

## Orientation and effects of channel H<sub>2</sub>O and CO<sub>2</sub> in cordierite

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### Abstract

Channel CO<sub>2</sub> and H<sub>2</sub>O has been re-introduced into a channel evacuated Mg-cordierite



from White Well, Australia. At 600°C and pressures up to 6 kbar, a maximum of 3.13 wt.% CO<sub>2</sub> and 2.6 wt.% H<sub>2</sub>O re-enter cordierite's channels. The H<sub>2</sub>O molecule orients preferentially with its H–H vector parallel to *c* (optic orientation: *c* = *X*, *b* = *Y*, *a* = *Z*). With increased H<sub>2</sub>O content, the refractive indices  $\gamma$  and  $\beta$  increase more strongly than  $\alpha$  whereas the *a*<sub>0</sub> cell edge decreases, *b*<sub>0</sub> perhaps increases slightly, and *c*<sub>0</sub> first increases and then levels off. The linear CO<sub>2</sub> molecule orients chiefly parallel to cordierite's *a*-axis so that, as CO<sub>2</sub> content increased, refractive index  $\gamma$  increased more strongly than did  $\alpha$  and  $\beta$ . As a result,  $2V_x$  increased with CO<sub>2</sub> content so as to exceed 90°. Simultaneously, *c*<sub>0</sub> increased, *a*<sub>0</sub> decreased, and *b*<sub>0</sub> remained constant. The distortion index  $\Delta$  for these crystals changed with H<sub>2</sub>O and CO<sub>2</sub> content.

### Introduction

Cordierites,  $Na_{x+2y}(Mg,Fe,Mn)_{2-y}^{VI}(Al_{4-x}Be_xSi_5)^{IV}O_{18} \cdot n[H_2O,CO_2]$ , consist structurally of six-membered rings of corner-sharing tetrahedra (T<sub>2</sub> in Fig. 1) stacked along the *c*-axis. These are linked (laterally and vertically) by corner-sharing with other tetrahedra (T<sub>1</sub>) to form a framework structure (Gibbs, 1966). The ring-stacking produces channels along the *c*-axis that pinch to 'bottlenecks' (~2.5 Å in diameter) or swell to large 'cages' whose maximum dimensions, which occur in the plane parallel to (001), are approximately 5.4 Å along *b* and 6.0 Å along *a*. In low cordierite, Al concentrates into the two equivalent T<sub>26</sub> tetrahedra in the Al<sub>2</sub>Si<sub>4</sub>O<sub>18</sub> ring so that, instead of possessing a 6-fold axis, as for the Si<sub>6</sub>O<sub>18</sub> ring in beryl, the Al<sub>2</sub>Si<sub>4</sub>O<sub>18</sub> ring possesses a 2-fold axis (like that drawn perpendicular to the lightly stippled plane in Fig. 1). This leads to the orthorhombic symmetry (*Cccm*, *Z* = 4, *a* > *b* > *c*) thus far exhibited by cordierites from metamorphic and igneous rocks (Selkregg and Bloss, 1980; Wallace and Wenk, 1980). High-cordierite or indialite, found where shale and sandstone fused near a

burning coal seam (Miyashiro and Iiyama, 1954), appears to be hexagonal, the rapid cooling apparently preventing Al from concentrating in T<sub>26</sub>. Such also seems the case for cordierites from meteorites or those synthesized by recrystallization from a glass (Schreyer and Schairer, 1961), by growth in a flux (Lee and Pentecost, 1976), and by hydrothermal methods (Schreyer and Yoder, 1964). Additional heating of such synthetics leads to Si/Al ordering accompanied by twinning along (110) and (310) (Putnis, 1980; Armbruster and Bloss, 1981).

In natural cordierites the structural channels accommodate numerous species of which CO<sub>2</sub> and H<sub>2</sub>O can be recognized by IR spectroscopy, and He and Ar by mass spectrometry (Damon and Kulp, 1958; Beltrame *et al.*, 1976). Hydrocarbons and other gases may also be present in small concentrations not revealed in IR spectra. Na<sup>+</sup> is an important channel occupant that centers in the bottlenecks (Meagher, 1967) and provides charge balance if Be<sup>2+</sup> substitutes for Al<sup>3+</sup> (Povondra and Langer, 1971). Small amounts of Fe in the channels may cause the pleochroism of some cordierites (Goldman *et al.*, 1977), but firm data in support of this are not yet at hand.

The degree of variation of the cordierite lattice from hexagonal symmetry is customarily measured

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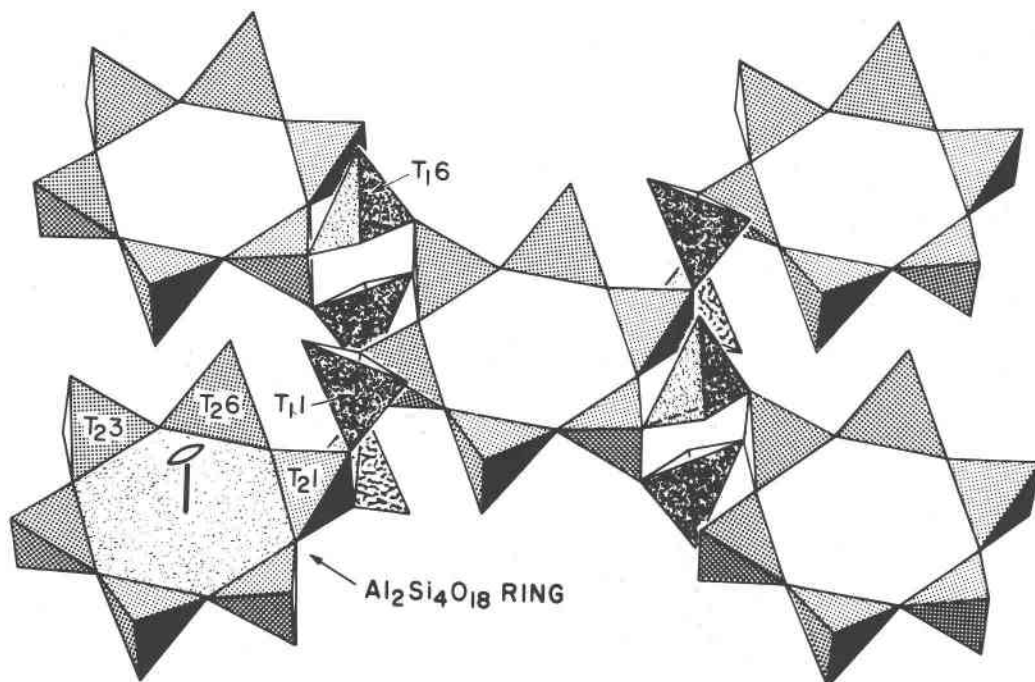


Fig. 1. Part of the cordierite framework (after Gibbs, 1966). The rings consisting of six (unstippled) tetrahedra are linked to each other, both laterally and vertically, by sharing corners with the darkly stippled tetrahedra. At the lower left, the *c*-axis (9.3 Å repeat) is shown centered in one of the Al<sub>2</sub>Si<sub>4</sub>O<sub>18</sub> rings. For orthorhombic cordierites, tetrahedral sites T<sub>2.1</sub>, T<sub>2.6</sub> and T<sub>2.3</sub> (symbolism of Meagher and Gibbs, 1977) are non-equivalent because Al concentrates in T<sub>2.6</sub>. Also, T<sub>1.1</sub> and T<sub>1.6</sub> are non-equivalent because Al concentrates in T<sub>1.1</sub>. In consequence the *c*-axis is 2-fold.

by the distortion index  $\Delta$ , which Miyashiro (1957) defined (for CuK $\alpha$  radiation) as

$$\Delta = 2\theta_{131} - (2\theta_{511} + 2\theta_{421})/2.$$

The assumption was then made that  $\Delta$  provided a reliable measure of Al/Si ordering for natural cordierites. However, systematic optical and X-ray studies of natural and heated cordierites (Selkregg and Bloss, 1980) has shown this to be unlikely for natural cordierites, excluding indialites. Instead, the  $\Delta$  index depended on Na-content, Fe- and Mn-substitution for Mg, and on water content. Although degrees of Si/Al ordering would without doubt influence  $\Delta$ , cordierites from metamorphic and igneous rocks appear to have cooled so slowly that, thus far, all crystal structure analyses (Gibbs, 1966; Cohen, *et al.*, 1977; Hochella *et al.*, 1979; Wallace and Wenk, 1980) have disclosed nearly perfect Si/Al ordering.

In studies of natural cordierites, the tendency has been to overlook all channel occupants except H<sub>2</sub>O. However, CO<sub>2</sub> also seems important (Suknev *et al.*, 1971; Schreyer *et al.*, 1979; Hörmann *et al.*, 1980) and leads to specific distortions of the cordierite

lattice and modifications of its optical properties (Armbruster and Bloss, 1980). The present paper examines the effects of the orientation of H<sub>2</sub>O and CO<sub>2</sub> on the crystal lattice and optical properties of a cordierite with a well ordered Si/Al distribution.

### Experimental

A Mg-cordierite from White Well, Australia (Pryce, 1973) with well ordered Si/Al distribution (Cohen *et al.*, 1977; Hochella *et al.*, 1979) was crushed to an average grain size of about 200  $\mu$ m. Transparent, inclusion-free grains were hand picked under a binocular microscope and tempered for one day between 1200°C and 1300°C in a SiC-furnace in order to expel the volatile channel occupants. Electron microprobe and coulombmetrical analyses for H<sub>2</sub>O and CO<sub>2</sub> established this cordierite's composition to be



prior to heating. After heating, through loss of the volatile channel occupants (and perhaps some Na),

it became



Samples (~50 mg) of this tempered cordierite were sealed in platinum tubes and held for 4 weeks at 600°C under H<sub>2</sub>O pressures of 0.5, 2, 4, 5 and 6 kbar. Others were similarly treated but in a CO<sub>2</sub> atmosphere (from decomposition of Ag<sub>2</sub>C<sub>2</sub>O<sub>4</sub>). Annealing times of 4 weeks were needed to achieve optical homogeneity within the crystal grains of the sample. The samples were subsequently checked by IR spectroscopy between 5000 cm<sup>-1</sup> and 400 cm<sup>-1</sup> using the KBr-powder technique. H<sub>2</sub>O and CO<sub>2</sub> were thus easily detected. H<sub>2</sub>O displays symmetric and asymmetric stretching modes between 3570 cm<sup>-1</sup> and 3700 cm<sup>-1</sup> (Farrell and Newnham, 1967; Goldman *et al.*, 1977). Gaseous CO<sub>2</sub> possesses an asymmetric stretching vibration at 2349 cm<sup>-1</sup>. Hence, absorption at or near 2354 cm<sup>-1</sup> for beryl (Wood and Nassau, 1967) and for cordierite (Farrell and Newnham, 1967) was attributed by them to CO<sub>2</sub> trapped in the channels. For the optical studies, single grains from each run were mounted on a goniometer head and oriented on the spindle stage. Refractive indices were measured by the double variation ( $\lambda$ , T) method and  $2V$  was calculated from extinction data at 400, 666 and 900 nm applying the computer program EXCALIBUR. A precise description of the optical methods used in this paper is given elsewhere (Selkregg and Bloss, 1980; Bloss, 1981). The precision of the refractive indices is within 0.0005 and  $2V$  within 0.5°. For the same crystal grains, cell dimensions were measured by the back-reflection Weissenberg method. For at least 60 indexed reflections (CuK $\alpha_1$ , CuK $\alpha_2$  and CuK $\beta$ ),  $2\theta$  values were submitted to the least square program of Burnham (1962, 1965), as revised by L. Finger, which corrects for film shrinkage and absorption. Water in the samples was quantitatively determined by the equation of Medenbach *et al.* (1980) which relates change in average refractive index (upon heating to complete dehydration) to water content for Mg-cordierite. CO<sub>2</sub> released from the sample at 1300° C was determined using the same equipment and method as described in detail by Johannes and Schreyer (1980). These CO<sub>2</sub> and H<sub>2</sub>O analyses lead to standard deviations below 5% of the total gas analyzed. The orientation of CO<sub>2</sub> in natural cordierites was determined by polarized IR spectroscopy from oriented single crystal slabs kindly provided by D. S. Goldman (and described by Goldman *et al.*, 1977). For these single crystal

slabs, we compared the relative intensity of the asymmetric stretching absorption of CO<sub>2</sub> at (*ca.*) 2350 cm<sup>-1</sup> along the *a*, the *b*, and the *c* axis.

### Results

The White Well cordierite, degassed by heating between 1200° and 1300°C in a slightly reducing atmosphere, developed oriented cracks but no color change or accompanying hematite reflections in its X-ray photographs. The cracks subsequently vanished for the H<sub>2</sub>O-treated crystals, probably because H<sub>2</sub>O enhances diffusion processes. By contrast, the CO<sub>2</sub> treatment caused only partial healing of the cracks. IR spectra confirmed H<sub>2</sub>O as the only volatile channel occupant in the H<sub>2</sub>O-treated cordierites and CO<sub>2</sub> as the only one in the CO<sub>2</sub>-treated ones. The H<sub>2</sub>O-containing cordierites show a strong absorption peak at 3690 cm<sup>-1</sup> and weaker ones at 3575 cm<sup>-1</sup> and 3630 cm<sup>-1</sup>. The CO<sub>2</sub>-containing cordierites exhibit strong peaks at about 2350 cm<sup>-1</sup>. There is no significant difference in the frequencies of the absorption bands for channel CO<sub>2</sub> and H<sub>2</sub>O observed by us from those observed by Farrell and Newnham (1967) or Goldman *et al.* (1977). Absorption bands attributed to AlO<sub>4</sub>, SiO<sub>4</sub>, and MgO<sub>6</sub> vibrations are consonant with the spectra observed by Langer and Schreyer (1969). Non-equilibrium conditions during hydration or carbonation were readily detected because the resultant crystals possessed higher refractive indices at their edges than at their cores. Moreover, such non-homogeneous crystals exhibited undulatory extinction because of the strong influence of H<sub>2</sub>O and CO<sub>2</sub> on the optic angle (Armbruster and Bloss, 1980). By contrast, crystals annealed at 600°C for 4 weeks displayed sharp extinction.

With increased gas in the channels, each refractive index increased significantly (Fig. 2, Table 1). As channel H<sub>2</sub>O increased,  $\beta$  increased more rapidly than  $\gamma$  and much more so than  $\alpha$ . Hence,  $2V_x$  decreases sharply. As channel CO<sub>2</sub> increased,  $\gamma$  increased much more rapidly than  $\beta$  and  $\alpha$ . Hence,  $(\gamma - \alpha)$  and  $2V_x$  increased strongly with channel CO<sub>2</sub> as first noted by Armbruster and Bloss (1980). CO<sub>2</sub> and H<sub>2</sub>O each affect the cell dimensions (Fig. 3, Table 2). Channel H<sub>2</sub>O decreases  $a_0$  but it slightly increases  $b_0$  and  $c_0$  so that cell volume remains constant (within the range of error). As a result, the distortion index  $\Delta$ , calculated by the simplified equation of Selkregg and Bloss (1980), namely

$$\Delta = 1.094 (a_0 - b_0 \sqrt{3})$$

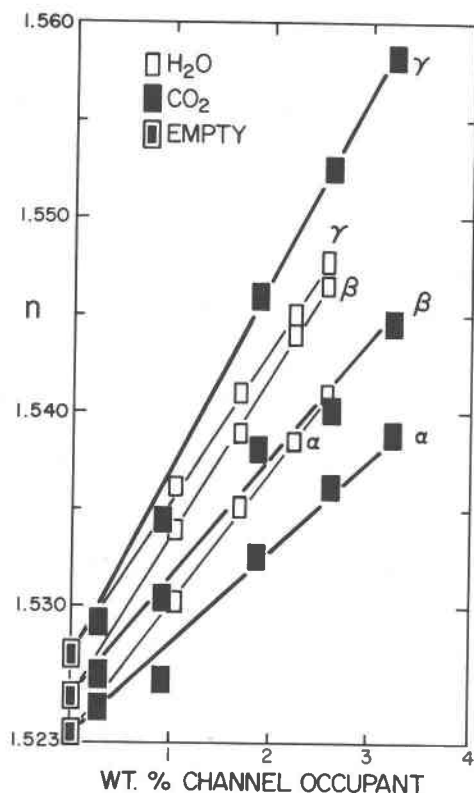


Fig. 2. The increase in refractive index of Mg-cordierite that results if H<sub>2</sub>O occupies the channels (light lines, open rectangles) or if CO<sub>2</sub> occupies the channels (bold lines, solid rectangles). The edges of the rectangles correspond to two estimated standard deviations. The data for the channel evacuated cordierite are represented by concentric open and solid rectangles.

decreased as H<sub>2</sub>O content increased. With increased channel CO<sub>2</sub>, *a*<sub>0</sub> first decreases but, beyond 2 weight percent CO<sub>2</sub>, it increases; *b*<sub>0</sub> seemed unaffected by CO<sub>2</sub> but *c*<sub>0</sub> increased strongly. Consequently, Δ decreased (up to 2 weight percent CO<sub>2</sub>), then increased slightly.

IR spectra for slabs cut from a crystal of natural cordierite confirmed the conclusions of Farmer (1974) that the linear CO<sub>2</sub> molecule aligns mainly in the (001) plane and, within this plane, chiefly along *a*.

Hartshorne and Stuart (1970, p. 137) state the Lorentz-Lorenz equation for molecular refractivity *R*<sub>M</sub> to be

$$R_M = \frac{n^2 - 1}{n^2 + 2} \cdot \frac{M}{\rho} = \frac{4}{3} \pi N P$$

where *n*, *M*, *ρ*, *N* and *P* respectively represent refractive index, molecular weight, density, the

Avogadro number, and the electron polarizability for a molecular compound, where *P* is the polarizability averaged over all possible orientations of the electric vector. Although the equation is not strictly applicable to anisotropic crystals, Hartshorne and Stuart note that a principal refractive index (*ε*, *ω*, *α*, *β*, or *γ*) may be substituted for *n*. In such case *P* reasonably approximates the crystal's polarizability along the principal vibration direction associated with this principal index and may thus be symbolized *P*<sub>ε</sub>, *P*<sub>ω</sub>, *P*<sub>α</sub>, *P*<sub>β</sub> or *P*<sub>γ</sub>, as the case may be. For example,

$$P_\epsilon = \frac{\epsilon^2 - 1}{\epsilon^2 + 2} \cdot \frac{M}{\rho} \cdot \frac{3}{4\pi N}$$

Actually, *M*/*ρ* represents the molar volume *V*<sub>M</sub>. For a crystal with known unit-cell volume (*V*<sub>uc</sub>), if *Z* represents the formula units encompassed by *V*<sub>uc</sub>, then its molar volume can be calculated since

$$V_M = \frac{N V_{uc}}{Z}$$

Hence the preceding equation can be rewritten as, for example,

$$P_\epsilon = \frac{3}{4\pi} \cdot \frac{\epsilon^2 - 1}{\epsilon^2 + 2} \cdot \frac{V_{uc}}{Z}$$

For mineralogists this is convenient since *V*<sub>uc</sub> and *Z* are more likely to be accurately known than density or molecular weight. The units for *P* will be those used for *V*<sub>uc</sub> (usually Å<sup>3</sup>).

For the cordierites here studied, polarizabilities *P*<sub>α</sub>, *P*<sub>β</sub>, and *P*<sub>γ</sub> were calculated using their principal

Table 1. Optical data of H<sub>2</sub>O and CO<sub>2</sub> treated White Well cordierite at 600°C.

Pressure (kbar)	Gas (wt %)	α	β	γ	$\bar{n}$	(γ - α)	2 <i>V</i> <sub>x</sub> *
Channel-evacuated							
--	--	1.5235	1.5254	1.5275	1.52546	0.0040	87.0
H <sub>2</sub> O re-introduced into channels							
0.5	1.05	1.5303	1.5340	1.5362	1.53350	0.0059	64.6
2.0	1.68	1.5351	1.5389	1.5410	1.53833	0.0059	56.6
4.0	2.24	1.5386	1.5440	1.5450	1.54253	0.0064	50.0
6.0	2.56	1.5409	1.5465	1.5477	1.54503	0.0068	44.1
CO <sub>2</sub> re-introduced into channels							
0.5	0.27	1.5238	1.5260	1.5287	1.52616	0.0049	86.8
2.0	0.91	1.5263	1.5305	1.5345	1.53043	0.0082	93.8
4.0	1.88	1.5326	1.5381	1.5459	1.53887	0.0133	105.0
5.0	2.61	1.5362	1.5410	1.5525	1.54323	0.0163	111.1
6.0	3.13	1.5388	1.5446	1.5582	1.54720	0.0187	112.4

\* Average of measurements at 400, 666, and 900 nm

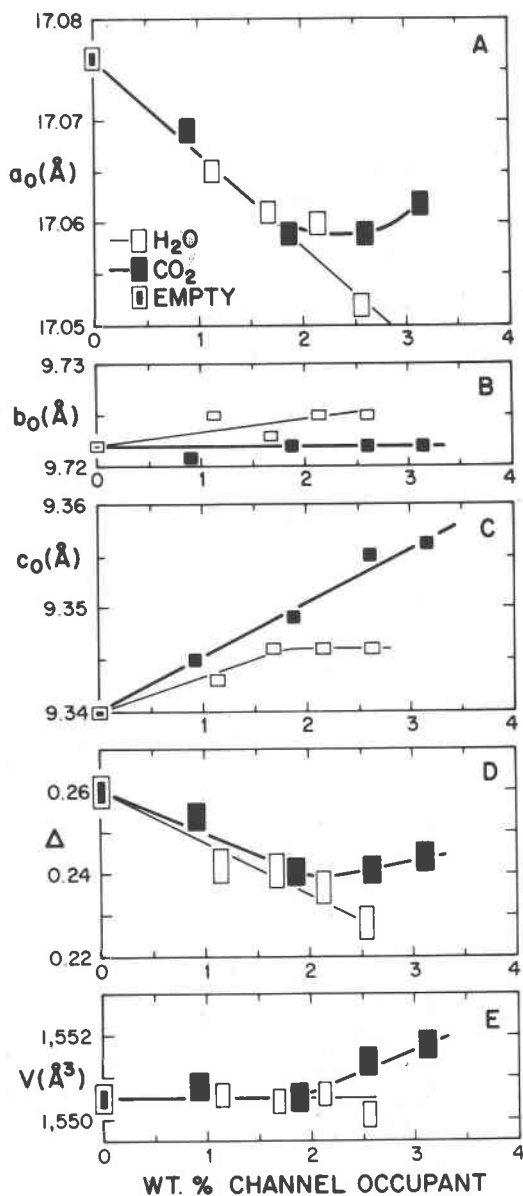


Fig. 3. Variation in unit cell edges (A, B and C), distortion index  $\Delta$  (D), and unit cell volume (E) as amount of channel occupant increases. The rectangles have the same significance as in Figure 2.

indices for sodium light ( $\alpha_D$ ,  $\beta_D$ , and  $\gamma_D$ ),  $V_{uc}$ , and  $Z$  (Table 3). These polarizabilities increase systematically with increased weight per cent of channel H<sub>2</sub>O or CO<sub>2</sub> (Fig. 4). The trends indicate that as channel H<sub>2</sub>O increases, cordierite's polarizabilities (and associated refractive indices) increase at the greatest rate for light vibrating parallel to  $Y$ , at a lesser rate for  $Z$ , and at the least rate for  $X$ . With increase in channel CO<sub>2</sub>, the polarizabilities (and associated

indices) increase most rapidly for  $Z$  ( $= a$  axis), least for  $X$  ( $= c$  axis) and, for  $Y$  ( $= b$  axis), at a slightly greater rate than that for  $X$ . The marked increase in polarizability along  $Z$ , as CO<sub>2</sub> increases, underscores (1) the ready polarizability of the linear CO<sub>2</sub> molecule along its length and, as IR spectra indicate, (2) the tendency of this molecule to align chiefly parallel (or subparallel) to the  $a$  axis in cordierite. This causes  $(\gamma-\alpha)$  and  $2V_x$  to increase with channel CO<sub>2</sub> content.

The linearity of the trends in Figure 4 reinforce confidence in the measurement of the refractive indices and in the determinations of H<sub>2</sub>O and CO<sub>2</sub>. However, specimens with a relatively low content of H<sub>2</sub>O and CO<sub>2</sub> seem at variance with these trends. Such variance may result (1) from experimental errors in determining CO<sub>2</sub> or H<sub>2</sub>O, or (2) from a possible tendency of channel CO<sub>2</sub> or H<sub>2</sub>O molecules to orient somewhat differently at low concentrations as compared to high.

### Discussion

From IR spectra of 8 natural cordierites, Goldman *et al.* (1977) confirmed two major types of orientation for the H-H vector for channel water—H-H parallel  $c$  (type I) or H-H parallel  $b$  (type II)—and noted that absence of channel cations, especially Na, favors type I. Tsang and Ghose (1972) investigated a Mg-rich cordierite from Madagascar by NMR technique and found only type I water. The relatively high distortion index ( $\Delta = 0.24$ ) calculated from the cell dimensions, suggests this specimen to be Na-poor (Selkregg and Bloss, 1980).

Table 2. Cell dimensions, volume and distortion index of H<sub>2</sub>O and CO<sub>2</sub> treated White Well cordierite at 600°C\*

Pressure (kbar)	Gas (wt %)	$a_0$ [Å]	$b_0$ [Å]	$c_0$ [Å]	$V$ [Å <sup>3</sup> ]	$\Delta$
Channel-evacuated						
--	--	17.076 <sup>(2)</sup>	9.722 <sup>(1)</sup>	9.340 <sup>(1)</sup>	1550.5 <sup>(3)</sup>	0.260 <sup>(3)</sup>
H <sub>2</sub> O re-introduced into channels						
0.5	1.05	17.065	9.725	9.343	1550.6	0.242
2.0	1.68	17.061	9.723	9.346	1550.4	0.241
4.0	2.24	17.060	9.725	9.346	1550.6	0.237
6.0	2.56	17.052	9.725	9.346	1550.1	0.228
CO <sub>2</sub> re-introduced into channels						
0.5	0.27	--	--	--	--	--
2.0	0.91	17.069	9.721	9.345	1550.8	0.254
4.0	1.88	17.059	9.722	9.349	1550.5	0.241
5.0	2.61	17.059	9.722	9.355	1551.1	0.241
6.0	3.13	17.062	9.722	9.356	1551.8	0.244

\* Numbers in parentheses in the table represent the estimated standard deviation in terms of the least decimal place cited for the value to the left and for all other values in the column below. The precisions achieved in determining the gas contents are discussed in the text.

Table 3. Polarizabilities for the H<sub>2</sub>O and CO<sub>2</sub> treated White Well cordierite

Gas (wt %)	P <sub>α</sub> (Å <sup>3</sup> )	P <sub>β</sub> (Å <sup>3</sup> )	P <sub>γ</sub> (Å <sup>3</sup> )	P <sub>̄</sub> (Å <sup>3</sup> )
Channel-evacuated				
--	28.291	28.377	28.472	28.380
H <sub>2</sub> O re-introduced into channels				
1.05	28.601	28.767	28.866	28.745
1.68	28.813	28.983	29.077	28.958
2.24	28.973	29.215	29.259	29.149
2.56	29.067	29.317	29.370	29.251
CO <sub>2</sub> re-introduced into channels				
0.27	28.310	28.410	28.532	28.420
0.91	28.424	29.613	29.793	28.610
1.88	28.702	28.949	29.298	28.984
2.61	28.875	29.090	29.602	29.190
3.13	29.005	29.264	29.868	29.380

The increase in refractive indices as channel H<sub>2</sub>O is reintroduced into channel-evacuated White Well cordierite (Table 3) is consonant with type I water. For this orientation, vectors drawn from the hydrogens to the nearest oxygens of the Al<sub>2</sub>Si<sub>4</sub>O<sub>18</sub> rings would represent directions of maximum (but weak) polarization of these oxygens by the hydrogens. Such directions, being more nearly parallel than perpendicular to (001), would hence favor, as Table 3 shows, a greater rate of increase of the indices β and γ than of α with increased channel water (since

$X = c$ ). The long distance (3.4 Å) between the water oxygen and the nearest oxygen of the Al<sub>2</sub>Si<sub>4</sub>O<sub>18</sub> ring precludes the existence of hydrogen bonding because, when hydrogen bonds exist, the associated O–O distances usually equal 2.7–2.8 Å (Vinogradov and Linell, 1971). Moreover, Langer and Schreyer (1976) point out that the slight energy decrease of about 80 cm<sup>-1</sup> (0.2 kcal) for the IR-stretching modes of H<sub>2</sub>O in cordierite compares to those for H<sub>2</sub>O vapor, this again indicating unlikelihood of strong hydrogen bonds, these latter usually being in the range of 5 kcal.

The linear CO<sub>2</sub> molecule (4.96 Å) has been observed with its elongation mostly normal to the *c* axis (Farrell and Newnham, 1967) and only to a smaller degree parallel to *c* (Farmer, 1974). Our IR results confirm this. Such observations of the orientation of CO<sub>2</sub> (and H<sub>2</sub>O) in natural specimens may be complicated because specific CO<sub>2</sub> and H<sub>2</sub>O orientations may result from interaction with other channel occupants (or each other). On the other hand, the small shifts of the CO<sub>2</sub> bands in IR spectra for cordierite as compared to those of CO<sub>2</sub> in the gaseous state suggest only minor interaction with the silicate framework. The marked increase in the refractive index γ, as channel CO<sub>2</sub> is reintroduced into a channel-evacuated White Well crystal (Table 3), suggests that CO<sub>2</sub> is dominantly aligned parallel to *a* (= *Z*). The refractive index α displays the least

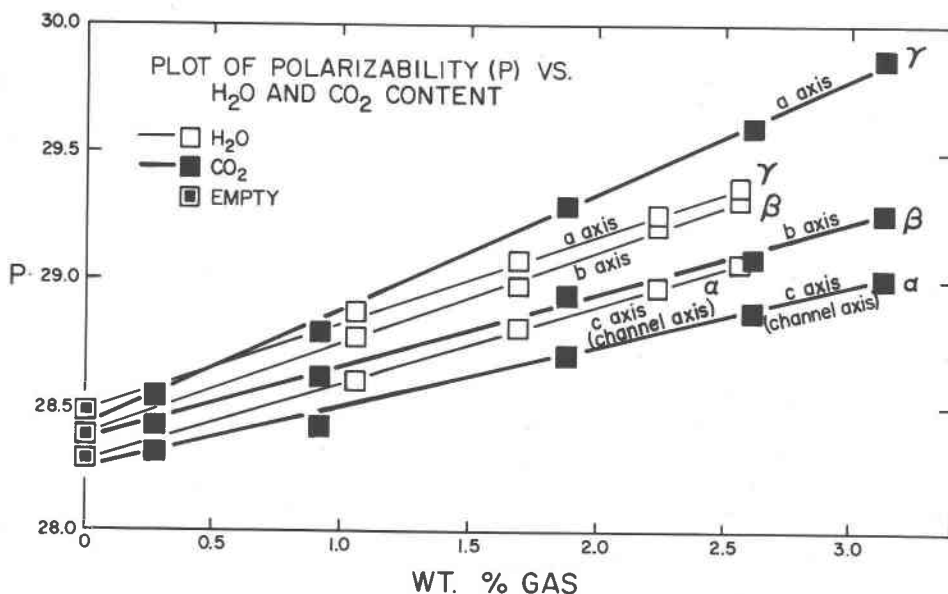


Fig. 4. Variation of the Lorentz-Lorenz polarizability  $P$  (in Å<sup>3</sup>) for White Well cordierite crystals as their channels are increasingly occupied by H<sub>2</sub>O (hollow squares) and by CO<sub>2</sub> (solid squares). Estimated standard deviations are not implied by the size of the squares.

rate of increase with channel CO<sub>2</sub>. This suggests that CO<sub>2</sub> is dominantly aligned perpendicular to *c* (since  $X = c$ ), even when not parallel to *a*, and thus agrees with the IR results.

Elongation of the CO<sub>2</sub> molecule parallel to *a* (type *a* orientation) avoids strong interaction with the oxygens that form the channel walls and with any other CO<sub>2</sub> molecules in the cavities above or below (Fig. 5A). The most probable position for CO<sub>2</sub>, if oriented parallel to *c* (type *c* orientation), is with its carbon centered within the Al<sub>2</sub>Si<sub>4</sub>O<sub>18</sub> ring (Fig. 5B) so that its oxygen atoms extend into the cages. In such case, however, only every second cavity could be occupied by CO<sub>2</sub> because  $c_0/2$  equals 4.7 Å whereas the length of CO<sub>2</sub> is 4.96 Å. Type *c* orientation might be favored by a low degree of channel filling, perhaps coupled with presence of cations (Na<sup>+</sup>, etc.) in adjacent available sites.

With increased re-introduction of channel CO<sub>2</sub>,  $a_0$  initially decreased but then increased,  $b_0$  remained constant, and  $c_0$  increased significantly (Table 2, Fig. 3). With increased channel H<sub>2</sub>O,  $a_0$  decreased,  $b_0$  perhaps increased slightly, and  $c_0$  initially increased but then levelled off. Channel H<sub>2</sub>O and channel CO<sub>2</sub> both decreased the distortion index  $\Delta$  (Table 2).

Under the same *P-T* conditions, fewer CO<sub>2</sub> molecules than H<sub>2</sub>O molecules enter cordierite's channels. At 600°C and 6 kbar, approximately twice as many H<sub>2</sub>O molecules as CO<sub>2</sub> molecules enter (so that the number of oxygens entering is about the same in each case). Presumably, the elongate CO<sub>2</sub> molecule enters the channel like a needle entering a tube of slightly larger diameter. After reaching a cavity, the CO<sub>2</sub> molecule must rotate 90° to become parallel to *a*, its favored orientation. With heating to expel CO<sub>2</sub> from the channels, any CO<sub>2</sub> molecule aligned along *a* must rotate 90° in order to diffuse outward along the channels or else decompose to smaller breakdown products.

Puzzingly, at 600°C and 5 kbar, we observed 2.61 wt.% CO<sub>2</sub> to enter the White Well cordierite whereas Johannes and Schreyer (1980) observed that, under these same *P-T* conditions only 1 wt.% CO<sub>2</sub> entered finely powdered synthetic Mg-cordierite. In each case the same equipment and methods were used to determine CO<sub>2</sub>. Possible synthetic cordierites possess a greater frequency of channel offsets or dislocations (by twin, domain, or grain boundaries). Also, Johannes and Schreyer (1980), using two different types of synthetic Mg-cordierites, noted a pronounced effect of the starting material

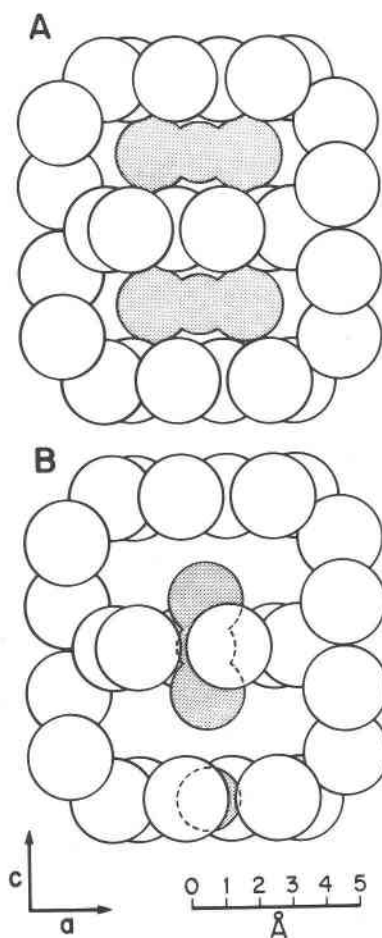


Fig. 5. A portion of the cordierite structure, viewed down the *b*-axis, showing only those oxygen atoms that form a single channel in the cordierite framework. Three "bottlenecks" are thus viewed edgewise, and midway between the centers of each two bottlenecks are the large open cages. (A) A CO<sub>2</sub> molecule (shaded dark) is shown in its preferred position (parallel to the *a*-axis) in the cage. (B) A CO<sub>2</sub> molecule oriented parallel to the *c*-axis, an orientation considerably less frequent than in (A). Centered in the lowest ring is an almost completely obscured sodium atom (shaded where seen).

on the amount of CO<sub>2</sub> (but not H<sub>2</sub>O) that entered. For natural cordierites, on the other hand, the effect of channel Na or of Fe content upon entry of CO<sub>2</sub> is not as yet known. A catalytic interaction between Na and CO<sub>2</sub> might favor incorporation of CO<sub>2</sub> into the channels. These conjectures aside, the internal consistency of our optical data, which yield Lorentz-Lorenz polarizabilities that plot quite linearly relative to wt.% CO<sub>2</sub> (Fig. 4), corroborate a value of 2.61 wt.% CO<sub>2</sub> for the grain held at 600°C, 5 kbar in a CO<sub>2</sub> atmosphere.

Knowledge of the effect of CO<sub>2</sub> and H<sub>2</sub>O on the

optical properties of cordierites may enable us (Armbruster and Bloss, in preparation) to estimate the absolute and/or relative amounts of H<sub>2</sub>O and CO<sub>2</sub> in natural cordierites. Ultimately, it is hoped, such data may provide estimates of the composition of the fluid phase with which a cordierite was last in equilibrium.

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