1 Revision 1

2 Effect of hydration on the single-crystal elasticity of pyrope at high

3 pressure and high temperature conditions: implications for the

4 **Earth's upper mantle**

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15 ABSTRACT:

16 The elasticity of single-crystal hydrous pyrope with ~ 900 ppmw H₂O has been 17 derived from sound velocity and density measurements using in situ Brillouin light 18 spectroscopy (BLS) and synchrotron X-ray diffraction (XRD) in the diamond anvil 19 cell (DAC) up to 18.6 GPa at room temperature and up to 700 K at ambient pressure. 20 These experimental results are used to evaluate the effect of hydration on the 21 single-crystal elasticity of pyrope at high pressure and high temperature (P-T)22 conditions to better understand its velocity profiles and anisotropies in the upper 23 mantle. Analysis of the results shows that all of the elastic moduli increase almost 24 linearly with increasing pressure at room temperature, and decrease linearly with 25 increasing temperature at ambient pressure. At ambient conditions, the aggregate 26 adiabatic bulk and shear moduli (K_{S0} , G_0) are 168.6(4) GPa and 92.0(3) GPa,

27 respectively. Compared to anhydrous pyrope, the presence of ~900 ppmw H₂O in 28 pyrope does not significantly affect its K_{S0} and G_0 within their uncertainties. Using the 29 third-order Eulerian finite-strain equation to model the elasticity data, the pressure 30 derivatives of the bulk $[(\partial K_S / \partial P)_T]$ and shear moduli $[(\partial G / \partial P)_T]$ at 300 K are derived 31 as 4.6(1) and 1.3(1), respectively. Compared to previous BLS results of anhydrous pyrope, an addition of ~900 ppmw H₂O in pyrope slightly increases the $(\partial K_S / \partial P)_T$, but 32 33 has a negligible effect on the $(\partial G/\partial P)_{T}$ within their uncertainties. The temperature 34 derivatives of the bulk and shear moduli at ambient pressure are $(\partial K_S / \partial T)_P = -0.015(1)$ 35 GPa/K and $(\partial G/\partial T)_P$ =-0.008(1) GPa/K, which are similar to those of anhydrous 36 pyrope in previous BLS studies within their uncertainties. Meanwhile, our results also 37 indicate that hydrous pyrope remains almost elastically isotropic at relevant high P-T38 conditions, and may have no significant contribution to seismic anisotropy in the 39 upper mantle. In addition, we evaluated the seismic velocities (V_P and V_S) and the 40 $V_{\rm P}/V_{\rm S}$ ratio of hydrous pyrope along the upper mantle geotherm and a cold subducted 41 slabs geotherm. It displays that hydrogen has also no significant effect on the seismic 42 velocities and the $V_{\rm P}/V_{\rm S}$ ratio of pyrope at the upper mantle conditions.

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Keywords: Hydrous pyrope, Single-crystal elasticity, High pressure and high
temperature, Brillouin light scattering, Upper Mantle

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55 INTRODUCTION

56 Silicate garnet is an important constituent in the Earth's upper mantle and transition 57 zone (e.g. Anderson 1989; Anderson and Bass 1984; Bina 2013; Duffy and Anderson 58 1989; Fan et al. 2009, 2011, 2015b, 2015c, 2017a; Frost 2008; Ita and Stixrude 1992; 59 McDonough and Sun 1995). Mantle compositional models such as pyrolite and 60 piclogite contain ~15% and ~22% of garnet in the upper mantle, respectively (e.g. 61 Bass and Anderson 1984; Ringwood 1975). The percentage can increase to ~40% or 62 even more in the transition zone because pyroxenes progressively dissolve into garnet 63 with increasing pressure (Li B W et al. 2018), forming majorite-garnet solid solutions 64 (Herzberg and Gasparik 1991; Ringwood 1991). Garnet is also one of the important 65 mineral for eclogite (e.g. Kimura et al. 2013; Liu 1980; Xu et al. 2019), formed by 66 high-pressure metamorphism of basalt or gabbro at subduction zones (Poli and 67 Schmidt 2002; Ringwood 1982). Although most natural garnets are complex solid 68 solutions, the most significant component of mantle garnets is its Mg end-member 69 pyrope (Mg₃Al₂Si₃O₁₂) (Rickwood et al. 1968; Sinogeikin and Bass 2002). Therefore, 70 pyrope or pyrope-rich garnet is an important mantle mineral, irrespective of what 71 compositional model of Earth's mantle is assumed (Ringwood 1975; Ita and Stixrude 1992). 72

73 In addition, previous studies have revealed that hydrogen could be incorporated 74 into nominally anhydrous minerals (NAMs) as structurally bound hydroxyl defects 75 (e.g. Ingrin and Skogby 2000; Skogby 2006; Smyth 1987). NAMs in the Earth's 76 mantle thus have significant implications for the Earth's deep water cycle (e.g. 77 Bolfan-Casanova et al. 2000; Hirschmann 2006; Hirschmann and Kohlstedt 2012; 78 Smyth and Jacobsen 2006). Water can incorporate in garnets as OH- defects 79 associated with charge balancing or oxidation-reduction reactions, or it may substitute 80 Si in the hydrogarnet substitution (Lu and Keppler 1997; Mookherjee and Karato 81 2010; Withers et al. 1998). Pyrope is a well-known hydrous-bearing NAMs phase in 82 the upper mantle (e.g. Ackermann et al. 1983; Rossman et al. 1989). Natural

pyrope-rich garnets from ultrahigh-pressure metamorphic rocks and kimberlite xenoliths generally contain tens to hundreds of ppmw H₂O (e.g. Aines and Rossman 1984a, 1984b; Bell and Rossman 1992a, 1992b; Beran and Libowitzky 2006; Li H Y et al. 2018). Moreover, experiments on water solubility in garnets also indicated that pyrope and pyrope-rich garnets could dissolve certain amounts of hydrogen, ranging from a few hundred to ~ 1000 ppmw H₂O (Geiger et al. 1991; Lu and Keppler 1997; Mookherjee and Karato 2010; Withers et al. 1998).

90 The accurate knowledge about the elastic property of pyrope or pyrope-rich garnet 91 is critical for deducing seismic velocities and density profiles and further constructing 92 reliable mantle mineralogy models (e.g. Bass and Anderson 1984; Bass et al. 2008; 93 Duffy and Anderson 1989; Weidner and Wang 2000). Up to now, numerous equation 94 of state studies of pyrope using XRD technique at high pressure/high temperature 95 have been reported (e.g. Leger et al. 1990; Levien et al. 1979; Sato et al. 1978; Wang et al. 1998; Zhang et al. 1998; Zou et al. 2012a). Additionally, the elastic properties of 96 97 pyrope at ambient and high pressure/temperature conditions have also been 98 investigated using theoretical calculations (e.g. Li et al. 2011; Hu et al. 2016). 99 Moreover, the adiabatic bulk and shear moduli of polycrystalline pyrope have been 100 reported up to 24 GPa and 1700 K by ultrasonic interferometry technique (e.g. 101 Chantel et al. 2016; Chen et al. 1997, 1999; Gwanmesia et al. 2006, 2007; Sumino 102 and Nishizawa 1978; Suzuki and Anderson 1983; Zou et al. 2012b). On the other 103 hand, BLS is another common technique to measure the elasticity of minerals (e.g. Bass and Zhang 2015; Speziale et al. 2014). It has tremendous advantages in deriving 104 105 the complete set of elastic moduli in single-crystal minerals at extremely high P-T 106 conditions (e.g. Fan et al. 2015a; Mao et al. 2015; Murakami et al. 2007; Yang et al. 107 2015; Zhang et al. 2016). There have been a number of single-crystal elasticity studies 108 on pyrope at ambient conditions (e.g. Leitner et al. 1980; O'Neill et al. 1991), high 109 pressure (e.g. Conrad et al. 1999; Sinogeikin and Bass 2000), and high temperature 110 (e.g. Sinogeikin and Bass 2002) conditions using BLS technique. Recently, Lu et al.

111 (2013) measured the single-crystal elasticity of Fe-bearing pyrope at high P-T112 conditions using BLS technique, which provided the detailed description of Fe effect 113 on the elastic moduli of pyrope.

114 Furthermore, the effect of hydrogen on the elasticity of other major mantle minerals 115 (e.g. olivine, wadsleyite, and ringwoodite) has been studied extensively at ambient 116 and high pressure/high temperature conditions (e.g. Inoue et al. 1998; Jacobsen et al. 117 2004, 2008; Mao et al. 2008, 2010, 2011, 2012; Sinogeikin et al. 2003; Wang et al. 118 2003, 2006). However, there is only one study on the acoustic velocities and 119 single-crystal elastic moduli of hydrous pyrope (~180 ppmw) using BLS technique at 120 ambient conditions (O'Neill et al. 1991). The effect of hydration on the acoustic 121 velocities and elastic moduli of pyrope at high P-T conditions remains unavailable, even though it is highly desirable to use its single-crystal elasticity for understanding 122 123 the geodynamic processes of the upper mantle. Up to now, only Fan et al. (2017b) conducted the high P-T equation of state study of hydrous pyrope using 124 125 synchrotron-based XRD technique.

In this study, we measured the acoustic velocities (V_P and V_S) of single-crystal hydrous pyrope at high pressures up to ~18.6 GPa and high temperatures up to 700 K using BLS technique, and derived its full set of single-crystal elastic moduli at high P-T conditions. Based on our results, we further evaluated the effects of hydrogen on the elastic moduli, sound velocities, and elastic anisotropies of pyrope. Finally, we applied our results to discuss the hydrogen effect on the velocity profile and V_P/V_S ratio of pyrope in the Earth's upper mantle.

133 EXPERIMENTAL METHODS

The single-crystal hydrous pyrope was synthesized in a multi-anvil pressure apparatus (YJ-3000T), at the Institute of Geochemistry, Chinese Academy of Sciences, Guiyang, China. More detailed information about the sample synthesis and sample characterization were presented elsewhere by Fan et al. (2017b). Here we briefly report results from Electron probe microanalysis (EPMA), Fourier transform infrared

139 (FTIR), and XRD. EPMA results show that our sample is homogeneous with a 140 chemical formula as Mg_{3.006}Al_{1.995}Si_{3.005}O₁₂. Hydrogen concentrations were 141 determined by FTIR spectroscopy and the absorption bands were readily attributed to 142 structural bonded hydroxyl groups in pyrope (Geiger et al. 1991; Mookherjee and 143 Karato 2010; Withers et al. 1998). The water content in our sample was determined to 144 be ~900(± 100) ppmw using the formula of Bell et al. (1995). Meanwhile, the XRD pattern of our sample confirmed a cubic structure with lattice parameter a=11.460(3)145 Å at ambient conditions, yielded the unit-cell volume $V_0=1505.24(8)$ Å³ and density 146 147 ρ =3.557(4) g/cm³. The unit-cell volume of our hydrous pyrope at ambient conditions 148 is ~0.15% higher than anhydrous pyrope (Du et al. 2015; Zhang et al. 1998), which 149 agrees with previous studies for other mantle minerals (e.g. Holl et al. 2008; Smyth et 150 al. 2003; Smyth and Jacobsen 2006; Ye et al. 2010, 2012).

151 Pyrope has a cubic structure with only three independent elastic moduli (C_{11} , C_{12} , 152 and C_{44}), and therefore, a single crystallographic orientation is sufficient to constrain 153 all three of them using BLS measurements. The crystallographic plane of sample 154 piece is (0.34, -0.53, 0.92) determined by single-crystal XRD at beamline 13-BMD of 155 the GeoSoilEnviroConsortium for Advanced Radiation Sources (GSECARS) of 156 Advanced Photon Source (APS), Argonne National Laboratory (ANL). We 157 double-sides polished our sample pieces to ~20-30 µm thickness with successively 158 finer grits down to a final 3M diamond lapping film of 1 µm grainsize. The thinly 159 polished platelet was then cleaved into several square pieces of the desired size (~150 160 μm) for high-pressure/high temperature measurements.

High-pressure BLS combined with XRD measurements were conducted on the single-crystal hydrous pyrope in a short symmetrical DAC at 13-BMD beamline of APS. An incident X-ray beam of 0.3344 Å wavelength focused to a $3\times7 \ \mu\text{m}^2$ area was used to determine the unit-cell volume of crystal in the DACs. Round Re gasket of 250 µm thick and 3 mm in diameter was pre-indented to ~55 µm thickness using a pair of 500 µm culet size diamond anvils. Subsequently, a cylindrical 300 µm

diameter hole was drilled in the pre-indented area as the sample chamber. A 167 168 single-crystal platelet with a diameter of ~150 µm was then placed into the sample 169 chamber, together with some ruby spheres of approximately 5 µm in diameter as the 170 pressure indicator (Mao et al. 1986) for neon gas loading as well as for high-pressure 171 experiments. The neon pressure medium was loaded into the sample chamber using 172 the gas-loading system at GSECARS of APS (Rivers et al. 2008). Pressures were 173 measured from ruby fluorescence spectra, while pressure uncertainties were 174 calculated using multiple measurements before and after collection of BLS spectra for each pressure point. The XRD spectra were used to determine density at each pressure 175 176 before and after BLS measurements (Table 1).

177 High-temperature BLS experiments were also performed at 13-BMD beamline of 178 APS. A single-crystal hydrous pyrope (~150 µm) was loaded into an 179 externally-heated DAC (EHDAC), which was equipped with an alumina ceramic heater coiled with two Pt wires of 200 µm in diameter and 48 cm in length (Fan et al. 180 181 2019; Kantor et al. 2012; Lu et al. 2013; Mao et al. 2015; Yang et al. 2014, 2016). Re 182 was used as the gasket material and pre-indented to \sim 55 µm thickness using a pair of 183 diamond anvils with 500 µm culet size and then a 300 µm diameter sample chamber 184 was drilled at the center of pre-indentation. The single-crystal hydrous pyrope sample 185 was sealed in the sample chamber. An R-type thermocouple was attached to one of 186 diamond surface approximately 500 µm away from its culet and clad with a ceramic 187 adhesive (Resbond 920) for temperature measurements. To minimize temperature 188 instability for each heating run, we firstly heated the sample chamber to a given 189 temperature and then kept it at this temperature for at least 30 minutes. Temperatures 190 of the sample chamber were actively stabilized within ± 1 K using the 191 temperature-power feedback program with a remotely controlled Tektronix Keithley 192 DC power supply during the experiments (Sinogeikin et al. 2006). Single-crystal 193 XRD patterns of hydrous pyrope before and after BLS measurements were also 194 collected to determine the lattice parameters and densities of the sample at high

195 temperatures (Table 2). Temperatures were increased every 100 K from room 196 temperature (300 K) to maximum temperature (700 K), and then decreased to room 197 temperature to check for possible changes in the lattice parameters and elastic moduli 198 at ambient conditions. From Table 2, we can find that the lattice parameters and 199 elastic moduli at ambient conditions of our hydrous pyrope are highly consistent 200 before and after heating. In addition, we did not measure the water content of our 201 sample after the high-temperature BLS measurements due to the relatively small size 202 and thickness of our sample. However, the previous study has determined the water 203 content of the hydrous garnet after heating up to ~ 1273 K and indicated the water 204 loss of hydrous sample less than 10% (Dai et al. 2012). The maximum experimental 205 temperature in this study (~ 700 K) is significantly lower than that in previous study 206 (~ 1273 K) (Dai et al. 2012). Therefore, we infer that there is no obvious water loss 207 during the high temperature BLS measurements in this study, which is also consistent 208 with the previous study that demonstrated that the intrinsic hydrogen loss in hydrous pyrope occurs at temperatures of \geq 500 °C (Bell et al. 1995). 209

210 The Brillouin system at 13-BMD beamline was equipped with a Coherent Verdi V2 211 solid-state laser with a wavelength of 532 nm, a Perkin-Elmer photomultiplier 212 detector (model: MP983), and a JRS six-pass tandem Fabry-Pérot interferometer (Lu 213 et al. 2013; Yang et al. 2014). BLS spectra were collected in the symmetric forward 214 scattering geometry with an external scattering angle of 50°, which was calibrated 215 using the elastic moduli of standard silicate glass, distilled water, and single-crystal 216 MgO (Ostwald et al. 1977; Polian et al. 2002; Sinogeikin and Bass 2000). The laser 217 beam focused on the sample position was approximately 15 µm in diameter. The 218 acoustic velocities ($V_{\rm P}$ and $V_{\rm S}$) of our sample were derived from analysis of the 219 measured Brillouin frequency shift as follows:

$$220 V_{P,S} = \frac{\lambda_0 \Delta v_B}{2 \sin\frac{\theta}{2}} (1)$$

where $V_{P,S}$ is the acoustic velocities, λ_0 is the incident laser wavelength, Δv_B is the Brillouin frequency shift, and θ is the external scattering angle.

223 **RESULTS AND DATA ANALYSES**

224 BLS and XRD spectra of single-crystal hydrous pyrope sample are collected up to 225 ~18.6 GPa at room temperature in 2-3 GPa pressure interval and up to 700 K at room 226 pressure in 100 K temperature interval. For all of the Brillouin spectra, one 227 quasi-longitudinal and one quasi-transverse acoustic mode are observed. Typical 228 Brillouin spectra at high pressure/high temperature conditions are shown in Figure 1. 229 The measured frequency shifts are converted to velocities along the horizontal axis 230 using equation (1). Most spectra show strong $V_{\rm P}$ and $V_{\rm S}$ peaks with high 231 signal-to-noise ratios except for some crystallographic directions where $V_{\rm P}$ peaks are 232 weakly observable (Fig. 1). Brillouin signals of neon pressure medium are also 233 observed at pressures below ~8 GPa, but they are too weak to be seen when the 234 pressures are increased above 8 GPa. For each platelet at each given P-T conditions, 235 Brillouin spectra are collected in 19 different crystallographic directions from 0 to 180° 236 of the azimuthal angle at every 10° (Fig. 2). The variation in measured $V_{\rm P}$ and $V_{\rm S}$ as a 237 function of azimuthal angle are not observed outside experimental uncertainties, 238 indicating that our hydrous pyrope is almost elastically isotropic at ambient and high pressure/high temperature conditions (Fig. 2). Furthermore, both V_P and V_S of hydrous 239 240 pyrope increase with increasing pressure, and decrease with increasing temperature.

Individual elastic moduli (C_{ij}) of single-crystal hydrous pyrope at each given pressure/temperature conditions (Tables 1 and 2) are obtained by fitting the measured spatial dispersion (velocity vs. orientation) of V_P and V_S to Christoffel's equation using non-linear least square method (Every 1980):

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$$|C_{ijkl}n_jn_l - \rho V_{P,S}^2 \delta_{ik}| = 0$$
 (2)

where C_{ijkl} are the elastic constant in full suffix notation, n_j and n_l are the direction cosines of the phonon along the propagation direction, ρ is the density at each pressure/temperature condition, δ_{ik} is the Kronecker delta function. The

249 root-mean-square (RMS) deviation of the fitting are about 20 m/s, indicating excellent 250 agreement between measured and calculated sound velocities at ambient and high 251 pressure/temperature conditions (Fig. 2). Previous studies also indicated that the 252 single-crystal elastic moduli of pyrope could be calculated by averaging the measured 253 acoustic velocities (e.g. Lu et al. 2013; Sinogeikin and Bass 2000). From Table S1, 254 we notice that the single-crystal elastic moduli of hydrous pyrope derived from the 255 non-linear least-squares fitting are indistinguishable within their uncertainties from 256 those calculated values by averaging the measured acoustic velocities assuming that 257 pyrope is elastically isotropic. All of the individual elastic moduli (C_{11} , C_{12} , and C_{44}) 258 for hydrous pyrope increase smoothly with increasing pressure, and decrease with 259 increasing temperature (Fig. 3).

260 Using the derived individual elastic moduli (C_{11} , C_{12} , and C_{44}) of hydrous pyrope, 261 the adiabatic bulk and shear moduli (K_S and G) are calculated according to the 262 Voigt-Reuss-Hill averages (Hill 1952). The aggregate adiabatic bulk (K_{S0}) and shear 263 moduli (G_0) of hydrous pyrope at ambient conditions are 168.6(4) and 92.0(3) GPa, 264 respectively. The pressure derivatives of elastic moduli at 300 K (Tables 3 and 4) are 265 obtained by fitting the elastic moduli at high pressure using the third-order Eulerian 266 finite-strain equation (Figs. 3a and 4a) (Birch 1978). The pressure derivatives of the 267 individual (C_{ii}) and aggregate (K_S and G) elastic moduli at room temperature are 268 derived to be $(\partial C_{11}/\partial P)_{\rm T}=6.2(1),$ $(\partial C_{12}/\partial P)_{\rm T}=3.7(1),$ $(\partial C_{44}/\partial P)_{\mathrm{T}}=1.5(1),$ 269 $(\partial K_{\rm S}/\partial P)_{\rm T}=4.6(1)$, and $(\partial G/\partial P)_{\rm T}=1.3(1)$, respectively. Due to the limited temperature 270 range for high-temperature data, a linear equation is applied to obtain the temperature 271 derivatives of elastic moduli (Figs. 3b and 4b). The temperature derivative of 272 individual and aggregate elastic moduli at ambient pressure (Tables 3 and 4) are 273 derived be $(\partial C_{11}/\partial T)_{\rm P} = -0.028(1)$ GPa/K, $(\partial C_{12}/\partial T)_{\rm P} = -0.009(1)$ to GPa/K. 274 $(\partial C_{44}/\partial T)_{P}=-0.006(1)$ GPa/K, $(\partial K_{S}/\partial T)_{P}=-0.015(1)$ GPa/K, and $(\partial G/\partial T)_{P}=-0.008(1)$ 275 GPa/K, respectively. The aggregate V_P and V_S of our hydrous pyrope at high 276 pressure/temperature conditions (Fig. 5) are calculated using the following equations:

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$$V_{\rm P} = \sqrt{\frac{\kappa_{\rm S} + \frac{4G}{3}}{\rho}}$$
(3)
278
$$V_{\rm S} = \sqrt{\frac{G}{\rho}}$$
(4)

279 **DISCUSSION**

280 Hydrogen effect on the elasticity of pyrope at high *P-T* conditions

281 In order to understand the effect of hydrogen on the elasticity of pyrope, we 282 compare our results with literature values for pyrope obtained from BLS and 283 ultrasonic interferometer measurements. Tables 3 and 4 show a complete list of 284 individual (C_{ii} s) and aggregate (K_{S0} and G_0) elastic moduli obtained in the present 285 study for hydrous pyrope along with those previous studies of anhydrous pyrope. 286 Compared to anhydrous pyrope (e.g. Sinogeikin and Bass 2000, 2002a) (Table 3), the 287 addition of ~ 900 ppmw H₂O in our pyrope sample has negligible effects on the 288 individual elastic moduli C_{ij} values at ambient conditions within experimental 289 uncertainties. Our aggregate bulk, K_{S0} , and shear moduli, G_0 , at ambient conditions of 290 hydrous pyrope are 168.6(4) and 92.0(3) GPa, respectively (Table 4). The presence of 291 ~900 ppmw H₂O in our hydrous pyrope also does not affect the values of K_{S0} and G_0 292 within experimental uncertainties compared to anhydrous pyrope (Table 4). This is 293 consistent with the conclusion from the study by O'Neill et al. (1991), who deduced 294 that there might have no discernable effect of hydrogen on the elastic properties of 295 pyrope at ambient conditions, though the water solubility in their hydrous pyrope is 296 significantly smaller (~180 ppmw H₂O) than our hydrous pyrope (~900 ppmw H₂O).

Our results also show that hydration has no visible effect on the pressure and temperature derivatives of C_{ij} s in pyrope within experimental uncertainties (Table 3). This is clearly different from the effect of composition (e.g. Fe and Ca) on the temperature derivatives of C_{ij} s in pyrope-rich garnet, where a distinctly larger temperature derivative of C_{12} and lower temperature derivative of C_{11} and C_{44} for (Fe,Ca)-bearing pyrope-rich garnet reported by Lu et al. (2013).

303 The $(\partial K_{\rm S}/\partial P)_{\rm T}$ of our hydrous pyrope is 4.6 (Fig. 4 and Table 4), which is slightly 304 higher than most results of anhydrous pyropes $[(\partial K_S/\partial P)_T=4.1-4.51]$ (Table 4), except 305 a distinctly lower value ($(\partial K_{\rm S}/\partial P)_{\rm T}=3.2$) reported by Conrad et al. (2000) and higher 306 value ($(\partial K_{\rm S}/\partial P)_{\rm T}$ =5.3) reported by Chen et al. (1999). Moreover, the $(\partial G/\partial P)_{\rm T}$ of our 307 hydrous pyrope $((\partial G/\partial P)_T = 1.3)$ is indistinguishable from most previous studies 308 $((\partial G/\partial P)_T = 1.3 - 1.5)$ on anhydrous pyrope within experimental uncertainties, except for 309 two slightly higher values ($(\partial G/\partial P)_T$ =1.66 and 1.7) reported from ultrasonic 310 interferometry experiments by Gwanmesia et al. (2006) and Chantel et al. (2016), 311 respectively. Furthermore, our derived $(\partial K_{\rm S}/\partial T)_{\rm P}$ and $(\partial G/\partial T)_{\rm P}$ of hydrous pyrope are 312 indistinguishable from previous BLS values for anhydrous pyrope within their 313 uncertainties, but their absolute values appeared to be slightly lower than most of 314 those from ultrasonic interferometry measurements except a consistent $(\partial G/\partial T)_{\rm P}$ 315 absolute value reported by Chantel et al. (2016) (Table 4). Therefore, we conclude 316 that the presence of ~900 ppmw H₂O in our pyrope slightly enhances the $(\partial K_S / \partial P)_T$, 317 but does not distinctly affect the K_{S0} , G_0 , $(\partial G/\partial P)_T$, $(\partial K_S/\partial T)_P$, and $(\partial G/\partial T)_P$ within 318 their uncertainties.

319 Hydrogen effect on the elastic anisotropy of pyrope at high *P-T* conditions

320 Elastic wave anisotropy is a critical feature of the upper mantle (e.g. Karato 1998). 321 One advantage of using single-crystal samples in BLS is that we can obtain the full 322 elastic moduli and put a constraint on the elastic anisotropy. The elastic anisotropy of 323 minerals expresses the difference in stiffness of materials in different crystallographic 324 directions, which can provide insights into seismic anisotropy and can be an indicator 325 of the mechanical stability of materials (e.g. Hu et al. 2016; Sinogeikin and Bass 326 2000). Thus, knowledge of elastic anisotropy for hydrous pyrope at high P-T may 327 shed light on understanding the seismic anisotropy within the Earth's upper mantle.

328 For the cubic pyrope, the elastic anisotropy factor (*A*) can be expressed as (Karki et329 al. 1997; Sinogeikin and Bass 2000):

330
$$A = \frac{2C_{44} + C_{12}}{C_{11}} - 1 \tag{5}$$

331 where A indicates the deviation from elastic isotropy, with A=0 for an elastically 332 isotropic material. Analysis of these parameters using our data show that the absolute A values of our hydrous pyrope slightly decreases with increasing pressure and 333 334 temperature, where A are -0.026 at ambient pressure, 0.001 at 18.53 GPa, and -0.019 335 at 700 K. Moreover, our hydrous pyrope at ambient conditions has slightly higher 336 absolute A value than anhydrous pyrope (A=-0.006 to -0.008) (e.g. Sinogeikin and 337 Bass, 2000, 2002a; O'Neill et al. 1991). However, all these absolute A values are still 338 pretty small and very close to zero. Thus, our results indicate that hydrogen does not 339 have a significant effect on the elastic anisotropy of pyrope. Hydrous pyrope remains 340 elastically isotropic at high pressure/high temperature conditions compared to the 341 other (olivine, orthopyroxene, and clinopyroxene) major minerals in the upper mantle 342 (e.g. Mao et al. 2015; Sang et al. 2014; Zhang et al. 2016). Therefore, pyrope may not 343 have a significant contribution to seismic anisotropy in the upper mantle, at least 344 when its water content is less than 900 ppmw.

345 IMPLICATION

346 Hydrogen effect on the velocity profiles of pyrope in the Earth's upper mantle

347 The NAMs in the Earth's deep mantle may serve as a large internal reservoir of 348 water and are important for understanding the evolution and dynamics of Earth's 349 interior (e.g. Bell and Rossman 1992a; Beran and Libowitzky 2006; Ohtani 2005, 350 2015). Hydration of NAMs has been proposed to correlate with the observed velocity 351 anomalies in the Earth's mantle (e.g. Nolet and Zielhuis 1994; Ohtani 2005; Song et 352 al. 2004; van der Meijde et al. 2003). With the obtained elastic moduli of hydrous 353 pyrope at high P-T conditions in this study, we evaluate the effect of hydration on the 354 sound velocities of pyrope at upper mantle conditions.

355 The presence of ~900 ppmw water in pyrope lowers its V_P and V_S by ~0.7% at 356 ambient conditions (Sinogeikin and Bass 2000). Furthermore, for a better

357 understanding of the hydrogen influence on the velocity behavior of pyrope, we have 358 modeled the velocity profiles of hydrous pyrope along the upper mantle geotherm 359 (Katsura et al. 2010) and a cold subducted slabs geotherm (Eberle et al. 2002) using 360 the updated high-P/T elasticity results. Our modeling here is limited to the 361 upper-mantle region ranging from 200 km to 400 km depth because of the much more 362 complex mineralogical, geochemical, and seismic heterogeneities above 200 km depth (e.g. Jordan 1975; Grand and Helmberger 1984). The modeled results are then 363 364 compared with the velocity profiles of anhydrous pyrope (Sinogeikin and Bass 2000, 365 2002a) and (Fe, Ca)-bearing pyrope-rich garnet (Lu et al. 2013). Briefly, the 366 third-order Eulerian finite-strain equation and the third-order Birch-Murnaghan 367 equation of state (Birch 1978) are used to obtain the $K_{\rm S}$, G, $V_{\rm P}$ and $V_{\rm S}$ of relevant 368 minerals by extrapolating the experimentally-derived elastic moduli and their P-T derivatives to relevant P-T conditions (see Lu et al. 2013 for details). By allowing 369 370 elastic parameters to vary within their plausible ranges, the uncertainties $(\pm 1\sigma)$ of 371 extrapolation results can also be estimated.

Figure 6 shows the calculated velocity-depth relationships of hydrous pyrope along 372 373 with those of anhydrous phase. Because of larger pressure derivative of $K_{\rm S}$ and similar 374 temperature derivative of $K_{\rm S}$, the $V_{\rm P}$ of hydrous pyrope increases more rapidly with 375 depth than that of anhydrous pyrope. Moreover, the $V_{\rm P}$ of hydrous pyrope crosses and 376 exceeds that of anhydrous phase at ~200 km and ~270 km depth along the upper 377 mantle geotherm and the cold subducted slabs geotherm, respectively. However, 378 considering the error bars presented in Fig. 6, the $V_{\rm P}$ profiles are indistinguishable 379 between hydrous and anhydrous pyrope at 200-400 km depth. Similarly, due to the 380 small effects of hydrogen on the pressure derivative of G, the difference in the $V_{\rm S}$ 381 between hydrous and anhydrous pyrope is also within the uncertainties over the 382 200-400 km depth.

383 On the other hand, hydrogen can enhance the anelasticity (e.g. Karato 1995) and 384 may further change the $V_{\rm P}$ and $V_{\rm S}$. Combining the anelastic effect, the change of $V_{\rm P}$

385 and $V_{\rm S}$ associated with ~900 ppmw H₂O in pyrope at the Earth's upper mantle may be 386 greater than that observed here. However, although up to several hundred ppmw H_2O 387 have been found in some natural garnets (e.g. Aines and Rossman 1984a, 1984b; Li H 388 Y et al. 2018), most of the mantle-derived garnets typically contain <100 ppmw H₂O 389 (e.g. Beran and Libowitzky 2006; Ohtani 2015). Combined with the limited effect of 390 \sim 900 ppmw H₂O on the velocities of pyrope, we thus infer that hydrogen has no 391 significant effect on the velocity profiles of pyrope at upper mantle conditions. 392 Nevertheless, compared to the effect of hydrogen on the velocities of pyrope, the 393 effect of compositions (e.g. Fe and Ca) is significant. Accordingly, the velocity 394 reduction produced by the effect of Fe and Ca is $\sim 2-3\%$ either along the upper mantle 395 geotherm or along the cold subducted slabs geotherm (Fig. 6). Therefore, the 396 elasticity studies of hydrous (Fe, Ca)-bearing pyrope-rich garnet at high P-T 397 conditions are needed to provide a more comprehensive understanding of the coupled 398 effect of compositions (e.g. Fe and Ca) and hydration on the elasticity and velocity 399 profiles of garnet and then the Earth's upper mantle.

400 Hydrogen effect on the $V_{\rm P}/V_{\rm S}$ ratio of pyrope in the Earth's upper mantle

401 The $V_{\rm P}/V_{\rm S}$ ratio has been proposed as one of the possible indicators to determine the 402 composition in the deep Earth (e.g. Anderson and Bass 1984; Duan et al. 2018; Li and 403 Neuville 2010; Mao et al. 2010), such as the silica content of continental crust 404 (Christensen 1996). It has been widely used to infer the thermal and compositional 405 state of the upper mantle (e.g. Afonso et al. 2010; Chou et al. 2009; Lee 2003; Niu et 406 al. 2004; Speziale et al. 2005). Here, we have investigated the effect of hydration on 407 the $V_{\rm P}/V_{\rm S}$ ratio of pyrope. The $V_{\rm P}/V_{\rm S}$ ratio of the hydrous pyrope at ambient conditions 408 is 1.78, which is the same with the anhydrous pyrope (1.78) (Sinogeikin and Bass 2000). The V_P/V_S ratio increases with pressure at an average rate of 5.03×10^{-3} GPa⁻¹. 409 410 but it remains constant with increasing temperature from 300 K to 700 K. At the depth 411 of 400 km, the $V_{\rm P}/V_{\rm S}$ ratio of the hydrous pyrope increases to 1.86 along the upper 412 mantle geotherm and to 1.85 along the cold subducted slabs geotherm, which are $\sim 0.5\%$

higher than that of the anhydrous phase. A similar increase in the $V_{\rm P}/V_{\rm S}$ ratio caused 413 414 by hydration was observed for ringwoodite (Jacobsen and Smyth 2006; Sinogeikin et 415 al. 2003). However, this contrasts with the behavior of olivine for which a $\sim 0.7\%$ 416 decreased in the $V_{\rm P}/V_{\rm S}$ ratio was observed in the hydrous olivine relative to anhydrous 417 phase (Mao et al. 2010; Zha et al. 1996). Figure 7 also shows a comparison of the 418 $V_{\rm P}/V_{\rm S}$ ratio for hydrous and anhydrous pyrope with the (Fe,Ca)-bearing pyrope-rich 419 garnet. We notice although all garnets exhibit the increased V_P/V_S ratio with 420 increasing depths, the hydrous pyrope has the highest value throughout the upper 421 mantle depths. However, considering the uncertainties of the $V_{\rm P}/V_{\rm S}$ ratio presented in 422 Fig. 7, the variation of the $V_{\rm P}/V_{\rm S}$ ratio among these garnet samples should be limited 423 at 200-400 km depth. This confirms the results of the previous study, which indicated 424 that the variation of mineral composition has only a weak effect on the $V_{\rm P}/V_{\rm S}$ ratio of 425 the upper mantle (Duan et al. 2018). Finally, we infer that the hydrogen has also no 426 significant effect on the $V_{\rm P}/V_{\rm S}$ ratio of pyrope at upper mantle conditions, especially 427 for the limited hydration level (<100 ppmw H₂O) of mantle-derived garnets.

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750 Figure Captions

Figure 1. Representative Brillouin spectra of single-crystal hydrous pyrope at 18.6 GPa and 300 K (a), and 1 atm and 700 K (b). Open circles: experimental data; solid lines: fitted V_P and V_S peaks, respectively. The average collection time was ~ 40 min and ~ 20 min for each spectrum of high-pressure and high-temperature measurements, respectively. The (0.34, -0.53, 0.92) crystallographic plane of single-crystal hydrous pyrope sample was used for both BLS experiments. Experimental uncertainties are smaller than the symbols. (Color online.)

758

Figure 2. V_P and V_S velocities of single-crystal hydrous pyrope as a function of the azimuthal angel measured at 18.6 GPa and 300 K (a), and 1 atm and 700 K (b). Open circles: experimental data; solid lines: modeled results. Error bars are smaller than the symbols when not shown. (Color online.)

763

Figure 3. Single-crystal individual elastic moduli (C_{11} , C_{12} , and C_{44}) of hydrous pyrope as a function of pressure (a) and temperature (b) compared with previous study of anhydrous pyrope. Solid symbols represent our experimental data; solid lines are modeled results using the third-order finite-strain equation fitting (a) or a linear fitting (b); dashed lines represent the experimental data of anhydrous pyrope (Sinogeikin and Bass 2000, 2002). Error bars are smaller than the symbols when not shown. (Color online.)

771

772	Figure 4. Adiabatic bulk (K_S) and shear modulus (G) of hydrous pyrope as a function
773	of pressure (a) and temperature (b) compared with previous study of anhydrous
774	pyrope. Solid symbols represent our experimental data; solid lines are modeled results
775	using the third-order finite-strain equation fitting (a) or a linear fitting (b); dashed
776	lines represent the experimental data of anhydrous pyrope (Sinogeikin and Bass 2000,
777	2002). Error bars are smaller than the symbols when not shown. (Color online.)
778	
779	Figure 5. Aggregate compressional (V_P) and shear velocity (V_S) of hydrous pyrope as
780	a function of pressure (a) and temperature (b) compared with previous study of
781	anhydrous pyrope. Solid symbols represent our experimental data; solid lines are
782	modeled results using the third-order finite-strain equation fitting (a) or a linear fitting
783	(b); dashed lines represent the experimental data of anhydrous pyrope (Sinogeikin and
784	Bass 2000, 2002). Error bars are smaller than the symbols when not shown. (Color
785	online.)
786	
787	Figure 6. Modeled velocities of pyrope garnets in the Earth's upper mantle along the
788	upper mantle geotherm and cold subducted slabs geotherm. Red lines: hydrous pyrope
789	(this study; Fan et al. 2017b); blue lines: anhydrous pyrope (Sinogeikin and Bass,
790	2000, 2002; Zou et al. 2012a); black lines: (Fe,Ca)-bearing pyrope (Lu et al. 2013;
791	Thieblot et al. 1998). Error bars represent the propagated uncertainties ($\pm 1\sigma$). (Color

792 online.)

793

Figure 7. Comparison of modeled V_P/V_S ratio of hydrous pyrope with anhydrous pyrope and (Fe,Ca)-bearing pyrope in the Earth's upper mantle along the upper mantle geotherm and cold subducted slabs geotherm. Error bars represent the propagated uncertainties ($\pm 1\sigma$). (Color online.)

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Р	Density	C_{11}	C_{12}	C_{44}	Kv	K _R	Ks	$G_{ m V}$	$G_{ m R}$	G	V_{P}	$V_{\rm S}$	$V_{\rm P}/V_{\rm S}$	AV
(GPa)	(g/cm^3)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(km/s)	(km/s)		
0.0001	3.557(2)	294.5(5)	105.7(6)	90.5(4)	168.6(4)	168.6(4)	168.6(4)	92.1(2)	92.0(3)	92.0(3)	9.05(1)	5.09(1)	1.78(1)	-0.026(1)
0.9(1)	3.577(2)	301.7(5)	109.9(6)	91.4(4)	173.8(3)	173.8(3)	173.8(3)	93.1(2)	93.2(3)	93.2(3)	9.13(1)	5.10(1)	1.79(1)	-0.030(1)
3.4(2)	3.629(3)	316.8(6)	119.2(7)	94.3(5)	185.1(4)	185.1(4)	185.1(4)	96.0(2)	96.1(2)	96.1(2)	9.29(1)	5.15(1)	1.81(1)	-0.028(1)
5.8(2)	3.677(3)	331.6(6)	126.3(5)	97.7(4)	194.7(5)	194.7(5)	194.7(5)	99.6(2)	99.7(2)	99.7(2)	9.44(1)	5.21(1)	1.81(1)	-0.030(1)
7.5(1)	3.709(2)	339.7(7)	133.5(6)	100.6(5)	202.2(4)	202.2(4)	202.2(4)	101.6(2)	101.6(2)	101.6(2)	9.54(1)	5.23(1)	1.82(1)	-0.015(1)
10.1(2)	3.757(2)	355.2(7)	142.6(7)	103.5(5)	213.5(3)	213.5(3)	213.5(3)	104.6(2)	104.6(3)	104.6(3)	9.69(1)	5.28(1)	1.84(1)	-0.016(1)
12.3(2)	3.795(3)	368.5(8)	150.6(7)	106.8(5)	223.2(3)	223.2(3)	223.2(3)	107.6(1)	107.7(2)	107.7(2)	9.83(1)	5.33(1)	1.85(1)	-0.012(1)
15.5(3)	3.849(3)	387.4(7)	162.5(6)	110.7(6)	237.5(4)	237.5(4)	237.5(4)	111.4(2)	111.4(2)	111.4(2)	10.01(1)	5.38(1)	1.86(1)	-0.009(1)
18.6(2)	3.899(3)	404.6(8)	174.8(8)	115.2(6)	251.4(5)	251.4(5)	251.4(5)	115.1(1)	115.1(2)	115.1(2)	10.19(1)	5.43(1)	1.88(1)	0.0010(2)

Table 1 Densities, elastic moduli and aggregate velocities of hydrous pyrope at high pressure and ambient temperature

Numbers in parenthesis represent standard deviations.

		,		66 6		5	1 2 1		1	8	1		
Density	C_{11}	C_{12}	C_{44}	$K_{\rm V}$	$K_{\rm R}$	Ks	$G_{ m V}$	G_{R}	G	$V_{\rm P}$	$V_{\rm S}$	$V_{\rm P}/V_{\rm S}$	
(g/cm^3)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(km/s)	(km/s)		

Table 2 Densities, elastic moduli and aggregate velocities of hydrous pyrope at ambient pressure and high temperature

Т	Density	C_{11}	C_{12}	C_{44}	$K_{\rm V}$	$K_{\rm R}$	$K_{\rm S}$	$G_{ m V}$	G_{R}	G	$V_{\rm P}$	V_{S}	$V_{\rm P}/V_{\rm S}$	AV
(K)	(g/cm^3)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(GPa)	(km/s)	(km/s)		
300 ^a	3.557(3)	294.8(6)	105.5(4)	90.7(5)	168.6(5)	168.6(5)	168.6(5)	92.2(3)	92.3(4)	92.3(4)	9.05(1)	5.09(1)	1.78(1)	-0.027(1)
300	3.557(2)	294.5(5)	105.7(6)	90.5(4)	168.6(4)	168.6(4)	168.6(4)	92.0(2)	92.1(3)	92.0(3)	9.05(1)	5.09(1)	1.78(1)	-0.026(1)
400	3.546(3)	291.8(4)	104.9(4)	89.9(2)	167.2(3)	167.2(3)	167.2(3)	91.3(2)	91.3(2)	91.3(2)	9.03(1)	5.07(1)	1.78(1)	-0.024(1)
500	3.534(2)	289.1(5)	104.1(6)	89.3(3)	165.8(3)	165.8(3)	165.8(3)	90.6(1)	90.6(2)	90.6(2)	9.00(1)	5.06(1)	1.78(1)	-0.022(1)
600	3.523(3)	286.6(5)	103.2(5)	88.7(4)	164.3(3)	164.3(3)	164.3(3)	89.9(2)	89.9(2)	89.9(2)	8.98(1)	5.05(1)	1.78(1)	-0.021(1)
700	3.511(3)	283.4(5)	102.3(6)	87.9(4)	162.7(4)	162.7(4)	162.7(4)	89.0(2)	89.0(3)	89.0(3)	8.95(1)	5.03(1)	1.78(1)	-0.019(1)

Numbers in parenthesis represent standard deviations.

T

^{a)} Represents measurement at temperature decreased from 700 K to room temperature.

Table 3 Single-crystal elastic moduli and their pressure and temp	erature derivatives of hydrous	pyrope at ambient conditions	s in comparison to
previous studies ^a			

I											
References	Composition	Method	C_{11}	C_{12}	C_{44}	$(\partial C_{11}/\partial P)_{\mathrm{T}}$	$(\partial C_{12}/\partial P)_{\mathrm{T}}$	$(\partial C_{44}/\partial P)_{\mathrm{T}}$	$(\partial C_{11}/\partial T)_{\rm P}$	$(\partial C_{12}/\partial T)_{\rm P}$	$(\partial C_{44}/\partial T)_{\rm P}$
			(GPa)	(GPa)	(GPa)				(GPa/K)	(GPa/K)	(GPa/K)
This study	Hydrous Prp100 °	BLS	294.5(5)	105.7(6)	90.5(4)	6.2(1)	3.7(1)	1.5(1)	-0.028(1)	-0.009(1)	-0.006(1)
O'Neill et al. (1991)	Hydrous Prp100 ^d	BLS	296.2(5)	111.1(6)	91.6(3)	b	b	b	b	b	b
Leitner et al. (1980)	Prp100	BLS	295(2)	117(1)	90(3)	b	b	b	b	b	b
Sinogeikin and Bass (2000)	Prp100	BLS	297(3)	108(2)	93(2)	5.8(4)	3.2(4)	1.3(3)	b	b	b
Sinogeikin and Bass (2002)	Prp100	BLS	298(3)	107(2)	93(2)	b	b	b	-0.031(3)	-0.006(2)	-0.007(2)
Lu et al. (2013)	Prp68Alm24Grs5	BLS	290(2)	106(2)	92.2(6)	6.0(1)	3.5(1)	1.2(1)	-0.021(4)	-0.0163(5)	-0.003(1)

Numbers in parenthesis represent standard deviations. Prp: Pyrope; Alm: Almandine; Grs: Grossular; BLS: Brillouin Light Scattering.

^{a)} Only Brillouin scattering results are listed for pyrope garnet.

^{b)} The value is not available in the text.

^{c)} ~900 ppmw H₂O.

^{d)} ~180 ppmw H₂O.

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Table 4 Bulk and shear	r moduli and	their pressure	and	temperature	derivatives	of hydrous	pyrope a	t ambient	conditions	in	comparison	to
previous studies ^a												

References	Composition	Methods	$K_{\rm S0}$	G_0	$(\partial K_{\rm S}/\partial P)_T$	$(\partial G/\partial P)_T$	$(\partial K_{\rm S}/\partial T)_P$	$(\partial G/\partial T)_P$
			(GPa)	(GPa)			(GPa/K)	(GPa/K)
This study	Hydrous Prp100	BLS	168.6(4)	92.0(3)	4.6(1)	1.3(1)	-0.015(1)	-0.008(1)
O'Neill et al. (1991)	Hydrous Prp100	BLS	172.8(3)	92.0(2)	b	b	b	b
Leitner et al. (1980)	Prp100	BLS	177(1)	89(1)	b	b	b	b
Conrad et al. (2000)	Prp100	BLS	172.7 °	92 °	3.2 °	1.4 °	b	b
Sinogeikin and Bass (2000)	Prp100	BLS	171(2)	94(2)	4.1(3)	1.3(2)	b	b
Sinogeikin and Bass (2002)	Prp100	BLS	171(2)	94(2)	b	b	-0.014(2)	-0.009(1)
Lu et al. (2013)	Prp68Alm24Grs5	BLS	168(2)	92(1)	4.4(1)	1.2(1)	-0.017(1)	-0.005(1)
Chen et al. (1999)	Prp100	UI	171(20)	92(1)	5.3(4)	1.6(2)	b	b
Gwanmesia et al. (2006)	Prp100	UI	175(2)	91(1)	3.9(3)	1.7(2)	-0.018(2)	-0.010(1)
Zou et al. (2012)	Prp100	UI	170.0(2)	93.2(1)	4.51(2)	1.51(2)	-0.0170(1)	-0.0107(1)
Chantel et al. (2016)	Prp100	UI	172(2)	89.1(5)	4.38(8)	1.66(5)	-0.018(2)	-0.008(1)

Numbers in parenthesis represent standard deviations.

^{a)} Only Brillouin scattering and ultrasonic interferometry results are listed for pyrope garnet.

^{b)} The value is not available in the text.

^{c)} The uncertainty is not available in the text.



Fig. 1



Fig. 2









Fig. 6

