Calculations of fluid-ternary solid solution equilibria: An application of the Wilson equation to fluid-(Fe,Mn,Mg)TiO₃ equilibria at 600 °C and 1 kbar

Y. SHIBUE

Geoscience Institute, Hyogo University of Teacher Education, Hyogo 673-1494, Japan

ABSTRACT

The Wilson equation (Wilson 1964) is applied to (Fe,Mn,Mg)TiO₃ solid solutions for obtaining the mixing properties of the ternary solid solution at 600 °C and 1 kbar. The present study utilizes data on cation exchange between (Fe,Mn)Cl_{2(aq)} and (Fe,Mn)TiO₃, between (Mn,Mg)Cl_{2(aq)} and (Mn,Mg)TiO₃, and between (Fe,Mg)Cl_{2(aq)} and (Fe,Mg)TiO₃ (Kubo et al. 1992). The molar excess Gibbs energy (G^{ex}) is the following: G^{ex} (kJ/mol) = $-7.260[X_{FeTiO_3}ln(X_{FeTiO_3} + 1.314X_{MnTiO_3} + 0.962X_{MgTiO_3})]$ + $X_{MnTiO_3}ln(0.585X_{FeTiO_3} + X_{MnTiO_3} + 0.393X_{MgTiO_3}) + X_{MgTiO_3}ln(0.406X_{FeTiO_3} + 0.371X_{MnTiO_3} + X_{MgTiO_3})]$, where X stands for the mole fraction of the subscripted component. The predicted compositions of (Fe,Mn,Mg)Cl_{2(aq)} fluids in equilibrium with the ternary solid solutions are in good agreement with the experimental values.

INTRODUCTION

The Margules equation has been used by many investigators to describe the mixing properties of mineral solid solutions (e.g., Ganguly and Saxena 1987). By means of the Margules equation, the molar excess Gibbs energy (G^{ex}) of a ternary solid solution is written as follows (e.g., Mukhopadhyay et al. 1993):

$$G^{\text{ex}} = X_1 X_2 (X_2 W_{12} + X_1 W_{21}) + X_1 X_3 (X_3 W_{13} + X_1 W_{31}) + X_2 X_3 (X_3 W_{23} + X_2 W_{32}) + X_1 X_2 X_3 C_{123}$$
(1)

where *X* and *W* stand for the mole fraction of the subscripted component and the Margules parameter for the subscripted pair, respectively. The term C_{123} is the ternary interaction parameter, which is independent of the composition and the binary Margules parameters. If the ternary system consists of only symmetric regular binaries, the ternary interaction parameter is equal to zero (Mukhopadyay et al. 1993). Otherwise, there are no reasons for assuming $C_{123} = 0$ without experimental confirmation.

As an alternative method for describing the mixing properties of multicomponent solutions, the Wilson equation (Wilson 1964) has been used for real solution mixtures. The molar excess Gibbs energy of a ternary solution is written as follows:

$$G^{\text{ex}} = -RT \sum_{i=1}^{3} X_i \ln \left(\sum_{j=1}^{3} \Lambda_{ij} X_j \right)$$
(2)

where R is the universal gas constant, T is the absolute temperature, and Λ_{ij} stands for the Wilson parameter for the pair of components i and j. By definition, Λ_{ij} is equal to 1 in Equation 2.

By an appropriate differentiation of Equation 2, the activity coefficient of component i (γ_i) is expressed as follows:

$$\ln \gamma_{i} = 1 - \ln \left(\sum_{j=1}^{3} X_{j} A_{ij} \right) - \sum_{k=1}^{3} \left(\frac{X_{k} A_{ki}}{\sum_{i=1}^{3} X_{j} A_{kj}} \right).$$
(3)

The activity coefficients of the components in the ternary solution can be computed by combining the Wilson parameters for the three binary solutions whose components make up the ternary system. Thus, the main advantage of the Wilson equation is its potential applicability to the ternary solution without the ternary interaction parameter. When the number of the components is n, the equations for G^{ex} and activity coefficients can be obtained by substituting n for the limit sum 3 in Equations 2 and 3. Therefore, we can compute activity coefficients for solutions of arbitrary number of components by combining the Wilson parameters for binary mixtures.

The major drawback of the Wilson equation is its inapplicability to partially miscible solutions. Renon and Prausnitz (1968) modified the Wilson equation by introducing additional parameters to make it applicable to those mixtures. Applications of their equation (the NRTL equation) to partially miscible multicomponent mixtures might require more complicated calculations than those of the Wilson equation.

This study focuses on an application of the Wilson equation to mineral solid solutions. As an illustrative example, this study analyzes the experimental results on fluid– (Fe,Mn,Mg)TiO₃ equilibria (Kubo et al. 1992). Although there are many studies on cation exchange reactions between minerals and aqueous solutions, only a few studies deal with both binary and ternary solid solutions. The reason for choosing fluid–(Fe,Mn,Mg)TiO₃ equilibria is the availability of such experimental data. Based on the results of the three binary cation exchange experiments, the present study obtains the mixing properties of the ternary solid solution. The present study computes the compositions of (Fe,Mn,Mg)Cl_{2(aq)} in equilibria

0003-004X/99/0009-1375\$05.00

^{*}E-mail:yshibue@sci.hyogo-u.ac.jp

with (Fe,Mn,Mg)TiO₃ solid solutions and compares these computed results with the experimental data. This study also compares the Wilson equation with the Margules equation in regard to the predictability of the fluid-phase compositions.

A subsidiary purpose of the present study was to use the nonlinear programming technique of Britt and Luecke (1973) for obtaining the Wilson parameters and Gibbs energy of reactions. Although the computational technique is more complicated than ordinary linear least-square regression, the results of the calculation obey the law of mass action.

THE WILSON EQUATION

The Wilson equation is based on the Flory-Huggins equation (Flory 1941; Huggins 1941) and the concept of "local volume fraction" (Wilson 1964), both of which have been only rarely referred to in the literature of geology, mineralogy, and geochemistry. The derivation of the Flory–Huggins equation was given in detail by Flory (1944) and is not repeated here. Only the resultant Flory–Huggins equation for a ternary mixture is shown below. Thereafter, this paper derives the Wilson equation by introducing the concept of "local volume fraction" following the method of Prausnitz et al. (1986).

According to the Flory–Huggins equation, the molar Gibbs energy of mixing (G^{mixing}) is expressed as follows:

$$G^{\text{mixing}} = \mathbf{R}T(X_1 \ln \phi_1 + X_2 \ln \phi_2 + X_3 \ln \phi_3) \tag{4}$$

where X_1 , X_2 , and X_3 stand for the mole fractions of the solvent, solute 1, and solute 2, respectively, and ϕ_1 , ϕ_2 , and ϕ_3 are the volume fractions of the solvent, solute 1, and solute 2, respectively.

Concept of "local volume fraction"

In a ternary solution being considered here, the focus is on a central molecule of type i. The probability of finding a molecule of type j, relative to finding a molecule of type i, adjacent to the central molecule is expressed in terms of the mole fraction and two Boltzmann factors:

$$\frac{X_{ji}}{X_{ii}} = \frac{X_{j} \exp\left(\frac{-\lambda_{ji}}{RT}\right)}{X_{i} \exp\left(\frac{-\lambda_{ii}}{RT}\right)}$$
(5)

where X_{ji} and X_{ii} stand for the mole fractions of components j and i around the central molecule of type i. The terms λ_{ji} and λ_{ii} are related to the potential energies of a j–i and an i–i pair, respectively. The ternary interaction parameter does not appear in Equation 5. It should be noted that Equation 5 considers only the nearest-neighbor molecules around the central molecule in a spherical liquid-like environment. The equation then assumes that the lattice sites on which the molecules are located could be displaced by the interaction energies between molecules.

Wilson (1964) defined the local volume fraction of component i (ϕ_{ii}) as follows:

$$\phi_{ii} = \frac{V_i X_{ii}}{V_1 X_{1i} + V_2 X_{2i} + V_3 X_{3i}}$$
(6)

where V_1 , V_2 , and V_3 stand for the molar volumes of the subscripted components in their pure state. It should be noted that $\phi_{11} + \phi_{22} + \phi_{33}$ is not always equal to 1 (Prausnitz et al. 1986). As a result, the local volume fraction is not equivalent to the volume fraction and should be considered as a conceptual quantity in the microscale environment. The Wilson parameter (Λ_{ij}) is defined as follows:

$$\Lambda_{ij} = \frac{V_j}{V_i} \exp\left(\frac{\lambda_{ii} - \lambda_{ij}}{RT}\right).$$
(7)

By the above definition, Λ_{ii} is equal to 1. By combining Equations 5 to 7, ϕ_{ii} is expressed as follows:

$$\phi_{ii} = \frac{X_{ii}}{\Lambda_{i1}X_1 + \Lambda_{i2}X_2 + \Lambda_{i3}X_3}.$$
(8)

Wilson (1964) substituted the local volume fractions of the components, instead of their volume fractions, into the Flory–Huggins equation (Eq. 4). After rearrangement, G^{mixing} can be expressed as follows:

$$\begin{aligned} G^{\text{mixing}} &= -RT\{[X_1 \ln(\Lambda_{11}X_1 + \Lambda_{12}X_2 + \Lambda_{13}X_3)] + [X_2 \ln(\Lambda_{21}X_1 + \Lambda_{22}X_2 + \Lambda_{23}X_3)] + [X_3 \ln(\Lambda_{31}X_1 + \Lambda_{32}X_2 + \Lambda_{33}X_3)] - (X_1 \ln X_1 + X_2 \ln X_2 + X_3 \ln X_3)\}. \end{aligned}$$

Therefore, the molar excess Gibbs energy (G^{ex}) can be derived as follows:

$$G^{\text{ex}} = G^{\text{mixing}} - RT(X_1 \ln X_1 + X_2 \ln X_2 + X_3 \ln X_3) \quad (10)$$

which is equivalent to Equation 2.

Activity coefficients of the individual components can be obtained from the partial derivatives of the excess Gibbs energy with respect to the number of moles of that component as follows:

$$\ln \gamma_{i} = \frac{1}{RT} \left(\frac{\partial n G^{ex}}{\partial n_{i}} \right)_{T,P,n_{|s|}}$$
(11)

where *n* and n_i stand for the total number of moles of all the components and the number of moles of component i, respectively. Using Equations 9 and 10, the activity coefficient of component i can be expressed as Equation 3.

APPLICATIONS TO FLUID-(FE,MN)TIO₃, FLUID-(MN,MG)TIO₃, AND FLUID-(FE,MG)TIO₃ EQUILIBRIA

Kubo et al. (1992) carried out cation exchange experiments on the following systems at 600 °C and 1 kbar by using the method of Uchida et al. (1989):

(A) $FeTiO_3 + MnCl_{2(aq)} = MnTiO_3 + FeCl_{2(aq)}$

(B) $MnTiO_3 + MgCl_{2(aq)} = MgTiO_3 + MnCl_{2(aq)}$

(C) $MgTiO_3 + FeCl_{2(aq)} = FeTiO_3 + MgCl_{2(aq)}$

Total molarity of metal chloride in the aqueous phase was 1.0 M in all their experiments. Only (Fe,Mn,Mg)TiO₃ was observed as the solid product.

By using the Wilson equation for describing the activity coefficients for the components of the solid phase, the Gibbs energy of reaction A can be written as follows:

$$\frac{\Delta G_{\rm A}^{0}}{RT} = -\ln \mathbf{K} = -\ln \frac{X_{\rm MnTiO_3} \gamma_{\rm MnTiO_3} Y_{\rm Fe}}{X_{\rm FeTiO_3} \gamma_{\rm FeTiO_3} Y_{\rm Mn}} = -\ln \frac{X_{\rm MnTiO_3} Y_{\rm Fe}}{X_{\rm FeTiO_3} Y_{\rm Mn}} + \ln \frac{\Lambda_{\rm MnFe} X_{\rm FeTiO_3} + X_{\rm MnTiO_3}}{X_{\rm FeTiO_3} + \Lambda_{\rm FeMn} X_{\rm MnTiO_3}} + \frac{\Lambda_{\rm MnFe} X_{\rm FeTiO_3} - X_{\rm FeTiO_3}}{X_{\rm FeTiO_3} + \Lambda_{\rm FeMn} X_{\rm MnTiO_3}} + \frac{X_{\rm MnTiO_3} - \Lambda_{\rm FeMn} X_{\rm MnTiO_3}}{\Lambda_{\rm MnFe} X_{\rm FeTiO_3} + X_{\rm MnTiO_3}}$$

$$(12)$$

where K stands for the equilibrium constant and Y_{Fe} and Y_{Mn} designate the ratios of the molalities of the subscripted species to the total molality of metal chloride in the fluid phase. The *Y* values are expressed as follows:

$$Y_{\rm Fe} = \frac{m_{\rm FeCl_2(aq)}}{m_{\rm FeCl_2(aq)} + m_{\rm MnCl_2(aq)} + m_{\rm MgCl_2(aq)}}$$
(13)

$$Y_{\rm Mn} = \frac{m_{\rm MnCl_2(aq)}}{m_{\rm FeCl_2(aq)} + m_{\rm MnCl_2(aq)} + m_{\rm MgCl_2(aq)}}$$
(14)

 ΔG^0 for Reaction B or C can be obtained by changing the subscripts appropriately. For example, Y_{Mg} denotes the following ratio.

$$Y_{\rm Mg} = \frac{m_{\rm MgCl_2(aq)}}{m_{\rm FeCl_2(aq)} + m_{\rm MnCl_2(aq)} + m_{\rm MgCl_2(aq)}}$$
(15)

The present study ignores the ionic species and assumes that the ratios of activity coefficients of the neutral aqueous species are equal to 1. Kubo et al. (1992) reported the concentrations of the aqueous species in units of molarity. Because the present study calculates the ratios of concentrations, the conversion factor of molar scale into molal scale does not appear in Equations 13 to 15.

Values of ΔG^0 for reactions A–C and the Wilson parameters were obtained by using the following constraints:

$$\Delta G_{\rm C}^0 = -\left(\Delta G_{\rm A}^0 + \Delta G_{\rm B}^0\right) \tag{16}$$

$$\Lambda_{\rm MgFe} = \frac{\Lambda_{\rm MgMn} \Lambda_{\rm MnFe} \Lambda_{\rm FeMg}}{\Lambda_{\rm MnMg} \Lambda_{\rm FeMn}}$$
(17)

The derivation of Equation 17 was given in Hala (1972). By imposing the above constraints on the calculations, the calculated ΔG^0 values as well as the Wilson parameters are internally consistent. For the calculation of all the values, the nonlinear programming technique derived by Britt and Luecke (1973) was used in the present study. Solid and fluid compositions were also computed for each experimental data point. This study considers that the standard errors of measurements of both compositions are 0.005 in terms of mole fraction or the ratio of molality. Substitutions of the calculated compositions of the solid and fluid phases and the Wilson parameters into Equation 12 make the resultant lnK values become constant. The same is true for reactions B and C. Detailed calculation procedures are shown in the Appendix. The results of the calculations of $\Delta G^{0}/RT$ and the Wilson parameters are listed in Table 1. The experimental and computed compositions are summarized in Tables 2–4. From Equation 12 and the equivalent expressions for reactions B and C, we get the following equations:

$$\frac{Y_{\text{Fe}}}{Y_{\text{Mn}}} = \exp\left[\left(-\frac{\Delta G_{\text{A}}^{0}}{\text{R}T}\right) - \ln\left(\frac{X_{\text{MnTiO}}, \gamma_{\text{MnTiO}}}{X_{\text{FeTiO}}, \gamma_{\text{FeTiO}}}\right)\right]$$
(18)

$$\frac{Y_{\rm Mn}}{Y_{\rm Mg}} = \exp\left[\left(-\frac{\Delta G_{\rm B}^0}{{\rm R}T}\right) - \ln\left(\frac{X_{\rm MgTiO}, \gamma_{\rm MgTiO_1}}{X_{\rm MnTiO_1} \gamma_{\rm MnTiO_1}}\right)\right]$$
(19)

$$\frac{Y_{\rm Mg}}{Y_{\rm Fe}} = \exp\left[\left(-\frac{\Delta G_{\rm C}^0}{RT}\right) - \ln\left(\frac{X_{\rm FeTiO_1}\gamma_{\rm FeTiO_1}}{X_{\rm MgTiO_1}\gamma_{\rm MgTiO_1}}\right)\right]$$
(20)

If we specify the solid compositions, activity coefficients of the components in the solid phase are determined by Equation 3. By using the activity coefficients at specific solid compositions and the ΔG^0 value for the relevant reaction, the fluidphase compositions can be obtained from the left-hand sides of Equations 18–20. Then cation exchange isotherms for reactions A, B, and C are computed from Equations 18, 19, and 20, respectively (Figs. 1–3). The computed isotherms show good agreement with the experimental results of Kubo et al. (1992). Activities of end-member components are plotted against the mole fractions for the three binary solid solutions (Fig. 4). The (Fe,Mn)TiO₃ and (Fe,Mg)TiO₃ solid solutions show almost ideal behavior whereas (Mn,Mg)TiO₃ shows distinct nonideality.

It is possible to compute the ΔG^0 value and the Wilson parameters for each reaction and to compare the results with the internally consistent values in Table 1. Calculations on each reaction gave internally inconsistent results. Furthermore, the iterative computations often showed divergences due to the nonlinearity of the Wilson equation. Therefore, this study did not attempt to extract thermodynamic quantities from the individual exchange reactions.

Calculations of fluid-(Fe,Mn,Mg)TiO₃ equilibria

By substituting the Wilson parameters into Equation 2, the molar excess Gibbs energy of $(Fe,Mn,Mg)TiO_3$ solid solution at 600 °C and 1 kbar is expressed as follows.

$$\begin{split} G^{\text{ex}}(\text{kJ/mol}) &= -7.260[X_{\text{FeTiO}_3} \ln(X_{\text{FeTiO}_3} + 1.314X_{\text{MnTiO}_3} + 0.962X_{\text{MgTiO}_3}) \\ &+ X_{\text{MnTiO}_3} \ln(0.585X_{\text{FeTiO}_3} + X_{\text{MnTiO}_3} + 0.393X_{\text{MgTiO}_3}) \\ &+ X_{\text{MgTiO}_3} \ln(0.406X_{\text{FeTiO}_3} + 0.371X_{\text{MnTiO}_3} + X_{\text{MgTiO}_3})] \end{split}$$

Table 1. △G⁰ for reactions A–C and Wilson parameters for (Fe,Mn)TiO₃, (Mn,Mg)TiO₃, and (Fe,Mg)TiO₃ solid solutions

 $\begin{array}{l} \text{Reaction A: } \text{FeTiO}_3 + \text{MnCl}_{2(aq)} = \text{MnTiO}_3 + \text{FeCl}_{2(aq)} \\ \text{Reaction B: } \text{MnTiO}_3 + \text{MgCl}_{2(aq)} = \text{MnTiO}_3 + \text{MnCl}_{2(aq)} \\ \text{Reaction C: } \text{MgTiO}_3 + \text{FeCl}_{2(aq)} = \text{FeTiO}_3 + \text{MgCl}_{2(aq)} \\ \end{array}$

$$\frac{\Delta G_{\rm A}}{RT} = 1.462(0.042), \frac{\Delta G_{\rm B}}{RT} = -0.975(0.036), \frac{\Delta G_{\rm C}}{RT} = -0.487(0.035)$$

 $\Lambda_{MnFe} = 0.585(0.220), \ \Lambda_{FeMn} = 1.314(0.370), \ \Lambda_{MgMn} = 0.371(0.087)$

 $\begin{array}{l} \Lambda_{\text{MnMg}} = 0.393(0.073), \ \Lambda_{\text{FeMg}} = 0.962(0.213), \ \Lambda_{\text{MgFe}} = 0.406(0.138) \\ \hline \textit{Notes:} \ \textit{Values in parentheses indicate the standard errors.} \ \textit{RT} \ is \ \textit{7.260} \\ \textit{(kJ/mol) at } 600 \ ^{\circ}\text{C.} \end{array}$

$$\ln \gamma_{\text{FeTiO}_{3}} = 1 - \ln \left(X_{\text{FeTiO}_{3}} + X_{\text{MnTiO}_{3}} \Lambda_{\text{FeMn}} + X_{\text{MgTiO}_{3}} \Lambda_{\text{FeMg}} \right) - \frac{X_{\text{FeTiO}_{3}}}{X_{\text{FeTiO}_{3}} + X_{\text{MnTiO}_{3}} \Lambda_{\text{FeMn}} + X_{\text{MgTiO}_{3}} \Lambda_{\text{FeMg}}} - \frac{X_{\text{MnTiO}_{3}} \Lambda_{\text{MnFe}}}{X_{\text{FeTiO}_{3}} \Lambda_{\text{MnFe}} + X_{\text{MnTiO}_{3}} + X_{\text{MgTiO}_{3}} \Lambda_{\text{MnMg}}} - \frac{X_{\text{MgTiO}_{3}} \Lambda_{\text{MgFe}}}{X_{\text{FeTiO}_{3}} \Lambda_{\text{MgFe}} + X_{\text{MnTiO}_{3}} \Lambda_{\text{MgMn}} + X_{\text{MgTiO}_{3}}}$$
(22)

$$\ln \gamma_{\rm MnTiO_3} = 1 - \ln \left(X_{\rm FeTiO_3} \Lambda_{\rm MnFe} + X_{\rm MnTiO_3} + X_{\rm MgTiO_3} \Lambda_{\rm MnMg} \right) - \frac{X_{\rm FeTiO_3} \Lambda_{\rm FeMn}}{X_{\rm FeTiO_3} + X_{\rm MnTiO_3} \Lambda_{\rm FeMn} + X_{\rm MgTiO_3} \Lambda_{\rm FeMg}} - \frac{X_{\rm MnTiO_3}}{X_{\rm FeTiO_3} \Lambda_{\rm MnFe} + X_{\rm MnTiO_3} + X_{\rm MgTiO_3} \Lambda_{\rm MnMg}} - \frac{X_{\rm MgTiO_3} \Lambda_{\rm MgFe}}{X_{\rm FeTiO_3} \Lambda_{\rm MgFe} + X_{\rm MnTiO_3} \Lambda_{\rm MgMn} + X_{\rm MgTiO_3}}$$
(23)

$$\ln \gamma_{\text{MgTiO}_{3}} = 1 - \ln \left(X_{\text{FeTiO}_{3}} \Lambda_{\text{MgFe}} + X_{\text{MnTiO}_{3}} \Lambda_{\text{MgMn}} + X_{\text{MgTiO}_{3}} \right)$$

$$- \frac{X_{\text{FeTiO}_{3}} \Lambda_{\text{FeMg}}}{X_{\text{FeTiO}_{3}} + X_{\text{MnTiO}_{3}} \Lambda_{\text{FeMg}}} + X_{\text{MgTiO}_{3}} \Lambda_{\text{FeMg}}}$$

$$- \frac{X_{\text{MnTiO}_{3}} \Lambda_{\text{MnMg}}}{X_{\text{FeTiO}_{3}} \Lambda_{\text{MnFe}} + X_{\text{MnTiO}_{3}} + X_{\text{MgTiO}_{3}} \Lambda_{\text{MnMg}}}$$

$$- \frac{X_{\text{MgTiO}_{3}}}{X_{\text{FeTiO}_{3}} \Lambda_{\text{MgFe}} + X_{\text{MnTiO}_{3}} \Lambda_{\text{MgMn}} + X_{\text{MgTiO}_{3}}}$$

$$(24)$$

By using Equation 3, activity coefficients of $FeTiO_3$, $MnTiO_3$, and $MgTiO_3$ in the ternary solid solution are expressed as follows:

On the basis of the above expressions of activity coefficients, the compositions of $(Fe,Mn,Mg)Cl_{2(aq)}$ in equilibria with $(Fe,Mn,Mg)TiO_3$ were computed. The solid compositions were taken from the experimental results of Kubo et al. (1992). The method of Shallcross et al. (1988) was used for the present calculation. Let us define Z values as follows.

$$Z_1 = \frac{Y_{\text{Fe}}}{Y_{\text{Mn}}}, \ Z_2 = \frac{Y_{\text{Mn}}}{Y_{\text{Mg}}}, \ Z_3 = \frac{Y_{\text{Mg}}}{Y_{\text{Fe}}}$$

From the solid compositions, activities of the end-member components can be computed. Thus, Z_1 , Z_2 , and Z_3 values are obtained from Equations 18–20. If we use Z_1 and Z_2 values, Y values are calculated as follows:

$$Y_{\rm Fe} = \frac{Z_1 Z_2}{Z_1 Z_2 + Z_2 + 1} \tag{25}$$

$$Y_{\rm Mn} = \frac{Z_2}{Z_1 Z_2 + Z_2 + 1} \tag{26}$$

$$Y_{\rm Mg} = \frac{1}{Z_1 Z_2 + Z_2 + 1} \tag{27}$$



FIGURE 1. The cation exchange isotherm computed in this study (solid line) and the experimental data points (Kubo et al. 1992) for the reaction $FeTiO_3+MnCl_{2(aq)} = MnTiO_3+FeCl_{2(aq)}$. X_{Fe} and Y_{Fe} stand for the mole fraction of $FeTiO_3$ in the solid phase and the molality ratio defined by Equation 13.

 TABLE 2.
 Experimental and calculated results on the cation exchange reaction $FeTiO_3 + MnCl_{2(aq)} = MnTiO_3 + FeCl_{2(aq)}$

Experimental results*		Calculated results†	
X _{FeTiO3}	Y _{Fe}	$\overline{X_{\text{FeTIO}_3}}$	Y _{Fe}
0.928	0.752	0.939	0.748
0.891	0.624	0.893	0.623
0.864	0.550	0.859	0.553
0.964	0.828	0.962	0.828
0.898	0.632	0.897	0.632
0.944	0.740	0.937	0.742
0.485	0.184	0.485	0.183
0.325	0.086	0.318	0.104
0.404	0.137	0.402	0.141
0.617	0.259	0.613	0.264
0.521	0.159	0.502	0.192
0.493	0.163	0.483	0.182
0.182	0.053	0.182	0.054
0.484	0.150	0.471	0.175
0.557	0.180	0.537	0.213
0.723	0.311	0.698	0.337
0.286	0.093	0.286	0.092
0.835	0.512	0.836	0.511

Other solutions of *Y* values are possible when we use the set of Z_1 and Z_3 or the set of Z_2 and Z_3 . Therefore, three sets of *Y* values are possible. The weighted average values of fluid-phase compositions can be obtained from the following equations:

$$Y_{\rm Fe} = \frac{X_{\rm MnTiO_3} Z_1 Z_2}{Z_1 Z_2 + Z_2 + 1} + \frac{X_{\rm FeTiO_3} Z_1}{Z_1 Z_3 + Z_1 + 1} + \frac{X_{\rm MgTiO_3}}{Z_2 Z_3 + Z_3 + 1} \quad (28)$$
$$Y_{\rm Mn} = \frac{X_{\rm MnTiO_3} Z_2}{Z_1 Z_2 + Z_2 + 1} + \frac{X_{\rm FeTiO_3}}{Z_1 Z_3 + Z_1 + 1} + \frac{X_{\rm MgTiO_3} Z_2 Z_3}{Z_2 Z_3 + Z_3 + 1} \quad (29)$$

$$Y_{\rm Mn} = \frac{X_{\rm MnTiO_3}Z_2}{Z_1Z_2 + Z_2 + 1} + \frac{X_{\rm FeTiO_3}}{Z_1Z_3 + Z_1 + 1} + \frac{X_{\rm MgTiO_3}Z_2Z_3}{Z_2Z_3 + Z_3 + 1} \quad (30)$$

 TABLE 3.
 Experimental and calculated results on the cation exchange reaction $MnTiO_3 + MgCl_{2(aq)} = MgTiO_3 + MnCl_{2(aq)}$

Experimental results*		Calculated results†	
X_{MnTiO_3}	Y _{Mn}	X _{MnTiO3}	$Y_{ m Mn}$
0.480	0.716	0.480	0.722
0.456	0.692	0.450	0.715
0.528	0.745	0.530	0.735
0.707	0.775	0.705	0.781
0.136	0.559	0.140	0.555
0.581	0.761	0.584	0.748
0.339	0.666	0.334	0.680
0.046	0.410	0.064	0.404
0.605	0.769	0.609	0.754
0.663	0.780	0.666	0.769
0.147	0.524	0.127	0.538
0.914	0.846	0.894	0.872
0.295	0.669	0.296	0.666
0.081	0.484	0.093	0.478
0.168	0.553	0.151	0.568
0.320	0.649	0.311	0.672
0.393	0.648	0.379	0.695
0.908	0.861	0.897	0.875
0.813	0.768	0.793	0.813
0.743	0.797	0.744	0.794
0.356	0.706	0.361	0.690
* Kubo et al.	(1992).		

† Calculated results of this study

TABLE 4. Experimental and calculated results on the cation exchange reaction MgTiO₃ + FeCl_{2(aq)} = FeTiO₃ + MgCl_{2(aq)}

Experimental results*		Calculated results†		
X _{MgTiO3}	Y _{Mg}	X _{MgTiO3}	$Y_{\rm Mg}$	
0.053	0 183	0.059	0 180	
0 122	0 248	0 100	0.261	
0.040	0.118	0.035	0.120	
0 151	0.318	0 143	0.324	
0.055	0.185	0.060	0.183	
0.359	0.548	0.375	0.524	
0.262	0.452	0.267	0.446	
0.167	0.328	0.155	0.338	
0.168	0.357	0.170	0.355	
0.346	0.507	0.347	0.505	
0.525	0.595	0.515	0.612	
0.922	0.922	0.925	0.919	
0.911	0.904	0.910	0.905	
0.792	0.787	0.785	0.796	
0.807	0.772	0.787	0.797	
0.474	0.605	0.482	0.592	
0.943	0.944	0.946	0.941	
0.594	0.611	0.571	0.647	
0.715	0.762	0.724	0.750	
0.343	0.456	0.321	0.487	
0.198	0.378	0.194	0.382	
0.457	0.607	0.471	0.585	
0.716	0.713	0.701	0.734	
0.608	0.623	0.586	0.657	
0.452	0.565	0.448	0.571	
0.361	0.493	0.351	0.508	
0.650	0.722	0.661	0.706	
* Kubo et al	(1992)			

† Calculated results of this study.

The weights correspond to the mole fractions of the components common to Z_i and Z_j .

Computed fluid compositions and the experimental results of Kubo et al. (1992) are compared in Figure 5. The agreement is generally good. Although the computed compositions deviate from the experimental compositions in the region of low Y_{Fe} , the deviations are within ±0.08 in terms of the Y value for any component (Table 5). It can be concluded that fluid



FIGURE 2. The cation exchange isotherm computed in this study (solid line) and the experimental data points (Kubo et al. 1992) for the reaction $MnTiO_3+MgCl_{2(aq)} = MgTiO_3+MnCl_{2(aq)}$. X_{Mn} and Y_{Mn} stand for the mole fraction of MnTiO₃ in the solid phase and the molality ratio defined by Equation 14.



FIGURE 3. The cation exchange isotherm computed in this study (solid line) and the experimental data points (Kubo et al. 1992) for the reaction $MgTiO_3$ +FeCl_{2(aq)} = FeTiO_3+MgCl_{2(aq)}. X_{Mg} and Y_{Mg} stand for the mole fraction of MgTiO₃ in the solid phase and the molality ratio defined by Equation 15.

compositions in equilibria with the ternary solid solutions may be predicted from the binary cation exchange experiments.

COMPARISON WITH THE MARGULES EQUATION

In the introduction, it was argued that the Margules equation without the ternary interaction parameter is not accurate for describing the mixing property of the ternary solid solution consisting of asymmetric binaries. The accuracy of the compo-





FIGURE 4. Activity (a) vs. composition (X) relations of (Fe,Mn)TiO₃, (Mn,Mg)TiO₃, and (Fe,Mg)TiO₃ solid solutions at 600 $^{\circ}$ C and 1 kbar.



FIGURE 5. Fluid compositions in equilibria with (Fe,Mn,Mg) TiO_3 solid solutions. Fluid compositions tied by lines correspond to the same solid compositions. The experimental compositions of solid and fluid phases are taken from Kubo et al. (1992). The tie lines are omitted when the computed compositions closely agree with the experimental compositions.

sitions of the fluid phase predicted by the Margules equation and by the Wilson equation can now be compared.

By an appropriate differentiation of Equation 1, activity coefficients of components 1 and 2 in a binary solid solution are written as follows (e.g., Mukhopadhyay et al. 1993).

 $\mathbf{R}T\mathbf{ln}\gamma_1 = [W_{12} + 2(W_{21} - W_{12})X_1]X_2^2 \tag{31}$

$$\mathbf{R}T \ln \gamma_2 = [W_{21} + 2(W_{12} - W_{21})X_2]X_1^2 \tag{32}$$

From the experimental results listed in Tables 2–4, ΔG^{0} values for reactions A–C and the Margules parameters were computed for the three binary solid solutions. The constraints (Eqs. 16 and 17) were not used in the calculation. The computation of three unknowns for each reaction was carried out with the nonlinear programming technique, following the procedure shown in Appendix.

The results of the calculation are listed in Table 6. By setting $C_{123} = 0$, activity coefficients of the components in the ternary solid solution of (Fe,Mn,Mg)TiO₃ were computed on the basis of the experimental compositions of the solid phase (Table 5). Activity coefficients of the three components can be written as follows (Mukhopadhyay et al. 1993):

TABLE 5. Experimental and calculated results on the cation exchange reactions between (Fe,Mn,Mg)TiO₃ and (Fe,Mn,Mg)Cl_{2(a0)}

Experimental results*		Calculated results†			
X _{FeTiO3}	$X_{\rm MnTiO_3}$	Y _{Fe}	Y _{Mn}	$Y_{\rm Fe}$	$Y_{\rm Mn}$
0.919	0.056	0.704	0.228	0.699	0.225
0.924	0.027	0.762	0.108	0.738	0.118
0.367	0.317	0.095	0.590	0.129	0.599
0.916	0.054	0.709	0.214	0.693	0.217
0.498	0.062	0.230	0.257	0.296	0.253
0.424	0.315	0.119	0.580	0.153	0.593
0.428	0.270	0.115	0.544	0.162	0.558
0.340	0.315	0.099	0.592	0.119	0.601
0.504	0.318	0.168	0.569	0.189	0.596
0.489	0.199	0.185	0.464	0.208	0.486
0.376	0.189	0.122	0.461	0.159	0.495
0.230	0.293	0.027	0.580	0.080	0.607
0.297	0.206	0.069	0.510	0.120	0.530
0.397	0.429	0.087	0.654	0.131	0.672
0.278	0.218	0.045	0.522	0.110	0.545
0.301	0.495	0.045	0.667	0.091	0.704
0.576	0.221	0.233	0.510	0.247	0.503
0.489	0.480	0.180	0.763	0.180	0.763
0.564	0.409	0.206	0.740	0.225	0.723
0.248	0.031	0.132	0.223	0.182	0.195
0.782	0.043	0.501	0.174	0.516	0.167
0.219	0.660	0.076	0.779	0.062	0.792
0.159	0.118	0.087	0.492	0.080	0.466
0.249	0.159	0.084	0.458	0.111	0.494
* Kubo et	al (1992)				

+ Calculated results of this study

TABLE 6. ΔG⁰ for reactions A–C and Margules parameters* for (Fe,Mn)TiO₃, (Mn,Mg)TiO₃, and (Fe,Mg)TiO₃ solid solutions

$\frac{\Delta G_{\rm A}^0}{RT} = 1.513(0.058), \frac{\Delta G_{\rm B}^0}{RT} = -0.908(0.043), \frac{\Delta G_{\rm C}^0}{RT} = -0.427(0.037)$
$W_{\text{FeMn}} = 0.337(0.577), W_{\text{MnFe}} = 3.254(1.661), W_{\text{MnMg}} = 11.549(1.105),$
$W_{MgMn} = 10.119(0.866), W_{MgFe} = 5.015(0.956), W_{FeMg} = 6.227(0.707)$
Notes Values in perentheses indicate the standard errors DT is 7.00

Notes: Values in parentheses indicate the standard errors. RT is 7.260 (kJ/mol) at 600 °C.

* Calculated from the results on each binary reaction. Reactions A, B, and C are indicated in Table 1 and in the text.

- $$\begin{split} \text{RTIn}\gamma_{\text{FeTiO}_3} &= W_{\text{FeMn}}X_{\text{MnTiO}_3}^2(1-2X_{\text{FeTiO}_3}) + 2W_{\text{MnFe}}X_{\text{FeTiO}_3}X_{\text{MnTiO}_3} \\ & (1-X_{\text{FeTiO}_3}) \\ & 2W_{\text{MnMg}}X_{\text{MnTiO}_3}X_{\text{MgTiO}_3}^2 2W_{\text{MgMn}}X_{\text{MnTiO}_3}^2X_{\text{MgTiO}_3} \end{split}$$
 - + $2W_{MgFe}X_{FeTiO_3}X_{MgTiO_3}(1 X_{FeTiO_3}) + W_{FeMg}X_{MgTiO_3}^2$ (1 - $2X_{FeTiO_3}$) (33)
- $RT\ln\gamma_{\text{MnTiO3}} = 2W_{\text{FeMn}}X_{\text{FeTiO3}}X_{\text{MnTiO3}}(1 X_{\text{MnTiO3}}) + W_{\text{MnFe}}X^{2}_{\text{FeTiO3}}$ $(1 2X_{\text{MnTiO3}})$
 - + $W_{MnMg}X^2_{MgTiO_3}(1 2X_{MnTiO_3}) + 2W_{MgMn}X_{MnTiO_3}X_{MgTiO_3}$ (1 - X_{MnTiO_3})

$$-2W_{\text{MgFe}}X^{2}_{\text{FeTiO}_{3}}X_{\text{MgTiO}_{3}} - 2W_{\text{FeMg}}X_{\text{FeTiO}_{3}}X^{2}_{\text{MgTiO}_{3}}$$
(34)

$$\begin{aligned} RT \ln \gamma_{MgTiO_3} &= -2W_{FeMn}X_{FeTiO_3}X^2_{MnTiO_3} - 2W_{MnFe}X^2_{FeTiO_3}X_{MnTiO_3} \\ &+ 2W_{MnMg}X_{MnTiO_3}X_{MgTiO_3}(1 - X_{MgTiO_3}) + W_{MgMn}X^2_{MnTiO_3} \\ &(1 - 2X_{MgTiO_3}) \\ &+ W_{MgFe}X^2_{FeTiO_3}(1 - 2X_{MgTiO_3}) + 2W_{FeMg}X_{FeTiO_3}X_{MgTiO_3} \\ &(1 - X_{MgTiO_3}) \end{aligned}$$

Compositions of $(Fe,Mn,Mg)Cl_{2(aq)}$ in equilibria with $(Fe,Mn,Mg)TiO_3$ were computed by the same method described

before (see Equations 28 to 30). After the computation of the fluid phase compositions, the sum of the distances were obtained between the experimental compositions and the calculated compositions in the ternary system. The overall fit to the experimental data using the Wilson equation is better than with the Margules equation having no ternary interaction parameter (Σ distances = 0.00938 vs. 0.01271). Thus, the prediction of the ternary solid solution–fluid equilibria by the Wilson equation is more accurate than that by the Margules equation on the basis of data on binary cation exchange reactions.

REFERENCES CITED

- Britt, H.I. and Luecke, R.H. (1973) The estimation of parameters in nonlinear, implicit models. Technometrics, 15, 233–247.
- Flory, P.J. (1941) Thermodynamics of high polymer solutions. Journal of Chemical Physics, 9, 660–661.
- (1944) Thermodynamics of heterogeneous polymers and their solutions. Journal of Chemical Physics, 12, 425–438.
- Ganguly, J. and Saxena, S.K. (1987) Mixtures and Mineral Reactions, 291 p. Springer–Verlag, Berlin.
- Hala, E. (1972) Note to Bruin–Prausnitz one–parameter local composition equation. Industrial and Engineering Chemistry Process Design and Development, 11, 638.
- Huggins, M.L. (1941) Solutions of long chain compounds. Journal of Chemical Physics, 9, 440.
- Kubo, T., Uchida, E., Furukawa, Y., and Imai, N. (1992) Experimental study on ion exchange equilibria between solid solution (Fe²⁺,Mn²⁺,Mg²⁺)TiO₃ and aqueous (Fe²⁺,Mn²⁺,Mg²⁺)Cl₂ solution. Kobutsugaku Zasshi (Journal of Mineralogical Society of Japan), 21, 59-67 (in Japanese with English abstract).
- Mukhopadhyay, B., Basu, S., and Holdway, M.J. (1993) A discussion of Margules– type formulation for multicomponent solutions with a generalized approach. Geochimica et Cosmochimica Acta, 57, 277–283.
- Prausnitz, J.M., Lichtenthaler, R.N., and Azevedo, E.G. (1986) Molecular Thermodynamics of Fluid-phase Equilibria. 2nd. ed., 600 p. Prentice Hall, New York.
- Renon, H. and Prausnitz, J.M. (1968) Local compositions in thermodynamic excess functions for liquid mixtures. American Institute of Chemical Engineers Journal, 14, 135–144.
- Shallcross, D.C., Herrmann, C.C., and McCoy, B.J. (1988) An improved model for the prediction of multicomponent ion exchange equilibria. Chemical Engineering Science, 43, 279–288.
- Uchida, E., Gima, M., and Imai, N. (1989) Experimental studies on ion exchange equilibria between minerals and aqueous chloride solution in the system CaWO₄-FeWO₄-MnWO₄ under supercritical condition. Geochemical Journal, 23, 339– 347.
- Wilson, G.M. (1964) Vapor–liquid equilibrium. XI. A new expression for the excess free energy mixing. Journal of American Chemical Society, 86, 127–130.

MANUSCRIPT RECEIVED MAY 18, 1998 MANUSCRIPT ACCEPTED MAY 3, 1999

PAPER HANDLED BY J. WILLIAM CAREY

APPENDIX

Britt and Luecke (1973) developed a nonlinear programming technique for obtaining the parameters in various types of regression equations. Their method was used in the present study to obtain the Gibbs energy of reactions, the Wilson parameters, and the compositions of solid and fluid phases.

To avoid lengthy notations and equations, the following abbreviations are used hereafter: $B_1 = \frac{\Delta G_n^{A}}{RT}, B_2 = \Lambda_{MnFe}, B_3 = \Lambda_{FeMn}, B_4 = \frac{\Delta G_n^{B}}{RT}, B_5 = \Lambda_{MgFe}, B_6 = \Lambda_{MnMg}, B_7 = \Lambda_{FeMg}.$

Using the constraint equations shown in the text (Eqs. 16 and 17), the other two unknowns ($\Delta G^0/RT$ for reaction C and Λ_{MgFe}) were obtained from the above seven parameters. In this study, the parameter F_i was defined as follows for all the data points relating to reaction A:

where X_i and Y_i stand for the calculated values of X_{FeTiO_3} and Y_{Fe} , respectively, for the i-th experimental data point. Similarly, by denoting X_i and Y_i as the calculated values of X_{MnTiO_3} and Y_{Mn} for the i-th experimental data point, F_i for reaction B is defined as follows.

$$F_{i} = B_{4} + \ln\left[\frac{(1-X_{i})Y_{i}}{X_{i}(1-Y_{i})}\right] - \ln\left[\frac{B_{5}X_{i}+1-X_{i}}{X_{i}+B_{6}(1-X_{i})}\right] - \frac{(B_{6}-1)X_{i}}{X_{i}+B_{6}(1-X_{i})} - \frac{(1-X_{i})(1-B_{5})}{B_{5}X_{i}+1-X_{i}}$$
(A2)

For reaction C, the following expression can be derived through the constraint equations.

$$F_{i} = -(B_{1} + B_{4}) + \ln\left[\frac{(1 - X_{i})Y_{i}}{X_{i}(1 - Y_{i})}\right] - \ln\left[\frac{B_{7}X_{i} + 1 - X_{i}}{X_{i} + \left(\frac{B_{2}B_{5}B_{7}}{B_{3}B_{6}}\right)(1 - X_{i})}\right] - \frac{\left[\left(\frac{B_{2}B_{5}B_{7}}{B_{3}B_{6}}\right) - 1\right]X_{i}}{X_{i} + \left(\frac{B_{2}B_{5}B_{7}}{B_{3}B_{6}}\right)(1 - X_{i})} - \frac{(1 - X_{i})(1 - B_{7})}{B_{7}X_{i} + 1 - X_{i}}$$
(A3)

where X_i and Y_i stand for the calculated values of X_{MgTiO_3} and

 Y_{Mg} for the i-th experimental data point. In the present study, calculations were carried out under the condition that all of the F_i values should be 0, which is equivalent to the condition that lnK for each reaction becomes constant. In other words, the computed results satisfy the law of mass action.

Standard errors of the measurements of both *X* and *Y* values are assumed to be 0.005 (= σ). Then this study minimizes the value of *Q*, which is defined as follows:

$$Q = \frac{1}{2} \sum_{i} \sigma^{-2} \left[\left(X_{i}^{\text{meas}} - x_{i} \right)^{2} + \left(Y_{i}^{\text{meas}} - y_{i} \right)^{2} \right]$$
(A4)

where the superscript, meas, stands for the measured value. The calculated values for solid and fluid compositions are designated as x and y, respectively.

We define L_i , which corresponds to the weight of the i-th experimental datum, as follows.

$$I_{i}^{(j)} = \sigma^{2} \left[\left(\frac{\partial F_{i}^{(j)}}{\partial X_{i}^{(j)}} \right)^{2} + \left(\frac{\partial F_{i}^{(j)}}{\partial Y_{i}^{(j)}} \right)^{2} \right]$$
(A5)

The superscript (j) designates the value at the j-th iteration. Also a matrix C(u,v) and a function E_i defined below are obtained.

$$C(\mathbf{u},\mathbf{v}) = \sum_{i} \frac{\left(\frac{\partial F_{i}^{(j)}}{\partial B_{u}^{(j)}}\right) \left(\frac{\partial F_{i}^{(j)}}{\partial B_{v}^{(j)}}\right)}{L_{i}^{(j)}}$$
(A6)

$$E_{i}^{(j)} = \frac{F_{i}^{(j)} + \frac{\partial F_{i}^{(j)}}{\partial X_{i}^{(j)}} (X_{i}^{meas} - X_{i}^{(j)}) + \frac{\partial F_{i}^{(j)}}{\partial Y_{i}^{(j)}} (Y_{i}^{meas} - Y_{i}^{(j)})}{L_{i}^{(j)}}$$
(A7)

Then true parameters for B_i values, b_i (i = 1...7), are obtained from the following simultaneous linear equation.

 $\left| \sum_{i} \left(\partial F_{i}^{(j)} \right)_{F_{i}^{(j)}} \right|$

After solving Equation A8, the solid and fluid compositions are computed by the following equations.

$$X_{i}^{(j)} - x_{i} = X_{i}^{(j)} - X_{i}^{\text{meas}} + \lambda_{i}^{(j)} \sigma^{2} \left(\frac{\partial F_{i}^{(j)}}{\partial X_{i}^{(j)}} \right)$$
(A9)

$$Y_{i}^{(j)} - y_{i} = Y_{i}^{(j)} - Y_{i}^{meas} + \lambda_{i}^{(j)} \sigma^{2} \left(\frac{\partial F_{i}^{(j)}}{\partial Y_{i}^{(j)}} \right)$$
(A10)

where

After b_i , x_i , and y_i values are obtained, those are substituted into Equations A1–A3 and A5–A7 as B_i , X_i , and Y_i . Then the newly computed F_i , E_i , and C(u,v) are substituted into Equations A8–A11. Successive iterations are repeated until all of the unknowns and the compositions of solid and fluid phases are converged. The following convergence criteria were adopted as follows:

 $|B_{u}^{(j+1)} - B_{u}^{(j)}| \le 0.00005 \ (u = 1 \dots 7), \ |X_{i}^{(j+1)} - X_{u}^{(j)}| \le 0.0001, \ |Y_{i}^{(j+1)} - Y_{i}^{(j)}| \le 0.0001.$

$$\lambda_{i}^{(j)} = \frac{F_{i}^{(j)} + \left(\frac{\partial F_{i}^{(j)}}{\partial X_{i}^{(j)}}\right) \left(X_{i}^{\text{meas}} - X_{i}^{(j)}\right) + \left(\frac{\partial F_{i}^{(j)}}{\partial Y_{i}^{(j)}}\right) \left(Y_{i}^{\text{meas}} - Y_{i}^{(j)}\right)}{L_{i}^{(j)}} - \frac{\left(\frac{\partial F_{i}^{(j)}}{\partial B_{i}^{(j)}}\right) \left(B_{2}^{\text{meas}} - b_{2}\right) + \left(\frac{\partial F_{i}^{(j)}}{\partial B_{3}^{(j)}}\right) \left(B_{3}^{\text{meas}} - b_{3}\right) + \left(\frac{\partial F_{i}^{(j)}}{\partial B_{4}^{(j)}}\right) \left(B_{4}^{\text{meas}} - b_{4}\right)}{L_{i}^{(j)}} \left(B_{5}^{\text{meas}} - b_{5}\right) + \left(\frac{\partial F_{i}^{(j)}}{\partial B_{5}^{(j)}}\right) \left(B_{7}^{\text{meas}} - b_{7}\right)}{L_{i}^{(j)}} \left(A11\right)$$