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3	Study on structure variations of incommensurately modulated labradorite feldspars
4	with different cooling histories
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# 26 Abstract

The incommensurately modulated structures of three intermediate plagioclase feldspars with 27 28 compositions of  $\sim An_{51}$  are determined by single crystal X-ray diffraction analyses. The samples selected cover a range of different cooling rate, from relatively fast to extremely slow. The 29 structures show various ordering states that are directly correlated with the cooling histories of 30 the samples. The slowest cooled sample shows an e1 structure with strong density modulation, 31 32 along with nanoscale exsolution lamellae. The fastest cooled sample displays an e2 structure, without second-order satellite reflections (*f*-reflections) and density modulation. The sample with 33 intermediate cooling rate shows a less ordered  $e_1$  structure with weak density modulation, but 34 the modulation period and orientation are the same as in  $e^2$  structure. The comparison of the 35 structures with the same composition reveals the ordering process and phase transitions during 36 the cooling of plagioclase within the compositional range of Bøggild intergrowth. New 37 parameters from modulation waves can be used for quantifying the ordering state of plagioclase 38 39 feldspars. Proposed phase relationship and T-T-T diagram for ~An<sub>51</sub> plagioclase feldspars are illustrated for explaining the relationship among  $C\overline{1}$ , e1 and e2 structures and relative cooling 40 41 rates of their host rocks.

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43 Keywords: labradorite, incommensurate, modulated structure, density modulation, gabbro, e-

44 plagioclase, cooling rate, ordering state, exsolution lamellae, intermediate plagioclase

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# 47 Introduction

Plagioclase feldspar ( $Ca_xNa_{1-x}Si_{3-x}Al_{1+x}O_8$ ), a coupled solid solution between albite (NaAlSi<sub>3</sub>O<sub>8</sub>) 48 49 and anorthite (CaAl<sub>2</sub>Si<sub>2</sub>O<sub>8</sub>), is the most abundant group of minerals in the Earth's crust. The intermediate labradorite feldspars (An<sub>50</sub>~An<sub>70</sub>) are common in mafic rocks. Plagioclase with 50 composition between ~An<sub>25</sub> to ~An<sub>75</sub> often show satellite diffractions (*e*-reflections) surrounding 51 the absent *b*-reflections, indicating an incommensurately modulated structure (Bown and Gay 52 53 1959; Ribbe 1983a; Smith and Brown 1988). The aperiodic structure of intermediate plagioclase 54 at low temperature and its complicated subsolidus phase relations have puzzled mineralogists 55 and petrologists for more than 70 years. Many structure models of the aperiodic structure have been proposed and refuted since the discovery of aperiodic plagioclase (Chao and Taylor, 1940). 56 Recent Z-contrast imaging study indicates polarity of the structure in a bytownite sample by Xu 57 (2015), and is confirmed by a high quality structure refinement of a metamorphic labradorite 58 59 sample by Jin and Xu (2017). The structure of e-plagioclase is composed of periodic I1 like domains connected by inversion twin boundaries, instead of  $I\overline{1}$  domains related by antiphase 60 boundaries. 61

62 Incommensurately modulated structures are quite common in nature, and have been an important subject of crystallography for a long time. The modulated structure consists of a lattice periodic 63 64 basic structure and a periodic modulation (Janssen et al. 2007). The structure of each unit cell of the basic lattice structure is modified by the modulation function from one to the next. The 65 66 structure is incommensurate when the ratio between the modulation period and the basic unit cell length is an irrational number. As the modulation function would never meet the basic lattice 67 structure at the same point, the whole structure become aperiodic. The aperiodicity would be 68 reflected in reciprocal space by satellite reflections that cannot be indexed with three integer 69 numbers. An extra dimension has to be introduced for describing incommensurately modulated 70 71 structure, where each parameter of the structure (occupancies, coordinates) is described by a periodic function. Characterization of incommensurately modulated structure in (3+1)D space is 72 a fully developed theory (van Smaalen 2007; Wagner and Schoenleber 2009), and powerful and 73 user-friendly software is available (Petříček et al. 2014). 74

Despite recent studies, many questions about the phase relations remain unsolved. The
 incommensurately modulated plagioclase, or *e*-plagioclase, has been categorized into two types,

e1 and e2, based on the existence of second order satellite reflections, or f-reflections (Ribbe, 77 78 1983b; Smith, 1984; Carpenter, 1994; McConnell, 2008). In the diffraction pattern of an el structure, not only can the *e*-reflections surrounding the absent *b*-positions be observed, but also 79 80 second order satellite reflections, *f*-reflections appear around the main *a*-reflections; whereas in 81 e2 structure, only the first order satellites, e-reflections are present. Almost all previous samples studied by single crystal X-ray diffraction were e1 with observable f-reflections. Little is known 82 83 about the other category, the  $e^2$  structure, since only one structure refinement was done on this 84 structure (Steurer and Jagodzinski 1988). The e1 structure is normally believed to only exist in 85 the Ca-rich plagioclase (Ribbe, 1983b; Smith, 1984; Smith and Brown, 1988), but Xu et al. (2016) discovered a plagioclase sample with a composition of  $\sim An_{45}$  that displays e1 structure, 86 which makes the case more complicated. It was proposed that the  $C\bar{1} \rightleftharpoons e2$  transition might be 87 thermodynamically continuous, even though the transition temperature is too low to make it 88 reversible in the lab (Carpenter, 1986). This makes the relationship between e1 and e2 structure 89 90 more problematic: is the difference between these two structures large enough to create a 91 miscibility gap between them, or are they essentially the same phase, with only gradually 92 decreasing intensities of satellite reflections?

93 The structure refinement of the labradorite from a metamorphic rock (Jin and Xu, 2017) reveals the ordering pattern of e-plagioclase in another dimension, which opens the door to a whole new 94 95 level of understanding the ordering states of plagioclase. Instead of using only average Al 96 occupancy of T<sub>1</sub>o site, the modulation of T-O bond distances, Ca-Na occupancies and density 97 modulation could all be used for quantifying the ordering state of the structure with much higher 98 sensitivity and accuracy. However, an individual structure of this composition can only provide 99 us the stability state of e-plagioclase, but nothing about the phase transition and ordering 100 mechanism. As the diffusion between Al and Si within the crystal lattices is extremely difficult (Grove et al., 1984; Liu and Yund, 1992), which makes experiments for these phase transitions 101 102 impractical in the laboratory, such information can only be preserved in plagioclase of igneous origin with a composition within the range of Bøggild intergrowth. In this paper, three carefully 103 chosen igneous plagioclase samples with modulated structure are analyzed with single crystal X-104 105 ray diffraction. The comparison between the samples with similar composition but different 106 cooling histories would provide important information about the phase transitions as well as

107 ordering mechanism of intermediate plagioclase, which are critical for mapping out the phase108 relations around the composition range of Bøggild intergrowth.

# **110** Samples and Experiments

111 Three plagioclase samples of igneous origin with similar composition ( $\sim An_{51}$ ) but different 112 cooling histories were analyzed in this study. All samples were carefully characterized using 113 selected-area electron diffraction (SAED), dark-field TEM imaging, and electron microprobe 114 analysis (EMPA) before they were used for single crystal X-ray diffraction analyses. Details 115 about the compositions of the samples are listed in Supplementary Table. K-feldspar component 116 in the samples are low ( $\sim 2$  mole% of Or).

Sample Dul-15-8B is a melanogranophyre intrusion in the anorthositic layer of the Duluth 117 118 Igneous Complex. The sample was collected from a roadcut on Skyline Parkway on the south side of Twin Ponds swimming area (Stop 1-8 in Miller's (2009) field guide). The intrusion is 119 120 considered to have formed from late-stage melts fluxing through the intercumulus pore spaces of the partially molten anorthositic series of the Duluth Layered Complex (Miller, 2009). The 121 122 intrusion is very small, with a diameter of only about 1 km at the current level of exposure, indicating a relatively fast cooling rate with respect to the main body of the Layered Complex 123 (Miller, 2009). The TEM image of the sample shows a homogeneous crystal with no exsolution 124 lamellae. The electron diffraction pattern shows definite yet weak e-reflections, but no f-125 126 reflections are observed (Fig. 1). The plagioclase in the sample shows a relatively large variation 127 in composition from crystal to crystal, from  $\sim An_{50}$  to  $\sim An_{57}$ .

Sample SK90-12 is from Lower Zone C (top of lower zone) of the Skaergaard layered intrusion, 128 129 which is similar to the Duluth igneous complex in many respects (Taylor, 1964). The size of the intrusion is about 10 km across (Nielsen, 2004; Winter, 2010). Plagioclase crystals in this layer 130 are quite homogeneous from core to rim with a composition of  $\sim An_{50}$  (Toplis *et al.*, 2008). No 131 132 exsolution lamellae are observed by TEM (Fig. 2). First-order satellite reflections (*e*-reflections) are sharp. Second-order satellite reflections (f-reflections) are weak in SAED patterns (Fig. 2). 133 134 Plagioclase from the Skaergaard intrusion are classic samples and have been subjected to extensive analyses for their incommensurate structures (Gay, 1956; Carpenter, 1986, 1994). No 135 136 sign of exsolution was reported in any previous work.

137 Sample R2923 is an iridescent labradorite crystal from the collection of Smithsonian National138 Museum of Natural History, Washington, DC. The sample shows red to dark orange iridescent

139 color, indicating nanoscale exsolution lamellae resulted from extremely slow cooling. TEM images (Fig. 3) show almost periodic exsolution lamellae with alternating thickness. Based on 140 141 EDS results, the wide lamellae are more Ca-rich than the thin lamellae. The compositions are about An54 and An48 respectively. The minor change of modulation direction across adjacent 142 143 lamellae indicates the difference in structure should be relatively small. The electron diffraction 144 pattern (Fig. 3B) also shows only sharp round satellites instead of splitting peaks. Iridescent labradorite feldspars are only found in slowly cooled gabbroic or anorthositic massifs (Smith and 145 Brown, 1988). As the diffusion between Al and Si is extremely sluggish at low temperature 146 (Grove et al., 1984; Liu and Yund, 1992), the rather coarse exsolution lamellae in sample R2923 147 148 prove extremely slow cooling of the host rock.

The composition of the samples was analyzed by electron microprobe around the areas on the thin section with thick albite twin lamellae for selecting crystals for X-ray diffraction. Sample Dul-15-8B and SK90-12 were analyzed with a CAMECA SXFive-FE at 15kV and 10nA beam current with a 10-micron beam size. Sample R2923 were analyzed with a CAMECA SX51 microprobe at 15 kV and 20 nA beam current with a 20-micron beam size. Data for Si, Al, Ca, Na, K and Fe were analyzed using plagioclase and iron oxide standards (*Supplementary Table*).

The TEM samples were prepared by crushing small plagioclase grains selected from area close to the points analyzed by EMPA in ethanol with an agate mortar. Drops of the suspension were placed on lacey-carbon coated Cu grids, and the small flaky crystals are left on the Cu grid after the ethanol evaporates. TEM images and selected-area electron diffraction pattern (SAED) were taken using a Philips CM200UT TEM. Compositions of the samples were examined by attached X-ray energy-dispersive spectroscopy (EDS). The EDS results are calibrated with the results from electron microprobe analyses.

162 The crystals for single crystal X-ray diffraction were picked after microprobe analyses. Areas 163 with composition around  $An_{50-52}$  were selected. Crystals are washed with acetone to clean off the 164 crystalbond from the thin section and any attached small grains, and then mounted on glass fibers 165 with epoxy glue. The X-ray diffraction data were collected on a Bruker Quazar APEXII single 166 crystal diffractometer with Mo K $\alpha$  I $\mu$ S source at 100K. 4  $\omega$  runs and 1  $\varphi$  run were programmed 167 for each sample with a scan width of 0.5°. Exposure times of 65s, 95s, and 65s were applied for 168 Dul-15-8B, SK90-12 and R2923 respectively. The instrument was running at power of 50kV and 169 0.6mA. The detector was at a distance of 5cm from the crystal. The X-ray diffraction data for 170 sample Dul-15-8B was analyzed again on a Bruker D8 VENTURE X-ray diffractometer, equipped with an air-cooled PHOTON II detector using CMOS technology. The X-ray source 171 172 was an Incoatec Mo IµS 3.0 microfocus tube coupled to a multilayer mirror optic. Data were 173 collected at 100K using an Oxford Cryosystems Cryostream 800 low-temperature apparatus to 174 intensify high-angle satellites. This new instrument is applied specifically for sample Dul-15-8B sample to confirm the nonexistence of *f*-reflections (second order satellites). More experimental 175 176 details are provided in Table 1.

Unit cell parameters were calculated and refined using APEX3 software. The structure was
solved with a charge-flipping algorithm (Oszlanyi and Suto, 2004, 2005) using SUPERFLIP
(Palatinus and Chapuis, 2007). The refinement of the average structure and modulated structure
was done with JANA2006 (Petříček *et al.*, 2014). The 3-D crystal structure was visualized by
VESTA (Momma and Izumi, 2011).

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# **183** Structure Refinement

#### 184 Sample Dul-15-8B

After the data were collected on Bruker Quazar APEXII with 65s exposure time, no *f*-reflections were observed. To confirm the extinction of second order satellites, the same crystal was analyzed again on a Bruker D8 VENTURE diffractometer for 90s exposure. With much stronger beam intensity, a more sensitive detector and longer exposure time, still no *f*-reflections were detected. Thus, the second-order satellites were confirmed to be absent. Therefore, no second order component should be allowed in the modulation function.

The modulated structure of Dul-15-8B was refined in the unconventional X1( $\alpha\beta\gamma$ )0 superspace group setting, with c~14Å subcell and a centering condition of (½ ½ 2 0), (0 0 ½ 2), (½ 2 0 ½) (Boysen and Kek, 2015; Jin and Xu, 2017). Symbol ( $\alpha\beta\gamma$ ) indicates the modulation is not constrained by the symmetry of the structure thus should have components along all three unit cell axes (*q*-vector in *Table 1*). The model used for refinement is the same as the metamorphic labradorite (987L) refinement by Jin and Xu (2017), with M2 only occupied by Ca and M1 containing both Ca and Na. The potassium component is not considered in the refinement. The total occupancy of the atoms in M1 and M2 sites are constrained to 1. Since no second-order satellites were observed, all atomic modulation functions are described by simple harmonic functions, including the atomic displacement parameters (ADPs or temperature parameters). The resulting atomic occupancies and positions are listed in *Table 2*. Detailed information about the crystal can be found in *Supplementary Data*.

The Al occupancies of tetrahedral sites are assigned according to the <T-O> bond distances (Å),
following the equation (Kroll and Ribbe, 1983):

$$Occ(Al) = 0.25(1 + n_{An}) + (< T_i - 0 > - \ll T - 0 \gg)/k$$

The value of k is estimated to be 0.135Å, which is conventionally used for albite and plagioclase feldspars (Angel *et al.*, 1990; Jin and Xu, 2017). The resulting occupancies of all tetrahedral sites lie in the reasonable range (between 0 and 1). The modulation of  $\langle T-O \rangle$  bond distances is plotted in Fig. 4A. The average  $\langle T-O \rangle$  bond distances for T<sub>1</sub>o, T<sub>1</sub>m, T<sub>2</sub>o and T<sub>2</sub>m are 1.6852Å, 1.6555Å, 1.6527Å and 1.6581Å, corresponding to Al occupancies of 0.55, 0.33, 0.31 and 0.35 respectively.

211 The displacement modulation of the M site shows a similar pattern as the structure of the metamorphic labradorite sample (987L), with the M2 site (Ca) almost stationary and the M1 site 212 213 (Ca+Na) moving back and forth (Jin and Xu, 2017). The displacement happens mainly within the **b-c** plane as the component along *a*-axis is very minor (Fig. 5A). The amplitude of M1 214 displacement, however, is obviously smaller compared to the structure of sample 987L (Jin and 215 Xu, 2017). The occupational modulation of M site is not as dramatic as in the metamorphic 216 217 labradorite structure (Jin and Xu, 2017), as shown in Fig. 6A. The variation (max-min in Table 2) of occupational modulation of Na in M1 site is 0.07, whereas in 987L the maximum difference in 218 Na occupancy is 0.35. 219

As explained by Jin and Xu (2017), the structure with no second-order satellites (*f*-reflections) would not display a density modulation. As shown in Fig. 7A, where all Na occupancies of all 8 symmetrically equivalent M sites are plotted, the average occupancy shows a virtually constant value. Therefore, the structure is chemically homogenous with only ordering within each subcell. For comparison, the average <T-O> bond distances are also plotted in Fig. 8A. Even though the individual average distance shows some variation (less than 0.005), the overall average bond distance of all the T-sites, which corresponds to the chemical composition (Kroll and Ribbe,1983), is basically constant.

228 A section of the modulated structure consisting of 3x7 anorthite unit cells along *b*- and *c*-axis is presented in Fig. 9. The structure can still be described as periodic domains with polarity related 229 230 by  $I\overline{I}$  inversion twin boundaries, except the inversion boundaries have the same composition as the I1-like domains. The ordering of Si-A1 at inversion twin boundaries shows an anorthite-like 231 pattern, with Al-rich sites surrounded by 4 Si-rich sites and vice versa. The ordering pattern of 232 233 the *I*1-like domains is unique and different from albite and anorthite structures. The Al occupancy of T<sub>1</sub>o sites in the *I*1-like domains is not higher than in the inversion twin boundaries, 234 235 whereas in low-temperature labradorite structure (Jin and Xu, 2017), the T<sub>1</sub>o site in the *I*1-like domains (rich in Na) is higher than the inversion twin boundaries, even though the total Al 236 237 component is lower than that at the inversion twin boundaries.

The M site configuration is always complicated in plagioclase structure. The correlation between 238 239 the split distance of M site and the surrounding tetrahedral framework was observed in the low temperature labradorite structure (Jin and Xu, 2017). As shown in Fig. 10, the M site lies almost 240 241 exactly on the plane defined by the T<sub>2</sub> sites. The split direction of the M site is very close to the diagonal line connecting T<sub>2</sub> sites. The split distance of M site is plotted against the diagonal 242 distance between T<sub>2</sub>oz and T<sub>2</sub>mo and shown in Fig. 11A. The two values are correlated with each 243 other, as the M site is more split where the diagonal T<sub>2</sub> distance is larger. The modulation 244 245 function of  $T_1 \circ O_A 1 - T_1 m$  angle is provided in Fig. 12A. The two angles are tilting towards 246 opposite (inward vs outward), except for when t=0 and t=0.5.

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#### 249 Sample SK90-12

With 95 seconds of exposure time, 284 *f*-reflections were observed in the X-ray diffraction data of Sample SK90-12, which is fewer than the number observed in low-temperature labradorite (sample 987L, 60s exposure) (Jin and Xu, 2017). The structure of Sample SK90-12 is refined in exactly the same model as sample 987L (Jin and Xu, 2017) with the same constrains. Therefore, all parameters are refined with a sum of harmonic functions of first and second order. The result atomic occupancies and positions are listed in *Table 3*. Complete structural information can be
found in *Supplementary Data*.

257 The M site modulation shows very similar pattern to the Dul-15-8B sample, but with an obvious tendency towards the more ordered structure as low-temperature labradorite (Fig. 5B) (Jin and 258 Xu, 2017). The modulation of Na in the structure, modified by the second-order component of 259 260 the modulation function, shows two separate and definite maxima within one period (Fig. 6B). 261 Although the variation of the occupancy modulation is only slightly larger than that in sample Dul-15-8B and much smaller than that in the low-temperature labradorite (Jin and Xu, 2017). By 262 averaging all 8 symmetrically equivalent Na occupancies within one unit cell (c~14Å), a density 263 modulation with variation of  $\sim 8$  mole % of anorthite is displayed (Fig. 7B). 264

265 The <T-O> bond distance modulation pattern also lies somewhere between sample Dul-15-8B 266 and the low-temperature labradorite (987L) (Fig. 4B). The average <T-O> bond distances for T<sub>1</sub>o, T<sub>1</sub>m, T<sub>2</sub>o and T<sub>2</sub>m are 1.6876Å, 1.6547Å, 1.6504Å and 1.6575Å, corresponding to Al 267 occupancies of 0.56, 0.32, 0.29 and 0.34 respectively. The modulated bond distance function of 268  $<T_1$ o-O> and  $<T_1$ m-O> is starting to broaden between t=0.25 and t=0.5. The variation of bond 269 270 distance modulation is also bigger than that of Dul-15-8B. The  $\langle T_1 \circ O \rangle$  bond, for example, has a variation of ~0.1Å between the maximum and minimum values, whereas the variations in Dul-271 15-8B and low-temperature labradorite (987L) are ~0.08Å and ~0.12Å respectively (Jin and Xu, 272 273 2017). The difference is even more obvious when we average all the symmetrically equivalent sites within one unit cell (Fig. 8B). The average  $\langle T_2 - O \rangle$  distances and the overall average 274 275 obviously show two maxima and two minima in one modulation period. The averages resemble 276 the pattern in the low-temperature labradorite sample 987L (Jin and Xu, 2017) but are different from the pattern in sample Dul-15-8B. The correlation between the split distance of M site and 277 the diagonal distance of  $T_2$  sites are shown in Fig. 11B. Both the variation of the M1-M2 278 279 distance and the diagonal T<sub>2</sub> distance are larger than in Dul-15-8B. The modulation of bond 280 angles also shows larger amplitude as plotted in Fig. 12.

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# 284 Sample R2923

Iridescent color in labradorite feldspars are generated by nanoscale semiperiodic exsolution 285 lamellae within the crystal. The thickness, or period of the lamellae is at the same scale with the 286 287 wavelength of visible light. Therefore, light of certain color/wavelength would be reflected by 288 the exsolution lamellae boundaries and interfere with each other, and display beautiful iridescent 289 colors. Thicker exsolution lamellae would generate iridescent color with longer wavelength. 290 Almost all previous labradorite samples used for single crystal X-ray analyses display blue iridescent color (Horst et al., 1981; Yamamoto et al., 1984; Boysen and Kek, 2015). This was 291 292 considered a serious problem (Ribbe, 1983a), as the reflections from the two components overlap 293 and therefore any structure refined is a composite of two structures of different compositions. The difference of the *e*-fringes across the exsolution lamellae is measured to be at least 10 294 295 degrees in a sample with blue iridescent color (McConnell, 1974; Wenk and Nakajima, 1980; Ribbe, 1983a). However, all igneous labradorite within the range of Bøggild intergrowth that are 296 cooled extremely slowly would display exsolution lamellae. Therefore, the structure of sample 297 298 R2923 was refined as a single phase, simply for the purpose of comparison.

As observed with TEM (Fig. 2), even though the exsolution lamellae are clearly seen, the 299 300 difference of *e*-fringes across the exsolution lamellae is very minor, indicating a minor difference 301 in composition and structure between the two phases. And the structure is dominated by the Ca-302 rich phase, unlike the sample with blue iridescent color where the two phases are almost of the 303 same proportion. Therefore, the labradorite with red iridescent color should be a better subject 304 for single crystal diffraction, and the refinement result should be more reliable than that of a sample with blue iridescent color. With an R value of 0.07 for f-reflections (second-order 305 306 satellites) and 0.03 for all reflections, the degree of fit of this refinement exceeds all previous 307 iridescent labradorite structures, which suggests that the resulting structure should be a good 308 approximation of both exsolution lamellae.

There were less than 300 third-order satellites observed during the data collection (Fig. S1 in Supplementary Material). This is very likely to have resulted from the suppressed thermal vibration and enhanced intensities by low temperature data collection. But considering they are weak and not from a pure phase, and also for the convenience of comparison, the third order satellites were not used in the refinement. The structure was refined in the same model with the

same conditions as the low-temperature labradorite sample (987L), except only the first-order 314 harmonic wave was used for the atomic displacement parameters (ADPs or temperature 315 parameters) modulation to avoid unrealistic results. The chemical composition based on the 316 refined M site occupancies is about  $An_{49}$ , which is obviously lower than the microprobe analysis 317 result. This disagreement is hardly surprising as the reflections used for refinement are simply a 318 sum integration of overlapped reflections from two different lamellae phases. Nonetheless, to get 319 a better idea of the structure, we constrained the total occupancy of Ca in M site to be 0.525 320 321 based on the microprobe analyses (Supplementary Table). The atomic occupancies and positions are listed in Table 4. Complete information of the modulated structure can be found in 322 323 Supplementary Data.

The resulting structure of R2923 is very similar to the low-temperature labradorite (sample 987L) 324 (Jin and Xu, 2017). The modulation orientations and periods of the two samples are basically the 325 same. The modulation of M site also matches that in the metamorphic labradorite (sample 987L) 326 327 (Jin and Xu, 2017) very well. Both the occupancy (Fig. 6C) and displacement (Fig. 5C) modulation are very similar to the modulation in sample 987L, even the electron density contour 328 (Fig. 5C) matches with 987L surprisingly well. And by averaging all 8 symmetrically equivalent 329 Na occupancies within one unit cell (c~14Å), a density modulation with variation of ~17 mole% 330 of anorthite is displayed (Fig. 7C). 331

The <T-O> bond distance modulation is also indistinguishable from the 987L sample (Fig. 4C 332 and Fig. 8C) (Jin and Xu, 2017). The modulation functions are obviously flattened in certain part 333 334 of the modulation period (Fig. 4C), which distinctly distinguish the structure from sample Dul-15-8B. Based on the same function, the average  $\langle T-O \rangle$  bond distances for T<sub>1</sub>o, T<sub>1</sub>m, T<sub>2</sub>o and 335 T<sub>2</sub>m are 1.6935Å, 1.6516Å, 1.6508Å and 1.6571Å, corresponding to Al occupancies of 0.61, 336 0.29, 0.29 and 0.34 respectively. The correlation between M site split distance and diagonal T<sub>2</sub> 337 338 distance displays an elongated loop compare to that of Dul-15-8B and SK90-12 (Fig. 11C). The modulation of T<sub>1</sub>o-O<sub>A</sub>1-T<sub>1</sub>m angle also shows larger amplitude compared to Dul-15-8B and 339 340 SK90-12 Fig. 12C.

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# 344 Discussion

345 The *e*-plagioclase feldspars were previously divided into two categories based on the existence of f-reflections, namely e1 and e2 (Ribbe, 1983b; Smith and Brown, 1988; Carpenter, 1994). The 346 compositional difference was believed to be the reason causing this difference, as *f*-reflections 347 barely appear in Na-rich e-plagioclase but are very common in Ca-rich e-plagioclase (Bown and 348 349 Gay, 1959; Smith and Brown, 1988). Many previous refinements of the structure only considered first-order satellites (e-reflections) and neglected the weak f-reflections (Horst et al., 1981; 350 351 Yamamoto et al., 1984), and thus failed to resolve any difference between e1 and e2 structures. The discovery of f-reflections in a Na-rich plagioclase ( $\sim An_{45}$ ) raised the complexity of the case, 352 353 and revealed cooling rate as an important factor in addition to composition in determining the structure of e-plagioclase (Xu et al., 2016). Because the satellites in Na-rich plagioclase samples 354 355 are usually weak, which makes data collection and structure refinement very difficult, only one e2 structure (An38) refinement has been previously published (Steurer and Jagodzinski, 1988). 356 357 Nonetheless, since there is no evidence suggesting otherwise, we are categorizing the structure of 358 sample Dul-15-8B as e2, which is the same as most Na-rich e-plagioclase feldspars.

359 With a sufficiently long X-ray beam exposure time for sample Dul-15-8B, we can confirm the extinction of second order satellites (*f*-reflections) and the absence of density modulation, which 360 means the difference between e1 and e2 is distinct and they should be categorized as two 361 separate modulated structures. On the other hand, the structure refinement does show great 362 similarities between two phases, which made the differences hard to uncover when f-reflections 363 364 were neglected. Other than the absence of density modulation, there are a few noticeable differences in the  $e^2$  structure from  $e^1$ . The amplitudes of the modulated parameters are 365 systemically smaller than those of the e1 structure. For instance, the  $<T_10-O>$  bond distance 366 modulation of sample Dul-15-8B dropped nearly 30% in variation from the iridescent labradorite 367 368 (sample R2923). The average occupancies in tetrahedral sites are also different. Average Al occupancy in T<sub>1</sub>o site is 0.55 for sample Dul-15-8B, which is obviously smaller than 0.59 in low-369 370 temperature labradorite (sample 987L) (Jin and Xu, 2017) and 0.61 in the iridescent labradorite (sample R2923). 371

The modulation period is another important difference between e1 and e2 structures. With very similar chemical composition, the modulation period of sample Dul-15-8B is ~34Å, which is about 4Å longer than that of the iridescent labradorite sample R2923 and low-temperature labradorite (sample 987L) (Jin and Xu, 2017). This result confirms the discontinuity observed by Slimming (1976) in electron diffraction patterns of *e*-plagioclase feldspars with different compositions. The orientation of the modulation is also different from more ordered *e*plagioclase (*e*1), tilting towards the Ca-rich side on the trend (Fig. 13).

379 It is worth noting that even though sample SK90-12 starts to show weak f-reflections, the 380 modulation period and orientation are still close to sample Dul-15-8B (Fig. 13). This indicates the crystal was transformed from e2 into e1 phase, where density modulation is starting to form 381 on the foundation of the e2 modulation period and orientation. Carpenter plotted the modulation 382 383 period against the composition for igneous samples from Skaergaard and metamorphic samples 384 from Broken Hill (Figure 15 in Carpenter, 1994), based on data from Gay (1956) and Slimming (1976). The metamorphic samples show a discontinuous trend whereas the igneous samples 385 386 show a continuous one. As the SK90-12 sample in this study shows no evidence of exsolution, the divergence of the trends can be easily explained by the different trend of e1 and e2, instead of 387 388 introduced by Bøggild intergrowth. While the Skaergaard sample may start to develop  $e^{1}$ ordering from e2 structure, it still preserves the modulation period and orientation of e2. As 389 390 Carpenter (1994) proposed, once a given structure is established ( $e^2$  in this case), the energy 391 advantage of reorientation will probably be small compared to degree of structural reorganization 392 required. The metamorphic Ca-rich plagioclase samples, on the other hand, formed directly at the 393 temperature where  $e_1$  structure is stable, therefore the modulation period and orientation follows a different trend from e2. So the data points in Figure 15 of Carpenter (1994) should follow two 394 separate trends: where all the Skaergaard samples and Na-rich metamorphic samples follow the 395 396 nice e2 trend, and the Ca-rich metamorphic samples follow the e1 trend.

With the structure of e-plagioclase solved, there is no reason to still think that e-plagioclase is a metastable stage leading to an unmixing of the structure into albite and anorthite. First of all, the modulation period does not get longer as the cooling rate of e-plagioclase decreases, whereas in exsolution process, the exsolution lamellae tends to grow thicker through time; the modulation period of e-plagioclase gets longer as the Ca component of the composition increases, which also distinguish the ordering of *e*-plagioclase from an unmixing process; and finally, spinodal decomposition is a non-equilibrium process, and should produce a relatively randomly distributed wavelength, instead of the perfectly ordered structure with strictly constrained period and orientation we see in *e*-plagioclase structure. Therefore, there is no doubt that *e*-plagioclase structures are thermodynamically stable for intermediate plagioclase at low temperature.

407 All the parameters used to estimate the ordering state of intermediate plagioclase were based on the average structure, such as  $\langle T_1 o \cdot O \rangle$  bond distance, gamma angle of the lattice parameters 408 (Kroll, 1983; Smith and Brown, 1988). Although Al occupancy in T<sub>1</sub>o site can indicate ordering 409 sates of alkali feldspars, it does not work well for intermediate plagioclase like labradorite 410 feldspars based on their average structures. Intermediate plagioclase feldspars display structural 411 modulation that characterized by additional e- and f-reflections. As the three samples refined in 412 this paper clearly show an increasing ordering state, and the structures are described in great 413 detail, we should be able to extract some more accurate parameters to quantify the ordering states 414 of the plagioclase structures. Some potential parameters are listed in Table 5. Variation of Na 415 occupancy, variation of density modulation, and variation of  $\langle T_1 o - O \rangle$  bond distance modulation 416 of the three samples are plotted against the gamma angles in Fig. 14, where clear correlations 417 418 between the parameters are shown. The gamma angle is a simple way for qualifying the ordering 419 state of intermediate plagioclase feldspar. However, it cannot tell any detailed ordering state 420 within the structure, let alone the difference among  $e_1$ ,  $e_2$  and  $C_{\overline{1}}$ . The structural modulation on 421 the other hand, shows relatively large variation among different ordering states, and clearly shows the difference among the three labradorite feldspar phases. The calorimetry work by 422 Carpenter et al (1985) reveals the significant enthalpy change associated with the ordering of e-423 plagioclase. The enthalpy difference between a disordered plagioclase with  $C\overline{1}$  symmetry and a 424 425 highly ordered *e*-plagioclase can be up to  $\sim$ 3 kcal/mole. The ordering states of the samples used 426 in their paper were only qualitatively characterized by the intensity and sharpness of the satellite 427 reflections in the electron diffraction pattern, which are very difficult to compare with one 428 another due to the difference in thickness and orientation of the TEM sample. However, with 429 single crystal diffraction data, the ordering enthalpy and the modulated structure could be directly connected, leading us to a comprehensive and detailed picture of the thermodynamics of 430 431 plagioclase feldspar solid solution.

432 The discovery of an  $e^2$  structure with Ca-rich composition (~An<sub>52</sub>) raises another problem that 433 needs to be solved, as e2 structure does not appear at this composition in any previously proposed phase diagrams. It was reported that disordered  $C\overline{1}$  structure is stable at high 434 temperature for this composition (Smith and Brown, 1988; Carpenter, 1994), and homogenous 435 e1 structure is stable at low temperature (below ~550°C) (McConnell, 2008; Jin and Xu, 2017). 436 Disordered  $C\overline{1}$  structure transforms to  $e^2$  structure on the Na-rich side of Bøggild intergrowth 437 (Carpenter, 1994; McConnell, 2008). Therefore, the solvus for Bøggild intergrowth must be 438 shaped as a loop which closes at low temperature, and serves as a first-order phase transition 439 boundary between disordered  $C\overline{1}$  structure and  $e_1$  structure (McConnell, 2008; Jin and Xu, 2017; 440 Xu et al., 2016). In this scenario, e2 cannot be a stable phase at any temperature for this 441 composition, but is more likely to be a metastable phase that is kinetically preferable when 442 443 exsolution is suppressed. Exsolution requires interdiffusion between CaAl and NaSi, which is extremely sluggish (Grove et al., 1984; Liu and Yund, 1992). A schematic diagram showing the 444 topology of these phases is provided in Fig. 15. The phase boundary between disordered  $C\overline{1}$ 445

structure and e2 is extended as the dash line. A corresponding time-temperature-transformation 446 (TTT) diagram for plagioclase of  $\sim An_{51}$  composition is also illustrated in Fig. 16. If the crystal is 447 quenched from high temperature (an erupted volcanic rock, for instance), the disordered  $C\overline{1}$ 448 structure, which is stable only at high temperature, would be preserved (path A in Fig. 16). In a 449 quickly cooled host rock (as in sample Dul-15-8B), a disordered  $C\overline{1}$  structure will not exsolve 450 into  $e_1$  structure, but transforms to  $e_2$  instead (path **B** in Fig. 16). As the cooling process gets 451 slower (path C in Fig. 16), the e2 structure will end up supercooled in the area where e1 is stable, 452 and further ordering results in  $e_1$  structure while the  $e_2$  modulation period and orientation is still 453 preserved (as in sample SK90-12). As for an extremely slow-cooled sample (sample R2923), 454 exsolution would take over once it enters the two-phase area of  $C\overline{1}$  structure and e1 and form 455 Bøggild intergrowth (path **D** in Fig. 16). When the temperature gets below the loop solvus, the 456 Ca-poor lamellae would also transform from  $C\overline{1}$  structure to e1, and the crystal will try to 457 homogenize into one single e1 phase. However, the temperature then would be too low for the 458 459 process to go to completion, therefore the exsolution lamellae are preserved.

460 Carpenter (1986) observed *e* ordering above the reversed  $e \rightleftharpoons I\overline{1}$  transition with starting crystals of 461  $C\overline{1}$  structure. The structure will transform eventually to  $I\overline{1}$  structure if held long enough at the

same temperature. This was explained, reasonably, as a metastable structure appearing in the 462 stable field of  $I\overline{1}$  structure, but the metastable *e* structure was presumed to be *e*1 based on the 463 composition in his paper. However, since ordering of e1 structure requires atom exchange 464 outside a unit cell, it is unlikely that e1 is kinetically preferable to  $I\overline{1}$  structure. The big 465 miscibility gap between  $C\overline{1}$  and e1 phase also supports this, while no gap has been observed 466 between  $C\overline{1}$  and  $I\overline{1}$  (Carpenter and McConnell, 1984). By comparing the structure of  $C\overline{1}$ , e2 and 467  $I\overline{1}$ , it is obvious that the e2 structure is actually closer to  $C\overline{1}$  structure. Both  $C\overline{1}$  and e2 structure 468 are chemically homogenous without density modulation; transition from  $C\overline{1}$  to e2 only requires 469 470 total reordering of tetrahedral sites of the inversion twin boundaries and minor adjustment for the Il-like domains, whereas  $I\overline{1}$  requires reordering of tetrahedral sites of every  $C\overline{1}$  unit cell 471 throughout the crystal. The T site bond distance modulations in sample Dul-15-8B show a 472 considerable extent of ordering. However, the transition from  $C\overline{1}$  to e2 might be initiated by only 473 M site ordering, as the volcanic plagioclase sample from Hogarth Ranges already starts to show 474 475 weak e-reflections (Fitz Gerald et al. 1986). Therefore, e2 structure is very likely to be more kinetically favorable than  $I\bar{1}$  structure for a starting crystal with  $C\bar{1}$  structure at certain 476 temperature. As shown in Fig. 15, the extended boundary between  $C\overline{1}$  and  $e^2$  might be above the 477 boundary between  $I\overline{1}$  and e1. Nonetheless, it would require careful measurement of the 478 modulation period and orientation of the experiment product in Carpenter's (1986) paper to 479 480 confirm this theory.

481

# 482 Implications

The single crystal X-Ray analyses of e1 and e2 structures with similar compositions reveals the 483 complicated phase relations around Bøggild intergrowth. The e2 structure is confirmed to be a 484 distinctly different phase from  $e_1$ , but the two modulated structures are similar in many ways. 485 486 The e2 structure is characterized by the absence of second-order satellites (f-reflections), indicating a less ordered state compared to the e1 structure. It can also be described as periodic 487 Il-like domains related by  $I\overline{1}$  like inversion boundaries. However, with no density modulation, 488 489 the domains are of the same chemical composition as the inversion twin boundaries. The variations of modulated parameters in e2 structure are systematically smaller than e1 structure. 490

491 The average Al occupancy in  $T_1$  o site of *e*2 structure is smaller compared to *e*1 structure for the 492 same composition, and the modulation period is longer than in the *e*1 structure.

The structure of e2 is closer to  $C\overline{1}$  structure, for only the tetrahedral sites of inversion boundaries 493 need to be completely rearranged. This similarity makes it kinetically preferable for a crystal 494 cooled from a  $C\overline{1}$  structure. The e2 phase normally exists as a stable phase for Na-rich 495 plagioclase at a relatively low temperature. Nevertheless, it can also appear as a metastable phase 496 for Ca-rich plagioclase, as the energy barrier for the phase transition from  $C\overline{1}$  to e1 or even  $I\overline{1}$ 497 can be relatively big. A fast-cooled plagioclase feldspar (~An<sub>51</sub>) will not exsolve, but transforms 498 into the  $e^2$  structure instead. The supercooled  $e^2$  structure would eventually transform to  $e^1$  if a 499 long enough cooling time is allowed, but with the modulation period and orientation of  $e^2$ 500 preserved. Only extremely slowly cooled plagioclase would exsolve to e1 directly and follow the 501 same modulation period and orientation as in metamorphic rock. 502

The result of this research clears up the topology of the lower end of the solvus loop for Bøggild intergrowth. The case for the higher end of the loop could be more complicated. Meanwhile, the nature and exact position of the boundary between e1 and e2 are still not clear. The two structures behave differently, but the symmetry did not change at all from this phase transition. Further work needs to be done for plagioclase samples with various cooling history on both sides of the Bøggild miscibility gap to comprehend the phase relations between  $C\overline{1}$ ,  $I\overline{1}$ , e1 and e2.

509 The completed structure variations of plagioclase could also be employed to develop a geospeedometry for quantifying the cooling rates of igneous rocks (Fig. 16). The new parameters 510 such as amplitudes of modulation waves, are proved to be much more sensitive and accurate in 511 quantifying the ordering states of intermediate plagioclase structures, compared to the 512 conventionally used gamma angles or  $T_1o$  occupancies of the average structure. These 513 parameters could be calibrated either internally, based on the CaAl-NaSi interdiffusion rates in 514 515 plagioclase from experiment data, or externally using cooling rate of the same rock calculated by other methods. 516

517

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524

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# 622 Captions of Figures

623

Fig. 1 Dark-field image (left) and selected-area electron diffraction (SAED) pattern ([111] zoneaxis) of sample Dul-15-8B. The SAED pattern (right) shows weak and diffuse *e*-reflections and
no *f*-reflections.

627

Fig. 2 Dark-field image and SAED pattern of sample SK90-12. The image (left) shows homogenous *e*-fringes with no exsolution lamellae. The diffraction pattern (right) ( $[5\overline{1}1]$  zone axis) shows sharp *e*-reflections and weak *f*-reflections next to main reflections (indexed). A fast Fourier transform (FFT) pattern is inserted at the left of the TEM image.

632

Fig. 3 A. Bright-field TEM image of sample R2923 showing nearly periodic exsolution lamellae. The wide lamellae (~230 nm thick) are Ca-rich. B. SAED pattern ( $[5\overline{1}1]$  zone axis) of R2923, showing strong *e*- and *f*-reflections, some weak third order satellite reflections can also be observed. C. Dark-field TEM image showing e-fringes crossing the exsolution lamellae boundary, with a slight change in orientation (~2 degrees or less) and period. The difference in *e*reflections of different lamellae is indistinguishable in the diffraction pattern.

639

Fig. 4 <T-O> bond distances of individual tetrahedral sites for the three samples. **A.** Sample Dul-15-8B; **B.** Sample SK90-12; **C.** Sample R2923. The modulation function wave shows flat part (e.g.  $T_1$ m between t=0.2 and t=0.5;  $T_2$ m between t=0.9 and t=0.2).

643

Fig. 5 The displacement modulation along three axes and electron density contour variation for all three samples, the red line represents M2(Ca) and green line represents M1(Na+Ca), xi (i = 1, 2, 3, 4) are unit cell edge in (3+1)D space. A. Sample Dul-15-8B; B. Sample SK90-12; C. Sample R2923.

Fig. 6 Ca/Na occupancy modulation of M1(Ca+Na) and M2(Ca) site of three samples. A. Sample
Dul-15-8B; B. Sample SK90-12; C. Sample R2923.

651

Fig. 7 The average Na occupancy of 8 symmetrically equivalent M sites in one unit cell (c~14Å)
of all three samples. A. The average occupancy of sample Dul-15-8B is constant, indicating no
density modulation in the structure; B. Sample SK90-12 shows a density modulation of ~8 mole%

An; C. Sample R2923 shows a density modulation of ~17 mole% An.

656

Fig. 8 The average <T-O> bond distance of 8 symmetrically equivalent sites and total average of
all 32 T sites in one unit cell (c~14Å) for all three samples. A. Sample Dul-15-8B; B. Sample
SK90-12; C. Sample R2923.

660

Fig. 9 A section of the modulated structure consists of 3x7 anorthite cells along *b*- and *c*-axis.
The modulation wavefront is marked with red plane in the figure. Oxygen atoms are omitted.
Ca/Na atoms are shown as big spheres, with blue for Ca and yellow for Na occupancies. Al/Si
atoms are shown as small spheres; Si occupancies are shown with dark blue color.

665

Fig. 10 M site and surrounding tetrahedra projected along  $a^*$  (A) and  $b^*$  (B) direction. The split M site lies within the plane defined by T<sub>2</sub> sites, which is parallel to the (100) plane. The split direction of the M site (green arrow) is close to the line connecting two diagonal T<sub>2</sub>m-T<sub>2</sub>o sites (red arrow).

670

Fig. 11 The split distances between M1 and M2 sites are plotted against the distances between diagonal  $T_2$ oz and  $T_2$ mo site for all three samples. The points for t=0 and t=0.5 are marked in red and green respectively. **A.** Sample Dul-15-8B; **B.** Sample SK90-12; **C.** Sample R2923.

Fig. 12 The modulation function of  $T_1 \circ O_A 1 - T_1 m$  angle for all three samples. **A.** Sample Dul-15-8B; **B.** Sample SK90-12; **C.** Sample R2923.

677

Fig. 13 Stereonet projection of the q-vectors of the three samples. The *q*-vectors lie almost exactly on the [923] zone-axis, tilting away from ( $\overline{1}03$ ) as the anorthite composition increases. The stereonet is projected along the *b*-axis, the same orientation as in Smith and Brown (1988). Poles and traces of the two zone-axes [111] and [ $\overline{5}1\overline{1}$ ] used for dark-field images are also labeled in the figure.

683

Fig. 14 Variations of selected modulation functions plotted against gamma angle of the structuresrefined. Clear correlations between the parameters are shown.

686

Fig. 15 A schematic phase diagram around An<sub>51</sub> showing topological relations of  $C\overline{1}$ , e1 and e2phases. The gray area is where  $C\overline{1}$  and e1 coexist as two thermodynamically stable phases, which produce Bøggild intergrowth in extremely slow cooled igneous labradorite (which stayed within the gray area long enough). The dash line is extended boundary between  $C\overline{1}$  and e2, as e2can appear with Ca-rich composition as a metastable phase. The loop area is modified from McConnell (2008) considering compositions for two ends from Ribbe (1983a).

693

Fig. 16 A schematic time-temperature-transformation (TTT) diagram for plagioclase with composition of  $\sim An_{51}$ . Path A, B, C and D represent four host rock with gradually decreasing cooling rate.





Fig. 2



Fig. 3





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Fig. 5

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Fig. 6



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Fig. 7



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Fig. 14





Fig. 16



Dul-15-8B	SK90-12	R2923	
hemical formula $Ca_{0.524}Na_{0.476}Si_{2.476}Al_{1.524}O_8$		Ca <sub>0.525</sub> Na <sub>0.475</sub> Si <sub>2.475</sub> Al <sub>1.525</sub> O <sub>8</sub>	
270.6	270.2	270.6	
Triclinic			
$X\overline{1}(\alpha\beta\gamma)0^{a}$			
8.1625 (3), 12.8586 (2), 14.2038 (4)	8.1614 (1), 12.8545 (1), 14.2112 (2)	8.1658 (2), 12.8542 (1), 14.2096 (2)	
93.5540 (7), 116.2161(2), 90.159 (3)	93.5904 (6), 116.2504 (9), 90.0075 (13)	93.6093 (7), 116.2726 (11), 89.8088 (14)	
$\delta h = 0.05301(8)$ $\delta k = 0.05905(8)$ $\delta l = -0.20038(8)$	$\delta h = 0.05451(10)$ $\delta k = 0.05751(10)$ $\delta l = -0.20208(10)$	$\delta h = 0.06664(7)$ $\delta k = 0.04916(7)$ $\delta l = -0.22680(6)$	
33.95	34.07	30.75	
1334.12 (7)	1333.88 (3)	1334.28 (4)	
0.116×0.100×0.089	0.086×0.077×0.061	0.139×0.115×0.093	
1074	1072	1074	
2.694	2.691	2.694	
1.26	1.25	1.26	
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, , ,	· · · · · · · · · · · · · · · · · · ·	5931(1600 <i>a</i> +2779 <i>e</i> +1552 <i>f</i> )	
		0.047	
$\theta_{\text{max}} = 30.5, \ \theta_{\text{min}} = 2.2$	$\theta_{\rm max} = 30.5,  \theta_{\rm min} = 2.0$	$\theta_{\rm max} = 28.3,  \theta_{\rm min} = 1.9$	
0.715		0.667	
$h = -11 \rightarrow 11,$	$h = -11 \rightarrow 11,$	$h = -11 \rightarrow 10,$	
		$k = -17 \rightarrow 17,$ $l = -19 \rightarrow 19$	
	Ca <sub>0.524</sub> Na <sub>0.476</sub> Si <sub>2.476</sub> Al <sub>1.524</sub> O <sub>8</sub> 270.6Triclinic $X\overline{1}(\alpha\beta\gamma)O^a$ 8.1625 (3), 12.8586 (2), 14.2038 (4)93.5540 (7), 116.2161(2), 90.159 (3) $\delta h = 0.05301(8)$ $\delta k = 0.05905(8)$ $\delta l = -0.20038(8)$ 33.951334.12 (7)0.116×0.100×0.08910742.6941.26Mo Kα IµSGraphite100K495346091(2007a+4084e) <sup>b</sup> 4975(1949a+3026e)0.042 $\theta_{max} = 30.5, \theta_{min} = 2.2$ 0.715	$\begin{array}{c c c c c c c c c c c c c c c c c c c $	

## Table 1 Experimental details of the single crystal X-Ray diffraction analyses.

Refinement							
Refined on	F <sup>2</sup>						
$R[F^2 > 2\sigma(F^2)]$	0.029(0.023/0.048) <sup>c</sup>	0.033(0.030/0.034/0.132)	0.032(0.028/0.032/0.076)				
R(all)	0.037(0.024/0.072)	0.064(0.032/0.061/0.380)	0.042(0.028/0.039/0.142)				
GOF(obs)	2.67	1.62	1.93				
GOF(all)	2.43	1.12	1.66				
No. of parameters	385	640	604				
No. of constraints	150	250	239				
Weighting scheme	Weighting scheme based	Weighting scheme based on measured s.u.'s $w = 1/(\sigma^2(I) + 0.0004I^2)$					
$(\Delta/\sigma)_{max}$	0.030	0.037	0.021				

<sup>a</sup> Centering condition X:  $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}, 0), (\frac{1}{2}, \frac{1}{2}, 0, \frac{1}{2}), (0, 0, \frac{1}{2}, \frac{1}{2}), (0, 0, \frac{1$ 

<sup>b</sup> Letter *a* and *e* (and *f*, if present) refer to main reflections and first order satellites (and second order satellites)

<sup>c</sup> Values in parentheses refer to R for a and e (and f reflections, if present) respectively

Table 2 Atomic positions and occupancies in the modulated structure of sample Dul-15-8B. Note that all the fractional coordinates are based on  $c\sim$ 14Å unit cell, even though the number of independent atoms is the same as in albite cell ( $c\sim$ 7Å)

label	atom	Average	Max	Min	х	У	Z	Uequiv
	Occ.	Occ.	Occ.		,			
M1	Ca	0.133(5)	0.169	0.098	0.26625(11)	-0.01220(3)	0.08100(2)	0.0280(6)
IVIT	Na	0.476(3)	0.512	0.440	0.20025(11)	0.01220(3)	0.00100(2)	0.0200(0)
M2	Са	0.391(6)	0.458	0.324	0.27157(10)	0.03077(7)	0.04735(9)	0.0134(3)
T₁o	Si	0.4535	0.747	0.160	0.49373(4)	0.33644(2)	-0.10713(2)	0.00656(10)
110	Al	0.5465	0.840	0.253	0.49373(4)	0.55044(2)		
Tm	Si	0.6736	0.938	0.409	0 50208(4)	0.31684(2)	0.11573(2)	0.00687(10)
T₁m	Al	0.3264	0.591	0.062	0.50308(4)			
То	Si	0.6943	0.910	0.478	0.69514(4)	0.10875(2)	0.15822(2)	0.00679(10)
T <sub>2</sub> 0	Al	0.3057	0.522	0.090	0.68514(4)			
Tm	Si	0.6546	0.923	0.387	0.19096(4)	0 27804(2)	0 17951(2)	0.00656(10)
$T_2m$	Al	0.3453	0.613	0.077	0.18086(4)	0.37894(2)	0.17851(2)	0.00656(10)
O <sub>A</sub> 1	0	1			0.49740(12)	0.37132(6)	0.01064(6)	0.0160(3)
O <sub>A</sub> 2	0	1			0.57988(10)	-0.00812(6)	0.13892(6)	0.0101(2)
O <sub>B</sub> O	0	1			0.81245(11)	0.10394(6)	0.09430(7)	0.0146(3)
O <sub>B</sub> m	0	1			0.31592(11)	0.35283(7)	0.12231(8)	0.0202(3)
O <sub>c</sub> o	0	1			0.48665(11)	0.20938(6)	-0.14032(6)	0.0135(3)
O <sub>c</sub> m	0	1			0.51385(11)	0.18786(6)	0.10648(6)	0.0148(3)
O <sub>D</sub> o	0	1			0.30173(10)	0.39288(6)	0.30853(6)	0.0134(3)
O <sub>D</sub> m	0	1			0.69006(11)	0.36635(6)	0.21588(6)	0.0164(3)

label	atom	Average	Max	Min	х	У	Z	Uequiv	
		Occ.	Occ.	Occ.					
M1	Ca 1	0.129(9)	0.193	0.075	0.26698(15)	-0.01241(12)	0.08054(10)	0.0288(5)	
	Na	0.499(12)	0.596	0.392				0.0200(0)	
M2	Ca	0.371(8)	0.532	0.250	0.27103(14)	0.02970(9)	0.04829(9)	0.0173(3)	
T1o	Si	0.4397	0.830	0.113	0.49383(3)	0 22610/2)	0 10008(2)	0.00784(9)	
110	Al	0.5603	0.887	0.170	0.49565(5)	0.33619(2)	-0.10698(2)		
T1	Si	0.6830	0.960	0.312	0 50202(2)	0.31716(2)	0.11584(2)	0.00801(9)	
T1m	Al	0.3170	0.688	0.040	0.50302(3)				
<b>T</b> 2	Si	0.7149	0.931	0.425		0.40070(2)	0.45000(2)	0.00702(0)	
T2o	Al	Al 0.2851 0.575 0.069 0.68524(3)	0.68524(3)	0.10878(2)	0.15808(2)	0.00783(9)			
<b>T</b> 2	Si	0.6624	0.930	0.277	0.4000.4(2)		0 ( = 0 = 0 ( 0 )		
T2m	AI	0.3376	0.723	0.070	0.18094(3)	0.37915(2)	0.17853(2)	0.00760(9)	
0 <sub>A</sub> 1	0	1			0.49757(10)	0.37118(5)	0.01087(5)	0.0164(3)	
O <sub>A</sub> 2	0	1			0.58015(9)	-0.00768(5)	0.13883(5)	0.0109(2)	
O <sub>B</sub> O	0	1			0.81201(9)	0.10437(5)	0.09420(6)	0.0151(2)	
O <sub>B</sub> m	0	1			0.31603(10)	0.35282(6)	0.12255(6)	0.0201(3)	
O <sub>c</sub> o	0	1			0.48687(9)	0.20888(6)	-0.14006(5)	0.0145(2)	
0 <sub>c</sub> m	0	1			0.51433(9)	0.18826(6)	0.10685(5)	0.0160(2)	
O <sub>D</sub> O	0	1			0.30143(9)	0.39273(5)	0.30838(5)	0.0142(2)	
O <sub>D</sub> m	0	1			0.69009(9)	0.36674(5)	0.21611(6)	0.0168(2)	

Table 3 Atomic positions and occupancies in the modulated structure of sample SK90-12.

label atom	atom	Average	Max	Min	X	У	Z	Uequiv
	atom	Occ.	Occ.	Occ.	Х			
M1	Ca	0.171(2)	0.319	0.047	0.26670(13)	-0.01331(17)	0.08044(15)	0.0244(4)
	Na	0.475(3)	0.664	0.265	0.20070(13)	-0.01331(17)	0.00044(13)	0.0244(4)
M2	Ca	0.354(2)	0.685	0.048	0.27123(14)	0.02886(9)	0.04801(9)	0.0089(2)
То	Si	0.3944	0.898	0.021	0.49380(4)	0.33570(2)	-0.10684(2)	0.00632(11)
T <sub>1</sub> 0	Al	0.6056	0.979	0.102	0.49560(4)	0.33370(2)		
тm	Si	0.7054	0.994	0.213	0 50200(4)	0 21772(2)	0.11606(2)	0.00634(10)
T₁m	Al 0.2946 0.787 0.006 0.50308(4	0.50508(4)	0.31772(2)	0.11606(2)	0.00054(10)			
То	Si	0.7110	0.954	0.334	0 (0522(4)	0.10907(2)	0 15702/2)	0.00642(11)
T <sub>2</sub> O	2 <sup>0</sup> Al 0.2890 0	0.666	0.046	0.68532(4)	0.10907(2)	0.15793(2)	0.00642(11)	
Tm	Si	0.6642	0.944	0.160	0.40442(4)	0.37948(2)	0.17858(2)	0.00631(11)
T₂m	Al	0.3358	0.840	0.056	0.18113(4)			
$O_A 1$	0	1			0.49761(11)	0.37165(6)	0.01149(6)	0.0145(3)
O <sub>A</sub> 2	0	1			0.58061(10)	-0.00708(6)	0.13887(6)	0.0096(3)
O <sub>B</sub> o	0	1			0.81203(10)	0.10514(6)	0.09420(6)	0.0131(3)
$O_{\text{B}}m$	0	1			0.31686(11)	0.35299(7)	0.12310(7)	0.0180(3)
O <sub>c</sub> o	0	1			0.48713(10)	0.20793(6)	-0.13962(6)	0.0123(3)
0 <sub>c</sub> m	0	1			0.51516(10)	0.18927(6)	0.10754(6)	0.0136(3)
O <sub>D</sub> o	0	1			0.30094(10)	0.39289(6)	0.30824(6)	0.0124(3)
$O_D m$	0	1			0.68940(11)	0.36694(6)	0.21607(6)	0.0153(3)

Table 4 Atomic positions and occupancies in the modulated structure of sample R2923.

	Dul-15-8B	SK90-12	R2923	987L <sup>a</sup>
Composition	An <sub>52</sub>	An <sub>50</sub>	An <sub>52</sub>	An <sub>51</sub>
Gamma (°)	90.159	90.008	89.809	89.840 <sup>b</sup>
Modulation Period (Å)	33.95	34.07	30.75	30.15 <sup>c</sup>
Modulation	е2	<i>e</i> 1	e1+e1	<i>e</i> 1
Na occ. variation (peak-trough)	0.072	0.204	0.399	0.353
Density variation (mole% An)	0	8	17	17
Average <t<sub>1o - O&gt; (Å)</t<sub>	1.6852	1.6876	1.6935	1.6907
Average <t₁m -="" o=""> (Å)</t₁m>	1.6555	1.6547	1.6516	1.6508
Average <t<sub>2o - O&gt; (Å)</t<sub>	1.6527	1.6504	1.6508	1.6474
Average <t₂m -="" o=""> (Å)</t₂m>	1.6581	1.6575	1.6571	1.6562
$< T_1 o-O > variation (Å)$	0.0790	0.0928	0.1193	0.1170
$< T_1 m-O > variation (Å)$	0.0713	0.0836	0.1066	0.1066
<t<sub>2o-O&gt; variation (Å)</t<sub>	0.0584	0.0648	0.0837	0.0842
<t<sub>2m-O&gt; variation (Å)</t<sub>	0.0724	0.0844	0.1067	0.1090

## Table 5 Potential parameters for estimating the ordering state of e-plagioclase

<sup>a</sup> Metamorphic labradorite with no exsolution, data collected at room temperature (Jin and Xu, 2017).

<sup>b</sup> For comparison, the unit cell parameters (a, b, c, α, β, γ) of the same crystal were measured to be (8.1610(1) Å, 12.8481(1) Å, 14.2074(2) Å, 93.6469(6)°, 116.2616(9)°, 89.8744(13)°) at 100K.

<sup>c</sup> The *q*-vector (δ*h*, δ*k*, δ*l*) measured at 100K was (0.06667, 0.05477, 0.2293), resulting in a period of 30.29 Å.