

## The crystal structure of CaGeO<sub>3</sub> perovskite and the crystal chemistry of the GdFeO<sub>3</sub>-type perovskites

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### Abstract

A single-crystal X-ray study indicates that the perovskite form of CaGeO<sub>3</sub> is orthorhombic (*Pbnm*), and isotypic with GdFeO<sub>3</sub>, although it was previously reported as a cubic form; the cell dimensions are  $a = 5.2607(6)\text{\AA}$ ,  $b = 5.2688(10)$ ,  $c = 7.4452(15)$  and  $V = 206.36(6)\text{\AA}^3$  ( $Z = 4$ ;  $D_{\text{calc}} = 5.17\text{ g/cm}^3$ ). The crystal structure is close to that of SmAlO<sub>3</sub>. Although the unit cell of CaGeO<sub>3</sub> is pseudocubic, the structure itself is very distorted through the tilting and distortion of polyhedra. The oxygen polyhedra are less tilted and less distorted than those of other GdFeO<sub>3</sub>-type perovskites. The structural deformation of the GdFeO<sub>3</sub>-type perovskite is determined primarily by the size-ratio of two kinds of cation occupying A and B sites. Some structural characteristics such as O(2)–O(2)–O(2) and A–O(1)–B angles and bond-length distortions exhibit systematic relationships as a function of the observed tolerance factor which is newly defined here. A strong correlation between the Goldschmidt tolerance factor and the observed tolerance factor has made possible some predictions for GdFeO<sub>3</sub>-type perovskites.

### Introduction

The perovskite structure, with general formula ABO<sub>3</sub>, consists of a framework of B octahedra that share corners with each other and triangular faces with cuboctahedra containing A cations in twelve coordination. For compositions in which their constituent atoms are not of the ideal relative size, many distorted types of perovskite may replace the ideal structure. Some metasilicates and metagermanates are known to crystallize as perovskites in cubic (SrGeO<sub>3</sub>; above 50 kbar), hexagonal (BaGeO<sub>3</sub>; above 95 kbar), and orthorhombic phases (MgSiO<sub>3</sub>; above 300 kbar) (Shimizu *et al.*, 1970; Liu, 1976a; Yagi *et al.*, 1978; Ito and Matsui, 1978). The high-pressure transformation of a metasilicate to a perovskite form is geophysically important for interpreting seismic wave velocities in the Earth's lower mantle. Also, germanates are useful as structural analogs of common silicate minerals at high pressures because new dense phases frequently exist at much lower pressures than the corresponding isotypic silicates.

CaGeO<sub>3</sub> is a typical example of such germanates and is known to transform from the wollastonite through a garnet-like structure to the perovskite structure at more than 65 kbar and 900°C (Susaki and Akimoto, private comm.). CaGeO<sub>3</sub> has previously been indexed as cubic perovskite with  $a = 3.723\text{\AA}$  (Ringwood and Major, 1967), although Prewitt and Sleight (1969) reported a doubling of the unit cell ( $a = 7.448\text{\AA}$ ). In order to confirm the cell dimension and the space group of CaGeO<sub>3</sub> perovskite, we

examined a single crystal using X-ray diffraction techniques, to refine its crystal structure, and to make a systematic study of crystallographic correlations among the GdFeO<sub>3</sub>-type perovskites. A preliminary communication on this result has been reported (Sasaki *et al.*, 1981).

### Experimental

#### Sample

A polycrystalline specimen of CaGeO<sub>3</sub> perovskite was synthesized by hot-pressing CaGeO<sub>3</sub> wollastonite powder in squeezer solid-media apparatus for two hours at  $P = 100\text{ kbar}$  and at  $T = 1000^\circ\text{C}$  (see details in Liebermann *et al.*, 1977). After sintering at elevated pressure and temperature, the run was slowly cooled (20–60 minutes) to room temperature after which the pressure was released. Examination of the recovered specimen using a polarizing microscope and X-ray powder diffraction analyses confirmed that it was a single phase with the perovskite structure. A single crystal of parallelepiped shape and dimensions,  $0.14 \times 0.10 \times 0.07\text{ mm}$  was extracted from the polycrystalline aggregate and prepared for single-crystal X-ray diffraction study.

#### Space group determination

The unit cell and space group of CaGeO<sub>3</sub> perovskite were determined with the aid of precession and Weissenberg photographs and intensity data collected with a Picker four-circle diffractometer. The space group deter-

mination was more difficult than expected, because the unit cell is pseudocubic, dimensionally and in intensity distribution for strong reflections. The cell dimensions determined with the four-circle diffractometer (graphite monochromatized  $\text{MoK}\alpha_1$ ;  $\lambda = 0.70926\text{\AA}$ ) are as follows:

$$\begin{aligned} a &= 5.2607(6)\text{\AA} \\ b &= 5.2688(10) \\ c &= 7.4452(15) \\ V &= 206.36(6)\text{\AA}^3 \\ Z &= 4 \text{ (formula units/cell)} \\ D_{\text{calc}} &= 5.17 \text{ g/cm}^3 \end{aligned}$$

The cell dimensions have the relation  $a \cong b \cong \sqrt{2} a_p$  and  $c \cong 2a_p$  ( $a_p$ : a pseudocubic subcell parameter), maintaining the characteristic of  $\text{GdFeO}_3$ -type perovskites. In order to emphasize the difference from cubic perovskite, the precession photograph for  $hk1$  is shown in Figure 1. If this crystal were cubic with  $a = 3.7\text{\AA}$ , no diffraction spots would appear on photographs for  $l = 2n + 1$  on the orthorhombic cell. The intensity distribution indicates orthorhombic symmetry. The systematic absences are: no conditions on  $hkl$ ;  $k = 2n$  on  $0kl$ ;  $h + l = 2n$  on  $h0l$ ; no conditions on  $hk0$ ; ( $h = 2n$ ) on  $h00$ ; ( $k = 2n$ ) on  $0k0$ ; ( $l = 2n$ ) on  $00l$ . The possible space group is either  $D_{2h}^{16}$ - $Pbnm$  or  $C_{2v}^2$ - $Pbn2_1$ . On the basis of no evidence for noncentro-

symmetry and the results of least-squares refinement of the structure, the space group is most likely  $Pbnm$ . Although some weak reflections apparently violating the above determination could be detected, they were reflections caused by a small twinned part of the crystal, as will be described in the following section.

### Twining

Preliminary investigation using the four-circle diffractometer and long-exposure precession photographs showed the presence of two types of reflections inconsistent with the  $\text{GdFeO}_3$ -type perovskite of space group  $Pbnm$ . The first type is of reflections on the reciprocal points at half-integral coordinates, such as  $1/2\ 5/2\ 0$ ,  $1/2\ 5/2\ 2$ ,  $1/2\ 3/2\ 3$ ,  $1/2\ 1/2\ 4$ ,  $1/2\ 5/2\ 4$ ,  $3/2\ 3/2\ 6$ , and  $1/2\ 5/2\ 6$ . Another is of the  $031$  and  $013$  reflections that apparently violate the extinction rule for the  $b$ -glide plane. All of the above reflections, however, can be clearly explained by the existence of twinning related by the following transformation matrices:

$$\begin{pmatrix} a^{*'} \\ b^{*'} \\ c^{*'} \end{pmatrix} = \begin{pmatrix} 1/2 & -1/2 & 1 \\ 1/2 & -1/2 & -1 \\ 1/2 & 1/2 & 0 \end{pmatrix} \begin{pmatrix} a^* \\ b^* \\ c^* \end{pmatrix}$$

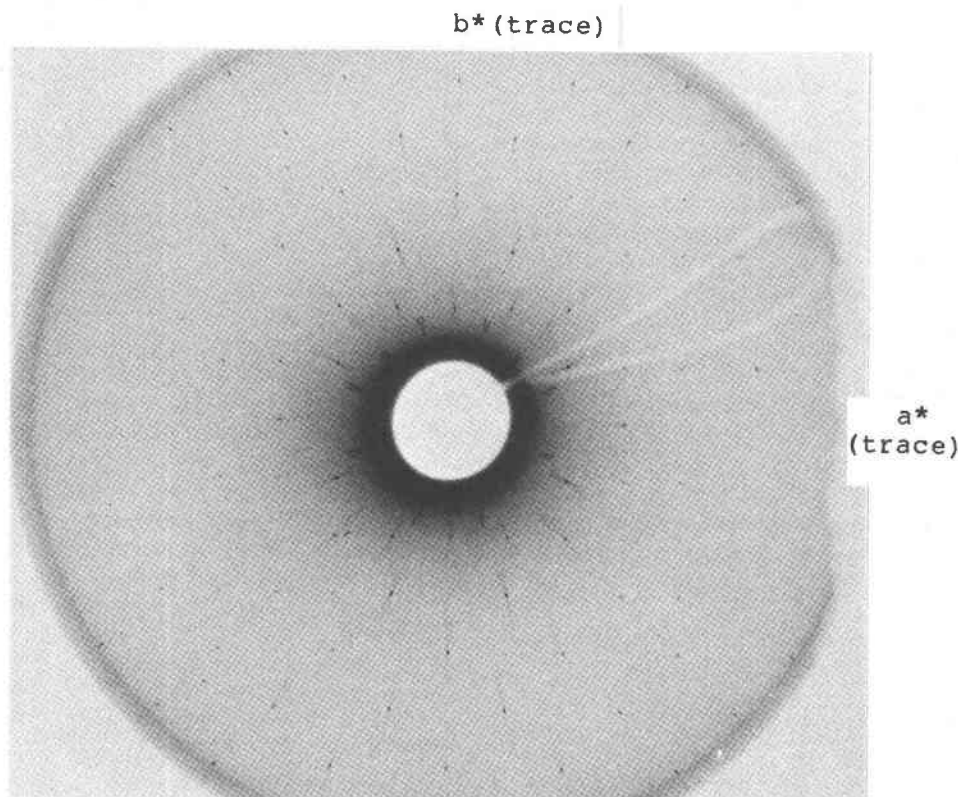


Fig. 1. Precession photograph of the  $hk1$  net for  $\text{CaGeO}_3$  perovskite, taken with  $\mu = 25^\circ$  (non-filtered  $\text{MoK}\alpha$  radiation; 35 kV and 15 mA for 336 hours). If this perovskite were cubic with  $a = 3.7\text{\AA}$ , no diffraction spots would appear for this net. The symmetry  $mmm$  (orthorhombic) is required for Laue group.

Table 1. Fractional atomic coordinates and thermal parameters ( $\times 10^4$ ) of CaGeO<sub>3</sub> perovskite obtained in least-squares refinements: (1) after and (2) before the correction of intensities diffracted by the twinned crystal. The  $\beta_{ij}$ 's are defined by  $\exp[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23})]$ .

(1) After correction:				
	Ca	Ge	0(1)	0(2)
x	-0.0051(3)	0	0.0606(9)	0.7157(6)
y	0.0283(2)	1/2	0.4916(8)	0.2830(6)
z	1/4	0	1/4	0.0320(5)
$\beta_{11}$	46(2)	21(1)	36(7)	49(5)
$\beta_{22}$	66(2)	25(1)	66(8)	48(5)
$\beta_{33}$	30(1)	10(1)	15(4)	24(2)
$\beta_{12}$	-9(2)	0(1)	4(6)	-16(5)
$\beta_{13}$	0	-0(1)	0	1(3)
$\beta_{23}$	0	-0(.4)	0	17(3)
B(equiv)	0.64(2)	0.24(1)	0.49(8)	0.53(5)
Ref. (meas)	689			
(used)	412			
R	0.035			
R(wt)	0.037			
(2) Before correction:				
	Ca	Ge	0(1)	0(2)
x	-0.0051(5)	0	0.0587(16)	0.7159(10)
y	0.0273(3)	1/2	0.4911(17)	0.2829(9)
z	1/4	0	1/4	0.0321(9)
$\beta_{11}$	46(4)	24(3)	42(17)	56(13)
$\beta_{22}$	68(4)	31(3)	118(23)	49(12)
$\beta_{33}$	24(4)	11(3)	22(17)	34(11)
$\beta_{12}$	-6(7)	-0(3)	2(17)	-20(11)
$\beta_{13}$	0	0(2)	0	11(10)
$\beta_{23}$	0	0(1)	0	1(10)
B(equiv)	0.60(6)	0.29(4)	0.75(27)	0.64(17)
Ref. (meas)	689			
(used)	343*			
R	0.035			
R(wt)	0.052			

\*The reflections greater than 4 $\sigma$  were used in refinement.

$$\begin{pmatrix} h' \\ k' \\ l' \end{pmatrix} = \begin{pmatrix} 1/2 & -1/2 & 1/2 \\ 1/2 & -1/2 & -1/2 \\ 1 & 1 & 0 \end{pmatrix} \begin{pmatrix} h \\ k \\ l \end{pmatrix}$$

There is no doubling of spots or peaks for any reflections that overlap. Since the reflections at half integral reciprocal points do not overlap each other, the comparison in intensities of such reflections gave roughly the volume ratio of 1:0.14(1).

### Data collection and structure refinement

The  $\omega$ -2 $\theta$  scan technique of the four-circle diffractometer was used to collect at room temperature the intensity data up to  $2\theta = 80^\circ$  in the bisecting mode. Each reflection was scanned at the appropriate speed with a maximum measurement time of four minutes; the scan width on  $2\theta$  ( $^\circ$ ) was  $2.0 + 0.7 \tan\theta$ . A standard reflection used was stable within  $\pm 1.5\%$  of the integrated intensity during data collection. Of a total of 689 reflections measured, the reflections less than three times the calculated standard deviation were omitted in refinements. A set of intensities was corrected for Lorentz and polarization factors and a prismatic absorption correction was applied using the program ACACA (Wuensch and Prewitt, 1965). The linear

Table 2. Interatomic distances ( $\text{\AA}$ ) and angles ( $^\circ$ ) in the Ge octahedron. Standard deviations are in parentheses. Square brackets, [ ], show the value calculated for the ideal cubic perovskite with the same polyhedral volumes.

Ge 1 - 0(1) i	x 2	1.889(.9)
- 0(2) v	x 2	1.898(3)
- 0(2) viii	x 2	1.889(3)
Mean <Ge - O>		1.892 [1.861]
Bond-length distortion, $\Delta_B$		0.005
0(1) i - 0(2) v	x 2	2.684(5)
- 0(2) iv	x 2	2.665(5)
- 0(2) i	x 2	2.671(5)
- 0(2) viii	x 2	2.677(5)
0(2) v - 0(2) viii	x 2	2.659(4)
- 0(2) iv	x 2	2.696(4)
Mean <O - O>		2.675 [2.632]
0(1) i - Ge 1	- 0(2) viii	90.3(2)
0(1) i -	- 0(2) v	90.3(2)
0(2) v -	- 0(2) viii	89.2(1)
0(2) i -	- 0(2) viii	90.8(1)

absorption coefficient is  $177.1 \text{ cm}^{-1}$  and the transmission factors for this crystal varied from 0.19 to 0.33.

All parameters were refined simultaneously using the full-matrix least-squares program, RADY, which is a modified version of RADIEL (Coppens *et al.*, 1979). The

Table 3. Interatomic distances ( $\text{\AA}$ ) and angles ( $^\circ$ ) with standard deviation in parentheses for cuboctahedron. Square bracket, [ ], shows the value calculated for the ideal cubic perovskite with the same polyhedral volume.

Ca 1 - 0(1) i	x 1	2.346(5)	
- 0(1) iv	x 1	2.465(5)	
- 0(1) iiii	x 1	2.929(5)	
- 0(1) v	x 1	2.849(5)	
- 0(2) i'	x 2	2.352(4)	
- 0(2) iv'	x 2	2.568(4)	
- 0(2) v	x 2	2.597(4)	
- 0(2) viii	x 2	3.069(4)	
Mean <M - O>		2.647 [2.632]	
Bond-length distortion, $\Delta_A$		9.40	
0(1) i - 0(1) v	x 2	3.303(7)	
0(1) v - 0(1) iiii	x 2	4.197(7)	
- 0(2) vii	x 4	2.684(5)	
- 0(2) ii'	x 4	2.665(5)	
0(1) iv - 0(2) iiii'	x 4	2.671(5)	
- 0(2) vi	x 4	2.677(5)	
0(2) i' - 0(2) ii'	x 2	3.246(5)	
0(2) viii - 0(2) viii	x 2	4.199(6)	
0(2) i' - 0(2) iv'	x 4	2.659(4)	
0(2) iiii' - 0(2) vi	x 4	2.696(4)	
0(2) i' - 0(2) iiii'	x 4	4.196(5)	
- 0(2) vii	x 4	4.571(5)	
- 0(2) v	x 2	3.256(4)	
0(2) iv' - 0(2) viii	x 2	4.251(4)	
0(1) i - Ca 1	- 0(1) iv	x 1	86.67(16)
0(1) iv -	- 0(1) iiii	x 1	101.83(15)
0(1) iiii -	- 0(1) v	x 1	93.18(13)
0(1) v -	- 0(1) i	x 1	78.31(15)
0(2) v -	- 0(2) vi	x 2	107.89(11)
0(2) iiii' -	- 0(2) iv'	x 2	78.41(12)
0(2) i' -	- 0(2) v	x 2	82.13(11)
0(2) iv' -	- 0(2) viii	x 2	97.51(10)
0(2) viii -	- 0(2) viii	x 2	86.34(10)
0(2) i' -	- 0(2) ii'	x 2	87.27(13)
0(2) v -	- 0(2) viii	x 2	55.23(9)
0(2) viii -	- 0(2) i'	x 2	57.85(10)
0(2) i' -	- 0(2) iv'	x 2	65.26(11)
0(2) iv' -	- 0(2) v	x 2	62.92(10)
0(1) i -	- 0(2) vii	x 2	57.42(9)
0(1) v -	- 0(2) vii	x 2	53.79(8)
0(1) v -	- 0(2) ii'	x 2	60.77(10)
0(1) iiii -	- 0(2) ii'	x 2	59.58(10)
0(1) iiii -	- 0(2) iiii'	x 2	57.56(10)
0(1) iv -	- 0(2) iiii'	x 2	64.06(11)
0(1) iv -	- 0(2) vi	x 2	63.80(9)
0(1) i -	- 0(2) vi	x 2	65.54(9)

Table 4. The metal-metal distances (Å) and Ge—O—Ge and O(2)—O(2)—O(2) angles (°), with standard deviations in parentheses.

Ge i - Ge i11	3.7227(4)
- Ge ii	3.7226(8)
Ca i - Ca i11	3.7606(19)
- Ca v	3.6853(19)
- Ca ii	3.7349( 8)
Ca i - Ge i	3.248(1)
- Ge i11	3.349(1)
- Ge v	3.204(1)
- Ge vi1	3.105(1)
Ge i - O(1) i - Ge ii	160.4(3)
- O(2) viii - Ge i11	159.0(2)
Ca i - O(1) i - Ge i	99.6(1)
Ca i11 - O(1) i - Ge i	90.0(1)
Ca i - O(2) i' - Ge v	97.3(1)
- O(2) iv' - Ge v	90.6(1)
- O(2) iv' - Ge vi1	86.7(1)
- O(2) v - Ge vi1	86.1(1)
- O(2) v - Ge i	91.2(1)
- O(2) viii - Ge i	77.9(1)
- O(2) viii - Ge i11	81.1(1)
- O(2) i' - Ge i11	103.8(1)
O(2) i - O(2) viii - O(2) i'	154.72(18)
O(2) viii - O(2) i' - O(2) iv'	105.10(14)
O(2) i' - O(2) i1' - O(2) vi1	100.18(14)

residuals of the function  $\sum w_i(|F_{\text{obs}}| - |F_{\text{calc}}|)^2$  were minimized with  $w_i$  kept as unity. Atomic scattering factor tables and anomalous dispersion coefficients were taken from *International Tables for X-Ray Crystallography, Vol. IV* (1974), Fukamachi (1971), and Tokonami (1964) and a secondary extinction correction was applied following Becker and Coppens (1974).

## Results and discussion

### Crystal structure of CaGeO<sub>3</sub>

The final positional and thermal parameters obtained are listed in Table 1<sup>1</sup>, in which we include data both before and after correction for the intensity of the twinned crystal. As seen from the table, the shift of atomic parameters due to the twinned crystal is relatively small. The interatomic distances and angles are also given in Tables 2, 3, and 4. Figure 2 shows a projection of the CaGeO<sub>3</sub> structure on (001) along with labels for the atoms related by symmetry operations. The atomic coordinates for this perovskite structure are similar to those of SmAlO<sub>3</sub> (Marezio *et al.*, 1972).

The cell distortion factor,  $d$ , is useful for estimating the departure from an ideal cubic model (for which  $d = 0$ ):

$$d = \{(a/\sqrt{2} - a_p)^2 + (b/\sqrt{2} - a_p)^2 + (c/2 - a_p)^2\} / 3a_p^2 \times 10^4,$$

where  $a_p = (a/\sqrt{2} + b/\sqrt{2} + c/2)/3$ . The value for

CaGeO<sub>3</sub> ( $d = 0.004$ ) is close to that for SmAlO<sub>3</sub> (0.003) and is much smaller than those for ScAlO<sub>3</sub> (5.64), GdTlO<sub>3</sub> (6.00), and GdFeO<sub>3</sub> (4.08). Although the cell for CaGeO<sub>3</sub> is dimensionally close to being cubic, the polyhedra are considerably distorted as in characteristic orthorhombic perovskites. The most prominent distortion from the ideal cubic perovskite is in the tilting of the polyhedra (Fig. 2). The O(2)<sup>i</sup>-O(2)<sup>viii</sup>-O(2)<sup>i'</sup> angle, which shows the degree of the tilting in the  $xy$  plane, is 154.7(2)° in contrast to 180° for a cubic perovskite. The value of O(2)<sup>i'</sup>-O(2)<sup>ii'</sup>-O(2)<sup>vii</sup> angle, corresponding to the tilting on the  $xz$  plane, is 100.2(1)° as opposed to 90° for a cubic perovskite. The polyhedra themselves have the following characteristics: (1) The ratio of  $(\langle M-O \rangle - \langle M-O \rangle_{\text{cubic}}) / \langle M-O \rangle$  shows that the oxygens around the Ge atom are more extended from the ideal cubic positions (0.016) than those around Ca (0.006), where  $\langle M-O \rangle_{\text{cubic}}$  was estimated by constraining a cubic cell to have the same volume as the actual unit cell (see Table 2). (2) When we define the bond-length distortion,  $\Delta$ , as  $1/n \sum \{(r_i - r) / r\}^2 \times 10^3$ , it is clear that the Ca cuboctahedron ( $\Delta_A = 9.4$ ) is distorted, whereas the Ge octahedron is close to a regular one ( $\Delta_B = 0.005$ ). The SmAlO<sub>3</sub> perovskite has similar values of the bond-length distortion:  $\Delta_A = 8.7$  and  $\Delta_B = 0.001$ . The distortions, however, are less than for those of most GdFeO<sub>3</sub>-type perovskites such as MgSiO<sub>3</sub> ( $\Delta_A = 27.31$ ,  $\Delta_B = 0.226$ ), ScAlO<sub>3</sub> ( $\Delta_A = 37.50$ ,  $\Delta_B = 0.034$ ) and GdFeO<sub>3</sub> ( $\Delta_A = 28.13$ ,  $\Delta_B = 0.032$ ). (3) The O—Ca—O angles, which in the ideal cubic perovskite are all 90°, vary from 107.9(1)° to 53.79(8)° (see Table 3).

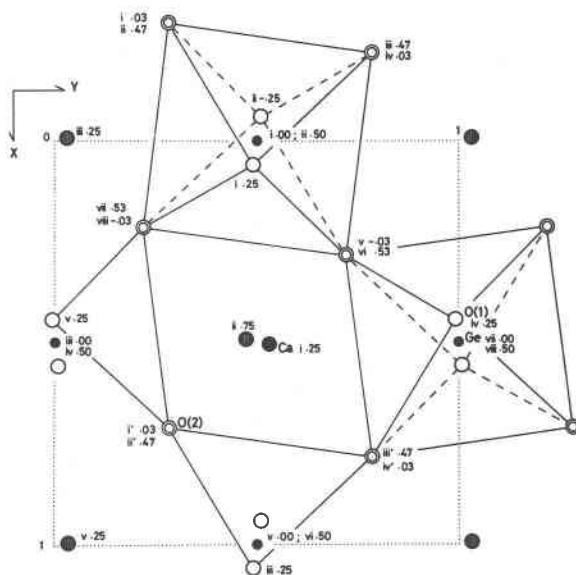


Fig. 2. Structure of CaGeO<sub>3</sub> perovskite projected on (001); Ca sites ruled, Ge sites solid, O(1) sites open, O(2) sites doubled. Each of the atoms is numbered with lower-case Roman numerals. Most atoms for which  $0.5 < z < 1.0$  have been omitted for clarity.

<sup>1</sup>To receive a copy of structure factor tables and Table 5, order document AM-83-233 from the Business Office, Mineralogical Society of America, 2000 Florida Ave. NW, Washington, D. C. 20009. Please remit \$1.00 in advance for microfiche.

Table 6. Comparison of tolerance factors and some structural characteristics such as B-O(1)-A and O(2)-O(2)-O(2) angles and bond-length distortions among orthorhombic GdFeO<sub>3</sub>-type perovskites, ABO<sub>3</sub>, based on single-crystal X-ray structure refinements.

	$\langle B-O \rangle$ (Å)	$\langle A-O \rangle$ (Å)	$t_{obs}$	$\Delta_B$	$\Delta_A$	B-O(1)-A 1-1-1 (°)	O(2)-O(2)-O(2) 1-viii-1' (°)	O(2)-O(2)-O(2) 1'-11'-vii (°)	$t_{IR}$	Reference
SeMgO <sub>3</sub>	2.150	2.808	0.924	0.374	58.49	112.41	119.52	114.88	0.883	1
ScAlO <sub>3</sub>	1.901	2.607	0.970	0.034	37.50	105.07( 5)	133.60( 4)	109.13( 4)	0.830	2
YAlO <sub>3</sub>	1.911	2.655	0.982	0.017	20.34	102.55(31)	145.89(68)	103.86(53)	0.884	3
SmAlO <sub>3</sub>	1.899	2.657	0.989	0.001	8.66	100.18(13)	155.41(21)	100.66(16)	0.906	4
MgSiO <sub>3</sub> *	1.790	2.476	0.978	0.226	27.31	103.8(15)	139.9(16)	107.7(15)	0.900	5
MgSiO <sub>3</sub> *	1.791	2.474	0.977	0.002	28.27	103.3(13)	138.1(16)	108.2(16)	0.900	6
CaTiO <sub>3</sub>	1.926	2.712	0.996	0.001	5.62	95.6	161.3	98.4	0.889	7
YTiO <sub>3</sub>	2.039	2.800	0.971	0.177	37.69	105.67(15)	134.65(18)	107.82(15)	0.826	8
LaTiO <sub>3</sub>	2.018	2.816	0.987	0.003	11.69	100.97(74)	151.62(57)	101.62(47)	0.874	8
NdTiO <sub>3</sub>	2.022	2.803	0.980	0.025	21.73	103.47(12)	143.72(14)	104.49(12)	0.857	8
SmTiO <sub>3</sub>	2.038	2.814	0.976	0.074	28.36	104.62(20)	139.65(24)	106.17(20)	0.847	8
GdTiO <sub>3</sub>	2.039	2.811	0.975	0.170	32.48	105.04(12)	137.52(14)	106.66(12)	0.838	8
InCrO <sub>3</sub>	1.987	2.708	0.964	0.096	38.40	105.18(16)	131.96(22)	110.39(19)	0.814	9
SeMnO <sub>3</sub>	2.235	2.894	0.916	0.424	64.51	112.76	117.16	115.86		1
LaMnO <sub>3</sub> #	2.018	2.818	0.987	3.75	16.44	101.15(16)	143.79(31)	101.89(26)	0.952	10
PrMnO <sub>3</sub> #	2.037	2.823	0.980	5.55	27.51	100.91(16)	133.45(29)	103.61(26)	0.871	11
YFeO <sub>3</sub>	2.015	2.771	0.972	0.081	34.26	104.67(29)	135.32(33)	108.09(31)	0.836	12
PrFeO <sub>3</sub>	2.008	2.791	0.983	0.008	17.67	102.25(11)	146.48(17)	103.75(13)	0.873	13
NdFeO <sub>3</sub>	2.011	2.790	0.981	0.007	20.24	102.89(11)	144.71(14)	104.53(12)	0.868	13
SmFeO <sub>3</sub>	2.013	2.785	0.978	0.040	24.83	103.48( 9)	141.46(11)	105.54( 9)	0.857	13
EuFeO <sub>3</sub>	2.013	2.783	0.978	0.038	26.85	103.64(11)	140.56(14)	105.82(12)	0.853	13
GdFeO <sub>3</sub>	2.012	2.781	0.977	0.032	28.13	103.87(16)	140.19(15)	105.83(13)	0.848	13
TbFeO <sub>3</sub>	2.012	2.775	0.975	0.047	30.17	104.03( 9)	138.42(11)	106.79( 9)	0.844	13
DyFeO <sub>3</sub>	2.012	2.773	0.975	0.059	32.02	104.15(12)	137.22(14)	107.10(12)	0.839	13
HoFeO <sub>3</sub>	2.010	2.768	0.974	0.054	33.36	104.36(13)	136.55(17)	107.44(14)	0.835	13
ErFeO <sub>3</sub>	2.011	2.765	0.972	0.038	34.85	104.83(12)	135.72(14)	107.83(12)	0.831	13
TmFeO <sub>3</sub>	2.011	2.763	0.972	0.029	35.87	104.67(13)	135.07(17)	108.26(15)	0.828	13
YbFeO <sub>3</sub>	2.011	2.759	0.970	0.029	37.61	104.69(17)	134.00(20)	108.61(16)	0.825	13
LuFeO <sub>3</sub>	2.010	2.755	0.969	0.030	38.62	105.01(15)	133.01(18)	109.30(15)	0.822	13
SeCoO <sub>3</sub>	2.148	2.805	0.923	0.688	58.30	112.39	121.00	114.15		1
TeCoO <sub>3</sub>	2.159	2.855	0.935	1.88	50.24	112.28	128.87	109.48		1
SeNiO <sub>3</sub>	2.113	2.768	0.926	0.285	55.72	111.74	121.59	113.64		1
SeCuO <sub>3</sub>	2.177	2.839	0.922	13.5	60.26	113.12	130.81	111.54		1
SeZnO <sub>3</sub>	2.160	2.815	0.922	0.775	59.14	111.99	120.63	114.51		1
CaGeO <sub>3</sub>	1.892	2.647	0.989	0.005	9.40	99.57(14)	154.72(18)	100.18(14)	0.923	14
SrZrO <sub>3</sub> #	2.091	2.919	0.987	0.006	10.65	101.18(31)?	152.92(35)	101.12(27)	0.887	15
BaCeO <sub>3</sub> #	2.241	3.128	0.987	0.044	10.57	101.13(62)	151.10(77)	102.97(76)	0.878	16
BaPrO <sub>3</sub> #	2.223	3.111	0.990	0.001	9.63	100.99(97)	154.6(12)	101.8(10)	0.886	16
NaTaO <sub>3</sub> #	1.978	2.765	0.988	0.001	8.23	99.51( 6)	155.81( 8)	99.43( 6)	0.894	17

$\langle B-O \rangle$ : Mean atomic distances for B site (VI coordinate).

$\langle A-O \rangle$ : Mean atomic distances for A site (XII coordinate).

$t_{obs}$ : Observed tolerance factor:  $t = \langle A-O \rangle / \sqrt{2} \langle B-O \rangle$ .

$t_{IR}$ : Tolerance factor calculated from the Shannon & Prewitt's ionic radii (VI and XII coordinates for B and A site, respectively).

$\Delta_B$ : Bond-length distortion for B-site (VI coordinate).

$\Delta_A$ : Bond-length distortion for A-site (XII coordinate).

\*: X-ray powder diffraction study.

#: Neutron powder diffraction study.

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The calculated powder diffraction pattern of CaGeO<sub>3</sub> perovskite for CuK $\alpha$  radiation ( $\lambda = 1.5418\text{\AA}$ ) is given in Table 5<sup>1</sup>. This pattern shows that a GdFeO<sub>3</sub>-type perovskite with a small unit cell distortion, such as that for CaGeO<sub>3</sub>, is difficult to distinguish from a cubic one using only X-ray powder data. Because all of the strong reflections in this pattern can be indexed on a pseudocubic cell, it is easy to understand why Ringwood and Major (1967) and Prewitt and Sleight (1969) reported the symmetry of CaGeO<sub>3</sub> to be cubic.

### Tilting of polyhedra in GdFeO<sub>3</sub>-type perovskites (Pbnm)

Glazer (1972; 1975) proposed a classification of perovskites based on tilting of octahedra and has suggested that the overall symmetry follows that of the tilts in spite of displacements and distortions. Although the tilting of the octahedra has been shown to have a large effect on lattice parameters, the description has been difficult because of different factors involved: for example, although SmAlO<sub>3</sub> and CaTiO<sub>3</sub> have similar tilting angles of octahedra, the ratios of cell constants are very different between these crystals:  $a/b = 5.291/5.290 (= 1.0)$  and  $c/\sqrt{2} = 5.285\text{\AA}$  for SmAlO<sub>3</sub>;  $a/b = 5.367/5.444$  and  $c/\sqrt{2} = 5.405\text{\AA}$  for CaTiO<sub>3</sub>. The following factor, newly defined here, makes it possible to describe the tiltings and/or distortions in a systematic way: we define the *observed tolerance factor*

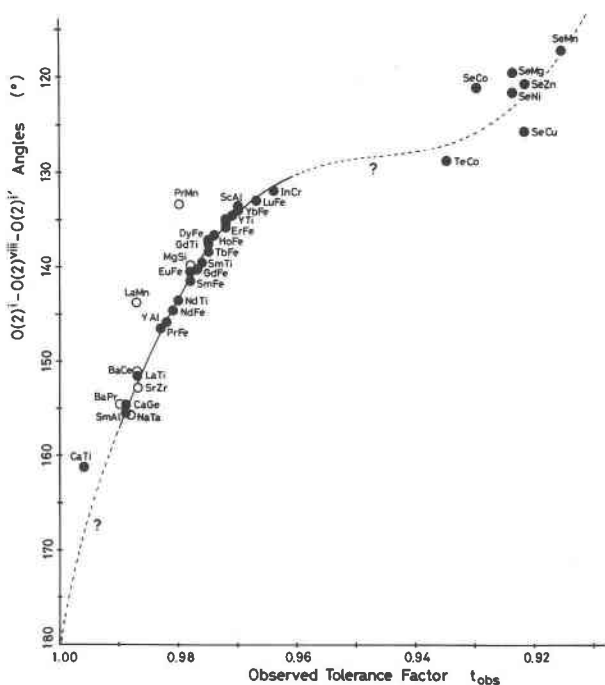


Fig. 3.  $O(2)^i-O(2)^{viii}-O(2)^{i'}$  angle vs.  $t_{obs}$  for the GdFeO<sub>3</sub>-type perovskites. Solid circles represent single-crystal X-ray results; open circles X-ray or neutron powder; shaded circle X-ray film with single crystal. The line is based on a polynomial equation (C Fit) shown in Table 7.

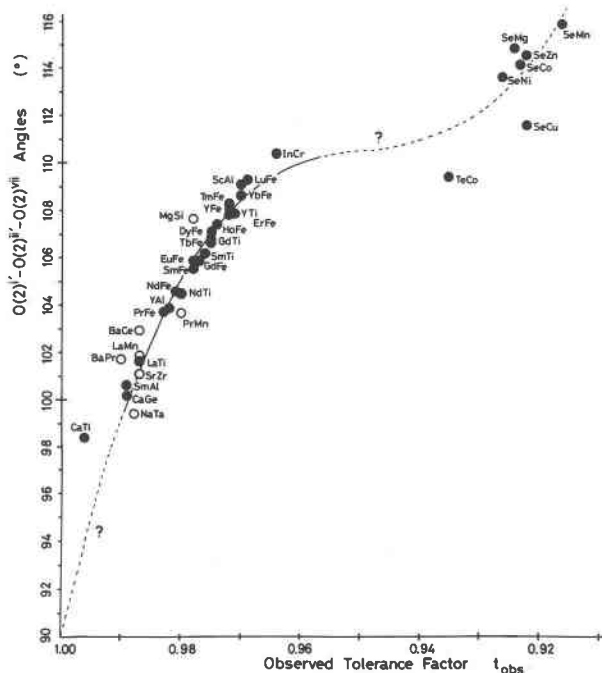


Fig. 4.  $O(2)^{i'}-O(2)^{ii'}-O(2)^{vii}$  angle vs.  $t_{obs}$  for the GdFeO<sub>3</sub>-type perovskites. Samples plotted are the same as those in Fig. 3. The line is based on a polynomial equation (C Fit) shown in Table 7.

based on the atomic parameters obtained from X-ray or neutron studies as follows:

$$t_{obs} = \langle A-O \rangle / \sqrt{2} \langle B-O \rangle,$$

where  $\langle A-O \rangle$  and  $\langle B-O \rangle$  are the mean interatomic distances with twelve and six coordination for A and B sites, respectively. The coordination numbers are fixed for all GdFeO<sub>3</sub>-type perovskites in order to maintain the same basis for comparison of structures with different coordinations. The  $t_{obs}$  values calculated from the atomic parameters are listed in Table 6.

The variations of angles of  $O(2A)^i-O(2)^{viii}-O(2)^{i'}$  [xy],  $O(2)^{i'}-O(2)^{ii'}-O(2)^{vii}$  [yz], and  $B^i-O(1)^i-A^i$  are shown in Figures 3, 4, and 5 as functions of  $t_{obs}$ . Some distinct features are evident among the GdFeO<sub>3</sub>-type perovskites: (1) The angles and bond length distortions concentrate clearly within two distinct groups  $t_{obs} = 0.96-0.99$  (group I) and  $0.93-0.92$  (group II, including selenides and tellurides) with a data gap in  $t_{obs}$  between the two groups. CaGeO<sub>3</sub> and SmAlO<sub>3</sub> are characterized as less-distorted perovskites ( $t_{obs} = 0.989$ ) and there is also a data gap between them and cubic perovskites ( $t_{obs} = 1.00$ ), except for CaTiO<sub>3</sub>. (2) There are systematic relationships between the angles showing the tilting of polyhedra and  $t_{obs}$ : (a) as  $t_{obs}$  decreases, the angle  $O(2)^i-O(2)^{viii}-O(2)^{i'}$  decreases (see Fig. 3); (b) as the  $t_{obs}$  decreases, the  $O(2)^{i'}-O(2)^{ii'}-O(2)^{vii}$  angle and the  $B^i-O(1)^i-A^i$  bond angle increase (see Fig. 4 and 5). (3) Table 7 shows that, in group I, the tilting of polyhedra versus  $t_{obs}$  can be interpolated

using the polynomial approximation of  $y = \sum a_j (1-t_{\text{obs}})^j \times 10^3$  ( $j = 0, 1, 2, \dots, n$ ; where  $n = 2$  or  $3$ ) (see "A Fit" in this table). (4) Also, these curves can be extended to fit over the full range as shown in "C Fit" of Table 7.

*The distortions of polyhedra in GdFeO<sub>3</sub>-type perovskites*

It is useful to plot distortions of polyhedra as well as tilts versus  $t_{\text{obs}}$ . Figures 6 and 7 show the correlation of bond-length distortions  $\Delta_A$  and  $\Delta_B$  versus  $t_{\text{obs}}$  for the A and B sites, respectively. Figure 6 shows that bond-length distortions for the larger A site have a systematic relationship with  $t_{\text{obs}}$ : (1) As the value of  $t_{\text{obs}}$  decreases, the distortions increase. (2) The trend can be approximated as a linear function within each of the groups, I and II. (3) The equations for distortions versus  $t_{\text{obs}}$  allow an interpolation within a group but does not permit an extrapolation. (4) The fit of the polynomial approximation, shown by the form of  $R_k = \sum \{|Y(\text{obs}) - Y(\text{calc})| / |Y(\text{obs})|\}_k$ , is not as good for distortions of polyhedra as it is for tilting (see Table 7). Although a comparison of bond-length distortion for the B site indicates similar characteristics to that of  $\Delta_A$ , there is no strong systematic relationship (Fig. 7).

*Application for germanates and silicates*

The systematic relationships observed in Figures 3 to 7 suggest that it might be possible to predict the deformation of perovskites containing various ions. For this purpose, we must be able to estimate the value of  $t_{\text{obs}}$  from known quantities such as ionic radii. In Figure 8, we

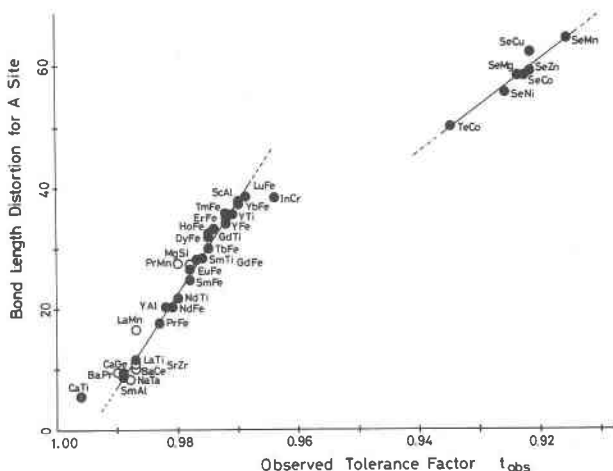


Fig. 6. Bond-length distortion for A site vs.  $t_{\text{obs}}$  for the GdFeO<sub>3</sub>-type perovskites. Samples plotted are the same as those in Fig. 3.

plot the Goldschmidt tolerance factor,  $t_{\text{IR}} = (r_A + r_O) / \sqrt{2(r_B + r_O)}$  versus  $t_{\text{obs}}$  for the perovskite crystals shown in Table 6, where  $r_A$ ,  $r_B$  and  $r_O (= 1.4 \text{ \AA})$  are the empirical ionic radii of the respective ions in the compound ABO<sub>3</sub> (Shannon and Prewitt, 1969; Shannon, 1976). Radii data for coordination numbers eight (A site) and six (B site) were used because values for twelve coordination for

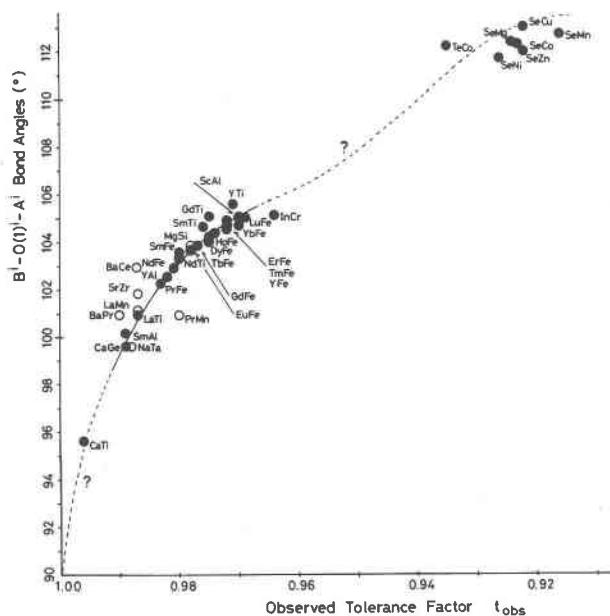


Fig. 5.  $B^i-O(1)^j-A^i$  angle vs.  $t_{\text{obs}}$  for the GdFeO<sub>3</sub>-type perovskite. Samples plotted are the same as those in Fig. 3. The line is based on a polynomial equation (C Fit) shown in Table 7.

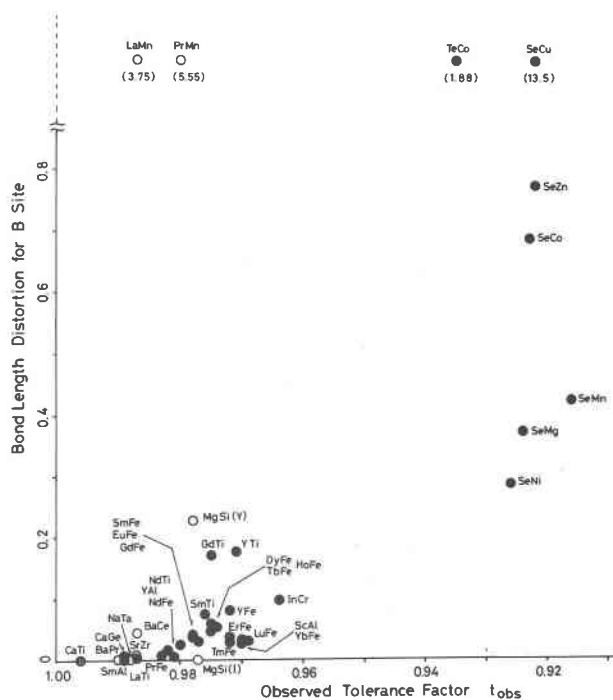


Fig. 7. Bond-length distortion for B site vs.  $t_{\text{obs}}$  for the GdFeO<sub>3</sub>-type perovskites. Samples plotted are the same as those in Fig. 3. For MgSiO<sub>3</sub>, Y and I represent values after Yagi *et al.* (1978) and Ito and Matsui (1978), respectively.

Table 7. The polynomial equations relating some bond angles and bond-length distortion versus  $t_{\text{obs}}$  using the mean squares method with least-squares calculations. The form is given by:  $Y = \sum a_j(1-t_{\text{obs}}) \times 10^3$  ( $j = 0, 1, 2, \dots, n$ , where  $n = 3, 4$ , or  $5$ ).

	$a_0$	$a_1$	$a_2$	$a_3$	$a_4$	$R_k$	Reference
<b>A Fit (Group I) (<math>0.96 &lt; t_{\text{obs}} \leq 0.99</math>)</b>							
$O(2)^i-O(2)^{viii}-O(2)^{i'}$	167.353	-0.88694	-0.031787	$0.80562 \times 10^{-3}$		0.009	Fig. 3
$O(2)^{i'}-O(2)^{ii'}-O(2)^{vii}$	94.153	0.62812	-0.004793			0.002	Fig. 4
$B^i-O(1)^i-A^i$	93.703	0.67376	-0.009870			0.002	Fig. 5
$\Delta_A$	15.340	-2.41560	0.206090	$-0.33625 \times 10^{-2}$		0.024	Fig. 6
<b>B Fit (Group I + Cubic) (<math>0.96 &lt; t_{\text{obs}} \leq 1.00</math>)</b>							
$O(2)^i-O(2)^{viii}-O(2)^{i'}$	180.001	-3.44519	0.149365	$-0.45824 \times 10^{-2}$	$0.57285 \times 10^{-4}$	0.008	Fig. 3
$O(2)^{i'}-O(2)^{ii'}-O(2)^{vii}$	89.998	1.50629	-0.069647	$0.20036 \times 10^{-2}$	$-0.22033 \times 10^{-4}$	0.002	Fig. 4
$B^i-O(1)^i-A^i$	90.019	1.20693	-0.033576	$0.33067 \times 10^{-3}$		0.002	Fig. 5
$\Delta_A$	-0.0015	0.68943	-0.013916	$0.31847 \times 10^{-2}$	$-0.69645 \times 10^{-4}$	0.023	Fig. 6
<b>C Fit (Groups I and II + Cubic) (<math>0.91 &lt; t_{\text{obs}} \leq 1.00</math>)</b>							
$O(2)^i-O(2)^{viii}-O(2)^{i'}$	179.956	-2.70904	0.048539	$-0.29879 \times 10^{-3}$		0.025	Fig. 3
$O(2)^{i'}-O(2)^{ii'}-O(2)^{vii}$	90.018	1.12058	-0.020962	$0.13474 \times 10^{-3}$		0.004	Fig. 4
$B^i-O(1)^i-A^i$	90.034	1.25746	-0.040800	$0.61616 \times 10^{-3}$	$-0.32303 \times 10^{-5}$	0.003	Fig. 5
$\Delta_A$	-0.2526	0.37869	0.066647	$-0.15607 \times 10^{-2}$	$0.98440 \times 10^{-5}$	0.036	Fig. 6

$t_{\text{obs}}$ : The observed tolerance factor

$R_k = \frac{\sum \{|Y(\text{obs}) - Y(\text{calc})| / |Y(\text{obs})|\}}{\sum \{|Y(\text{obs})|\}}$ ; For the calculation of  $O(2)^i-O(2)^{viii}-O(2)^{i'}$ ,  $Y = 180 - Y$ .

many ions are lacking. It is very significant that the GdFeO<sub>3</sub>-type perovskites with the same B ion can be characterized as a linear function in the  $t_{\text{IR}}-t_{\text{obs}}$  diagram, and that each function for the different B ions has almost the same slope. These characteristics make the estimation of  $t_{\text{obs}}$  possible for structures containing arbitrary kinds of ions.

Let us consider several germanates and silicates because of their geophysical importance. Average gradients of the  $t_{\text{obs}}-t_{\text{IR}}$  lines of titanates, aluminates, and orthoferites were assumed for germanates and silicates making use of the data for CaGeO<sub>3</sub> (this study) and MgSiO<sub>3</sub> (Yagi *et al.*, 1978). The estimated values of  $t_{\text{obs}}$ , which we shall call  $t'_{\text{obs}}$ , can be derived from the assumption that:

$$t'_{\text{obs}} = 0.274 t_{\text{IR}} + 0.736 \quad (1)$$

for germanates, and

$$t'_{\text{obs}} = 0.274 t_{\text{IR}} + 0.731 \quad (2)$$

for silicates (see Fig. 8).

SrGeO<sub>3</sub> is known to crystallize as a cubic perovskite (Shimizu *et al.*, 1970). The  $t'_{\text{obs}}$  value of 1.003 (SrGeO<sub>3</sub>) estimated from  $t_{\text{IR}} = 0.975$  is consistent with the observation that SrGeO<sub>3</sub> is a cubic perovskite. The value of  $t'_{\text{obs}} = 1.019$  for hexagonal BaGeO<sub>3</sub> (Shimizu *et al.*, 1970) may show that the orthorhombic GdFeO<sub>3</sub>-type perovskites can be separated crystallographically from hexagonal ones. On the other hand, for the silicate CaSiO<sub>3</sub>, which was reported to transform into a cubic perovskite at  $P = 160$  kbar and  $T = 25^\circ\text{C}$  (Liu and Ringwood, 1975), the value of  $t'_{\text{obs}} = 1.002$  ( $t_{\text{IR}} = 0.99$ ) is close to 1.0, in accord with the formation of cubic perovskite.

A cubic CdGeO<sub>3</sub> perovskite with  $a = 3.7\text{\AA}$  has been

reported to exist at 130 kbar and 900°C (Ringwood and Major, 1967). The value of  $t'_{\text{obs}} = 0.987$  ( $t_{\text{IR}} = 0.916$ ), determined from Equation (1), is close to those of CaGeO<sub>3</sub> and SmAlO<sub>3</sub>. This means that CdGeO<sub>3</sub> should crystallize in space group *Pbnm*, but with more distorted polyhedra than is found for CaGeO<sub>3</sub>. Because CaGeO<sub>3</sub> was previously reported incorrectly as cubic perovskite, a more detailed study of the symmetry for CdGeO<sub>3</sub> is required. If CdGeO<sub>3</sub> is orthorhombic, the  $t'_{\text{obs}}$  value predicts a crystal structure similar to that of LaTiO<sub>3</sub>:  $O(2)^i-O(2)^{viii}-O(2)^{i'} = 152^\circ$ ,  $O(2)^{i'}-O(2)^{ii'}-O(2)^{vii} = 102^\circ$ ,  $\text{Ge}^i-O(1)^i-\text{Cd}^i = 101^\circ$ , and  $\Delta_{A(=\text{Cd})} = 11$ . MnGeO<sub>3</sub> quenched from 250 kbar and 1400–1800°C has been found to be orthorhombic (Liu, 1976b). The value of  $t'_{\text{obs}} = 0.973$  ( $t_{\text{IR}} = 0.865$ ) supports this report and predicts the space group *Pbnm* (or *Pbn2<sub>1</sub>*). Liu (1977) reported MgGeO<sub>3</sub> as an orthorhombic perovskite, having space group *Pmmm* or *P222*. Examination of the  $t_{\text{obs}}-t_{\text{IR}}$  diagram shows that MgGeO<sub>3</sub> is within the stability of the GdFeO<sub>3</sub>-type perovskite ( $t'_{\text{obs}} = 0.966$ ;  $t_{\text{IR}} = 0.839$ ) and has similar characteristics to InCrO<sub>3</sub>. However, because the report for MgGeO<sub>3</sub> was based on powder data of five mixed phases, it is also recommended that the space group be reexamined.

It should be emphasized that the data are consistent with MgSiO<sub>3</sub> (Yagi *et al.*, 1978; Ito and Matsui, 1978) being a GdFeO<sub>3</sub>-type perovskite (see Figs. 3 to 7). In contrast, Madon, Bell, Mao, and Poirier (1980) assigned tetragonal symmetry to MgSiO<sub>3</sub> ( $a = b = 9.3$ ,  $c = 13.1\text{\AA}$ ), based on lattice parameters measured using electron diffraction patterns. Their observations, however, are not sufficient to confirm the symmetry and their conclusion may result from other factors such as twinning or pseudo-symmetry. It should be noted that an experiment using



only electron microscopy is not sufficient to determine the symmetry.

We may use the observed relationship to predict values of  $t'_{\text{obs}}$  for perovskites of some germanate and silicate compounds which have not yet been synthesized. For this purpose, the values of  $t'_{\text{obs}}$  associated with various cations were calculated for germanates and silicates from  $t_{\text{IR}}$  and are shown in Table 8. The compounds for which the  $t'_{\text{obs}}$  lies within the range for group I would probably crystallize as orthorhombic perovskites and have the tilting and distortion of oxygen polyhedra as estimated from Table 7: these are  $\text{FeGeO}_3$  ( $t'_{\text{obs}} = 0.969$ ),  $\text{CoGeO}_3$  (0.967), and  $\text{ZnGeO}_3$  (0.967) for germanates, and  $\text{MnSiO}_3$  ( $t'_{\text{obs}} = 0.985$ ),  $\text{FeSiO}_3$  (0.981),  $\text{CoSiO}_3$  (0.979), and  $\text{ZnSiO}_3$  (0.978) for silicates. On the other hand,  $\text{BaSiO}_3$  ( $t'_{\text{obs}} = 1.035$ ) and  $\text{SrSiO}_3$  (1.017) should not occur as the  $\text{GdFeO}_3$ -type perovskite. However, the above discussion illustrates only one geometrical aspect concerning crys-

Table 8. Estimated values of  $t_{\text{obs}}$  for some germanates and silicates.

	$t_{\text{IR}}$	$t'_{\text{obs}}$		$t_{\text{IR}}$	$t'_{\text{obs}}$
Germanates			Silicates		
$\text{MgGeO}_3$	0.839	0.966	$\text{CaSiO}_3$	0.990	1.002
$\text{MnGeO}_3$	0.865	0.973	$\text{MnSiO}_3$	0.927	0.985
$\text{FeGeO}_3$	0.850	0.969	$\text{FeSiO}_3$	0.911	0.981
$\text{CoGeO}_3$	0.843	0.967	$\text{CoSiO}_3$	0.904	0.979
$\text{ZnGeO}_3$	0.843	0.967	$\text{ZnSiO}_3$	0.904	0.979
$\text{SrGeO}_3$	0.975	1.003	$\text{SrSiO}_3$	1.045	1.017
$\text{CdGeO}_3$	0.916	0.987	$\text{CdSiO}_3$	0.982	0.986
$\text{BaGeO}_3$	1.033	1.019	$\text{BaSiO}_3$	1.108	1.035

tallization of orthorhombic perovskites; often other factors such as crystal-field effects determine what phases are stable. For example, although  $\text{FeSiO}_3$  is one of the possible orthorhombic perovskites (see Table 8), it is thought to be unstable with respect to mixed oxides (Yagi et al., 1978).

### Summary

$\text{CaGeO}_3$  perovskite is orthorhombic ( $Pbnm$ ) and isostructural with  $\text{GdFeO}_3$ . The oxygen polyhedra are less tilted and less distorted than those of most  $\text{GdFeO}_3$ -type perovskites. Comparison between structural characteristics and  $t_{\text{obs}}$  has made possible useful predictions for the existence of previously-unknown  $\text{GdFeO}_3$ -type perovskites. The following predictions are obtained for germanate and silicates: (1)  $\text{CdGeO}_3$  should crystallize as an orthorhombic perovskite rather than a cubic one; (2) the perovskite form of  $\text{MgSiO}_3$  probably has the  $\text{GdFeO}_3$ -type structure; (3) the space group,  $Pmmm$  (or  $P222$ ), reported for a  $\text{MgGeO}_3$  perovskite should be reexamined with the possibility of it being  $Pbnm$ ; (4) in the geometrical view of crystals, germanate and silicate perovskites containing Mg, Mn, Co, Fe, and Zn ions in the large A site should have the  $\text{GdFeO}_3$ -type structure.

### Acknowledgments

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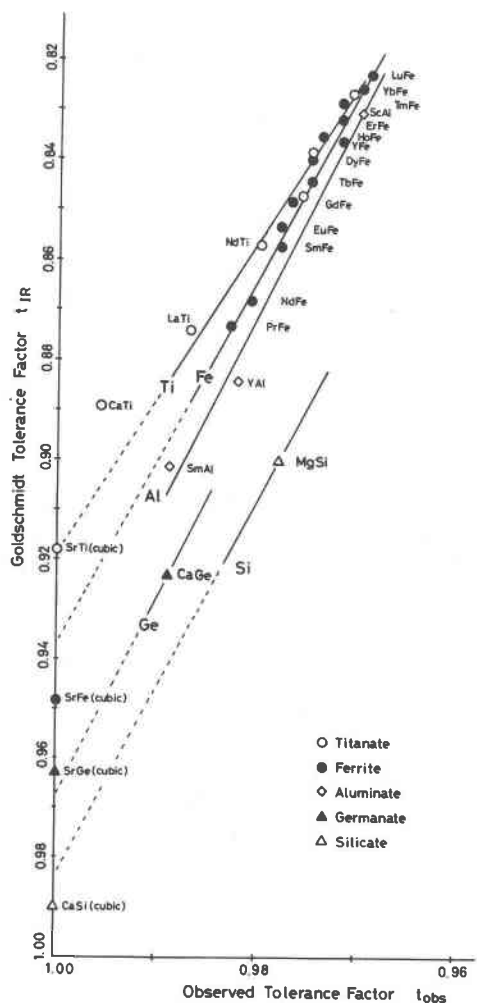


Fig. 8. The  $t_{\text{obs}}-t_{\text{IR}}$  diagram of some  $\text{GdFeO}_3$ -type perovskites such as orthoferrites, titanates, aluminates, germanate, and silicates.

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