANES spectrum on glass with FeO concentrations between 5 and 10 wt.%
402 (e.g. Cottrell and Kelley, 2011; Shorttle et al., 2015). Therefore, successive XANES
403 spectra, collected at a single sample position, may not show any measurable change if
404 damage begins to plateau before acquisition of the first spectrum is finished.

405

406 *Conditions XANES_{3, 4}. Successive Analysis at one position.* Using glass B11 again, we
407 show in Figure 4a and Table S3 how the centroids of successive analyses on the same
408 spot in glass B11 increase if we employ condition XANES₄, a very low flux achieved by
409 attenuating and defocusing the beam to 50x50µm. The centroid increases approximately
410 logarithmically in energy for each successive spectrum. Contrast this with successive
411 spots at the same incident flux, but with the beam focused to 2x2µm (XANES₃), in which
412 case the total number of photons delivered to a given area (photons/µm²) is higher after
413 one spectral acquisition than after successive analyses under condition XANES₄. The
414 centroid is commensurately higher, and plateaus after one acquisition. Consistent with
415 Cottrell et al., 2009, we observe no change in the centroid or pre-edge peak intensities

416 with successive analyses on anhydrous glass AII_0. Figure 4b reveals how the integrated
417 intensity (area) of the Gaussian extracted for the first multiplet ($I_{\text{Fe}^{2+}}$) decreases while the
418 integrated intensity of the second multiplet ($I_{\text{Fe}^{3+}}$) increases. Evolution of the spectra in
419 this manner is consistent with oxidation of Fe^{2+} to Fe^{3+} .

420

421 Observations of the raw spectra (Figure 5) confirm the spectral evolution seen in the
422 modelled parameters. Raw spectra collected sequentially at a given position on sample
423 B11, under condition XANES₄, show decreases in the intensity of the first pre-edge
424 multiplet and increases in intensity of the second pre-edge multiplet. In addition, we see
425 the main absorption edge move to higher energy as radiation dose accumulates. Spectral
426 changes are consistent with an increase in the average oxidation state of Fe with
427 increasing radiation dose.

428

429 *Beam Damage Persists Indefinitely*

430 To assess the longevity of beam damage, we mapped glass B11 approximately one year
431 after a session in which we investigated beam damage. In figure 6, image brightness
432 scales with the ratio between the Fe K α fluorescence peak amplitude at 7113.2 eV, the
433 energy maximum for the Fe^{3+} pre-edge multiplet peak, and at 7111.6 eV, the energy
434 maximum for the Fe^{2+} pre-edge multiplet peak. In the false-color inset, the bright white
435 areas correspond to previously analysed spots. From this map it is evident that areas
436 exposed to high radiation dose retain higher $\text{Fe}^{3+}/\text{Fe}^{2+}$ ratios than adjacent areas that were
437 not exposed to the X-ray beam at least one year after exposure. Beam damage persists
438 indefinitely.

439

440 *Coordination of Fe in Hydrous and Anhydrous Glasses*

441 Figure 7 compares the summed integrated pre-edge intensities ($I_{\text{Fe}^{2+}} + I_{\text{Fe}^{3+}}$) of hydrous
442 and anhydrous glasses as a function of their centroids – what Wilke et al. (2001) coined a
443 “variogram.” A negative non-linear correlation between coordination and oxidation state
444 is inferred for glasses as well as minerals with mixed geometries (e.g. Jackson et al.,
445 2005; Wilke et al., 2001). Because $I_{\text{Fe}^{2+}}$ and $I_{\text{Fe}^{3+}}$ are extremely sensitive to detector dead
446 time, edge-step normalization procedure, self-absorption etc., they are generally not
447 comparable between facilities and are even difficult to compare between analytical
448 sessions at the same facility (Cottrell et al., 2009). We only compare spectra acquired at
449 APS in Figure 7; these have been collected and processed under conditions as uniform as
450 can be practically achieved.

451

452 All glasses display typical pre-edge intensities indicative of 5-fold or mixed 4-fold and 6-
453 fold coordination (Wilke et al., 2001; Wilke et al., 2005). We see that hydrous glasses fall
454 along the same trend of increasing integrated pre-edge intensity with centroid energy
455 under both the high flux density condition XANES₃ and the low flux density condition
456 XANES₄. Covariation in the oxidation state and coordination of Fe in these glasses is
457 consistent with a shift in Fe coordination upon exposure to radiation only in proportion to
458 the shift in Fe oxidation state. This analysis therefore also points to actual oxidation of
459 Fe^{2+} to Fe^{3+} in the hydrous glasses.

460

461 The variogram also corroborates the site geometry analysis of the Mössbauer spectra,

462 which revealed higher Fe^{3+} isomer shift (and smaller quadrupole splitting) in hydrous B
463 glasses relative to anhydrous glasses (Figure S1), consistent with longer, weaker, O—
464 Fe^{3+} bonds, and Fe^{3+} in higher coordination in the hydrous glasses than the anhydrous AII
465 and LW glasses. In Figure 7, anhydrous glasses display higher integrated intensities
466 (lower average coordination) than hydrous glasses at any given centroid energy,
467 consistent with longer and weaker O— Fe^{3+} bonds in the hydrous glasses. Correcting for
468 over-absorption does not change this relationship.

469

470 **DISCUSSION**

471 *Beam Damage Mechanism*

472 The *Is* absorption edge and pre-edge peaks in *3d* transition metals, like Fe, are due to
473 electron transitions from *Is* to *3d* states (e.g. Henderson et al., 2014). Quadrupole (*Is3d*)
474 and dipole (*Is* to the *4p* character of the *3d* band) transitions are possible in silicate
475 glasses, owing to iron's mixed coordination. In XANES spectra, Fe^{2+} exhibits 2 or 3 pre-
476 edge peaks, and Fe^{3+} exhibits 1 or 2 pre-edge peaks, depending on coordination (Westre
477 et al., 1997). $I_{\text{Fe}^{2+}}$ and $I_{\text{Fe}^{3+}}$ increase as coordination number decreases (i.e. with higher
478 extents of *3d* + *4p* mixing) because of the added dipole contribution (Penner-Hahn, 2005;
479 Roe et al., 1984; Westre et al., 1997). The relative integrated intensities and energies of
480 the pre-edge peaks therefore depend on the metal site geometry: the relative proportions
481 of Fe^{2+} and Fe^{3+} and their coordination as tetrahedral (4-fold) or octahedral (6-fold)
482 (Henderson et al., 2014; Wilke et al., 2001). Thus, it is necessary to evaluate the degree
483 to which the observed beam-induced changes in the XANES spectra reflect changes in
484 effective Fe valence as compared to coordination, particularly since in glasses these

485 effects are coupled.

486

487 We considered the possibility that exposure to radiation in hydrous glasses, but not
488 anhydrous glasses, might increase $3d-4p$ hybridization of Fe^{3+} centers. For example,
489 $\text{Fe}^{3+}\text{—OH}^-$ bonds could respond differently to radiation than $\text{Fe}^{3+}\text{—O}$ bonds. Augmented
490 hybridization of Fe^{3+} centers in hydrous glasses would increase the dipole contribution to
491 the second multiplet, increase its integrated intensity, and increase the centroid, without
492 any change in the actual average oxidation state of Fe. Yet, several lines of evidence
493 argue against this possibility. First, upon exposure to a high photon flux density, the
494 XANES spectra of the hydrous glasses show that the Fe^{3+} multiplet integrated intensity
495 increases and the Fe^{2+} multiplet intensity concomitantly decreases (Figure 4). Moreover,
496 the white line shifts to progressively higher energy as well (Figure 5). These three
497 spectral changes together are much more consistent with oxidation of Fe^{2+} to Fe^{3+} than an
498 increase in $3d-4p$ hybridization, as the latter would result in neither a decrease in the
499 intensity of the first multiplet, nor a shift in the main edge.

500

501 Because no measurable change in the XANES spectra is observed as a function of dose in
502 spectra collected from anhydrous glasses (e.g. Figures 2, 4), we hypothesize that the
503 changes we observe in the spectra of hydrous B glasses with exposure to ionizing
504 radiation relates to their dissolved water contents. Here we use $\text{Fe}^{3+}/\Sigma\text{Fe}_{\text{Möss}}$ –
505 $\text{Fe}^{3+}/\Sigma\text{Fe}_{\text{XANES}}$ as a proxy for beam-induced Fe oxidation, where Mössbauer spectra
506 establish the “true” $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio, and the XANES centroid the “apparent” $\text{Fe}^{3+}/\Sigma\text{Fe}$
507 ratio (Table 3). Indeed, when we compare glass B8 to B16, glasses of similar $\text{Fe}^{3+}/\Sigma\text{Fe}$

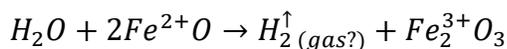
508 ratio, we see that the centroid of glass B16, with > 3.6 times more dissolved water, is
509 overestimated by ~70%, whereas the centroid of glass B8 is overestimated by ~27%
510 when exposed to the same radiation dose. Figure 8a shows, however, that total dissolved
511 H₂O in the glasses does not correlate with beam-induced Fe oxidation. For example, the
512 anhydrous XANES calibration of Cottrell et al. (2009) predicts the Fe³⁺/ΣFe ratio of
513 hydrous glass B17, with 4.67 wt.% total water, within error of its “true” value, while the
514 centroid of B11, with nearly indistinguishable dissolved water content, is overestimated
515 by >70% (Figure 4, 11). The latter overestimation of Fe³⁺/ΣFe ratio leads to
516 overestimation of *f*_{O₂} by two orders of magnitude.

517

518 Following this, it is also clear that when the initial oxidation state of Fe is high, the
519 spectral response to radiation (beam damage) is lessened (Figure 8b). Upon exposure to
520 the same high photon dose (condition XANES₁), the Fe³⁺/ΣFe ratios of hydrous glasses
521 with lower Fe³⁺/ΣFe ratios are greatly overestimated in comparison to glasses with higher
522 Fe³⁺/ΣFe ratios, which appear to be well predicted by the Mössbauer calibration (R²
523 linear correlation = 0.53; R² exponential correlation = 0.72, Figure 8b).

524

525 All the evidence suggests to us that the mechanism of Fe oxidation involves breaking of
526 O—H bonds by incident ionizing radiation, loss of hydrogen, and concomitant oxidation
527 of Fe²⁺ to Fe³⁺. One possible reaction, analogous to that proposed for radiation-induced
528 oxidation of S dissolved in hydrous silicate glasses by Klimm et al. (2012) is:



529 The loss of hydrogen gas is consistent with the longevity of the beam damage, as

530 discussed above and shown in Figure 6. This mechanism is also consistent with the
531 observed dependence on both the initial $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio and the dissolved water content of
532 the hydrous glass.

533

534 The hypothesized mechanism suggests that a fruitful metric for predicting the
535 susceptibility of hydrous glass to beam damage is the molar ratio of dissolved total water
536 to the ferric/ferrous ratio as determined by Mössbauer spectroscopy
537 ($\Phi = \text{XHO}_{1.5} * \text{XFeO} / \text{XFeO}_{1.5}$ ratio) (Table 1). We find that beam damage susceptibility,
538 Φ , correlates strongly with the extent to which high photon flux XANES₁ overestimates
539 the $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio ($R^2 = 0.92$, $n=14$, $b \equiv 0$ because anhydrous glasses show no damage)
540 (Figure 8c).

541

542 Inspection of the raw spectra of the B glasses under conditions XANES₃ and XANES₄
543 confirms that spectral changes increase as a function of dose and increasing Φ , with
544 anhydrous glasses showing no spectral change (Figure 9). Visual inspection also verifies
545 that the spectral features all evolve in a manner consistent with oxidation of ferrous iron
546 to ferric iron: $I_{\text{Fe}^{2+}}$ falls; $I_{\text{Fe}^{3+}}$ grows; and the K edge shifts to higher energy. Spectral
547 changes increase in proportion to Φ . We provide the raw spectra for Figure 9 in Table S4.

548

549 We tested this hypothesis by subjecting glass B10 to a 25 μm beam of high photon flux
550 of $\sim 1 \times 10^{12}$ photons/sec for ~ 12 minutes (Figure 10). We subsequently mapped the area
551 surrounding this spot on the sample via Fourier Transform Infrared spectroscopy (FTIR).
552 The FTIR data show that the area of the glass adjacent to the exposed spot (e.g., pixel 1,

553 Figure 10) also has significantly lower dissolved H₂O than areas of the glass distant from
554 exposure (e.g., pixel 2, Figure 10). We reduced individual FTIR spectra from pixels 1 and
555 2 by fitting the background of the total H₂O absorption band at 3500 cm⁻¹ with a spline
556 function that was then subtracted from the spectra before measuring the band height. This
557 analysis shows that pixel 2, away from the potential damage zone, has a total dissolved
558 H₂O content of 2.43 wt.%, which is identical to the reported H₂O content for this glass by
559 Botcharnikov et al. (2005). The dissolved H₂O content of pixel 1, however, is 900 ppm
560 lower (2.34 wt.%). Pixel 1 is not coincident with the location of the XANES beam at the
561 sample surface. Rather, the incident X-ray beam intersects the sample surface at a 45°
562 angle relative to the FTIR incident beam, which is normal to the sample surface. The high
563 flux X-ray beam travelled through the sample, from right to left in Figure 10, at 45°
564 angle, and thus interacted with the glass chip to the left of the beam cross-hair location.
565 We conclude from this assessment that exposure of this sample to a high-flux XANES
566 beam caused both oxidation of Fe and loss of H₂O from the glass, which supports the
567 proposed mechanism of Fe oxidation via hydrogen loss. We note that only a few percent
568 of the Fe fluorescence we measure comes from the sample surface. The geometrical
569 considerations we present above and the magnitude of the spectral changes and water loss
570 we observe require chemical changes to the bulk affected volume. Damage cannot be
571 limited to the sample surface.

572

573 Cryo-cooling has been shown to mitigate beam damage (e.g. Steinbrecht and Zierold,
574 2012; Corbett et al., 2007). Our finding that cryo-cooling exacerbated Fe²⁺ oxidation in
575 hydrous glass is an important and counterintuitive observation. We have shown that

576 hydrogen concentrations in the glass are lower in areas that have experienced oxidizing
577 beam damage, consistent with radiation-induced breaking of O—H bonds, and
578 preferential diffusion of hydrogen, relative to oxygen, away from the area. Because it is
579 the relative diffusivity of hydrogen and oxygen that matters, and diffusivity depends
580 strongly on temperature, we speculate that oxygen is able to diffuse farther from the
581 analysed spot in the room temperature case than the cryo-cooled case (whereas hydrogen
582 is always able to diffuse away). This may be why cryo-cooling is counterproductive
583 when attempting to mitigate radiation-induced oxidation of Fe in hydrous silicate glass.
584 Indeed, Gonçalves Ferreira et al. (2013) found that heating soda-lime glasses to 500°C
585 actually mitigated radiation-induced *reduction* of iron, emphasizing that the effect of
586 temperature on beam damage will depend on the nature and mechanism of beam damage
587 and is difficult to predict a priori.

588

589 **IMPLICATIONS**

590 *Prior XANES studies of hydrous glasses*

591 We have shown that synchrotron radiation causes progressive changes to the XANES
592 spectra of hydrous glasses, but not anhydrous glasses, as a function of radiation dose.
593 Spectral changes are consistent with preferential migration of hydrogen relative to
594 oxygen, and oxidation of iron, in the analysis spot. Our study reveals that oxidative beam
595 damage occurs in hydrous glasses even at the lowest flux densities achieved, but that
596 damage is not detectable outside of the error of anhydrous calibrations with exposure to
597 radiation doses typically found at bending magnet sources. Conversely, it would be
598 prudent to re-evaluate results from studies carried out on hydrous glasses at undulator

599 beam lines using a focused spot.

600

601 Mafic glasses may be synthesized experimentally with a wide range of Φ
602 ($X\text{HO}_{1.5} * X\text{FeO} / X\text{FeO}_{1.5}$) that may readily exceed those investigated in this study. Most
603 naturally occurring basaltic glasses, however, will have $0 < \Phi < 1$. For example, basaltic
604 glass inclusions in volcanic rocks from subduction zones have Φ ranging from ~ 0.34 to
605 ~ 1 ; natural basaltic glasses from sea floor volcanoes erupting in back arc basins have
606 ratios ranging from ~ 0.2 to ~ 0.6 ; and mid-ocean ridge basalts have ratios between ~ 0.05
607 to ~ 0.28 (ranges estimated from data in Brounce et al. 2014b; Cottrell and Kelley, 2011;
608 Kelley and Cottrell, 2012; Zhang et al., submitted). Based on the work presented here, all
609 submarine basalts and undegassed melt inclusions would suffer significant beam damage
610 under nominal acquisition conditions (i.e. a focused beam) at unattenuated undulator
611 beam lines. We have also shown, however, that a radiation dose of $\leq 6 \times 10^7$
612 photons/sec/ μm^2 mitigates (but does not prevent) beam damage on hydrous glasses with
613 Φ up to ~ 1 .

614

615 We hypothesize that previous studies conducted at undulator beam lines have not
616 detected beam damage in hydrous glasses because, under nominal operating conditions,
617 radiation damage is maximum or “complete” after acquisition of a single spectrum (~ 600
618 seconds for spectra acquired under condition XANES_1). Collection of multiple spectra at
619 a single position on the glass under condition XANES_1 will therefore return the same
620 centroid because damage is complete (“saturated”) before the pre-edge is scanned the
621 first time (e.g. Moussallam et al., 2016; Moussallam et al., 2014). For materials with

622 relatively low Φ , it may be sufficient at an undulator source to attenuate the beam to
623 prevent beam damage from exceeding detection limits (e.g. Shorttle et al., 2015). We
624 have shown here, however, that even at the lowest fluxes achievable at 13-ID-E,
625 attenuation is insufficient to prevent beam damage; the beam must also be defocused.

626

627 *Recommended analytical protocols for acquisition of Fe-XANES spectra on hydrous*
628 *glasses*

629 Synchrotron radiation causes progressive changes to the XANES spectra of hydrous
630 glasses as a function of radiation dose and Φ . Dose is a function of photon flux, beam
631 diameter, and exposure time, and can/should be calculated and reported for any
632 synchrotron experiment. Attenuating the beam, defocusing the spot, or both, may
633 mitigate beam damage. Using a cryostat is unlikely to mitigate beam damage; it may
634 even exacerbate damage. This study suggests that a radiation dose of $\leq 6 \times 10^7$
635 photons/sec/ μm^2 will allow the centroids of most terrestrial basalts to lie within analytical
636 error of anhydrous calibration curves. We emphasize that beam damage must be assessed
637 at each beamline and for the specific materials under analysis in a given study.

638

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657

658

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837
838
839
840 **FIGURE CAPTIONS**

841

842 **Figure 1.** XANES spectra for hydrous B glasses and anhydrous LW and AII glasses
843 under three radiation dose conditions. Left hand panels show full spectra; right hand
844 panels magnify the pre-edge region. Spectra are color-coded by relative Fe³⁺/ΣFe ratio as
845 determined by Mössbauer spectroscopy (black are the most reduced glasses, grading
846 through red and orange to yellow, which are the most oxidized glasses). Note the
847 smooth, monotonic, decrease in the intensity in the first pre-edge multiplet and increase
848 in intensity of the second pre-edge multiplet as a function of Fe³⁺/ΣFe ratio for the
849 hydrous samples under the bending magnet and low radiation dose conditions, and the
850 anhydrous glasses under all dose conditions. Spectra of hydrous B glasses under the high
851 radiation dose condition do not evolve systematically as a function of Fe³⁺/ΣFe ratio.

852

853 **Figure 2. (A)** $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios determined by Mössbauer spectroscopy versus $\text{Fe}^{3+}/\Sigma\text{Fe}$
854 ratios determined by XANES at a flux of 10^{6-7} photons/sec/ μm^2 using the anhydrous
855 basalt calibration curve of Cottrell et al. (2009). Filled circles: anhydrous basalts, Cottrell
856 et al. (2009). Open squares: hydrous basalts, Botcharnikov et al (2005). Filled squares
857 and filled triangles: hydrous basalts, this study. The anhydrous basalts, on which this
858 calibration is based, fall slightly off the 1:1 line because we have used the average
859 centroid of these glasses collected over a 10 year period (n=30-50 per glass) as reported
860 in Zhang et al. (submitted) to highlight the reproducibility of the measurement. **(B)** The
861 same analysis as in (A), but at a higher flux of 10^{10} photons/sec/ μm^2 . Vertical 1σ error
862 bars are smaller than the symbol sizes. Horizontal 1σ error bars are the errors on the
863 Mössbauer determinations. Dashed lines show the predictive capability of the model
864 using a leave-one-out cross validation method (Arlot and Celisse, 2010). The root mean
865 square uncertainty for the predicted $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratios of unknowns determined from the
866 basaltic calibration is $<\pm 0.01$ (1σ).

867

868 **Figure 3.** Time-dependence of the intensity of the second pre-edge multiplet ($\Delta\text{FeK}\alpha/\text{IO}$)
869 during exposure to synchrotron radiation under four beam conditions: high flux density
870 (filled squares) and low flux density (open circles); and at ambient temperature (red) and
871 with the sample cryogenically cooled to -195°C (blue). The intensity of the second pre-
872 edge multiplet is higher (i.e. more radiation-induced oxidative beam damage) when the
873 sample is cryogenically cooled under both beam conditions.

874

875 **Figure 4.** Evolution of spectral features as a function of spectrum number (proxy for

876 dose) for spectra acquired in succession in the same position on the sample for hydrous
877 glass B11 and anhydrous glass AII. **(A)** Evolution of centroids under two beam
878 conditions, high radiation dose condition XANES₃ ('+' and 'x' symbols) and low
879 radiation dose condition XANES₄ (circles and triangles). Centroids of successive spots on
880 anhydrous glass AII show no significant changes under either beam condition.

881 Application of a high radiation dose, typical of focused beams at undulator beamlines,
882 does not result in significant changes to the centroids of successive spectra on the same
883 spot of hydrous glass B11; however, centroids acquired using a high radiation dose are
884 higher than centroids collected on the same sample at low radiation dose. Significant
885 increases in the centroids of successive spots on hydrous glass B11 only become apparent
886 using a low radiation dose. Vertical 1σ error bars are smaller than the symbol sizes.

887 **(B)** Evolution of $I_{\text{Fe}^{2+}}$ and $I_{\text{Fe}^{3+}}$ with successive spectra taken at the same position under
888 low radiation dose condition XANES₄. Vertical 1σ error bars are smaller than the symbol
889 sizes.

890

891 **Figure 5. (A)** Evolution the background-subtracted pre-edge, acquired via successive
892 analyses in the same sample position, as a function of radiation dose, using condition
893 XANES₄. As the total dose increases, the intensity of the first pre-edge multiplet
894 decreases while the second increases. **(B)** The main edge moves to progressively higher
895 energy as dose increases.

896

897 **Figure 6.** The grey-scale image is a peak amplitude map collected at ~ 7113.2 eV (energy
898 of 2nd pre-edge multiplet) of hydrous glass B11. Image color is scaled to maximum

899 intensity, where white is high and black is low. The inset map is a peak amplitude map
900 collected at 7113.2/7111.6eV (peak amplitude of the 2nd pre-edge multiplet divided by
901 the peak amplitude of the 1st pre-edge multiplet). Image color is scaled to maximum peak
902 amplitude, where white is high and red is low. White areas show spots analysed >1 year
903 earlier.

904

905 **Figure 7.** A “variogram” showing the spectral characteristics of centroid versus
906 integrated pre-edge intensity (area) of Fe in end-member structural environments (large
907 empty circles) from Wilke et al., 2001 and glasses from this study. Anhydrous (black
908 circles) and hydrous glass (red and blue squares) spectral characteristics from this study
909 are consistent with iron in five-fold coordination or a mixture of four-fold and six-fold
910 coordination. Under the same radiation dose conditions, anhydrous glasses display higher
911 pre-edge intensities relative to hydrous glasses at a given centroid energy. Hydrous
912 glasses under high radiation dose conditions display higher pre-edge intensities than
913 under low radiation dose conditions; however, increases in pre-edge intensities are
914 proportional to increases in the centroid ($\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio). Vertical 1σ error bars are
915 smaller than the symbol sizes.

916

917 **Figure 8.** We use $\Phi = \text{Fe}^{3+}/\Sigma\text{Fe}_{\text{Möss}} - \text{Fe}^{3+}/\Sigma\text{Fe}_{\text{XANES}}$ as a proxy for beam-induced Fe
918 oxidation (where Mössbauer spectra establish the “true” $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio, and the XANES
919 centroid the “apparent” $\text{Fe}^{3+}/\Sigma\text{Fe}$ ratio). The molar concentration of **(A)** dissolved water
920 (single cation basis), $\text{XHO}_{1.5}$, **(B)** ferric iron (single cation basis), $\text{XFeO}_{1.5}$, and **(C)** the
921 ratio of dissolved water multiplied by the ferrous/ferric ratio, $\text{XHO}_{1.5} * \text{XFeO}/\text{XFeO}_{1.5}$,

922 versus $\%Fe^{3+}/\Sigma Fe_{Möss} - Fe^{3+}/\Sigma Fe_{XANES}$. The correlation in (C) is forced through the
923 origin, because we observe no oxidative beam damage in anhydrous glasses. 1σ error
924 bars are smaller than the symbol sizes.

925

926 **Figure 9.** Pre-edge spectra of one anhydrous and four hydrous glasses, with varying Φ
927 and under two radiation dose conditions. We display merged, successive spectra,
928 acquired under two different beam conditions: high radiation dose condition XANES₃
929 (red) and low radiation dose condition XANES₄ (blue). The black line shows the
930 difference between the spectra.

931

932 **Figure 10. (A)** FTIR transmission map of glass B10 surrounding an area purposely
933 exposed with a $2 \times 2 \mu\text{m}$ synchrotron radiation beam at a photon flux of $\sim 1 \times 10^{12}$
934 photons/sec for ~ 12 minutes. We show the location of the beam at the sample surface, but
935 the beam intersects the sample surface at a 45° angle and travels through the sample from
936 left to right. Color scale shows the intensity of the H_2O_{Total} absorbance band at 3500 cm^{-1}
937 after subtraction of a reference intensity at 2300 cm^{-1} . Each pixel is $20 \times 20 \mu\text{m}$. FTIR
938 spectra for pixels designated (1) and (2) are discussed in the text.

939

Table 1: Room Temperature Mössbauer hyperfine parameters, Fe³⁺/ΣFe ratios, and select compositional parameters

Sample	Fe ³⁺ /ΣFe * 100				spectral fitting parameters (this study)										B et al., 2005		Φ [^]
	B et al. 2005	1σ	this study	1σ	χ ²	Bkgrd.	Fe ³⁺ CS	Fe ³⁺ δCS	Fe ³⁺ QS	Fe ³⁺ δQS	Fe ²⁺ CS	Fe ²⁺ δCS	Fe ²⁺ QS	Fe ²⁺ δQS	FeO Total Wt.%	H2O Total Wt.%	
B11	19	2	18.7	1.45	1.0	1024543	0.46	0.19	0.93	0.14	0.95	0.12	1.98	0.44	12.49	4.82	1.05
B16	18	2	16.8	1.23	0.8	1427668	0.53	0.09	0.92	0.39	0.95	0.14	1.97	0.41	12.93	3.02	0.81
B12	11	2	12.3	1.81	1.2	598904	0.52	0.28	1.12	0.01	0.97	0.14	2.00	0.42	12.02	1.9	0.77
B13	10	2	12	1.41	1.0	1988275	0.48	0.34	1.00	0.02	0.94	0.03	1.91	0.48	11.39	1.71	0.72
B10	40	4	36.3	1.61	1.1	680688	0.50	0.02	0.90	0.38	0.91	0.28	2.02	0.12	12.61	2.43	0.24
B6	39	3	36.9	1.72	1.2	644411	0.44	0.02	0.91	0.46	0.96	0.12	1.93	0.38	12.94	2	0.19
B9	29	3	26.7	2.14	1.4	570588	0.55	0.19	0.82	0.00	0.91	0.28	1.99	0.10	12.62	1.14	0.19
B7	23	2	21.1	0.94	0.6	1068960	0.57	0.18	0.87	0.00	0.92	0.28	1.94	0.12	12.91	0.97	0.22
B8	16	2	17.8	1.15	0.8	905216	0.59	0.02	0.92	0.34	0.93	0.28	1.92	0.14	12.67	0.82	0.23
B17	60	2	55.6	2.03	1.4	392639	0.50	0.16	0.88	0.29	0.93	0.29	2.11	0.09	12.65	4.67	0.19
B21	49	2	45.7	1.58	1.1	735417	0.50	0.05	0.90	0.38	0.89	0.31	2.04	0.15	12.34	2.9	0.19
B19	52	2	49.3	1.68	1.1	628828	0.54	0.15	0.82	0.28	0.81	0.28	2.21	0.11	12.91	1.85	0.11
B22	28	2	25.3	0.38	0.3	5545128	0.58	0.20	0.82	0.00	0.89	0.28	1.98	0.11	12.77	0.94	0.17
B20	25	2	24.8	1.00	0.7	2254268	0.58	0.20	0.83	0.00	0.90	0.28	1.98	0.11	12.80	0.72	0.13
			C et al., 2009 ^v	1σ	χ ²	Bkgrd.	Fe ³⁺ CS	Fe ³⁺ δCS	Fe ³⁺ QS	Fe ³⁺ δQS	Fe ²⁺ CS	Fe ²⁺ δCS	Fe ²⁺ QS	Fe ²⁺ δQS			
All_0	-	-	14.7	0.38	0.55	5408375	0.44	0.09	1.31	0.47	0.95	0.15	1.88	0.44	-	-	-
All_05	-	-	19.4	0.63	0.63	7463896	0.50	0.17	1.27	0.46	0.95	0.16	1.88	0.42	-	-	-
All_-05	-	-	13.8	0.50	0.5	6400103	0.44	0.04	1.30	0.63	0.95	0.18	1.88	0.37	-	-	-
All_15	-	-	24.7	0.44	0.44	5948886	0.44	0.03	1.26	0.50	0.95	0.17	1.88	0.42	-	-	-
All_-15	-	-	9.3	0.60	0.6	694310	0.44	0.33	1.29	0.03	0.95	0.14	1.88	0.45	-	-	-
All_25	-	-	35.6	0.38	0.38	5951129	0.41	0.09	1.27	0.50	0.95	0.18	1.88	0.41	-	-	-
All_35	-	-	48.4	0.40	0.4	12486195	0.38	0.10	1.25	0.48	0.93	0.17	1.88	0.14	-	-	-
All_45	-	-	61.1	0.26	0.26	10318722	0.37	0.09	1.26	0.48	0.91	0.18	1.88	0.44	-	-	-
LW_0	-	-	16.2	0.55	0.38	7123408	0.44	0.14	1.23	0.55	0.95	0.16	1.90	0.42	-	-	-
LW_10	-	-	23.5	0.53	0.53	7922396	0.46	0.17	1.22	0.40	0.95	0.15	1.91	0.43	-	-	-
LW_-10	-	-	13.0	0.58	0.58	9201577	0.44	0.10	1.25	0.70	0.95	0.16	1.90	0.42	-	-	-
LW_20	-	-	30.3	0.14	1.4	2072205	0.30	0.12	1.22	0.48	0.86	0.10	1.91	0.47	-	-	-
LW_-20	-	-	8.8	0.53	0.53	6835643	0.44	0.03	1.28	0.03	0.95	0.15	1.90	0.42	-	-	-

^v Cottrell et al., 2009. Recommended Fe³⁺/ΣFe ratios, considering recoilless fraction, can be found in Zhang et al., submitted. Room Temperature values are used in this work for comparison with Botcharnikov et al., 2005.

[^] Φ = XH_{0.5}*XFeO/XFeO_{1.5} based on Fe³⁺/ΣFe ratios from this study and H₂O from Botcharnikov et al., 2005.

940

Table 2: Synchrotron radiation dose conditions

Flux Condition	spot size (μm)	photon flux (photons/sec)	flux density (photons/sec/ μm^2)	Gray/sec**	notes
XANES ₁	2x2	$5 \times 10^{10-11}$	$1 \times 10^{10-11}$	$9.0 \times 10^{10-11}$	nominal at 13-ID-D
XANES ₃	2x2	$1.5-8.6 \times 10^9$	4×10^8 to 2×10^9	$\sim 1.3 \times 10^7$	
XANES ₄	50x50	$1.5-8.6 \times 10^9$	8×10^5 to 3×10^6	$\sim 1.3 \times 10^4$	
XANES ₅	25x25	2×10^9	3×10^6	3.8×10^4	
XANES _{CRYO-HI}	38x28	$3-4 \times 10^7$	$3-4 \times 10^4$		monochromator = 7113.3eV
XANES _{CRYO-LO}	38x28	$2-4 \times 10^6$	$2-4 \times 10^3$		monochromator = 7113.3eV
X26A	9x5	3×10^9	6×10^7	7.7×10^5	

**Dose in Gray (Joules/kg/sec) is calculated for glass of basaltic composition through 1/e absorption length

Table 3: XANES centroids and integrated pre-edge intensities

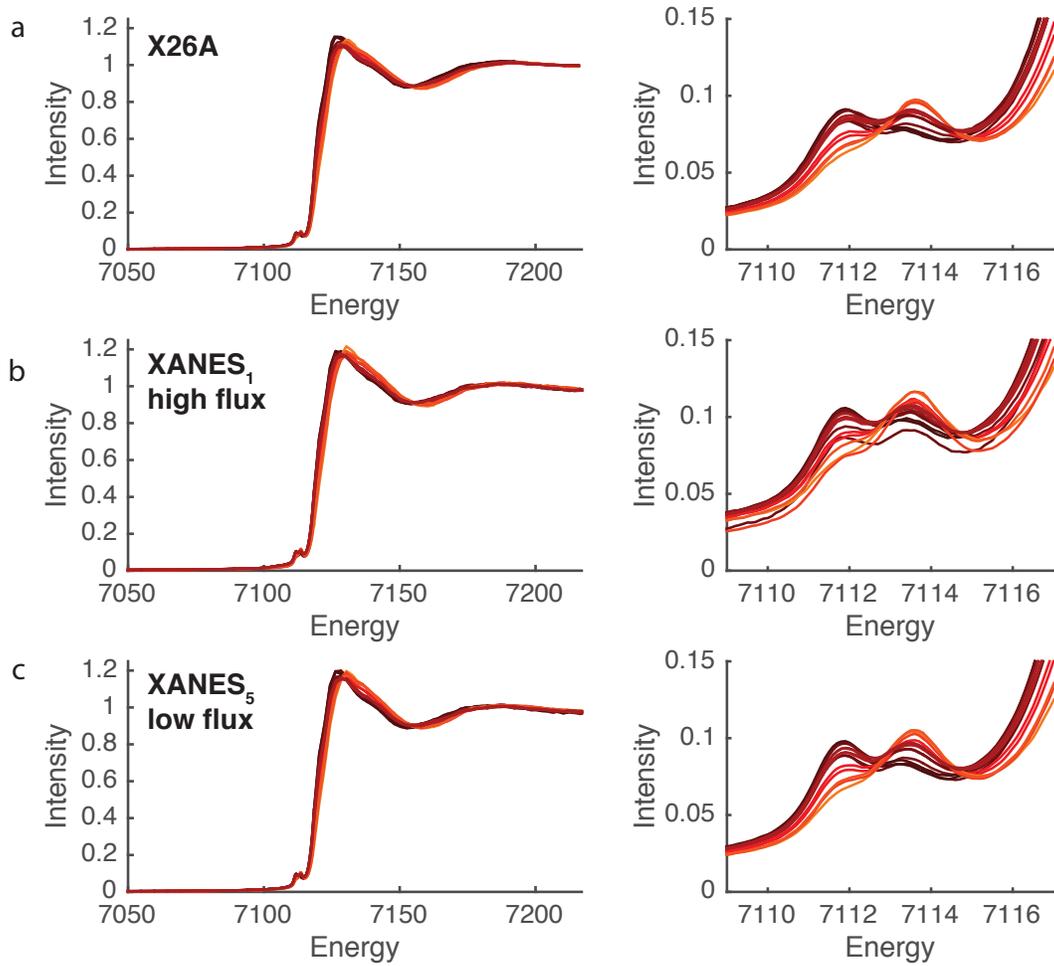
sample	X26A♦				XANES ₁				XANES ₅				
	5x10 ¹⁰ photons/um ²				3x10 ¹⁵ photons/um ²				2x10 ⁹ photons/um ²				
	centroid	σ	I1	I2	centroid	σ	I1	I2	centroid	σ	I1	I2	%Δ**
B11	7112.402	0.012	0.06	0.037	7112.654	0.015	0.053	0.066	7112.421	0.003	0.058	0.045	71.9
B16	7112.342	0.008	0.063	0.031	7112.578	0.007	0.057	0.057	7112.362	0.01	0.061	0.04	69.2
B12	7112.226	0.007	0.073	0.026	7112.4	0.016	0.07	0.046	7112.228	0.002	0.074	0.032	66.4
B13	7112.177	0.002	0.077	0.024	7112.328	0.011	0.073	0.038	7112.211	0.015	0.076	0.031	46.0
B10	7112.784	0.01	0.047	0.062	7112.86	0.018	0.045	0.082	7112.773	0.004	0.048	0.074	18.4
B6	7112.733	0.012	0.05	0.06	7112.793	0.017	0.049	0.076	7112.735	0.007	0.052	0.07	6.5
B9	7112.585	0.006	0.06	0.05	7112.633	0.006	0.06	0.066	7112.568	0.011	0.06	0.058	16.5
B7	7112.477	0.007	0.064	0.042	7112.505	0.009	0.066	0.054	7112.468	0.009	0.064	0.05	18.7
B8	7112.413	0.007	0.07	0.04	7112.451	0.026	0.066	0.05	7112.418	0.006	0.067	0.047	27.2
B17	7113.083	0.013	0.03	0.088	7113.108	0.016	0.029	0.106	7113.059	0.001	0.032	0.098	3.8
B21	7112.988	0.006	0.039	0.082	7113.029	0.014	0.039	0.104	7112.957	0.008	0.04	0.088	15.5
B19	7113.01	0.008	0.04	0.083	7112.99	0.011	0.042	0.097	7112.969	0.008	0.042	0.094	2.3
B22	7112.587	0.01	0.062	0.051	7112.613	0.008	0.06	0.064	7112.582	0.01	0.06	0.06	19.1
B20	7112.548	0.002	0.064	0.048	7112.551	0.009	0.064	0.058	7112.538	0.009	0.063	0.057	9.5
	5x10 ¹⁰ photons/um ²				4x10 ¹⁵ photons/um ²								
All0	7112.275	0.01	0.091	0.03	7112.271	0.009	0.079	0.036					
All05	7112.361	0.011	0.086	0.038	7112.368	0.002	0.08	0.046					
All-05	7112.219	0.01	0.09	0.029	7112.22	0.004	0.084	0.034					
All15	7112.52	0.013	0.078	0.047	7112.529	0.002	0.069	0.058					
All-15	7112.147	0.026	0.093	0.02	7112.134	0.007	0.087	0.025					
All25	7112.742	0.014	0.068	0.07	7112.76	0.009	0.061	0.079					
All35	7112.962	0.027	0.055	0.091	7112.974	0.002	0.052	0.106					
All45	7113.184	0.025	0.046	0.124	7113.186	0.003	0.042	0.14					
LW_0	7112.3	0.01	0.085	0.032	7112.302	0.01	0.079	0.04					
LW10	7112.434	0.013	0.078	0.042	7112.433	0.009	0.07	0.05					
LW-10	7112.154	0.024	0.095	0.023	7112.155	0.004	0.085	0.028					
LW20	7112.608	0.02	0.072	0.058	7112.62	0.002	0.069	0.069					
LW-20	7112.112	0.01	0.097	0.018	7112.128	0.002	0.092	0.027					

♦ centroids are 7-yr avgs (n=30-48) as reported in Zhang et al., submitted

** $\text{Fe}^{3+}/\Sigma\text{Fe}_{\text{Möss}} - \text{Fe}^{3+}/\Sigma\text{Fe}_{\text{XANES}}$

Doses reported in this table are per spectrum (photons/sec/μm²*total collection time)

Hydrous B Glasses



Anhydrous LW and All Glasses

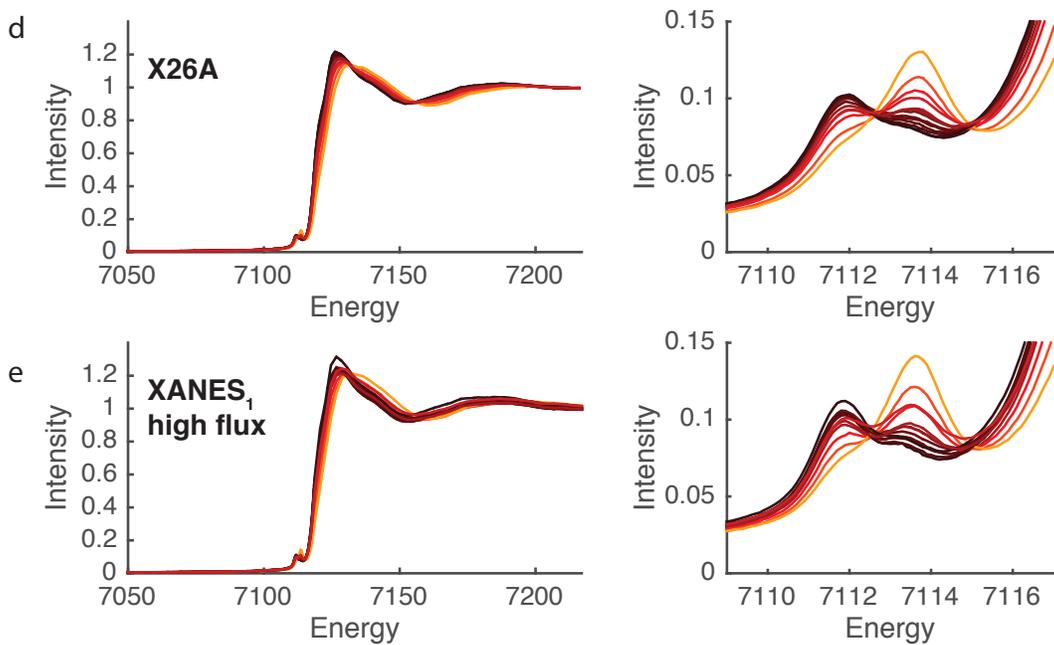


Figure 1

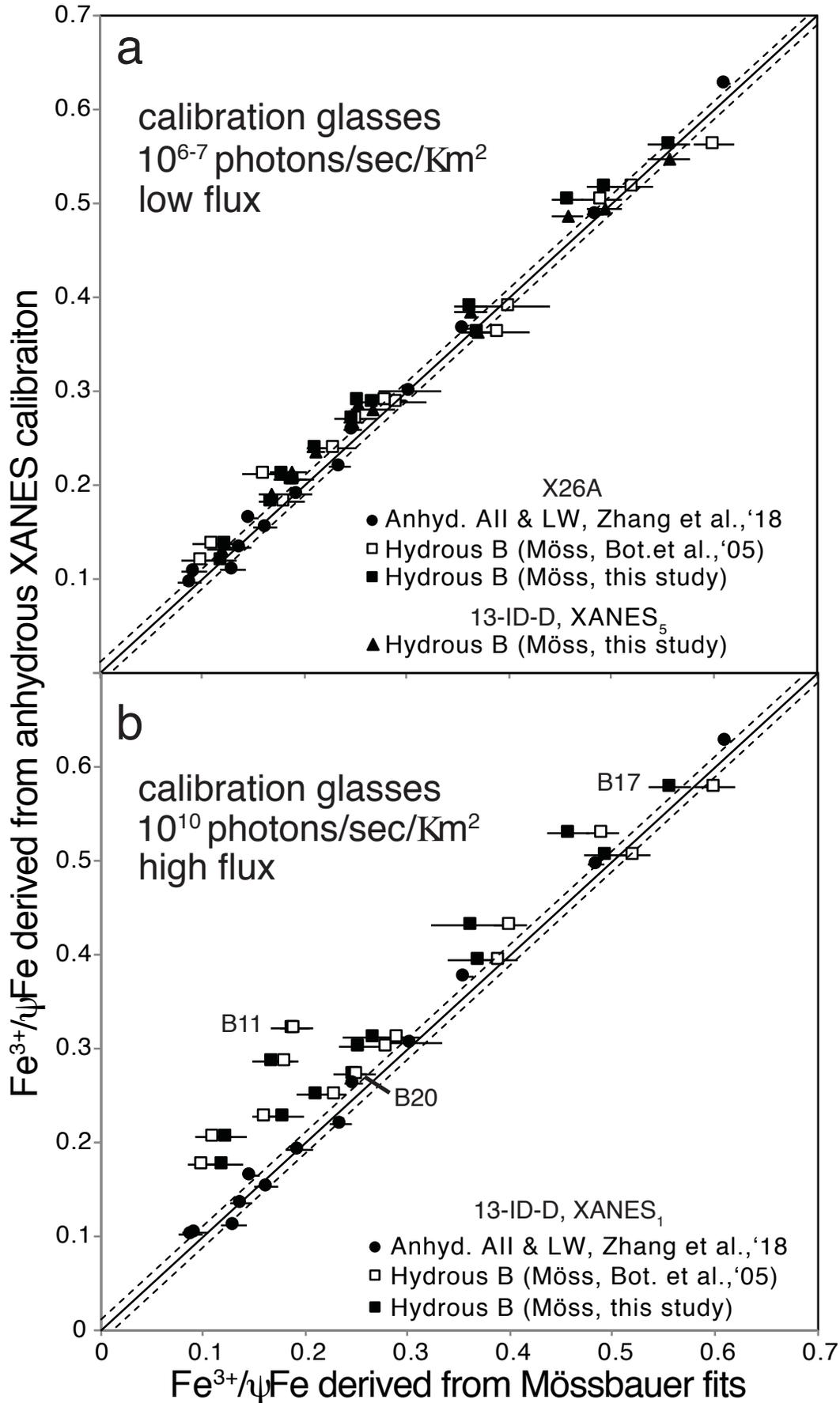


Figure 2

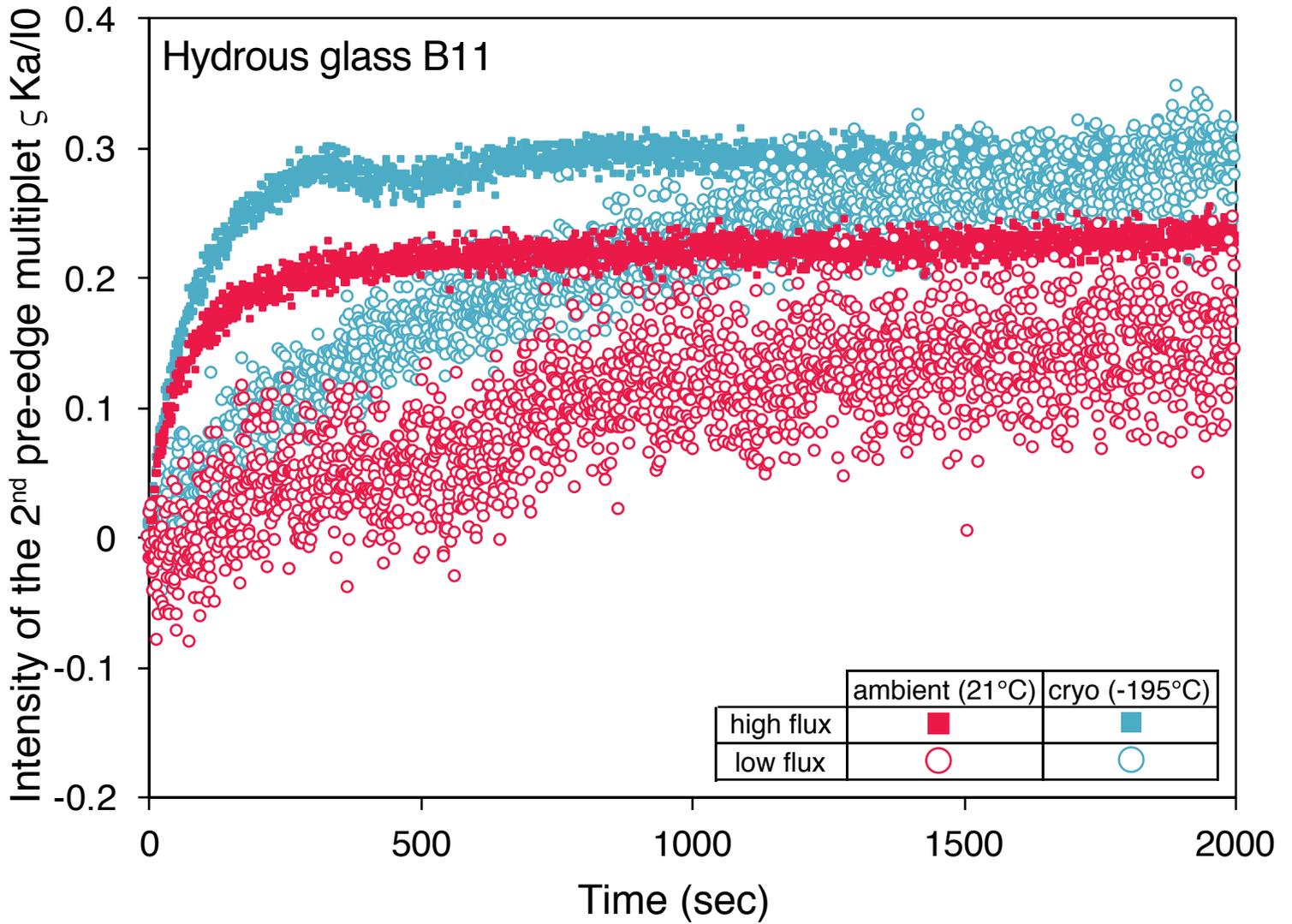


Figure 3

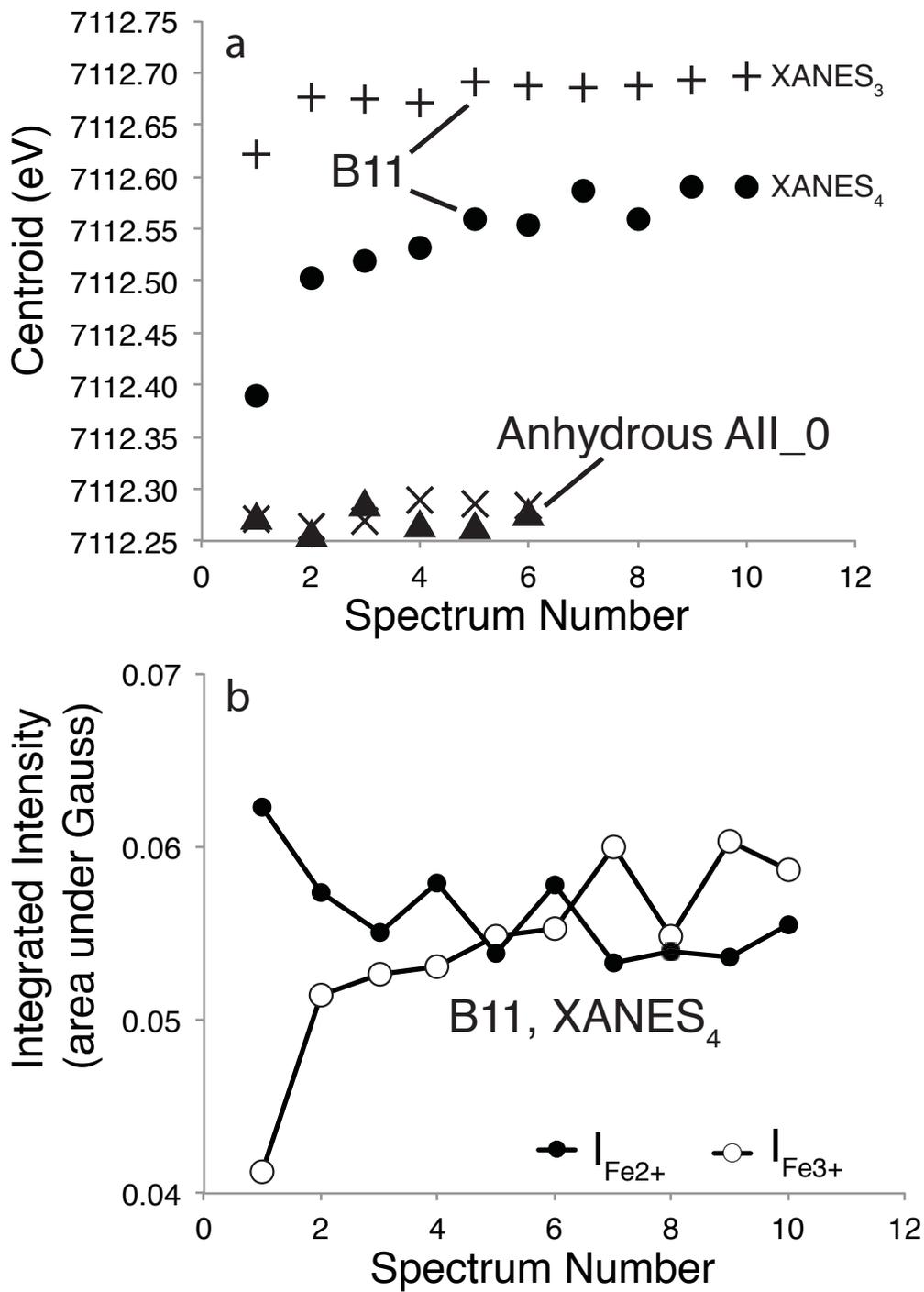


Figure 4

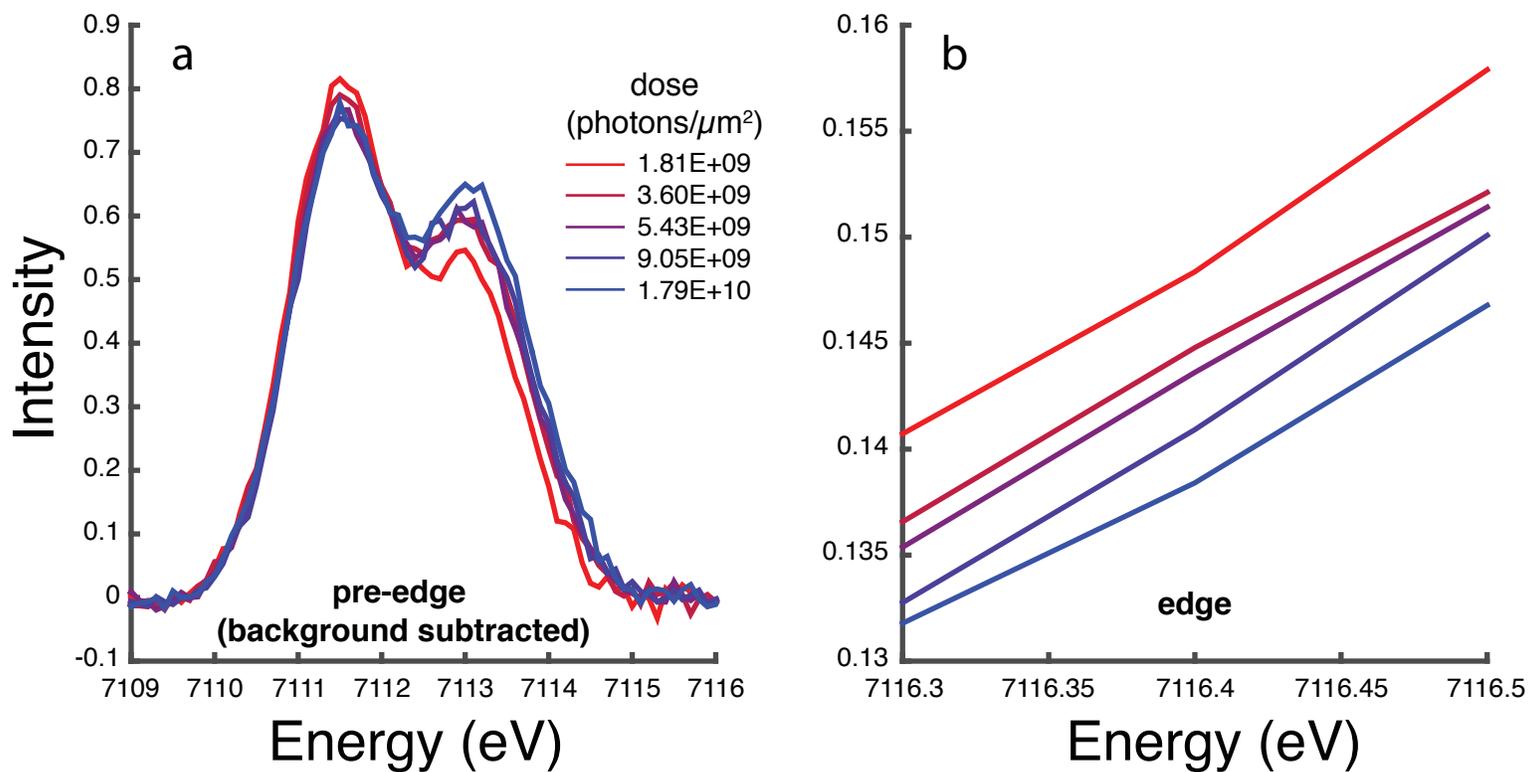


Figure 5

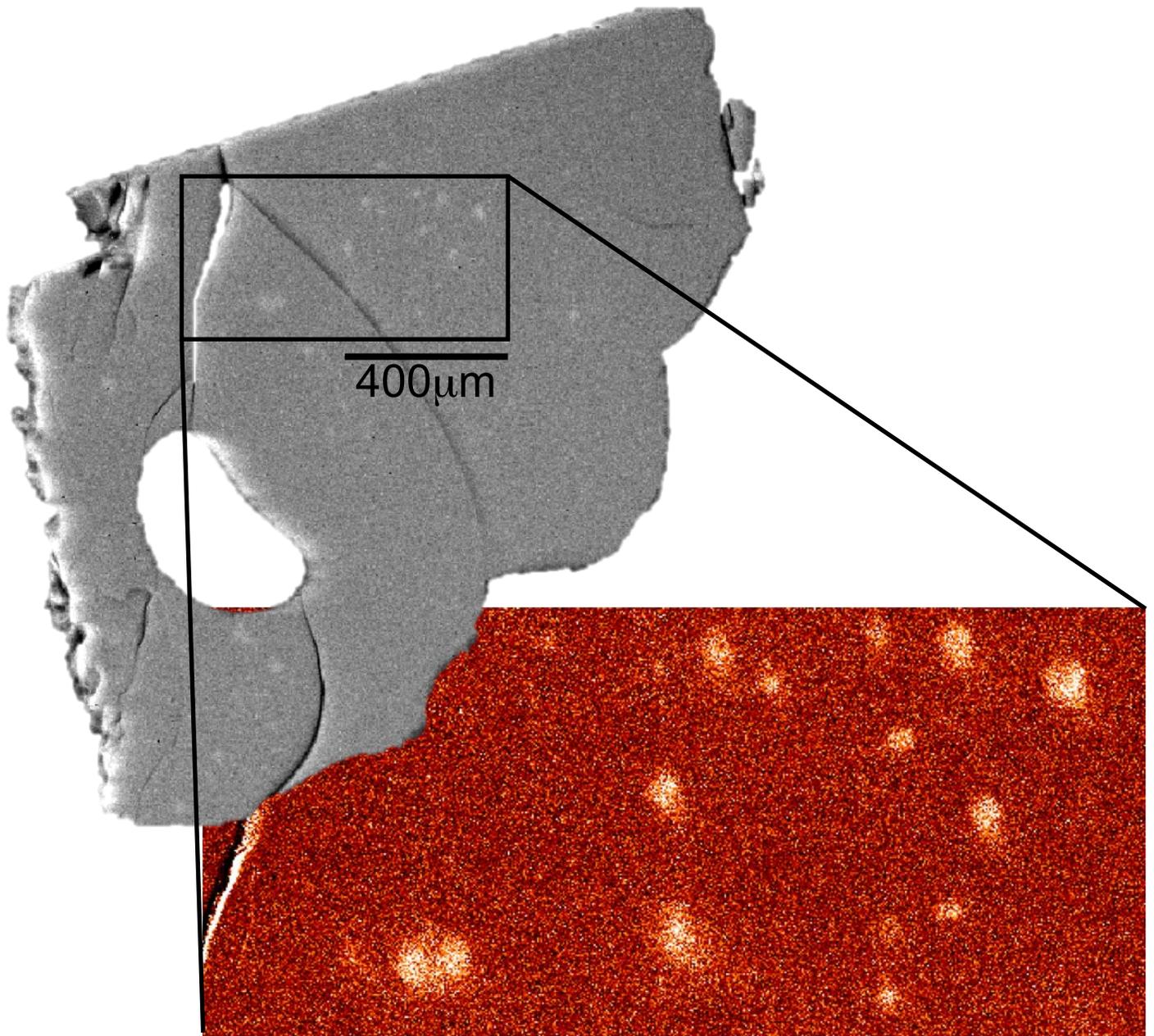


Figure 6

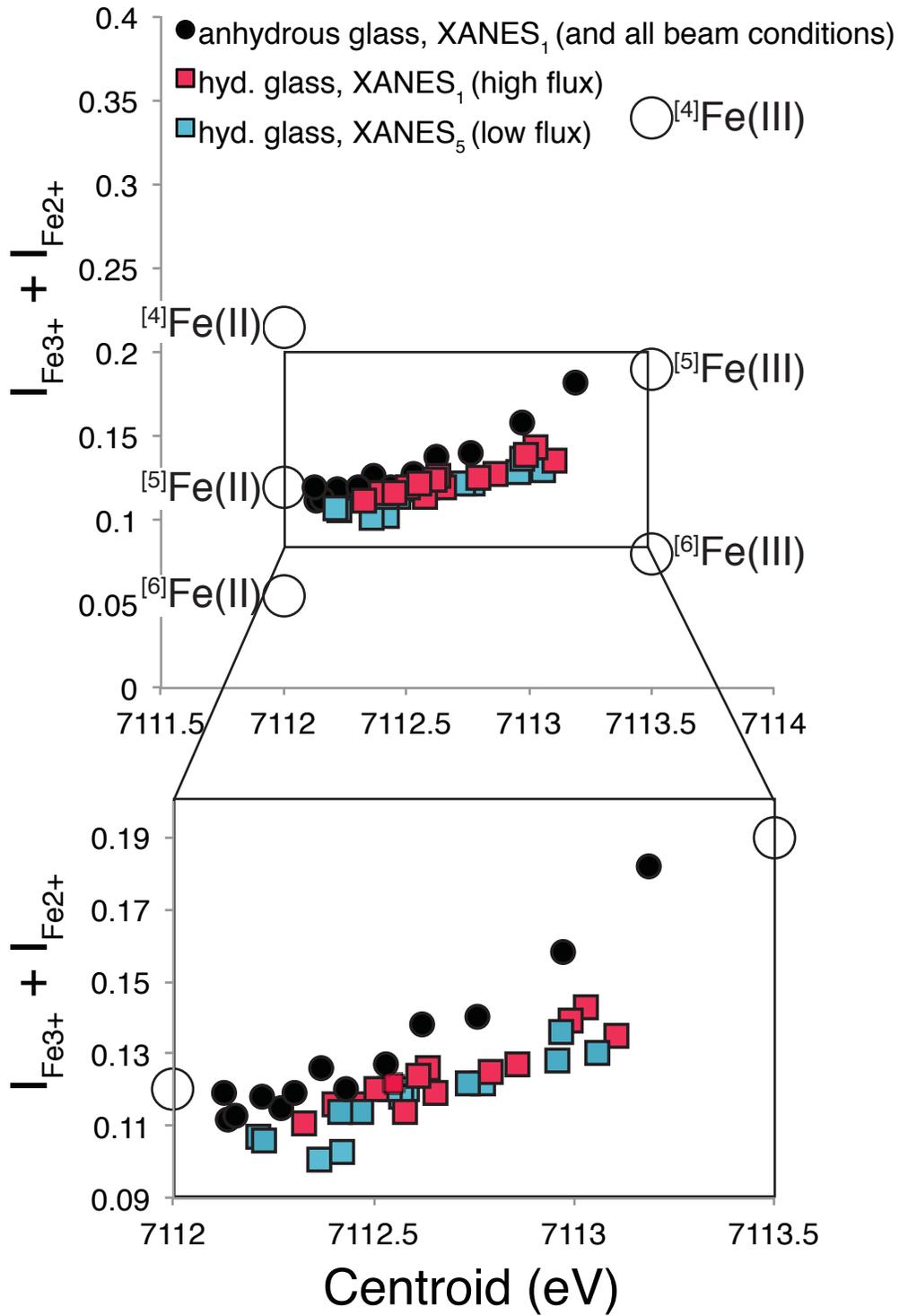


Figure 7

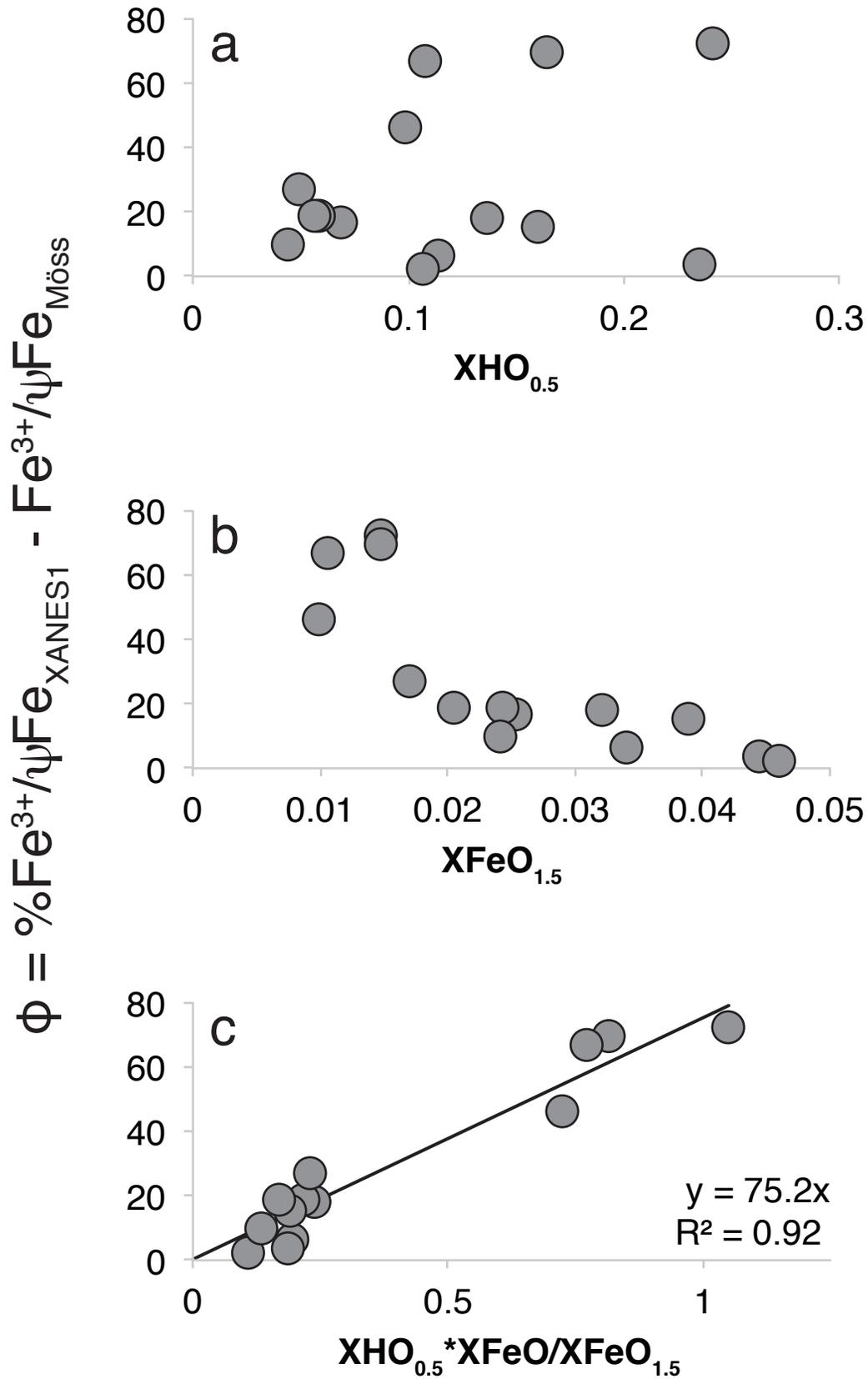


Figure 8

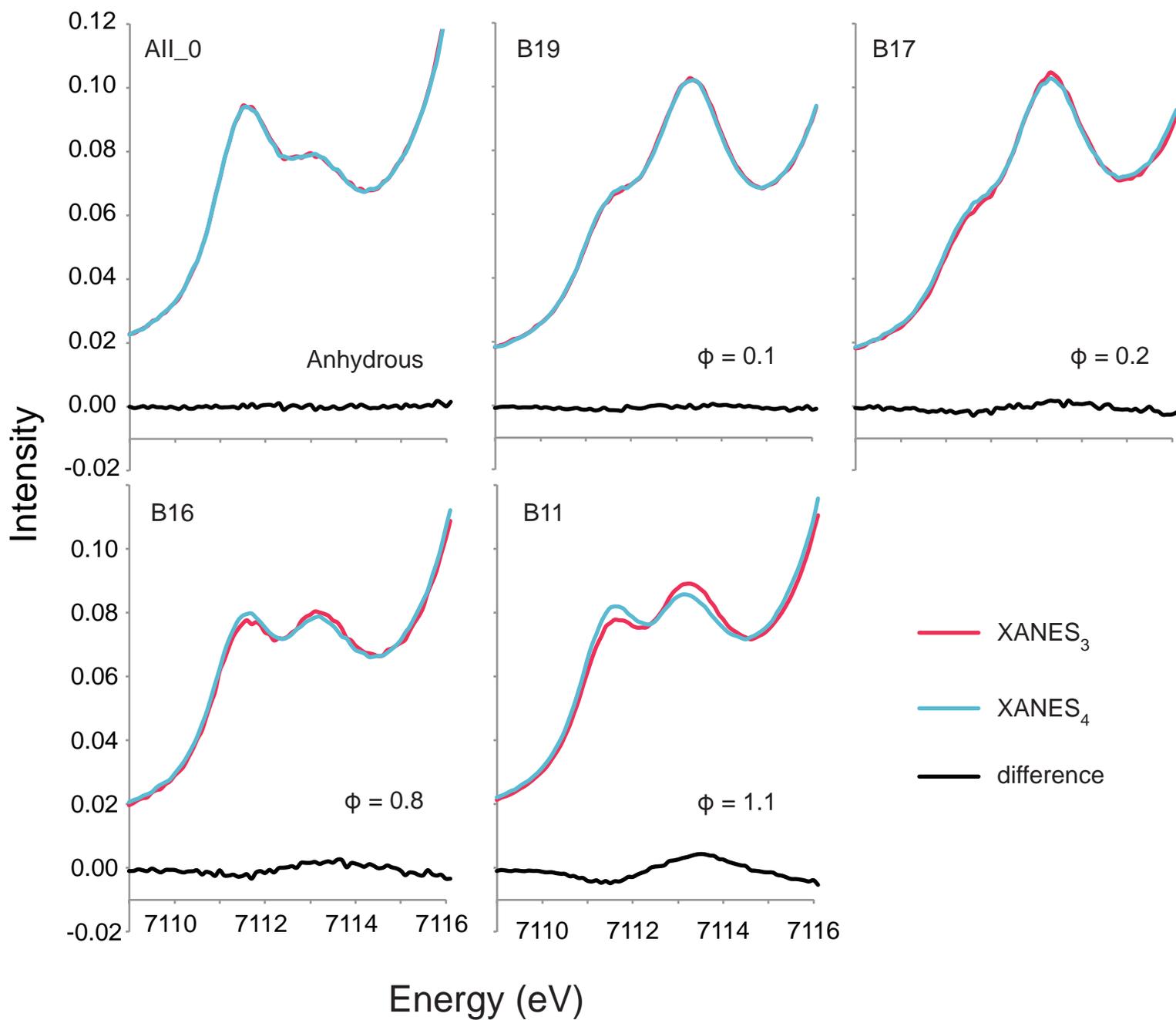


Figure 9

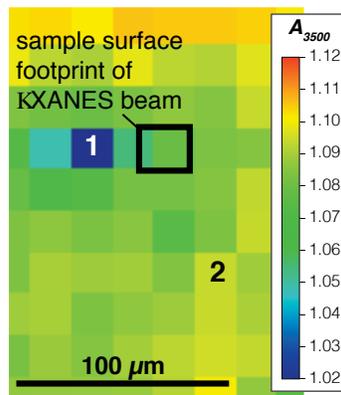


Figure 10