

## Supplementary Material:

### Crystal Chemistry and Thermodynamic Properties of Zircon Structure-Type Materials

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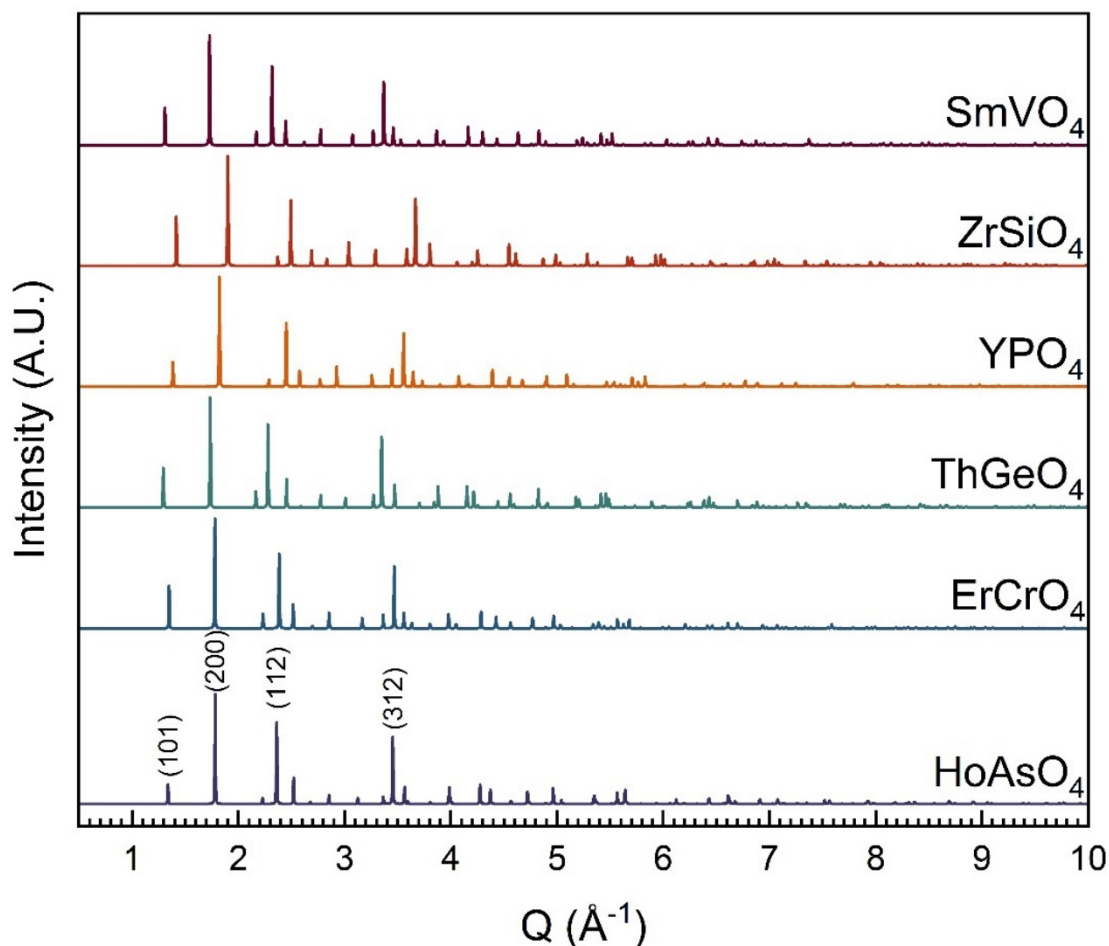
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## Methods of investigation into zircon-type compounds

**Powder X-ray diffraction** Studying the crystallographic structure with the help of powder X-ray diffraction, will lead to a pattern where the strongest diffraction peaks correspond, in descending order, to the following crystallographic planes: (200), (112), (312), and (101). The precise position of these XRD peaks will change as a result of the exact crystal chemistry of the zircon structure owing to changes in the size of the unit-cell (Figure S1). These shifts tend to be systematic across an entire series of compounds, such as with xenotimes (Ni et al. 1995; Ushakov et al. 2001a), resulting in a linear shift in unit cell parameters. This can be used further in both single substituted solid solutions, where only the lanthanide element is replaced (*i.e.*,  $\text{Er}_x\text{Yb}_{1-x}\text{PO}_4$ ) (Strzelecki et al. 2022) and in coupled substitution of solid solutions in which cations and anions are replaced, such as in phospho-silicates solid solutions ( $\text{Th}_x\text{Er}_{1-x}(\text{SiO}_4)_x(\text{PO}_4)_{1-x}$ ) (Mesbah et al. 2016; Shelyug et al. 2021), as a means of determining the chemical composition. The reason for this is that zircon-type solid solutions follow Vegard's law. However, according to the work of Marcial et al. (2021) there are unique circumstance for which Vegard's law is not applicable, *e.g.*, for the uranothorite solid solution that exhibits a negative deviation from Vegard's law.



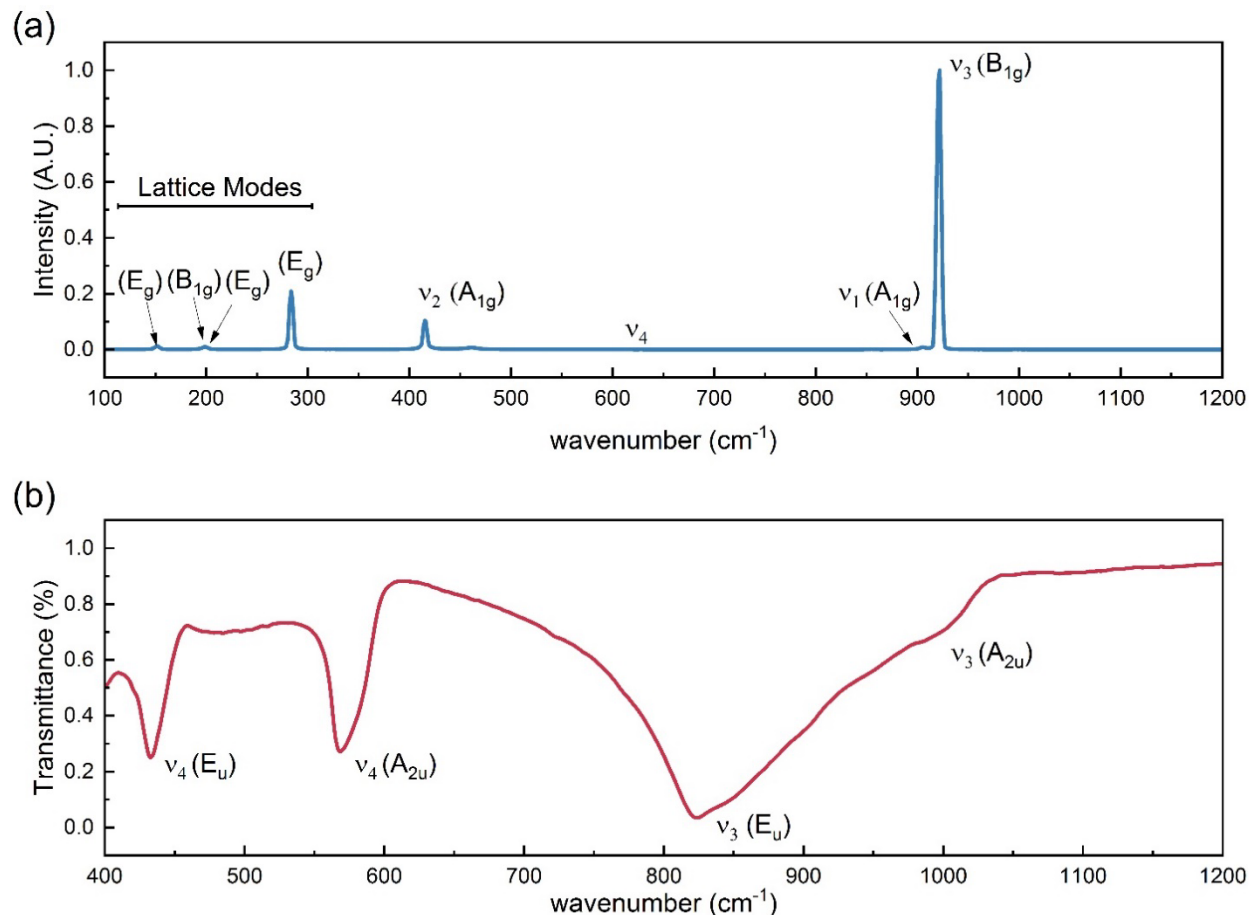
**Figure S1.** Calculated powder X-ray diffraction patterns for a variety of zircon structure materials, with the four strongest reflections labeled. Plots were calculated using the VESTA 3 software package (Momma and Izumi 2011) and crystallographic information files (CIFs) reported in the literature for each compound (Robinson et al. 1971; Ni et al. 1995; Mullica et al. 1996; Jiménez et al. 2002b; Schmidt et al. 2005; Achary et al. 2007).

**Raman.** As zircon structure-type phases belong to the  $D_{4h}$  point group, they have twelve active Raman vibrational modes, which can be determined by factor-group analysis (Dawson et al. 1971; Kolesov et al. 2001; Clavier et al. 2014). For Raman spectroscopy, seven ( $\Gamma_{\text{int}} = 2A_{1g} +$

$2B_{1g} + B_{2g} + 2E_g$ ) (Dawson et al. 1971; Kolesov et al. 2001) can be assigned to the internal vibrations of the  $TO_4$  tetrahedron, while the remaining five ( $\Gamma_{int} = 2B_{1g} + 3E_g$ ) (Dawson et al. 1971; Kolesov et al. 2001) can be assigned to the external vibrations of the  $TO_4$  tetrahedron (Dawson et al. 1971; Hoskin and Rodgers 1996; Kolesov et al. 2001). Due to the interaction of the  $TO_4$  tetrahedra with the  $MO_8$  dodecahedra, the tetrahedra cannot be considered as strictly independent units (Syme et al. 1977), and so there are no reported spectra with all twelve active Raman modes for zircon-type phases (Clavier et al. 2014). Instead, the visible Raman modes can be divided into three zones of interest: zone one,  $900 - 1200\text{ cm}^{-1}$ ; zone two,  $350 - 900\text{ cm}^{-1}$ ; and zone three,  $100\text{-}500\text{ cm}^{-1}$  (Figure S2a). In the first zone, the T-O stretching motions are observed and assigned as the symmetric stretching ( $\nu_1 - A_{1g}$ ) and the anti-symmetric stretching ( $\nu_3 - B_{1g}$ ) modes. In the second zone, the T-O bending motions are observed and assigned as the symmetric bending mode ( $\nu_2 - A_{1g}$ ) and antisymmetric bending mode ( $\nu_4$ ). The third zone corresponds to the lattice vibrational modes. There is still controversy over the assignment of some of the lattice motions as well as the antisymmetric bending modes ( $\nu_4$ ). For a more in-depth discussion, the reader may refer to the work of Nasdala et al. (2003). A number of the zircon structure materials containing f-elements, which have been studied by Raman, can either exhibit fluorescence (i.e.,  $U^{4+}$ ) or display photoluminescence (i.e.,  $REE^{3+}$ ) (Bauer et al. 2014; Clavier et al. 2014; Lenz et al. 2015, 2019; Strzelecki et al. 2021). There are a number of advantages of using these “artifacts” (Lenz et al. 2015, 2019, 2020).

**Fourier transform infrared spectroscopy (FTIR).** For FTIR, the  $I4_1/amd$  space group induces seven active FTIR vibrational modes, which was determined again through a factor-group analysis (Dawson et al. 1971; Clavier et al. 2014). Similar to Raman spectroscopy, these

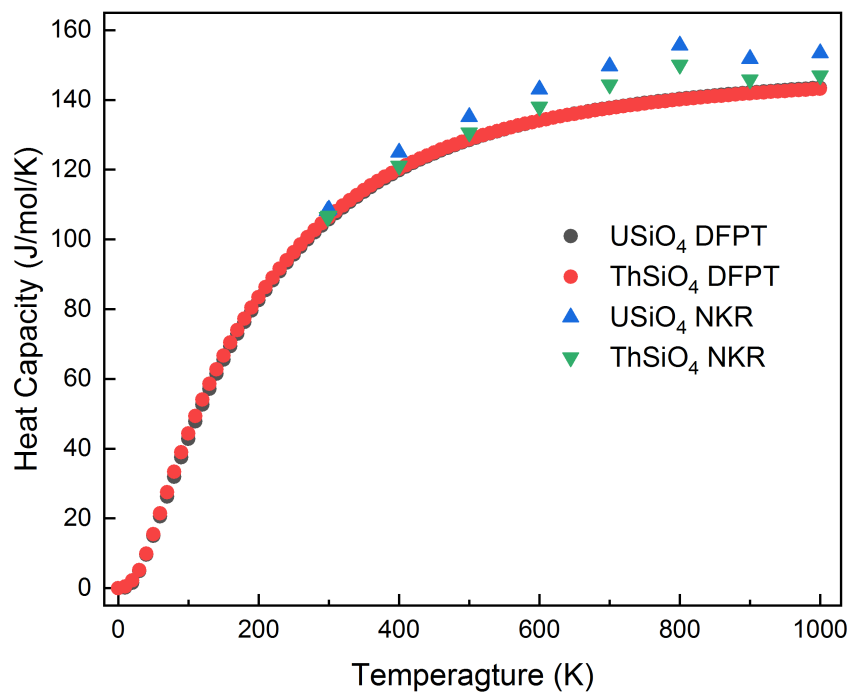
vibrational modes can be assigned to the internal and external vibrations of the  $\text{SiO}_4$  tetrahedron. Of the seven vibrational modes, several of them ( $\Gamma_{\text{int}} = 2A_{2u} + 2E_u$ ) can be assigned to internal vibrational modes and the others ( $\Gamma_{\text{ext}} = A_{2u} + 2E_u$ ) can be assigned to external vibrational modes. Again, like with Raman, the active modes in FTIR spectra can be grouped into three zones of interest: zone one,  $800 - 1200 \text{ cm}^{-1}$ ; zone two,  $400 - 800 \text{ cm}^{-1}$ , and zone three,  $100-400 \text{ cm}^{-1}$  (Figure S2b). As only the antisymmetric internal modes of the  $\text{TO}_4$  tetrahedron are visible by IR, this means that the doubly degenerated antisymmetric stretching mode ( $\nu_3 - E_u$ ) and the antisymmetric stretching mode ( $\nu_3 - A_{2u}$ ) would appear in the first zone. In the second zone, the antisymmetric deformation mode ( $\nu_4 - A_{2u}$ ) and the doubly degenerated antisymmetric deformation mode ( $\nu_4 - E_u$ ) are visible. In the third zone, the three external vibrational modes ( $\Gamma_{\text{ext}} = A_{2u} + 2E_u$ ) can be observed. However, unlike with Raman many of these vibrational modes are quite broad. The recent work reported by Clavier *et al.* (2014, 2018), on uranorthorite and gadolinium orthophosphate demonstrate how to differentiate the components assigned to symmetric and antisymmetric stretching vibration modes. One advantage of utilizing FTIR over Raman is when studying the simultaneous cationic and anionic solid solutions, such as with phosphor-silicate solid solutions  $(\text{Th}_x\text{Er}_{1-x})(\text{SiO}_4)_x(\text{PO}_4)_{1-x}$  (Mesbah et al. 2016), one can see the effects of stretching of both  $\text{TO}_4$  sites. For more in-depth discussion we invite the reader to the works of Nasdala et al. (2003).



**Figure S2.** Representative vibrational spectra, **(a)** Raman and **(b)** FTIR, for a zircon structure material (CeSiO<sub>4</sub>) (Strzelecki et al. 2021). The Raman was collected with a 532 nm laser and the FTIR was collected by Attenuated Total Reflectance.

**Enthalpy ( $\Delta H$ ).** There are direct and indirect experimental methods for determining the  $\Delta H_f$  value of a condensed phase material, such as calorimetric techniques (*i.e.*, acid, bomb, and oxide melt), and solubility measurements. In a direct calorimetric method, heat required for forming a material is measured. For indirect calorimetric methods, heat of a given reaction (*i.e.*, dissolution or leaching, including or not oxidative processes) of the interested material is measured and used

in thermochemical cycles based on Hess's law for deriving  $\Delta H_f$ . As previously stated in the introduction, zircon-type materials are highly refractory and resistant to alteration in many conditions, making the techniques based on acid or bomb calorimetry extremely difficult to implement. On the other hand, high temperature oxide melt drop solution calorimetry (Holm and Kleppa 1967; Navrotsky and Kleppa 1967), benefiting from the readily dissolution of refractory materials in high temperature oxide melts (lead borate,  $\text{PbO} \cdot \text{B}_2\text{O}_3$  and sodium molybdate,  $3\text{Na}_2\text{O} \cdot 4\text{MoO}_3$ ), is one principal method to determine  $\Delta H_f$  of zircon-type materials. In this technique, a small sample (~5 mg) is dropped into oxide melt of either  $\text{PbO} \cdot \text{B}_2\text{O}_3$  or  $3\text{Na}_2\text{O} \cdot 4\text{MoO}_3$  as hot solvent maintaining at 973 K or 1073 K (Navrotsky 1977a, 1977b, 1997, 2014). Another technique used to derive  $\Delta H_f$  is based on the vaporization using the mass spectrometric method in combination with Knudsen effusion (Costa et al. 2017; Kowalski et al. 2021). When studying xenotime and monazite based samples, cautious validation, however, may be needed for such refractory type materials, due to the possible formation of gaseous products during vaporization and the formation of oxyphosphates during the decomposition of the rare earth orthophosphates (Ushakov et al. 2001b).



**Figure S3.** Constant volume heat capacity evaluated by Neumann Kopp's rule and DFT calculations.



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