1	
Ŧ	

REVISION 2

2	A revisit to phase transition behavior of K-feldspar at
3	high-pressure and high-temperature: Implications on metastable
4	K-feldspar in cold subduction
5	Chengcheng He ^{1,2} · Heping Li ^{1*}
6	1 Key Laboratory of High-temperature and High-pressure Study of the Earth's
7	Interior, Institute of Geochemistry, Chinese Academy of Sciences, Guiyang, 550081,
8	China
9	2 University of Chinese Academy of Sciences, Beijing, 100049, China
10	* Corresponding author: <u>liheping@mail.gyig.ac.cn</u>
11	
12	
13	
14	
15	
16	
17	
18	
19	
20	

21

Abstract

22	Felsic gneiss is main type of ultrahigh-pressure metamorphic rocks in continental
23	subduction zones. As one of the important mineral of felsic gneiss, the phase behavior
24	of K-feldspar is of great significant to study subducting seismic activities,
25	crust-mantle interactions and plate tectonics in the Earth's interior. Thus, the phase
26	relations in K-feldspar are constrained by Raman spectroscopy combined with
27	externally-heating diamond anvil cell (DAC) under high-pressure (28 GPa),
28	high-temperature (400 °C), and simultaneous high pressure-temperature (P-T) (14
29	GPa, 430 °C) conditions. The phase diagram of K-feldspar at relatively lower
30	temperature has been complemented, and a more comprehensive P - T phase diagram
31	of K-feldspar is established along with previous results. K-feldspar undergoes a
32	reversible phase transition from $C-1$ to $P-1$ (metastable K-feldspar) symmetry at 10.3
33	GPa. Metastable K-feldspar is an intermediate phase from K-feldspar to K-holl-I
34	(KAlSi ₃ O ₈ with a hollandite-I structure) in extremely cold subduction slabs ($<2^{\circ}C/km$)
35	or the center of some old, cold, and rapidly subducting slabs. Metastable K-feldspar is
36	stable even at 11.4 GPa and 400°C in enriched hydroxyl group (OH) environments,
37	which shows greater high P - T stability than K-feldspar. So metastable K-feldspar
38	could withstand subduction-zone fluids low-temperature metamorphism and survive
39	to deeper than previously expected. This study enhances our understanding of the
40	formation paths and conditions of K-holl-I, the subducted depth of K-feldspar, the
41	effect of subduction-zone fluids on continental subduction, and provides a possible
42	reason for the origin of intermediate- and deep-focused seismicity.

- 43 Keywords: metastable K-feldspar, Raman spectroscopy, high pressure-temperature
- 44 conditions, continental subduction, subduction-zone fluids

46

Introduction

47	Feldspar minerals comprise about 50-60% of the Earth's crust and are widely
48	found in the surface of other planets (e.g., Moon, Mars, Venus) (Sprague et al. 1997;
49	Rogers and Nekvasil 2015; Hashimoto et al. 2009). K-feldspar (microcline, sanidine,
50	orthoclase; $KAlSi_3O_8$) is more stable than other two principal rock-forming
51	end-members (albite, anorthite), and has been regarded as a predominant potassic
52	phase in the Earth's crust and the uppermost mantle (e.g., Liu et al. 2010; Deng et al.
53	2011). In addition, K-feldspar could enter the deep Earth via subduction process,
54	crustal constituents are directly transferred from the descending slab into the mantle
55	wedge at different depths, leading to the crust-mantle interaction (Irifune et al. 1994;
56	Zhang et al. 2003; Zheng 2012). Furthermore, K-holl-I is regarded as a host phase for
57	large cations such as K, Na, Sr, Ba, and Pb in the mantle, which is of great
58	geochemical interest (e.g., Yamada et al. 1984). Thus, the stability of K-feldspar is of
59	great geophysical and geochemical significance to understand the dynamics of
60	subducting lithospheres, Earth's long-term thermal evolution, and the material
61	exchange of the crust-mantle system.

The framework of K-feldspar is based on the rings of four TO₄ tetrahedra (T = Si^{4+} , Al^{3+}) forming double crankshaft-like chains. All Si and Al atoms are bonded to four O atoms in the TO₄ tetrahedra, and potassium atoms occupy the large voids. K-feldspar dives into the Earth interior along with a series of structural and compositional changes. Phase equilibrium experiments confirmed that at temperatures above 700 °C, K-feldspar dissociated into a denser assemblage of K₂Si₄O₉ wadeite +

68	kyanite + coesite/stishovite (Wd + Ky + Coe/St) at 6-7 GPa, and latter the assemblage
69	recombined into K-holl-I at 9-10 GPa (e.g., Urakawa et al. 1994; Yagi et al. 1994; Liu
70	et al. 2010). K-holl-I transformed into K-holl-II (KAlSi $_3O_8$ with a hollandite-II
71	structure) at near 20 GPa, both ambient and elevated temperatures. K-holl-II is
72	unquenchable and stable over the entire range of mantle conditions (Sueda et al. 2004;
73	Nishiyama et al. 2005; Ferroir et al. 2006; Hirao et al. 2008); In addition, K-feldspar
74	melted to leucite (KA1Si ₂ 0 ₆) and silica-rich liquid at 2 GPa, 1150 °C (Lindsley 1966),
75	and transformed into metastable phase induced by an increase in the Al coordination
76	number at 10-13 GPa and ambient temperature (Pakhomova et al. 2020).
77	Metastability of subducted minerals could encourage subducting slabs
78	deceleration or stagnation (e.g., Van Mierlo et al. 2013; Agrusta et al. 2014). It is well
79	known that low-temperature parts in the Earth's subduction zones enable the
80	preservation of metastable phases at greater depths than those under equilibrium
81	conditions (e.g., Hogrefe et al. 1994; Kawakatsu and Yoshioka 2011). However,
82	previous high $P-T$ studies of K-feldspar are insufficient to be applied to in the
83	ultracold subducted environments, and the lower temperature parts (below 700 °C) of
84	the phase diagram has not be carefully investigated yet. Besides, Pakhomova et al.
85	(2020) studied the stability of metastable albite and anorthite by quenching
86	experiments, and proposed anorthite may exist at deeper depths than expected in the
87	Earth's interior. As one of the important mineral of felsic gneiss, the high P - T stability
88	of metastable K-feldspar is also worthy of attention. Herein, the phase relations at
89	relatively low temperature of K-feldspar are investigated firstly by Raman

90	spectroscopy in externally-heating DAC under high-pressure (28 GPa),
91	high-temperature (400 °C), and simultaneous high P-T (14 GPa, 430 °C) conditions.
92	We also presented the isothermal (γ_{iT}) mode Grüneisen parameters, isobaric (γ_{iP}) mode
93	Grüneisen parameters, and intrinsic anharmonic mode parameters of K-feldspar.
94	Besides, a series of ex situ multi-anvil experiments were performed to clarify the high
95	P-T stability of metastable K-feldspar. This study complements the gap in the phase
96	diagram of K-feldspar at relatively low temperature, and a more comprehensive P-T
97	phase diagram of K-feldspar is established along with previous studies. The results
98	provide new insights into the phase behavior of K-feldspar in subduction zone at
99	different thermal gradients (especially at the ultracold), and the possible seismicity
100	caused by their intruding plates (Zheng 2019).

- 101
- 102

Material and Methods

103 Sample synthesis and characterization

High-purity K-feldspar was synthesized by solid-state reaction. Reagent-grade 104 $KHCO_3$, Al_2O_3 and SiO_2 powders were mixed in a molar 2:1:6 ratio. The powder 105 mixtures were crushed with ethanol in an agate mortar for 1 h, and then placed in a Pt 106 crucible and heated up to 800 °C in a stove for 12 h to ensure a sufficient 107 decarbonation. The recovered mixture was ground finely and then pressed into pellets 108 109 with diameter of 6 mm and length of 3 mm under uniaxial pressure of 4 MPa. High *P-T* synthesis experiment was performed on DS 6×600 t cubic-anvil-type apparatus 110 at the Key Laboratory of High-Temperature and High-Pressure Study of the Earth's 111

112	Interior, Institute of Geochemistry, CAS, Guiyang, China. Figure 1 (a) shows a
113	schematic of the cross section of the experimental assemblage. A pellet was covered
114	by the silver foil of 0.025 mm thickness, pyrophyllite cube (32.5 mm \times 32.5 mm \times
115	32.5 mm) and h-BN were applied as the pressure medium, and a graphite was used as
116	the heater at 1.8 GPa, 900 °C for 2 h. The pressure was reduced to ambient conditions
117	after quenching. The reactant was confirmed as a K-feldspar single phase by both
118	powder X-ray and Raman spectrum, as Figures 1 (b-c) shown.

119

120 In situ high-pressure and high-temperature Raman experiments

121 The Raman spectra of synthetic K-feldspar under high-pressure (up to 28 GPa) and room-temperature, high-temperature (up to 400 °C) and ambient-pressure, and 122 simultaneous high P-T conditions (up to 14 GPa, 430 °C) were collected in DAC. 123 The sample chambers were firstly compressed up to target pressures, and then heated. 124 High-pressure conditions were generated by a pair of 400 mm diamond culets. The 125 126 250-µm-thick rhenium gaskets were pre-intended to dense disks of ~60 µm thickness, and then drilled a hole as the sample chamber with a diameter of $\sim 160 \ \mu m$. The 4:1 127 mixture of methanol-ethanol (ME), 16:3:1 mixture of methanol-ethanol-water (MEW), 128 helium, nitrogen or ultrapure water (H₂O) were selected as pressure-transmitting 129 130 mediums (PTMs). The gaseous PTMs were loaded by a high-pressure gas-loading 131 apparatus sealed at above 0.4 GPa. The gas-loading apparatus is developed in the 132 Sanchez Technologies and located at the Key Laboratory of High-Temperature and High-Pressure Study of the Earth's Interior, Institute of Geochemistry, CAS, Guiyang, 133

134	China. High-temperature environments were achieved by two external electric NiCr
135	resistance heating furnace in series around the diamonds. Temperatures were
136	measured by a R/S type platinum-rhodium and platinum thermocouple (Omega
137	technology company, America) with the precision of ± 1 °C in our experimental
138	temperature range. The thermocouple was used for temperature measurement and
139	stabilization through a feedback loop in the temperature controller (CN3251,Omega
140	technology company). Every time before the Raman spectrum was collected, we kept
141	the temperature at least 5 min to guarantee the stabilization and uniformity of
142	temperature in the sample chamber. The R1 line of tiny ruby (Cr ³⁺ doped α -Al ₂ O ₃)
143	spheres and 0-0 line of $SrB_4O_7:Sm^{2+}$ grains (synthetic method was described in
144	previous study) (Mao et al. 1986; Zhao et al. 2017) were selected to calibrate pressure
145	at high-pressure and high $P-T$ experiments, respectively. Compared to Ruby, the 0-0
146	line of SrB ₄ O ₇ :Sm ²⁺ presents very small temperature response, including frequency
147	shift and broadening. So it is the preferred pressure calibration under high $P-T$
148	conditions. The spectra of the calibrator were measured at each pressure before and
149	after sample data collection, and the average pressure values were taken. The Raman
150	spectra were collected at a confocal Renishaw in Via Raman Microscope in the Key
151	Laboratory of High-Temperature and High-Pressure Laboratory Institute of
152	Geochemistry, Chinese Academy of Sciences, Guiyang. The tilting and rotation of the
153	spectrometer relative to the incident laser beam were calibrated using a silicon wafer
154	with a static spectrum centered at 520 cm ⁻¹ . Samples were excited by an Argon ion
155	laser ($\lambda = 514.5$ nm) operated at 20 mW with a focused laser spot of ~5 μ m in

diameter through a SLM Plan 50 \times Olympus microscope objective. A spectrometer with a liquid nitrogen-cooled CCD detector was used to collect the Raman data. The accurate peak positions of the Raman spectra were fitted by *PeakFit* software.

159

160 Multi-anvil experiments

The high P-T conditions of quenching experiments were provided using a 161 1000-ton Kawai-type multi-anvil apparatus at the Key Laboratory 162 of High-Temperature and High-Pressure Study of the Earth's Interior, Institute of 163 164 Geochemistry, CAS, Guiyang, China. The sample assembly was displayed in Figure 165 S1. Eight tungsten carbide cubes with a truncation of 6 mm were used as second-stage anvils, and a Cr₂O₃-doped MgO octahedron with edges of length of 14 mm was used 166 as pressure medium. Starting materials were the high-purity K-feldspar synthesized in 167 this study. Graphite sleeve, platinum closure column, MgO tube, ZrO₂ tube and Al₂O₃ 168 plug were used as the heater, the sample capsule, the insulator, and the thermal 169 170 insulation sleeve, respectively. The temperature was measured using a 171 $W_{95}Re_5-W_{74}Re_{26}$ thermocouple inserted axially under the sample capsule. To remove absorbed water, all parts composed of MgO, Al₂O₃, or ZrO₂ were baked at 1000 °C 172 for 1 h before the experiments. The sample chambers of three quench experiments 173 were firstly compressed up to 15 GPa, and then heated to 500 °C, 650 °C, and 800 °C 174 175 at a rate of 20 °C/min, respectively. The duration time at each target temperature was at least 2 h before rapid quenching. 176

178

Results and Discussion

179	Raman spectrum of K-feldspar at ambient conditions is presented in Figure 1 (b).
180	12 modes of the 39 Raman modes $(20A_g + 19B_g)$ predicted by factor group analysis
181	(Tseng et al. 1995) were observed. The strongest triplet peaks at 451.98, 475.44, and
182	512.98 cm ⁻¹ (denoted as I_a , I_b , I_c modes) are ascribed to the ring-breathing of the
183	four-membered rings of tetrahedra. Three Raman peaks at 117.23, 156.34, and 195.19
184	(expressed as III_a , III_b , III_c modes) are assigned to cage-shear. Two peaks at 282.91
185	and 403.21 cm ⁻¹ (named as II_a , II_b modes) are due to the rotation-translation of the
186	four-membered rings. Two weak modes at 758.76, 803.59 cm ⁻¹ (marked as IV_a and IV_b
187	modes) are attributed to the deformation of the tetrahedra. The broad peaks with
188	position at around 1023.44 and 1118.38 cm ⁻¹ (labeled as V_a and V_b) belong to the
189	vibrational stretching modes of the tetrahedra. Most Raman modes coincide well with
190	previous studies, except for the modes of IV_a , IV_b , V_a and V_b due to the weak intensity
191	(Freeman et al. 2008; Bendel and Schmidt 2008; Peterson et al. 2009; Befus et al.
192	2018). The slight differences in wavenumber for observed modes may result from the
193	synthetic environment, peak overlap, and orientation dependence.

194

195 Raman spectroscopy of K-feldspar at high-pressure/high-temperature

High-pressure Raman spectra of K-feldspar were collected up to 28 GPa at room temperature, and were only studied in the spectral range 100-800 cm⁻¹ due to the intense Raman bands of MEW at 880-1200 cm⁻¹. The typical high-pressure Raman spectra are showed in Figure 2 (a). At high pressure, a general deterioration of the

200	Raman signal and amorphization effect can be observed. I_b , I_c , II_a , III_b modes can be
201	distinguished clearly within 10 GPa, while only the strongest mode I_c can be observed
202	clearly within the entire studied pressure range. All the Raman modes shift to higher
203	wavenumbers with increasing pressure because of the TO ₄ tetrahedra shrinkage and
204	the stronger electrostatic repulsion. The corresponding pressure shifts and the full
205	width at half maximum (<i>FWHM</i>) of I_c mode are plotted in Figures 2 (b-c), from which
206	two obvious inflection points at 6.5 GPa and 14.7 GPa can be observed, and another
207	obvious inflection point at 10.3 GPa can been seen in Figure 2 (c). The pressure
208	coefficient $(\partial v_i/\partial P)$ changes from 3.72(2) to 3.07(9) cm ⁻¹ GPa ⁻¹ at the first
209	discontinuity point (at 6.5 GPa). However, X-ray diffraction (XRD) studies reported
210	microcline experienced a continuous compression up to 7 GPa, then underwent an
211	elastic softening with an increase of the <i>a</i> -axis and a decrease of the <i>b</i> -axis at 7-9 GPa;
212	meanwhile, the α and β unit-cell angles rapidly increased while the γ unit-cell angle
213	decreased at pressures above 6.7 GPa (Allan and Angel 1997; Pakhomova et al. 2020).
214	Thus, the first discontinuity point at 6.5 GPa could not be ascribed to a phase
215	transition but the change in compressional mechanism at atomic scale, which result in
216	the elastic behaviour change. Similar discontinuity was also found in the isostructural
217	mineral, low albite, at 6.5 GPa (Aliatis et al. 2017). The wavenumbers of I_c mode
218	change linearly upon further compression from 6.1 to 14.2 GPa, wheras the FWHM
219	shows an abrupt change at ~10.3 GPa. Similar discontinuity was also found in low
220	albite at 8.5 GPa with the $\partial v_i / \partial P$ of the mode at 815 cm ⁻¹ (corresponding to IV_b modes)
221	showing a discontinuity (Aliatis et al. 2017). Pakhomova et al. (2020) reported a

222	displacive phase transition of K-feldspar from C-1 to P-1 (metastable phase)
223	symmetry, which was induced by an increase in the Al coordination number from four
224	to six upon 10.4-12.8 GPa. Thus, we infer the second discontinuity at 10.3 GPa is
225	relate to the phase transition from K-feldspar to metastable phase. The accurate phase
226	transition range is 9.7 to 10.8 GPa. An additional Raman experiment up to 13 GPa
227	demonstrated the reversibility of the phase transition (Figure S2 (a)). The third
228	discontinuity at 14.7 GPa is relate to the solidification of MEW which caused an
229	irreversible pressure-induced amorphization. The high-pressure behaviour of feldspars
230	is highly sensitive to deviatoric stress. Previous study reported metastable K-feldspar
231	was stable at least up to 27 GPa in hydrostatic conditions (Pakhomova et al. 2020).
232	While in this study, the $FWHM$ of I_c mode increased and the intensity decreased
233	rapidly at above 14.7 GPa. Upon further compression at above 23.1 GPa, sample was
234	fully amorphous and remained amorphous structure on recovery to ambient conditions.
235	Another high-pressure Raman experiment in H ₂ O up to 20.4 GPa indicated the
236	reversibility of the not fully amorphization (Figure S2 (b)). Previous results about the
237	pressure-induced phase transition of feldspars at room temperature reported a
238	displacive phase transition occurred for both anorthite and oligoclase at 2.2-3.9 GPa,
239	accompanied by the occurrence of a new peak near the Ic mode, which was mainly
240	caused by the reduction of Ca-O bond length (Xie et al. 2012; Pakhomova et al. 2020).
241	However, the phase transition behavior was not found in K-feldspar and albite during
242	the compressing processes (Aliatis et al. 2017; Pakhomova et al. 2020). We infer the
243	phase behavior at relative low pressure of feldspars is significantly affected by the

(Ca,Na,K)-O bond length and bond energy. Moreover, the three end-members all
experienced a coordination number increase of Al atoms to further adapt to the
high-pressure conditions (Pakhomova et al. 2020).

High-temperature Raman spectra of K-feldspar were collected up to 400 °C at 247 248 ambient pressure. The typical high-temperature Raman spectra are showed in Figure 3 (a), and the variation of Raman shift at different temperatures is plotted in Figure 3 (b). 249 The Raman shifts of all observed bands for K-feldspar continuously decrease with 250 increasing temperature, and no phase transition was observed during heating in this 251 252 study, which was consistent with previous studies (Henderson 1979; Hovis et al. 253 2010). Generally, the Raman peaks shift slightly to lower frequency regions with increasing temperature because of the longer bond length and weaker bond. Changes 254 of FWHM and signal intensity of Raman peaks under high-temperature conditions are 255 not great as that under high-pressure conditions (i.e., increase of *FWHM* is \sim 4.7 cm⁻¹ 256 for I_c mode in the investigated temperature range). 257

Combined the Raman spectra of K-feldspar under high-pressure and 258 high-temperature conditions, the pressure coefficients $(\partial v_i/\partial P)$ and the temperature 259 coefficients $(\partial v_i / \partial T)$ of observed Raman bands are listed in Table 1 and Figure S3. 260 261 The Raman modes could be seen approximately as continuous increases within ~ 10.3 GPa. The pressure coefficients of Raman vibrations of K-feldspar vary from 1.41 to 262 5.32 cm⁻¹GPa⁻¹, and the temperature coefficients vary from -0.17×10^{-2} to -0.82×10^{-2} 263 $cm^{-1}K^{-1}$. The modes relate to the rotations of the four non-equivalent tetrahedra and 264 the translation of K^+ cations (e.g., III_a , III_b , Aliatis et al. 2015) show the biggest 265

266 change in slope with pressure and temperature. The isothermal (γ_{iT}) and isobaric (γ_{iP}) mode Grüneisen parameters are important physics quantities to describe the elasticity 267 268 and anharmonicity under high P-T conditions, which can be calculated by the following equations and are listed in Table 1 (Gillet et al. 1989): 269 $\gamma_{iT} = (K_T / V_{i0}) (\partial V_i / \partial P))_T$ 270 (1) $\gamma_{iP} = -(1/v_{i0}\alpha)(\partial v_i/\partial T)_P$ (2)271 where v_i is the Raman shift of the *i*-th vibrational mode, v_{i0} is the frequency of mode v_i 272 273 at 0 GPa, K_T is the isothermal bulk modulus with value of 58.3 GPa (Allan and Angel 1997), α is the volume thermal expansion coefficient with value of $1.44 \times 10^{-5} \text{ K}^{-1}$ 274 (Henderson, 1979), defined as $\alpha = (1/V)(\partial V/\partial T)_P$. The calculated values of γ_{iT} 275 range from 0.26 to 1.96, and γ_{iP} change from 0.41 to 4.51. The average γ_{iT} value of 276 Raman modes of tetrahedra (0.32) are much smaller than that of the tetrahedral cages 277 (1.96), and the thermal expansion differs more with a smaller average γ_{iP} value of 278 tetrahedra (0.61) and a larger average γ_{iP} value of the tetrahedral cages (3.42). The 279 average γ_{iT} and γ_{iP} values are small for the TO₄ tetrahedra, which reflects their relative 280 incompressibility and weak expansivity. Hence, the weaker bonds (K-O) are affected 281 282 more by increasing pressure and temperature. It is obvious that the tetrahedra are 283 much more incompressible than the tetrahedral cages, which is consistent with other feldspars (e.g., oligoclase from Xie et al. 2012). The γ_{iP} are different with the γ_{iT} for 284 285 same mode, which means an intrinsic anharmonicity. The intrinsic anharmonic mode 286 parameter βi can be calculated by the following equation and are listed in Table 1: 287

288

 $\beta_i = \alpha(\gamma_{iT} - \gamma_{iP})$ (3)

the obtained β_i change from -0.46×10^5 to -0.16×10^{-5} K⁻¹, with an average of 289 $-0.26 \times 10^{-5} \text{ K}^{-1}$.

290

291 Raman spectroscopy of K-feldspar at simultaneous high *P-T* conditions

Raman spectra of K-feldspar in different PTMs were collected at pressures up to 292 14.7 GPa and temperatures up to 430 °C (Figure S4). During heating at a certain 293 294 initial pressure, the real pressure of sample chamber changed because of thermal 295 expansion and stress relief of experimental materials. Thus, we measured the pressure of sample chamber by the fluorescence spectra of SrB₄O₇:Sm²⁺ in real-time. The 296 representative Raman spectra of K-feldspar obtained under simultaneous high P-T 297 conditions are presented in Figure 4. Given the 4:1 ME mixture can conveniently 298 provide a hydrostatic condition at below 10.5 GPa (Klotz et al. 2009), and has 299 300 relatively simple ingredients, it was firstly selected as the PTM. At first initial pressure points of 1.8 GPa, with temperature increasing, no any meaningful variation 301 302 except pressure- and temperature-induced frequency shifts were observed (Figure S4 303 (a)). With further compressing and heating, the peak shape changed firstly at 400 $^{\circ}$ C and 5.9 GPa with a new peak appearing at \sim 430 cm⁻¹, and the Raman shift of mode at 304 ~ 290 cm⁻¹ clearly showed a discontinuous evolution (Figure S4 (b)). Then the 305 306 variation became apparent at 400 °C and 6.7 GPa with several new peaks at 206, 215, 357, 392, 430, and 465 cm⁻¹, while some new peaks disappeared after further heating 307 (Figure S4 (c)). We even observed some bubbles welling up in the sample chamber, 308 309 but it was hard to record and measure because of its quick breaking. Similar Raman 310 spectra variation were also observed at P-T range of 350-430 °C and 7.3-8.1 GPa (Figure S4 (d)), while surprisingly did not appear at 11.4-13.9 GPa (Figure S4 (e)). 311

312	The new peaks revealed at least one reaction occurred. Due to the signal-to-noise ratio
313	deterioration and the Raman peaks broadening under high P-T environments, it was
314	difficult to identify the new phase via the ambiguous peaks. Luckily, the new peaks
315	sharpened and the shape of the main strong peaks remain unchanged during cooling
316	and decompressing periods (Figure S5), which made the recovered sample easy to
317	distinguish, and could be referred to the high P-T new phase. Raman spectra of
318	samples recovered from high P - T to ambient conditions in different PTMs are shown
319	in Figure 5. The Raman spectrum of recovered sample from 430 °C, 6.7 GPa (ME) in
320	Figure 5 is virtually same as that in Figures S4 (b-d), and it is consistent well with the
321	characteristic peaks of an assemblage of coesite, stishovite and KHCO ₃ (Figure S6
322	(a)). Besides, we accidentally conducted another experiment at 4.7-5.9 GPa,
323	20-400 °C in the 16:3:1 MEW mixtures (Figure S4 (f)), and the Raman spectrum of
324	recovered sample can be indexed to coesite and muscovite (Figure S6 (b)). However,
325	because of the appearance of coesite, and the reaction did not occurred in the $P-T$
326	range of 1.8-3.2, 11.4-13.9 GPa and 20-400 °C (Figures S4 (a) and S4 (e)), it was
327	difficult to confirm whether the ME/MEW mixtures reacted with K-feldspar or with
328	its phase transition products, such as wadeite-type $K_2Si_4O_9$ (Wd). In order to further
329	study the stability of K-feldspar under high <i>P-T</i> conditions, we replaced the PTM. We
330	conducted two high P-T experiments in helium and in no PTM, respectively (Figures
331	S4 (g-h)), and found K-feldspar was stable at the experimental <i>P</i> - <i>T</i> range.
332	All up, pure K-feldspar is stable and does not decompose to an assemblage of Wd

333 + Ky + Coe/St within the P-T range (0-11.4 GPa, 20-430 $^{\circ}$ C) in this study.

334	Nonetheless, if fluid components added, the stability of K-feldspar is different. The
335	results of in situ Raman experiments under high P-T conditions in ME/MEW show:
336	K-feldspar (low Al coordination number) could be stable only below 4.5 GPa, 400 $^\circ$ C
337	in ME/MEW fluids. Once exceeding this limit, it transforms into the assemblage
338	containing coesite, then with pressure increasing, the transformation temperature
339	gradually drops to 350 °C (8.1 GPa). The phase transition of K-feldspar to the
340	assemblage containing coesite is irreversible and has negative Clapeyron slope.
341	However, after the structural transition of K-feldspar into metastable structure (high
342	Al coordination number) at above 10.3 GPa, metastable K-feldspar could remain
343	stable even at 11.4 GPa and 400 °C, showing a higher structural stability. Subduction
344	zone fluids greatly affect the mineral transformation, rock melting, element mobility,
345	and water (H2O) is an important fluid of subduction-zone rocks. For verifying
346	whether the higher structural stability of metastable K-feldspar still remain in the deep
347	Earth fluids, we conducted another two high P-T Raman experiments using H ₂ O as
348	the PTM. In H ₂ O, K-feldspar transformed into the assemblage containing coesite at
349	6.9 GPa and 350 °C (Figure S4 (i), Figure S6 (c)), while metastable K-feldspar could
350	remain stable at up to 10.4 GPa and 430 °C (Figure S4 (j)), almost same as in
351	ME/MEW. The decomposition of metastable K-feldspar into coesite mixture after
352	keep in 430 °C for 1 h may be ascribed to the pressure change. The pressure of sample
353	chamber decreased to 9.8 GPa because of thermal expansion and stress relief after 1 h,
354	so the metastable phase transformed back to K-feldspar, then K-feldspar was not
355	stable and decomposed. Therefore, metastable K-feldspar may be harder to undergo

356	metasomatism with subduction-zone fluids and can survive to deeper than previously
357	expected. In addition, the coesite mixture remain stable at 430 °C, 9.8-10.4 GPa after
358	50 min (Figure S7), which indicates the decomposition of K-feldspar in fluid
359	composition could occur at 4.5 GPa, 400 °C and maintained at a certain $P-T$ range
360	before the transition from K-feldspar to metastable phase. It is hard to identify
361	unambiguously the equilibrium phase(s) due to the kinetics of the phase transition.
362	But the stable existence of coesite with pressure and temperature increasing indicate
363	that the assemblage containing coesite is more stable than K-feldspar with low Al
364	coordination number at above 4.5 GPa, 400 °C.

365

366 Pressure-temperature phase diagram of K-feldspar

The phase diagram of K-feldspar under high P-T conditions constrained by 367 experimental data is shown in Figure 6 along with the literature data (Lindsley 1966; 368 Zhang et al. 1993; Urakawa et al. 1994; Yagi et al. 1994; Nishiyama et al. 2005; 369 Akaogi et al. 2004; Ferroir et al. 2006; Yong et al. 2006; Hirao et al. 2008; Chen et al. 370 2019; Guo et al. 2020; Pakhomova et al. 2020). The black solid lines represent the 371 phase boundaries determined by high P-T experiments or from previous phase 372 diagram, the dotted lines indicate the extrapolated phase boundaries by limited data. 373 We make following changes comparing to previous studies of the phase relations of 374 375 K-feldspar (Urakawa et al. 1994; Sueda et al. 2004; Nishiyama et al. 2005; Yong et al. 376 2006; Hirao et al. 2008; Liu et al. 2010): First of all, this study complements the blank of previous data under relatively low P-T conditions, and get more accurate P-T phase 377

378	boundaries between K-feldspar and Wd + Ky + Coe/St by Raman spectroscopy.
379	Secondly, previous phase diagrams revealed that $KAlSi_3O_8$ with a hollandite-I and -II
380	structures (K-holl) was stable under relatively high-pressure and low-temperature
381	conditions. However, former studies referred to the experimental data of e.g.,
382	Nishiyama et al. (2005), in which they used K-holl-I as starting sample rather than
383	K-feldspar. Since K-holl-I phase is quenchable, it can be detected under high-pressure,
384	low-temperature and even ambient conditions. Combined high-pressure stability of
385	K-feldspar in this study, the former data are not reasonable to interpret the phase
386	behavior of K-feldspar at this $P-T$ area in which K-holl phase has not formed yet. To
387	solve this problem, three quench experiments of K-feldspar in multi-anvil apparatus
388	were performed at 15 GPa, 500 °C, 15 GPa, 650 °C, and 15 GPa, 800 °C, respectively.
389	Raman analysis on the three quenching samples revealed the existence of K-feldspar,
390	K-feldspar + K-holl-I, and K-holl-I (with all peaks indexed well to that in Chen et al.
391	2019), respectively (Figure 7). For sufficient phase equilibrium, the quench
392	experiments at 15 GPa, 500 °C was kept for 12 h before rapid quenching, and the
393	characterized result at 15 GPa, 500 °C indicates K-feldspar is a stable phase at below
394	15 GPa, 500 °C. The coexistence of K-feldspar and K-holl-I at 15 GPa, 650 °C shows
395	metastable K-feldspar is the intermediate phase between K-feldspar and K-holl-I
396	under relatively high-pressure and low-temperature conditions. Thus, on the basis of
397	previous phase diagram, we re-predicted the phase boundary between K-feldspar and
398	K-holl by the results of three quenching experiments in this study, and only reserved
399	the data of Nishiyama et al. 2005 at above the phase boundary. Finally, the phase

400 boundary between K-holl-I and K-holl-II phase was re-constrained by the data of
401 Hirao et al. (2008) in which they used K-feldspar as the starting sample.

402

403

Implications

404	The phase diagram of K-feldspar in the ultracold and dry subducting slab, as
405	shown in Figure 8 (a), suggests that the pressure-induced phase transition of
406	K-feldspar at relatively low-temperature (< 650 °C) could be different from those
407	occurring at relatively high-temperature (> 650 °C). Instead of the K-feldspar→liquid
408	/ liquid + leucite in the normal or warmer mantle (> 5 °C/km), or K-feldspar \rightarrow Ky +
409	Wd + Coe/St in the ultracold subduction environments (~5 $^{\circ}C/km$) (Zheng and
410	Hermann 2014; Zheng 2019), K-feldspar firstly transforms into metastable phase at
411	above 10.3 GPa at relatively low-temperature. The metastable phase transition could
412	occur in extremely cold plates (<2 °C/km, like Tonga proposed by Syracuse et al.
413	2010) at depths of \sim 300 km. Additionally, because the slab center is much colder than
414	the margin in some cold, old, fast, and steep slabs (Ranalli et al. 2000; Kubo et al.
415	2009; Xu et al. 2020), metastable feldspar is also likely to preserve within the center
416	of some ultracold, old, and rapidly subducting slabs. The density is similar between
417	metastable K-feldspar and the mantle at the phase transition depth (~300 km)
418	(Pakhomova et al. 2020). Thus, metastable K-feldspar has possible influence on the
419	dynamics of subducted continental lithosphere, that is, affects buoyancy, hampers
420	further subduction, and leads to a deceleration and even a potential stagnation of slabs
421	at ~300 km, just like the effects of metastable pyroxene, olivine, and plagioclase on

422 the slab dynamics (e.g., Hogrefe et al. 1994; Kawakatsu and Yoshioka 2011; Agrusta

423 et al. 2014; Pakhomova et al. 2020).

The high P-T Raman experimental results have potential implications for the 424 effects of subduction-zone fluids on K-feldspar. Seismic studies have shown that a 425 amounts of water could be transported along the subducted slabs deep into the mantle 426 (Kawakatsu and Watada 2010). Although most of the water is released from the 427 oceanic crust at depths of 150-200 km (Iwamori 2007), under cold subduction 428 conditions, water could be stored in hydrous minerals, or enter deeper mantle as 429 430 nominally anhydrous minerals containing hydroxyl group (OH) (Wang et al. 2020). 431 Our study indicates that K-feldspar with high Al coordination number is relatively more stable than that with low coordination number in the deep Earth enriched OH 432 fluids (such as water). The phase relations of K-feldspar in the ultracold and hydrous 433 subducting slab are shown in Figure 8 (b). At temperature above 350-400 °C 434 subducting environments, K-feldspar with low coordination structure cannot resist 435 fluid metamorphism and rapidly decomposes into the coesite-containing assemblage 436 437 at depth above 135 km. An obvious dispersion of sample in ME/MEW/H₂O after the decomposition was observed in the experiment, so we infer low-temperature fluid 438 metamorphism of K-feldspar may potentially affect continental slab breakoff and 439 seismic activities. Besides, quartz is highly soluble in metamorphic fluids, so the 440 assemblage containing coesite may be difficult to deeper subduction in metamorphic 441 environments, which inhibits its further phase transition to form K-holl-I. Thus, the 442 well-known phase transitions of K-feldspar \rightarrow Ky + Wd + Coe/St \rightarrow K-holl-I maybe 443

only occur within dry and cold slabs. However, in the subduction zones where 444 metastable phases can be preserved, K-feldspar firstly transforms into metastable 445 phase which has higher structural stability to withstand fluid metamorphism and is 446 stable at least at close to mantle transition zone, where are enriched in OH. So 447 metastable K-feldspar could survive to deeper than previously expected, and more 448 likely further subducts and transforms into K-holl-I. The phase transitions of 449 K-feldspar \rightarrow metastable phase \rightarrow K-holl-I could occur in both dry and hydrous slabs. 450 K-holl-I is considered to be the most abundant phase in a continental crust 451 452 composition in the upper and lower mantle. Studies about the existence of K-holl-I 453 was mostly carried out in a high-pressure multi-anvil apparatus, in which samples were firstly compressed up to a target pressure and then heated up (e.g., Wu et al. 454 2009; Ishii et al. 2012). So the metastable phase transition process may has been 455 ignored. The phase behavior of metastable K-feldspar in ultracold subduction could 456 help us to explore the possible formation paths and conditions of K-holl-I, existing 457 depth of K-feldspar, and have further understanding of continental subduction. 458 459

460

Acknowledgement

We appreciate two anonymous reviewers for their constructive suggestions and comments, which helped improve the manuscript significantly. This work was financially supported by project of National Natural Science Foundation of China (U1812402).

466	
467	References
468	Agrusta, R., Van Hunen, J., and Goes, S. (2014) The effect of metastable pyroxene on
469	the slab dynamics. Geophysical Research Letters, 41(24), 8800-8808.
470	Akaogi, M., Kamii, N., Kishi, A., and Kojitani, A. (2004) Calorimetric study on
471	high-pressure transitions in KAlSi ₃ O ₈ . Physics and Chemistry of Minerals, 31(2),
472	85-91.
473	Aliatis, I., Lambruschi, E., Mantovani, L., Bersani, D., Andò, S., Diego Gatta, G.,
474	Gentile, P., Salvioli-Mariani, E., Prencipe, M., Tribaudino, and M., Lottici, P.P.
475	(2015) A comparison between ab initio calculated and measured Raman
476	spectrum of triclinic albite (NaAlSi ₃ O ₈). Journal of Raman Spectroscopy, 46(5),
477	501–508.
478	Aliatis, I., Lambruschi, E., Mantovani, L., Bersani, D., Gatta, G.D., Tribaudino, M.,
479	and Lottici, P.P. (2017) High-pressure Raman spectroscopy on low albite.
480	Physics and Chemistry of Minerals, 44(3), 213-220.
481	Allan, D.R., and Angel, R.J. (1997) A high-pressure structural study of microcline
482	(KAlSi ₃ O ₈) to 7 GPa. European Journal of Mineralogy, 9(2), 263–276.
483	Befus, K.S., Lin, J.F., Cisneros, M., and Fu, S. (2018) Feldspar Raman shift and
484	application as a magmatic thermobarometer. American Mineralogist, 103(4),
485	600-609.
486	Bendel, V., and Schmidt, B.C. (2008) Raman spectroscopic characterisation of
487	disordered alkali feldspars along the join KAlSi ₃ O ₈ -NaAlSi ₃ O ₈ : application to
488	natural sanidine and anorthoclase. European Journal of Mineralogy, 20(6),

489	1055-1065.
490	Chen, T., Gwanmesia, G.D., Ehm, L., Le Losq, C., Neuville, D.R., Phillips, B.L., Li,
491	B., and Liebermann, R.C. (2019) Synthesis and characterization of
492	polycrystalline KAlSi $_{3}O_{8}$ hollandite [liebermannite]: Sound velocities vs.
493	pressure to 13 GPa at room temperature. Comptes Rendus Geoscience, 351(2-3),
494	113–120.
495	Deng, L., Liu, X., Liu, H., and Zhang, Y. (2011) A first-principles study of the phase
496	transition from Holl-I to Holl-II in the composition KAlSi ₃ O ₈ . American
497	Mineralogist, 96(7), 974–982.
498	Ferroir, T., Onozawa, T., Yagi, T., Merkel, S., Miyajima, N., Nishiyama, N., Irifune, T.,
499	and Kikegawa, T. (2006) Equation of state and phase transition in KAlSi $_3O_8$
500	hollandite at high pressure. American Mineralogist, 91(2-3), 327-332.
501	Freeman, J.J., Wang, A., Kuebler, K.E., Jolliff, B.L., and Haskin, L.A. (2008)
502	CHARACTERIZATION OF NATURAL FELDSPARS BY RAMAN
503	SPECTROSCOPY FOR FUTURE PLANETARY EXPLORATION. The
504	Canadian Mineralogist, 46(6), 1477-1500.
505	Gillet, P., Guyot, F., and Malezieux, J.M. (1989) High-pressure, high-temperature
506	Raman spectroscopy of Ca ₂ GeO ₄ (olivine form): some insights on anharmonicity.
507	Physics of the Earth and Planetary Interiors, 58(2–3), 141–154.
508	Guo, X., Chen, S., Li, P., Zhang, Y., Wu, X., and Zhang, J. (2020) Phase transition of
509	sanidine (KAlSi ₃ O ₈) and its effect on electrical conductivity at pressures up to 11
510	GPa. Physics and Chemistry of Minerals, 47(4), 1–8.

511	Hashimoto, G.L., Roos-Serote, M., Sugita, S., Gilmore, M.S., Kamp, L.W., Carlson,
512	R.W., and Baines, K.H. (2009) Felsic highland crust on Venus suggested by
513	Galileo Near-Infrared Mapping Spectrometer data. Journal of Geophysical
514	Research: Planets, 114(5), 1–10.
515	Henderson, C.M.B. (1979) An Elevated Temperature X-ray Study of Synthetic
516	Disordered Na-K Alkali Feldspars. Contributions to Mineralogy and Petrology,
517	70(1), 71–79.
518	Hirao, N., Ohtani, E., Kondo, T., Sakai, T., and Kikegawa, T. (2008) Hollandite II
519	phase in KAlSi ₃ O ₈ as a potential host mineral of potassium in the Earth's lower
520	mantle. Physics of the Earth and Planetary Interiors, 166, 97–104.
521	Hogrefe, A., Rubie, D. C., Sharp, T. G., and Seifert, F. (1994) Metastability of
522	enstatite in deep subducting lithosphere. Nature, 372(6504), 351-353.
523	Hovis, G.L., Medford, A., Conlon, M., Tether, A., and Romanoski, A. (2010)
524	Principles of thermal expansion in the feldspar system. American Mineralogist,
525	95(7), 1060–1068.
526	Iwamori, H. (2007) Transportation of H ₂ O beneath the Japan arcs and its implications
527	for global water circulation. Chemical Geology, 239(3-4), 182-198.
528	Irifune, T., Ringwood, A.E., and Hibberson, W.O. (1994) Subduction of continental
529	crust and terrigenous and pelagic sediments: an experimental study. Earth and
530	Planetary Science Letters, 126, 351–368.
531	Ishii, T., Kojitani, H., and Akaogi, M. (2012) High-pressure phase transitions and
532	subduction behavior of continental crust at pressure-temperature conditions up to

533	the upper part of the lower mantle. Earth and Planetary Science Letters, 357-358,								
534	31-41.								
535	Kawakatsu, H., and Watada, S. (2010) Seismic Evidence for Deep-Water								
536	Transportation in the Mantle. Science, 316(June), 1468-1471.								
537	Kawakatsu, H., and Yoshioka, S. (2011) Metastable olivine wedge and deep dry cold								
538	slab beneath southwest Japan. Earth and Planetary Science Letters, 303(1-2),								
539	1-10.								
540	Klotz, S., Chervin, J.C., Munsch, P., and Le Marchand, G. (2009) Hydrostatic limits								
541	of 11 pressure transmitting media. Journal of Physics D: Applied Physics, 42(7).								
542	Kubo, T., Kaneshima, S., Torii, Y., and Yoshioka, S. (2009) Seismological and								
543	experimental constraints on metastable phase transformations and rheology of								
544	the Mariana slab. Earth and Planetary Science Letters, 287(1-2), 12-23.								
545	Lindsley, D.H. (1966). Melting relations of KAlSi ₃ O ₈ : effect of pressures up to 40								
546	kilobars. American Mineralogist, 51, 1793-1799.								
547	Liu, X., Hu Z., and Deng, L. (2010) Feldspars under conditions of high-temperature								
548	and high-pressure. Acta Petrologica Sinica, 26(12), 3641-3650.								
549	Mao, H.K., Xu, J., and Bell, P.M. (1986) Calibration of the ruby pressure gauge to								
550	800 kbar under quasi-hydrostatic conditions. Journal of Geophysical Research,								
551	91(B5), 4673.								
552	Nishiyama, N., Rapp, R.P., Irifune, T., Sanehira, T., Yamazaki, D., and Funakoshi, K.I.								
553	(2005) Stability and <i>P-V-T</i> equation of state of KAlSi ₃ O ₈ -hollandite determined								
554	by in situ X-ray observations and implications for dynamics of subducted								

555	continental crust material. Physics and Chemistry of Minerals, 32(8-9), 627-							
556	637.							
557	Pakhomova, A., Simonova, D., Koemets, I., Koemets, E., Aprilis, G., Bykov, M.,							
558	Gorelova, L., Fedotenko, T., Prakapenka, V., and Dubrovinsky, L. (2020)							
559	Polymorphism of feldspars above 10 GPa. Nature Communications, 11(1), 1-8.							
560	Peterson, J.W., O'Meara, T.A., Seymour, M.D., Wang, W., and Gu, B. (2009) Sorption							
561	mechanisms of cephapirin, a veterinary antibiotic, onto quartz and feldspar							
562	minerals as detected by Raman spectroscopy. Environmental Pollution, 157(6),							
563	1849-1856.							
564	Ranalli, G., Pellegrini, R., and D'Offizi, S. (2000) Time dependence of negative							
565	buoyancy and the subduction of continental lithosphere. Journal of Geodynamics,							
566	30(5), 539–555.							
567	Rogers, A.D., and Nekvasil, H. (2015) Feldspathic rocks on Mars: compositional							
568	constraints from infrared spectroscopy and possible formation mechanisms.							
5.00								
569	Geophysical Research Letters, 42, 2619–2626.							
569	Geophysical Research Letters, 42, 2619–2626. Sprague, A.L., Nash, D.B., Witteborn, F.C., and Cruikshank, D.P. (1997) Mercury's							
570 571	Geophysical Research Letters, 42, 2619–2626.Sprague, A.L., Nash, D.B., Witteborn, F.C., and Cruikshank, D.P. (1997) Mercury's feldspar connection mid-ir measurements suggest plagioclase. Advances in							
570 571 572	 Geophysical Research Letters, 42, 2619–2626. Sprague, A.L., Nash, D.B., Witteborn, F.C., and Cruikshank, D.P. (1997) Mercury's feldspar connection mid-ir measurements suggest plagioclase. Advances in Space Research, 19, 1507–1510. 							
570 571 572 573	 Geophysical Research Letters, 42, 2619–2626. Sprague, A.L., Nash, D.B., Witteborn, F.C., and Cruikshank, D.P. (1997) Mercury's feldspar connection mid-ir measurements suggest plagioclase. Advances in Space Research, 19, 1507–1510. Sueda, Y., Irifune, T., Nishiyama, N., Rapp, R.P., Ferroir, T., Onozawa, T., Yagi, T., 							
570 571 572 573 574	 Geophysical Research Letters, 42, 2619–2626. Sprague, A.L., Nash, D.B., Witteborn, F.C., and Cruikshank, D.P. (1997) Mercury's feldspar connection mid-ir measurements suggest plagioclase. Advances in Space Research, 19, 1507–1510. Sueda, Y., Irifune, T., Nishiyama, N., Rapp, R.P., Ferroir, T., Onozawa, T., Yagi, T., Merkel, S., Miyajima, N., and Funakoshi, K.I. (2004) A new high-pressure form 							
570 571 572 573 574 575	 Geophysical Research Letters, 42, 2619–2626. Sprague, A.L., Nash, D.B., Witteborn, F.C., and Cruikshank, D.P. (1997) Mercury's feldspar connection mid-ir measurements suggest plagioclase. Advances in Space Research, 19, 1507–1510. Sueda, Y., Irifune, T., Nishiyama, N., Rapp, R.P., Ferroir, T., Onozawa, T., Yagi, T., Merkel, S., Miyajima, N., and Funakoshi, K.I. (2004) A new high-pressure form of KAlSi₃O₈ under lower mantle conditions. Geophysical Research Letters, 							

577	Syracuse, E.M., van Keken, P.E., and Abers, G.A. (2010) The global range of
578	subduction zone thermal models. Physics of the Earth and Planetary Interiors,
579	183(1–2), 73–90.
580	Tseng, H., Heaney, P.J., and Onstott, T.C. (1995) Characterization of Lattice Strain
581	Induced by Neutron Irradiation. Physics and Chemistry of Minerals, 22, 399-
582	405.
583	Urakawa, S., Kondo, T., Igawa, N., Shimomura, O., and Ohno, H. (1994) Synchrotron
584	Radiation Study on the High-Pressure and High-Temperature Phase Relations of
585	KAlSi ₃ O ₈ . Physics and Chemistry of Minerals, 21(6), 387–391.
586	Van Mierlo, W. L., Langenhorst, F., Frost, D. J., and Rubie, D. C. (2013) Stagnation
587	of subducting slabs in the transition zone due to slow diffusion in majoritic
588	garnet. Nature Geoscience, 6(5), 400-403.
589	Wang, X., Chen, Q. F., Niu, F., Wei, S., Ning, J., Li, J., Wang, W., Buchen, J., and Liu,
590	L. (2020) Distinct slab interfaces imaged within the mantle transition zone.
591	Nature Geoscience, 13(12), 822-827.
592	Wu, Y., Fei, Y., Jin, Z., and Liu, X. (2009) The fate of subducted Upper Continental
593	Crust: An experimental study. Earth and Planetary Science Letters, 282(1-4),
594	275-284.
595	Xie, C., Du, J., Cui, Y., Chen, Z., Zhang, W., Yi, L., and Deng, L. (2012) Variation of
596	Raman Spectra of Oligoclase under $1.0 \sim 4.4$ GPa. Spectroscopy and Spectral
597	Analysis, 32(3), 691-694.
598	Xu, J., Fan, D., Zhang, D., Guo, X., Zhou, W., and Dera, P. K. (2020) Phase transition

599	of enstatite-ferrosilite solid solutions at high pressure and high temperature:
600	Constraints on metastable orthopyroxene in cold subduction. Geophysical
601	Research Letters, 47(12), 0–2.
602	Yagi, A., Suzuki, T., and Akaogi, M. (1994) High-Pressure Transitions in the System
603	KAlSi ₃ O ₈ –NaAlSi ₃ O ₈ . Physics and Chemistry of Minerals, 21, 12–17.
604	Yamada, H., Matsui, Y., and Ito, E. (1984) Crystal-chemical characterization of
605	KAlSi ₃ O ₈ with the hollandite structure. Mineralogical Journal, 12, 29–34.
606	Yong, W., Dachs, E., Withers, A.C., and Essene, E.J. (2006) Heat capacity and phase
607	equilibria of hollandite polymorph of KAlSi ₃ O ₈ . Physics and Chemistry of
608	Minerals, 33(3), 167–177.
609	Zhang, J., Ko, J., Robert, M.H., and Charles, T.P. (1993) High-pressure crystal
610	chemistry of KAlSi ₃ O ₈ hollandite. American Mineralogist, 78, 493–499.
611	Zhang, R.Y., Liou, J.G., Yang, J.S., and Ye, K. (2003) Ultrahigh-pressure
612	metamorphism in the forbidden zone: the Xugou garnet peridotite, Sulu terrane,
613	eastern China. Journal of Metamorphic Geology, 21(6), 539-550.
614	Zhao, C., Li, H., Wang, Y., Jiang, J., and He, Y. (2017) SrB ₄ O ₇ :Sm ²⁺ : an optical sensor
615	reflecting non-hydrostatic pressure at high-temperature and/or high pressure in a
616	diamond anvil cell. High Pressure Research, 37(1), 18-27.
617	Zheng, Y. F. (2012) Metamorphic chemical geodynamics in continental subduction
618	zones. Chemical Geology, 328, 5-48.
619	Zheng, Y.F., and Hermann, J. (2014) Geochemistry of continental subduction-zone
620	fluids. Earth, Planets and Space, 66(1), 93.

621 Zheng, Y.F. (2019) Subduction zone geochemistry. Geoscience Frontiers, 10(4),
622 1223–1254.



640



641

Figure 2. Representative Raman spectra of K-feldspar at various pressures and ambient temperature, displayed in the frequency range of 100-800 cm⁻¹ (a), the pressure dependence of the Raman shift (b), and the full width at half maximum (*FWHM*) for the *Ic* mode as a function of the pressure (c).



647

Figure 3. Representative Raman spectra of K-feldspar at various temperatures and ambient
pressure, displayed in the frequency range of 100-1200 cm⁻¹ (a), and the temperature dependence
of Raman bands of K-feldspar at ambient pressure (b).

651





668



669

Figure 5. Raman spectra of samples recovered from high *P*-*T* to ambient conditions in different

671 PTMs (the PTMs are given in brackets), displayed in the frequency range of 100-1200 cm⁻¹. Peaks

672 with black, red, blue, green, and purple asterisk are the characteristic Raman peaks of K-feldspar

673 (K-fsp), coesite (Coe), Muscovite (Mus), KHCO₃, and stishovite (St), respectively. The black and

674 green solid lines represent K-feldspar and the assemblage containing coesite, respectively.

675



677

678 Figure 6. Phase diagram of K-feldspar constrained by high P-T experiments. Squares, circles, pentagon, hexagon, erected triangles, right triangles, left triangles, rhombus and stars indicate the 679 680 P-T conditions where K-feldspar (K-fsp), wadeite + kyanite + coesite (Wd + Ky + Coe), wadeite + 681 kyanite + stishovite (Wd + Ky + St), liquid (Liq), leucite + liquid (Lc + Liq), liquid + kyanite + 682 coesite (Liq + Ky + Coe), hollandite (K-holl-I) and hollandite II (K-holl-II) are stable, respectively. 683 Meanwhile, adding different colors to distinguish them more clearly. The black solid lines 684 represent the phase boundaries determined by high P-T experiments or from previous phase 685 diagram, and the dotted lines indicate the extrapolated phase boundaries by limited data (Lindsley 686 1966; Zhang et al. 1993; Urakawa et al. 1994; Yagi et al. 1994; Nishiyama et al. 2005; Akaogi et al. 2004; Ferroir et al. 2006; Yong et al. 2006; Hirao et al. 2008; Chen et al. 2019; Guo et al. 2020; 687 688 Pakhomova et al., 2020). Green and blue lines denote the geothermal gradients at 2°C/km and 689 5°C/km, respectively, in subduction zones (Zheng and Hermann 2014; Zheng 2019).





Figure 7. Raman spectra of samples quenched from different *P*-*T* conditions.

693



Figure 8. A schematic diagram of the ultracold dry (a) and hydrous (b) subducting slab and the phase relations of K-feldspar. The phase transition of K-feldspar \rightarrow Ky + Wd + Coe/St \rightarrow K-holl-I at relatively high-temperature regions (> 650 °C) is from previous studies (e.g., Zheng and Hermann

698 2014; Zheng 2019), while the other phase transitions (K-feldspar→metastable phase→K-holl-I,

699 K-feldspar→Coesite-containing assemblage) are from this study. The figure is not drawn to scale.

700 K-fsp: K-feldspar; Ky + Wd + Coe/St: kyanite + wadeite + coesite/stishovite; K-holl-I:

701 K-hollandite-I.

Table 1. Constants determined in the expression: $v_P = v_{i0} + (\partial v_i / \partial P) \times P$, $v_T = v_{i0} + (\partial v_i / \partial T) \times T$, the

isothermal (γ_{iT}) and isobaric (γ_{iP}) mode Grüneisen parameters, and the intrinsic anharmonic mode

704 parameter (β_i) of K-feldspar.

Modes	$v_{i\theta}$ (cm ⁻¹)	$\frac{\partial v_i}{\partial P}$ (cm ⁻¹ /GPa)	Y iT	$v_{i\theta}$ (cm ⁻¹)	$(\partial v_i / \partial T) \times 10^2$ (cm ⁻¹ /K)	ү іР	β _i ×10 ⁵ (K ⁻¹)
III _a	-	-	-	124.28(13)	-0.82(3)	4.56(1)	-
III_b	157.99(13)	5.32(45)	1.96(2)	158.33(13)	-0.52(5)	2.28(2)	-0.46(2)
II_a	283.15(16)	1.41(28)	0.29(3)	284.27(16)	-0.17(1)	0.42(1)	-0.19(1)
Ia	-	-	-	455.42(10)	-0.73(6)	1.12(11)	-
I_b	474.68(10)	2.13(7)	0.26(7)	476.42(9)	-0.28(2)	0.41(5)	-0.22(3)
Ic	513.48(2)	3.56(9)	0.40(6)	515.22(2)	-0.38(1)	0.51(3)	-0.16(1)

705 Notes: The frequency at room pressure (v_{i0}) , pressure coefficients $(\partial v_i / \partial P)$, and temperature coefficients $(\partial v_i / \partial T)$

obtained in this study were used to calculate the mode Grüneisen parameters $(\gamma_{iT}, \gamma_{iP})$ and intrinsic anharmonic

707 mode parameter βi , using the isothermal bulk modulus (K_T) with values of 58.3(2.0) GPa (Allan and Angel 1997),

and the volume thermal expansion coefficient (α) with values of 1.44(8)×10⁻⁵ K⁻¹ (Henderson 1979).



721 Supplementary Information

- 722 for the manuscript "A revisit to phase transition behavior of K-feldspar at high-pressure and
- 723 high-temperature: Implications on metastable K-feldspar in cold subduction" by Chengcheng He
- 724 et al.
- 725
- 726
- 727 Supplementary



728

729 Figure S1. Schematic cross-section of cell assembly for quenching experiments.



732

Figure S2. Representative Raman spectra of K-feldspar measured with increasing pressure in nitrogen (a) and H_2O (b). The decompression spectra as shown in (a) is nearly same with the spectrum at ambient condition, which reveals the phase transition of K-feldspar into metastable

- phase under high was reversible. The spectra only displayed in the frequency range of 200-1200
- cm^{-1} because of the intense Raman bands at 100 and 200 cm⁻¹ of the pressure medium nitrogen.
- 738 Another decompression spectra as shown in (b) is also similar with the spectrum at ambient
- condition, which indicates the reversibility of the not fully amorphization.

740



741

742 Figure S3. Pressure dependence of the Raman bands of K-feldspar at ambient temperature.



744





746





748



749



751





754



Figure S4. Raman spectra of the K-feldspar up to ~14.7 GPa, 430 °C within various PTMs. Raman spectra of K-feldspar at high P-T in 4:1 ME mixture, 16:3:1 MEW mixture, helium, no PTM, and H₂O, respectively (a)-(e), (f), (g), (h), (i)-(j). N(1-7) are the new bands that may correspond to the assemblage containing coesite revealed in Raman spectra of the recovered

760 samples.



761

762 Figure S5. Raman spectra of K-feldspar during compressing (a), cooling (b), and decompressing

763 (c) process.

764



766





769 Figure S6. Determination of Raman spectra of recovered sample from 430 °C, 6.7 GPa (compressed in ME) (a), 770 430 °C, 5.9 GPa (compressed in MEW) (b) and 400 °C, 5.9 GPa (compressed in H₂O) (c) by comparing with 771 standard Raman spectrum of possible substances. The standard Raman spectra of KHCO3 was obtained by 772 measuring its high purity powder (99.99%, Alfa Aesar), and the others were download from RRUFF database 773 (https://rruff.info). The recovered sample from 430 °C, 6.7 GPa (ME) contain coesite, stishovite, and KHCO3. 774 The recovered sample from 430 °C, 5.9 GPa (MEW) contain coesite, muscovite, and KHCO3. The recovered 775 sample from 400 °C, 5.9 GPa (H2O) contain coesite and muscovite. The possible substances corresponding to the 776 sharp peaks at 900-1100 cm⁻¹ has not been determined yet, as shown in Figures (b-c).



777

Figure S7. Raman spectra of K-feldspar holding at 430 $^{\circ}$ C in H₂O for 1h.

779