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2 **Synthesis and crystal structure of LiNbO<sub>3</sub>-type Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub>: A possible indicator of shock conditions of**  
3 **meteorites**

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13  
14 **Abstract**

15 LiNbO<sub>3</sub>-type Mg<sub>2.98(2)</sub>Al<sub>1.99(2)</sub>Si<sub>3.02(2)</sub>O<sub>12</sub> (py-LN) was synthesized by recovering a run product from  
16 2000 K and 45 GPa to ambient conditions using a large volume press. Rietveld structural refinements were  
17 carried out using the one-dimensional synchrotron XRD pattern collected at ambient conditions. The unit-cell  
18 lattice parameters were determined to be  $a = 4.8194(3)$  Å,  $c = 12.6885(8)$  Å,  $V = 255.23(3)$  Å<sup>3</sup> with  $Z = 6$   
19 (Hexagonal,  $R3c$ ). The average A-O and B-O distances of the AO<sub>6</sub> and BO<sub>6</sub> octahedra have values similar to

20 those that can be obtained from the sum of the ionic radii of the averaged *A*- and *B*-site cations and oxygen  
21 (2.073 and 1.833 Å, respectively). The present compound has the *B*-site cations at the octahedral site largely  
22 shifted along the *c* axis compared with other LiNbO<sub>3</sub>-type phases formed by back-transition from perovskite  
23 (Pv)-structure, and as a result, the coordination number of this site is better described as 3 + 3. It appears  
24 therefore that the *B*-site cation in the octahedral position cannot be completely preserved during the back-  
25 transition because of the small size of Si and Al, which occupy usually a tetrahedral site at ambient conditions.  
26 The formation of py-LN can be explained by the tilting of BO<sub>6</sub> octahedra of the perovskite structure having the  
27 pyrope composition and formed at high *P-T* conditions. The tilting is driven by the decrease in ionic radius  
28 ratio between the *A*-site cation and oxygen during decompression. This also explains why there is no back-  
29 transition from the Pv-structure to the ilmenite-structure during decompression, since this is a reconstructive  
30 phase transition whose activation energy cannot be overcome at room temperature. Py-LN may be formed in  
31 shocked meteorites by the back-transformation after the garnet-bridgmanite transition, and will indicate shock  
32 conditions around 45 GPa and 2000 K.

33 **Key words:** bridgmanite, large volume press, Rietveld analysis, LiNbO<sub>3</sub>, high pressure

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## Introduction

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LiNbO<sub>3</sub> (LN)-structure (space group No. 161: *R3c*) crystallizes in the trigonal system and is very  
similar to the ilmenite (Ilm)-structure (space group: *R $\bar{3}$* ), which has the same crystal system. LN- and Ilm-type  
*ABO*<sub>3</sub> compounds consist of two kinds of octahedra, referred as AO<sub>6</sub> and BO<sub>6</sub> in this paper. The larger cations

39 are stored in the  $AO_6$  octahedra, whereas the smaller ones are in the  $BO_6$  octahedra. The two structures can be  
40 distinguished by the different arrangement of the  $AO_6$  and  $BO_6$  octahedra. The Ilm-structure has alternating  
41 layers of edge-sharing  $AO_6$  and  $BO_6$  octahedra normal to the  $c$  axis. On the other hand, each layer in the LN-  
42 structure consists of alternating edge-sharing  $AO_6$  and  $BO_6$  octahedra. The majority of LN-type  $ABO_3$   
43 compounds result from the back-transformation of perovskite (Pv)-type high-pressure polymorphs during  
44 decompression. The LN-type structure is non-centrosymmetric and therefore this kind of compounds is used as  
45 functional materials of nonlinear optics and piezoelectricity. Recently, it was suggested that the degree of  
46 polarity increases with the distortion of the  $BO_6$  octahedra (Inaguma et al., 2012). Therefore, efforts have been  
47 made to synthesize novel LN-type compounds, sometimes by applying high-pressure techniques, and to  
48 analyze their structures. Leinenweber et al. (1991) reported that  $MnSnO_3$  and  $FeTiO_3$  with the Ilm-type  
49 structure transform to the Pv-type structure at high pressure and temperature, and back-transform to the LN-  
50 type structure during decompression at room temperature. They interpreted this back-transition based on the  
51 tolerance factor for the Pv-structure,  $t$ , which is defined as  $t = (r_A + r_O) / \sqrt{2}(r_B + r_O)$ , where  $r_A$ ,  $r_B$ , and  $r_O$  are the  
52 ionic radii of the  $A$ - and  $B$ -site cations and oxygen, respectively, and concluded that Pv-structured compounds  
53 with  $t < 0.84$  transform to the LN-structure by pressure release.

54 In the field of Earth science, various kinds of  $ABO_3$  compounds have been investigated as analogue  
55 materials of  $MgSiO_3$  with the orthorhombic Pv-structure (bridgmanite, Brg), because Brg is considered to be  
56 the most abundant mineral in the Earth's lower mantle. Brg can incorporate  $Al^{3+}$  by two possible substitution  
57 mechanisms (a charge-coupled substitution ( $Mg + Si$  for  $2Al$ ) and an oxygen vacancy substitution ( $2Si + O$  for

58 2Al + one oxygen vacancy)). This incorporation shifts the stability field of Brg to higher pressures and changes  
59 its physical properties. For example, rigidity decreases and electrical conductivity increases (e.g. Kesson et al.,  
60 1995; Irifune et al., 1996; Xu et al., 1998; Walter et al., 2004; Murakami et al., 2012). It was also pointed out  
61 that the oxygen vacancy substitution may allow water incorporation (Navrotsky, 1999) although it is still  
62 unclear whether aluminous Brg can incorporate significant amounts of water. Therefore, understanding the  
63 stability and Al-substitution mechanism in Brg is one of the important subjects in Earth science.

64 Funamori et al. (1997) using a laser-heated diamond-anvil cell (LH-DAC) discovered that  
65  $(\text{Mg,Fe,Ca,Mn})_3\text{Al}_2\text{Si}_3\text{O}_{12}$  natural garnet directly transforms to the Pv-structure at high pressures, and then  
66 transforms to a metastable phase with the LN-structure during decompression. Miyajima et al. (1999)  
67 examined the Pv-LN transition using natural garnets with various compositions also using a LH-DAC. These  
68 authors suggested that Brg with  $\text{Al}_2\text{O}_3$  contents larger than 25 mol% transforms to the LN-structure during  
69 decompression. However, the structures of these LN-type compounds were not analyzed in detail because only  
70 small amount of the LN-type material was obtained together with amorphous phases, which were also formed  
71 by back-transformations in DACs. Recently, Ishii et al. (2016) succeeded in synthesizing a bulk sample of the  
72 LN-type phase with the pyrope composition ( $\text{Mg}_3\text{Al}_2\text{Si}_3\text{O}_{12}$ ) at pressures over 40 GPa using a multi-anvil  
73 press. Since this product is almost completely composed of the LN-type phase, it is suitable for the structural  
74 analysis by the Rietveld method.

75 In this paper, we report the crystal structure of the LN-type  $\text{Mg}_3\text{Al}_2\text{Si}_3\text{O}_{12}$  determined by X-ray  
76 powder diffraction and Rietveld analysis. The transformation mechanism of the Pv-LN back-transition and

77 implications for a possible natural occurrence of this compound are discussed.

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## Experimental methods

### 80 1. Sample synthesis and characterization

81 The  $\text{Mg}_3\text{Al}_2\text{Si}_3\text{O}_{12}$  with LN-type structure was synthesized via the following three steps. Firstly, a  
82 glass with a nominal composition of  $\text{Mg}_3\text{Al}_2\text{Si}_3\text{O}_{12}$  (py-glass) was produced by quenching in water a mixture  
83 of  $\text{MgO}$ ,  $\text{Al}_2\text{O}_3$  and  $\text{SiO}_2$  with a 3:1:3 molar ratio melted at 1950 K for 1 h. The py-glass was then converted to  
84 a sintered aggregate of Ilm-type  $\text{Mg}_3\text{Al}_2\text{Si}_3\text{O}_{12}$  (py-Ak) at 27 GPa and 1170 K for 1 h. The sintered py-Ak  
85 aggregate was then subjected to a press load of 15 MN and 2000 K for 3 h, quenched to ambient temperature,  
86 and then decompressed to ambient pressure. The recovered sample consisted of a sintered aggregate of LN-  
87 type  $\text{Mg}_3\text{Al}_2\text{Si}_3\text{O}_{12}$  (py-LN) compound. The high  $P$ - $T$  syntheses were conducted using a multi-anvil apparatus  
88 for ultra-high pressure generation, IRIS-15, at the Bayerisches Geoinstitut, University of Bayreuth, Germany.  
89 Details of the experimental setup are described in Ishii et al. (2016). Liu et al. (2016) reported that Brg with a  
90 pyrope composition (py-Brg) is stabilized at 45 GPa and 2000 K, and that py-LN formed by back-transition  
91 from py-Brg during decompression in a multi-anvil press installed at the BL04B1 in the synchrotron facility,  
92 Spring-8. Therefore, py-LN obtained in this study also has been likely formed by the back-transition from py-  
93 Brg during decompression from at least 45 GPa.

94 The phases present in the synthesis run products recovered from the second and the third final steps  
95 were identified using a laboratory micro-focused X-ray diffractometer (MF-XRD) (Bruker, D8 DISCOVER)

96 with a two-dimensional solid-state detector (VÅNTEC500) and a micro-focus source ( $I\mu S$ ) of Co-K $\alpha$  radiation  
97 operated at 40 kV and 500  $\mu A$ . The X-ray beam was focused to 50  $\mu m$  beam using a IFG polycapillary X-ray  
98 mini-lens. The run products were also examined in back-scattered electron (BSE) images by means of scanning  
99 electron microscopy (SEM) in order to determine whether minor phases were present. Compositional analysis  
100 was carried out using an electron microprobe with wavelength-dispersive spectrometers (JEOL, JXA-8200)  
101 and an accelerated voltage and probe current of 15 kV and 15 nA, respectively. Natural pyrope was used as  
102 standard material for Mg, Al and Si.

## 104 **2. Synchrotron X-ray diffraction and Rietveld refinement**

105 Angle-dispersive powder X-ray diffraction was conducted at the beam line BL10XU of the synchrotron  
106 radiation facility SPring-8, Japan. The py-LN aggregate was crushed to a powder. The powdered sample was  
107 loaded into a hole with a diameter of 300  $\mu m$  in a rhenium plate with a thickness of 20  $\mu m$ . X-ray beams  
108 monochromatized with a Si double monochromator were collimated to 10-20  $\mu m$  in diameter to irradiate the  
109 sample. The two-dimensional diffraction patterns were collected for 1 s on a flat panel detector (PerkinElmer).  
110 The X-ray wavelength ( $\lambda = 0.41429 \text{ \AA}$ ) and detector distance from the sample were calibrated using a CeO<sub>2</sub>  
111 standard. Five one-dimensional diffraction profiles of the sample were obtained by integrating the two-  
112 dimensional profile at different azimuth angles using the software “IPAnalyzer” (Seto et al., 2010). The five  
113 one-dimensional diffraction patterns were combined for the Rietveld refinement.

114 Rietveld analysis was performed using the RIETAN-FP/VENUS package (Izumi and Momma, 2007).

115 The VESTA software (Momma and Izumi, 2011) was used to draw crystal structures. The structure parameters  
116 of LN-type FeTiO<sub>3</sub> (Leinenweber et al., 1991) were used as an initial structure model. Unit-cell lattice  
117 parameters, scale factor, atomic coordinates and isotropic atomic displacement parameters (Table 1), and a  
118 split-type pseudo-Voigt profile fitting function formulated by Toraya (1990) were refined simultaneously. Since  
119 peaks of SiO<sub>2</sub> stishovite (St) were present in the synchrotron XRD, this phase was included in the Rietveld  
120 analysis, and the lattice parameters, the scale factor and the atomic coordinates were refined. A Legendre  
121 polynomial function with 12 parameters was used for fitting of the XRD background.

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## Results and Discussion

### 1. Phase identification and sample composition

124  
125 The compositions of py-glass, py-Ak, and py-LN were found to be Mg<sub>2.93(2)</sub>Al<sub>2.00(2)</sub>Si<sub>3.04(2)</sub>O<sub>12</sub>,  
126 Mg<sub>2.99(2)</sub>Al<sub>2.01(2)</sub>Si<sub>3.01(1)</sub>O<sub>12</sub>, and Mg<sub>2.98(2)</sub>Al<sub>1.99(2)</sub>Si<sub>3.02(2)</sub>O<sub>12</sub>, respectively, in agreement with the nominal  
127 composition and confirming that the high-pressure phases have a pyrope composition.

128 The MF-XRD patterns of py-Ak and py-LN are shown in Fig. 1. Although both phases have similar  
129 structural features, their diffraction patterns show clear differences. For instance, the 003 peak is present in the  
130 pattern of py-Ak but it is absent in the py-LN pattern, as required by the reflections conditions of the two  
131 different space groups ( $R\bar{3}$  and  $R3c$ , respectively) in which these two materials crystallize. All peaks can be  
132 identified as belonging to py-Ak (Fig. 1, bottom) or to py-LN (Fig. 1, top); however, some St and very small  
133 amounts of Al<sub>2</sub>O<sub>3</sub> corundum (Crn) grains were identified in the BSE images taken for py-LN. The amounts of

134 these minor phases appear to be below the detection limit of the Bruker diffractometer. These trace phases may  
135 have formed during heating to 2000 K for the synthesis of py-Brg. As mentioned in the experimental session, in  
136 the synchrotron XRD pattern of py-LN, only St was present but no peaks of Crn were observed.

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## 138 **2. Rietveld refinement of LiNbO<sub>3</sub>-type Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub>**

139 The results of Rietveld refinements of the synchrotron XRD pattern of py-LN are shown in Fig. 2.  
140 Because the X-ray scattering factors of Mg, Al and Si atoms are very similar, it is practically impossible to  
141 determine their site occupancies at the *A* and *B* sites. Thus, these were fixed by assuming 25% of Al at each  
142 site based on the compositional analysis and complete order of Mg-Si resulting in the following occupancies:  
143 Mg<sub>0.75</sub>Al<sub>0.25</sub> at the *A* site and Si<sub>0.75</sub>Al<sub>0.25</sub> at the *B* site. This assumption is based on the fact that no cation  
144 disorder has been reported either for Fe-bearing, Al-free or Fe-Al-bearing bridgmanites (e.g. Horiuchi et al.,  
145 1987; Vanpeteghem et al., 2006), thus the same cation distribution is likely retained in py-LN since it is formed  
146 by the back-transformation of py-Brg. During refinement, the isotropic atomic displacement parameter ( $U_{iso}$ ) of  
147 oxygen converged towards a much larger value than the one generally observed for several structures (i.e.  
148  $< \sim 0.013 \text{ \AA}^2$ ) (see Table 1). Therefore, we also performed refinements with the  $U_{iso}$  of oxygen fixed to the value  
149 obtained for LiNbO<sub>3</sub> by means of single-crystal structural refinement (Hsu et al., 1997). Both refinement  
150 strategies, i.e. with oxygen  $U_{iso}$  fixed or free to vary resulted in very similar crystal structures (see Table 1 and  
151 3). In both cases, the discrepancy indexes converged to satisfactory values of less than 5% (Table 1)

152 The difference between the Ilm- and LN-structures can be described in terms of their *c/a* ratio



153 (Funamori et al. 1997). The ranges of  $c/a$  ratios of Ilm- and LN-type phases are 2.76-2.87 and 2.58-2.69,  
154 respectively (e.g. Ko and Prewitt, 1988; Leinenweber et al., 1991; Funamori et al., 1997; Hsu et al., 1997;  
155 Inaguma et al., 2010; Akaogi et al., 2002). The present py-LN has a  $c/a$  ratio of 2.63, which is within the range  
156 of the LN-type phases defined by Funamori et al. (1997), and is smaller than that of py-Ak reported by Akaogi  
157 et al. (2002) ( $c/a = 2.80$ ). The unit-cell volumes and densities of the  $Mg_3Al_2Si_3O_{12}$  compounds are summarized  
158 in Table 2. There is a volume increase for the py-Brg to py-LN transition at ambient conditions of  $\sim 3.0\%$  and  
159 this value is in the range obtained for  $MnTiO_3$ ,  $MnSnO_3$ ,  $FeTiO_3$ ,  $ZnGeO_3$  and  $MgGeO_3$  compounds (2.7-  
160 3.6%) (Leinenweber et al., 1991). The density of py-LN lies between those of py-Ak and py-Brg, which may  
161 explain the formation of py-LN and not py-Ak by back-transformation from py-Brg.

162 The crystal structure of py-LN is shown in Fig. 3. The bond length, bond angles, and bond valence sum  
163 (BVS) (Brown and Altermatt, 1985) and effective coordination numbers ( $n_c$ ) (Hoppe, 1979) calculated from  
164 the bond lengths of the py-LN are summarized in Table 3. Neighboring  $AO_6$  octahedra ( $A = Mg_{0.75}Al_{0.25}$ ) and  
165 neighboring  $BO_6$  octahedra ( $B = Si_{0.75}Al_{0.25}$ ) are corner-sharing, whereas  $AO_6$  and  $BO_6$  octahedra are edge-  
166 sharing in the  $ab$ -plane and share a face along the  $c$ -axis, resulting in  $A$ - $B$ -empty site periodicity along the  $c$ -  
167 axis (Fig. 3). The average  $A$ -O and  $B$ -O distances of the  $AO_6$  and  $BO_6$  octahedra are 2.033-2.043 Å and 1.840-  
168 1.843 Å, respectively. The ionic radii of  $Mg^{2+}$ ,  $Al^{3+}$ ,  $Si^{4+}$ , and  $O^{2-}$  in six-fold coordination are  $r_{Mg} = 0.72$ ,  $r_{Al} =$   
169  $0.535$ ,  $r_{Si} = 0.4$  and  $r_O = 1.40$  Å, respectively (Shannon, 1976). The average  $A$ -O and  $B$ -O distances calculated  
170 using these ionic radii for  $0.75r_{Mg} + 0.25r_{Al} + r_O$  (2.073 Å) and  $0.75r_{Si} + 0.25r_{Al} + r_O$  (1.833 Å), respectively, are  
171 very similar to those obtained from the structural refinement, hence supporting the assumption of complete

172 order of  $\text{Mg}^{2+}$  and  $\text{Si}^{4+}$ . The BVS value of the *A* site (Table 3) is slightly larger than 2 as expected due to the Al  
173 substitution, whereas the BVS value of the *B* site is still close to the value of 4 in spite of the presence of  $\text{Al}^{3+}$ ,  
174 which has a lower valence than  $\text{Si}^{4+}$ . Moreover, the  $n_c$  value of the *B*-site cation (3.13-3.40) is much smaller  
175 than 6 (Table 3); this derives from the large difference between the two non-equivalent *B*-O lengths (0.427-  
176 0.505). The *B*-site cation is largely shifted from the position at the center of the octahedron along the *c* axis,  
177 reducing the coordination number to almost three. The  $n_c$  values of  $\text{Ti}^{4+}$  cations in LN-type  $\text{FeTiO}_3$  and  
178  $\text{MnTiO}_3$  compounds (5.0-5.2) which formed by back-transition from the corresponding Pv-type phases are  
179 significantly larger (e.g. Ko and Prewitt, 1988; Leinenweber et al., 1991; Hsu et al. 1997). This suggests that  
180 during the back-transformation from the py-Brg to the py-LN, the Si and Al cations tend to shift toward a  
181 tetrahedral coordination, given that this is their preferred coordination at ambient conditions. However, a  
182 complete reversion to such coordination is likely not possible at room temperature and the *B* site can therefore  
183 be seen as a 3 + 3 coordination site. The large off-set of the *B* atoms from the center of the octahedra is likely  
184 favoured by the acentric space group in which the py-LN phase crystallizes. However, additional property tests  
185 (e.g. ferroelectric hysteresis measurement) need to be conducted to characterise the dielectric properties of py-  
186 LN phase which likely may not present the ferroelectric behaviour of many other LN-type structures.

### 188 3. Back-transformation mechanism from py-Brg to py-LN

189 This study shows that the investigated compound has the LN-type structure. Liu et al. (2016) reported  
190 that Brg with 25 mol%  $\text{Al}_2\text{O}_3$  transforms to py-LN during decompression, whereas Brg with a lower  $\text{Al}_2\text{O}_3$

191 content can be recovered. Bridgmanite has the orthorhombic perovskite structure which can be view as a result  
192 from the distortion of the cubic Pv-structure due to the too small ratio in ionic radii between the A-site cation  
193 and the 12-coordinated anion. The ionic radius of the  $Mg^{2+}$  is too small for the oxygen packing in the ideal  
194 cubic Pv-structure. Distortion to the orthorhombic symmetry decreases the actual coordination number from  
195 twelve to eight + four, which changes the  $AO_{12}$  dodecahedral site to a  $AO_6$  prism site with two longer A-site  
196 cation-oxygen bonds (bicapped prism site) with consequent reduction of the polyhedral volume. The substitution  
197 of  $Mg^{2+}$  by  $Al^{3+}$  makes the average ionic radius of the A-site cation even smaller. Accordingly, bridgmanite  
198 with pyrope composition has likely a much more distorted orthorhombic structure than pure  $MgSiO_3$   
199 bridgmanite and can be stabilized only at higher pressures, because at these conditions the relatively larger  
200 compressibility of the oxygen anions with respect to the cations favor the substitution of smaller atoms at the A  
201 site. By decreasing pressure, the oxygen atoms expand more than the A-site cations and the resulting A-site  
202 volume becomes too large for the small Mg and Al cations. This likely hinders the possibility of recovering the  
203 orthorhombic Pv-structure containing a large amount of Al (25 mol%) outside its stability field and favors the  
204 transformation to the LN-structure because of the reduction of the coordination number of the A-site cation to  
205 six.

206 A comparison among the orthorhombic Pv-structure ( $MgSiO_3$  Brg), the LN-structure with pyrope  
207 composition and  $MgSiO_3$  akimotoite is given in Figs. 4a, 4b and 4c, respectively. The Pv- and LN-structures  
208 have a very similar  $BO_6$  octahedral framework. The transformation from the orthorhombic Pv-structure to the  
209 LN-structure is caused by the change in the tilting of the  $BO_6$  octahedra and by the shift of the A-site cations, as

210 shown by the thick yellow and thin red arrows in Fig. 4a. These changes correspond to the transformation from  
211 the [111] cubic closed-packing layer composed of oxygen and *A*-site cations in the orthorhombic Pv-structure  
212 to the [001] hexagonal closed-packing layer of oxygen with an interstitial *A*-site cations between two oxygen  
213 layers in the LN-structure. Red balls in Figs. 4a and 4b represent the oxygen atoms coordinating the *A*-site  
214 cation in the Pv- and LN-structures. Due to the back-transformation, the *A*-site cation shifts closer to the  
215 octahedral framework giving rise to the octahedral coordination observed in the LN-structure (Fig. 4b).

216 The octahedral tilting can be described for the two structures in terms of *B-O-B* angles. These angles  
217 are all the same in the LN-structure and their value (143.0-143.3°, Table 3) is smaller than the two *B-O-B*  
218 angles (146.5° in the *c* direction and 147.0° in the *b* direction) of the MgSiO<sub>3</sub> Brg structure (Horiuchi et al.,  
219 1982). However, as already mentioned, the Al substitution in bridgmanite may give rise to a larger  
220 orthorhombic distortion (similar to what observed for Fe-Al-bearing bridgmanites, Vanpeteghem et al., 2006),  
221 resulting in *B-O-B* angles smaller than those observed for pure MgSiO<sub>3</sub> bridgmanite. It is very likely, therefore,  
222 that the Brg to LN transformation results in an increase of the *B-O-B* angle and hence a smaller octahedral  
223 tilting.

224 Py-Brg does not transform to py-Ak, but to py-LN, even though the Ilm- and LN-structures are very  
225 similar. As shown in Fig. 4c, the Ilm-structure has alternating layers of corner-sharing AO<sub>6</sub> and BO<sub>6</sub> octahedra.  
226 This arrangement cannot be obtained simply by tilting of BO<sub>6</sub> octahedra and displacement of the *A*-site cations  
227 of the Pv-structure, but requires a reconstructive transition whose activation energy cannot be overcome during  
228 the pressure release at room temperature.

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## Implication

### 231 Possible existence of py-LN in shocked meteorites

232 A natural LN-type FeTiO<sub>3</sub> compound was recently discovered in a shocked gneiss from the Ries Crater,  
233 Germany (e.g. Dubrovinsky et al. 2009). Akaogi et al. (2016) constrained the shock pressure needed to obtain  
234 such material in the range of 14-28 GPa, based on the stability field of Pv-type FeTiO<sub>3</sub> which likely was  
235 formed during the shocking event and then back-transformed to the LN-type FeTiO<sub>3</sub> structure after the shock  
236 relaxation. Similarly, py-LN phases might form after quenching in shocked meteorites where the pressures  
237 may reach the values needed for the transition from garnet to Brg. The recovery of such phase after a shock  
238 event depends on its thermal stability, however the py-LN phase may be also a good indicator of the formation  
239 of the Brg phase with pyrope garnet composition and can be used to constrain the *P-T* conditions of the  
240 relatively heavy shock event around 45 GPa and 2000 K at least in a similar manner as amorphous phases with  
241 the (Mg,Fe)SiO<sub>3</sub> Brg composition have been used as evidence of the formation of Brg in meteorites (Sharp et  
242 al., 1997; Tomioka and Fujino, 1997).

243

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- 322
- 323

### Figure captions

324 **Figure 1.** MF-XRD patterns of akimotoite (bottom) with the pyrope composition (py-Ak) synthesized at 27  
325 GPa and 1170 K and LiNbO<sub>3</sub>-type phase (top) also with the pyrope composition (py-LN) obtained from  
326 recovering a sample from 2000 K and 45 GPa.

327

328 **Figure 2.** Rietveld refinement of LiNbO<sub>3</sub>-type Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub>. The XRD pattern was collected at atmospheric  
329 pressure and room temperature. Data points and solid lines represent the observed and the calculated profiles,  
330 respectively, and the residual curve is shown at the bottom. Bragg peak positions are shown by small ticks. The  
331 upper and lower ticks are for LiNbO<sub>3</sub>-type Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub> and rutile-type SiO<sub>2</sub> (stishovite), respectively.

332

333 **Figure 3.** Crystal structure of LiNbO<sub>3</sub>-type Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub> obtained using the oxygen-*U<sub>iso</sub>*-fixed model (note that  
334 this structure is identical to that obtained from the oxygen-*U<sub>iso</sub>*-free model). The solid line is the unit-cell.  
335 Occupancies of Mg and Al at the *A* site and Si and Al at the *B* site are shown with areas of different colors in  
336 each sphere. Smaller spheres are oxygen. BVS and *n<sub>c</sub>* represent the bond valence sum values and effective  
337 coordination numbers of each site, respectively (Table 3).

338

339 **Figure 4.** Structural relationship among (a) orthorhombic Pv-type MgSiO<sub>3</sub> (Brg), (b) LN-type Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub>  
340 (py-LN) and (c) Ilm-type MgSiO<sub>3</sub> (Ak). Thick yellow and thin red arrows indicate the tilting directions of the  
341 octahedra and the shifting directions of the *A*-site cations during the Brg to LN transition, respectively. The  
342 structure of Brg is viewed down the *a* direction, whereas, for comparison, py-LN and Ak are viewed in the

343  $[1\bar{1}2]_{\text{hex}}$  and  $[012]_{\text{hex}}$  plane. Red balls represent the oxygen atoms bonded to the A-site cations in Brg and py-LN.  
344 The two dashed white lines in Brg represent the two longer A-site cation-oxygen bonds of the  $AO_8$  polyhedron.  
345  $AO_6$  and  $BO_6$  octahedra are shown for Ak. Dashed lines in Ak indicate the directions of the alternating  $AO_6$  and  
346  $BO_6$  layers.  
347

348 **Table 1.** Structural parameters of LiNbO<sub>3</sub> (LN)-type Mg<sub>2.98(2)</sub>Al<sub>1.99(2)</sub>Si<sub>3.02(2)</sub>O<sub>12</sub> refined in space group *R3c* (No.  
 349 161)

350	Site	Wyckoff site	<i>g</i> (Mg or Si)	<i>g</i> (Al)	<i>x</i>	<i>y</i>	<i>z</i>	<i>U</i> <sub>iso</sub> (Å <sup>2</sup> )
351	<i>A</i> <sup>a</sup>	<i>6a</i>	0.75 <sup>b</sup>	0.25 <sup>b</sup>	0	0	0	0.0089(7)
352			0.75 <sup>b</sup>	0.25 <sup>b</sup>	0	0	0	0.0160(8)
353	<i>B</i> <sup>a</sup>	<i>6a</i>	0.75 <sup>b</sup>	0.25 <sup>b</sup>	0	0	0.2154(1)	0.0038(8)
354			0.75 <sup>b</sup>	0.25 <sup>b</sup>	0	0	0.2145(1)	0.0113(9)
355	O	<i>18b</i>	-	-	0.3473(5)	0.0681(4)	0.1035(3)	0.0067 <sup>c</sup>
356			-	-	0.3430(5)	0.0613(4)	0.1060(3)	0.0242(8)

357 For each atom the values obtained from the oxygen-*U*<sub>iso</sub>-fixed refinements are reported in the upper lines,  
 358 whereas those obtained from the oxygen-*U*<sub>iso</sub>-free refinements are reported in the lower line.

359 For the oxygen-*U*<sub>iso</sub>-fixed refinements we obtained unit-cell lattice parameters: *a* = 4.8194(3) Å, *c* =  
 360 12.6885(8) Å, *V* = 255.23(3) Å<sup>3</sup>, *Z* = 6, *V*<sub>*m*</sub> (molar volume) = 25.617(3) cm<sup>3</sup>/mol and *D* = 3.934(1) g/cm<sup>3</sup>.

361 Discrepancy factors: *R*<sub>wp</sub> = 4.108%, *R*<sub>e</sub> = 3.253%

362 LN phase: *R*<sub>B</sub> = 1.188%, *R*<sub>F</sub> = 0.506%

363 SiO<sub>2</sub> stishovite: *R*<sub>B</sub> = 1.420%, *R*<sub>F</sub> = 1.020%

364 For the oxygen-*U*<sub>iso</sub>-free refinements we obtained unit-cell lattice parameters: *a* = 4.8194(3) Å, *c* = 12.6887(8)  
 365 Å, *V* = 255.24(3) Å<sup>3</sup>, *Z* = 6, *V*<sub>*m*</sub> = 25.618(5) cm<sup>3</sup>/mol and *D* = 3.934(1) g/cm<sup>3</sup>.

366 Discrepancy factors: *R*<sub>wp</sub> = 3.615%, *R*<sub>e</sub> = 3.255%

367 LN phase:  $R_B = 0.838\%$ ,  $R_F = 0.396\%$

368 SiO<sub>2</sub> stishovite:  $R_B = 3.693\%$ ,  $R_F = 1.274\%$

369 
$$R_{wp} = \left\{ \frac{\sum_i w_i [y_i - f_i(x)]^2}{\sum_i w_i y_i^2} \right\}^{1/2}, R_B = \frac{\sum_K |I_0(\mathbf{h}_K) - I(\mathbf{h}_K)|}{\sum_K I_0(\mathbf{h}_K)}, R_F = \frac{\sum_K ||F_0(\mathbf{h}_K)| - |F(\mathbf{h}_K)||}{\sum_K |F_0(\mathbf{h}_K)|}, R_e = \left\{ \frac{N-P}{\sum_i w_i y_i^2} \right\}^{1/2}$$

370 where  $y_i$ ,  $w_i$  and  $f_i(x)$  are the intensity observed at step  $i$ , the statistical weight and theory intensity, respectively.

371  $I_0(\mathbf{h}_K)$ ,  $I(\mathbf{h}_K)$ ,  $F_0(\mathbf{h}_K)$  and  $F(\mathbf{h}_K)$  are the observed and calculated intensities and structure factors for reflection  $K$ ,

372 respectively.  $N$  and  $P$  are number of all data points and refined parameters, respectively.

373  $g(M)$ : site occupancy of M.

374 <sup>a</sup>A and B atom positions are occupied by Mg or Al and Si or Al, respectively.

375 <sup>b</sup>Site occupancies are fixed.

376 <sup>c</sup>Isotropic atom displacement parameter of oxygen is fixed to that of LiNbO<sub>3</sub> reported by Hsu et al. (1997).

377

378 **Table 2.** Volumes and densities of Mg<sub>3</sub>Al<sub>2</sub>Si<sub>3</sub>O<sub>12</sub> compounds.

379		$V_m$ (cm <sup>3</sup> /mol)	$V$ (Å <sup>3</sup> /formula unit)	$D$ (g/cm <sup>3</sup> )	Reference
380	Pyrope	113.13(3)	46.959(3)	3.564(1)	Geiger et al. (1987)
381	py-Ak	104.63(5)	43.437(5)	3.853(1)	Akaogi et al. (2002)
382	py-LN	102.47(5)	42.538(5)	3.934(1)	This study
383	py-Brg <sup>a</sup>	99.5(7)	41.3(7)	4.06(3)	Liu et al. (2016)

384 Py-Ak, py-LN and py-Brg represent ilmenite, LiNbO<sub>3</sub> and perovskite phases with pyrope composition,  
385 repectively.

386 <sup>a</sup>Unit-cell lattice parameters of Brg with Al<sub>2</sub>O<sub>3</sub> content (~29±5 mol%) from Liu et al. (2016).

387

388

389 **Table 3.** Interatomic distances and angles of LiNbO<sub>3</sub>-type Mg<sub>2.98(2)</sub>Al<sub>1.99(2)</sub>Si<sub>3.02(2)</sub>O<sub>12</sub>.

390	Bond length (Å)					
391	$A-O^i \times 3$	2.021(3)	$B-O^{iii} \times 3$	1.587(2)		
392		2.031(3)		1.629(3)		
393	$A-O^{ii} \times 3$	2.064(2)	$B-O^{vi} \times 3$	2.092(2)		
394		2.035(3)		2.056(3)		
395	Average	2.043	Average	1.840		
396		2.033		1.843		
397	BVS (A)	2.27	BVS (B)	4.23		
398		2.33		3.96		
399			<sup>IV</sup> BVS (B) <sup>a</sup>	3.66		
400				3.33		
401	$n_{c,A}$	5.98	$n_{c,B}$	3.13		
402		6.00		3.40		
403						
404	Bond angles (°)					
405	$O^i-A-O^{iii}$	82.3(1)	$O^{ii}-A-O^{iv}$	105.9(1)	$O-A-O^v$	160.8(2)
406		82.1(1)		106.5(1)		159.6(2)
407	$O^{vi}-B-O^{vii}$	102.3(2)	$O-B-O^i$	89.0(1)	$O-B-O^{iii}$	163.2(2)

408		101.0(2)		89.07(6)		164.9(2)
409	$A-O-B^v$	103.7(2)	$A-O-B$	83.27(1)	$A-O^{viii}-B^v$	86.6(1)
410		101.5(2)		83.40(8)		83.41(8)
411	$A-O^{vii}-A^{ix}$	117.6(1)	$B-O^{vii}-B^{vii}$	143.3(1)		
412		118.5(1)		143.0(2)		

---

413 Symmetry codes: (i)  $-y, x-y, z$ . (ii)  $-x+y+2/3, y+1/3$ . (iii)  $-x+y, -x, z$ . (iv)  $x+2/3, x-y+1/3, z+5/6$ . (v)  $-y+2/3,$   
414  $-x+1/3, z+5/6$ . (vi)  $x+1/3, x-y+2/3, z+1/6$ . (vii)  $-y+1/3, -x+2/3, z+1/6$ . (viii)  $-x+y+2/3, y+1/3, z+5/6$ . (ix)  $x+2/3,$   
415  $y+1/3, z+1/3$ .

416 The upper and lower lines reported for each site represent the oxygen- $U_{iso}$ -fixed and oxygen- $U_{iso}$ -free  
417 refinements, respectively.

418 BVS: bond valence sum value

419  $n_c$ : effective coordination number

420  $^{alV}BVS(B)$  represents the value calculated with three shorter and one longer oxygen- $B$  cation bond lengths.

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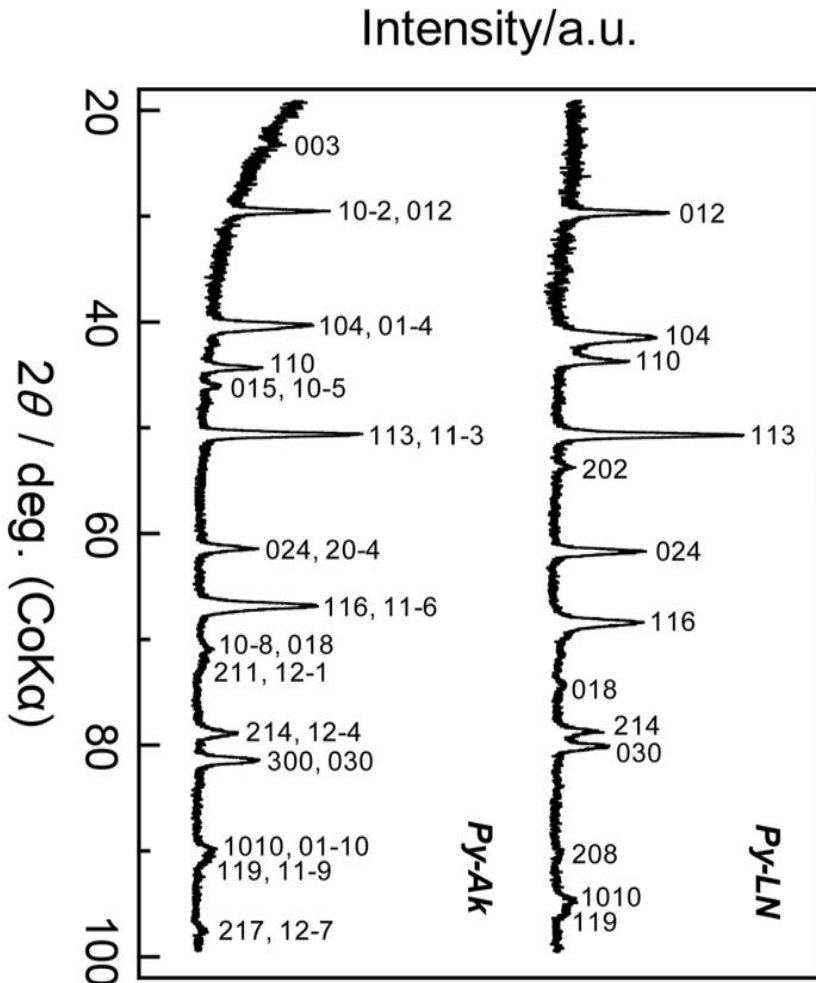


Fig. 1

Fig. 2

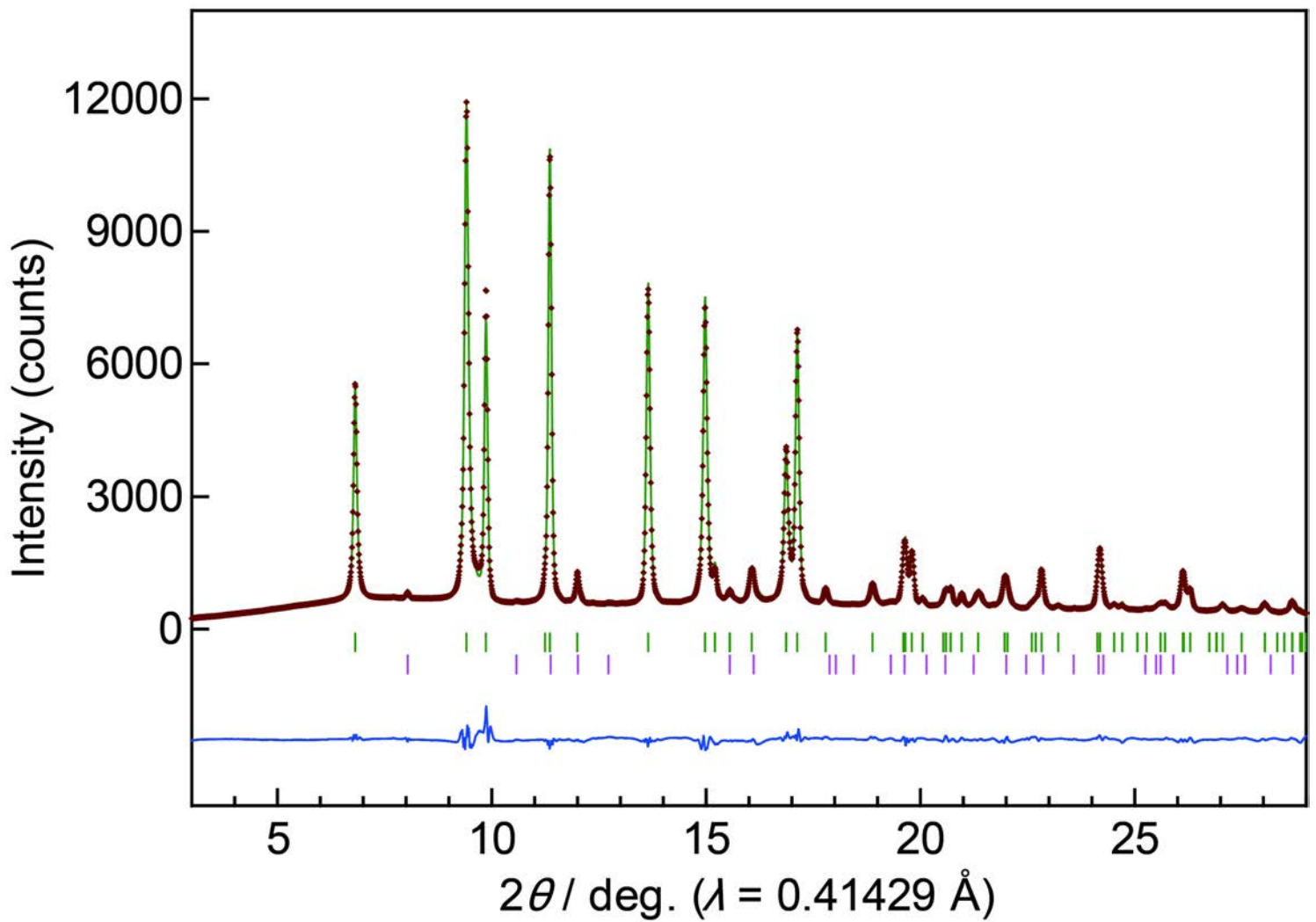


Fig. 3

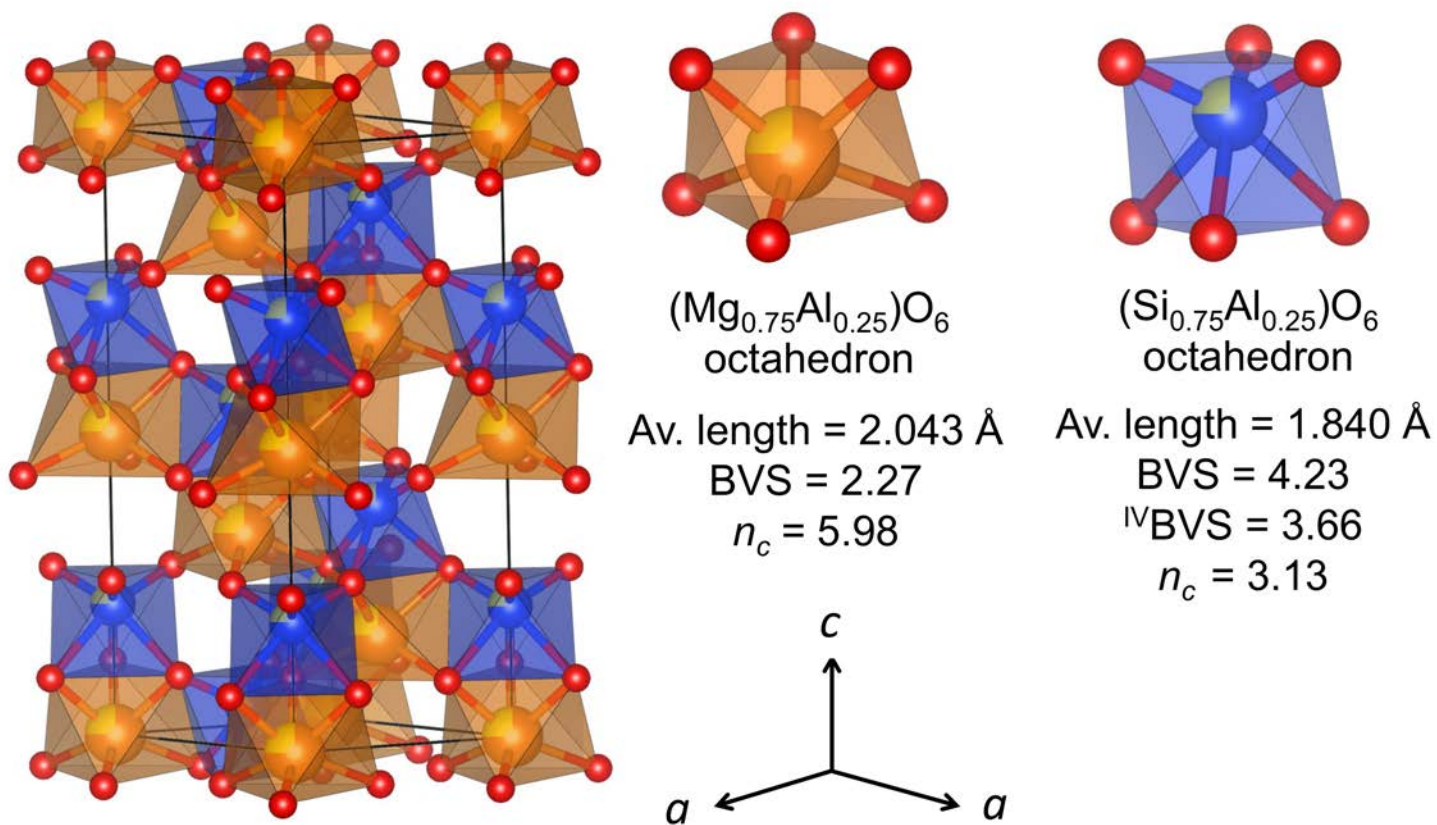


Fig. 4

