

Revision 1

Surface energy of fayalite and its effect on Fe-Si-O oxygen buffers and the olivine-spinel transition

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Abstract

The surface energy (hydrated surfaces) of fayalite (α -Fe₂SiO₄) was determined to be 2.47 ± 0.25 J/m² using high temperature oxide melt solution calorimetry. This is larger than the surface energy of magnetite (Fe₃O₄), but lower than that of forsterite (α -Mg₂SiO₄). The changes in the positions of the quartz/fayalite/magnetite (QFM) and quartz/iron/fayalite (QIF) buffers with particle size reduction were calculated. QFM is lowered in fO₂ by 3-7 log units as a function of temperature for 30 nm particles while QIF is raised by 1-2 log units. The estimated surface energy difference between olivine and spinel polymorphs decreases the pressure of the olivine-spinel transition in Fe₂SiO₄ by about 1 GPa.

INTRODUCTION

Oxidation - reduction equilibria of nanoscale iron oxides have been shown to depend significantly on particle size because the phases involved have different surface energies (Navrotsky et al. 2010). Although surface energies have been measured for hematite, maghemite - magnetite solid solutions, and ulvospinel (Majzlan et al. 2003; Lilova et al. 2014), few experimental data are available for silicates. Because fayalite (α -Fe₂SiO₄) is an important component of the rock-forming olivines, and is involved in a number of “redox buffers” that define oxygen fugacity in synthetic and petrologic systems, including important extraterrestrial materials, and participates in high pressure phase transitions, its surface energy is needed to understand the behavior of these reactions when small particles are involved. In addition, because there has been significant discussion of olivines as mantle water reservoirs (Férot and Bolfan-Casanova 2012), this work determines the enthalpy of the hydrated fayalite surface using high temperature oxide melt solution calorimetry of one bulk and several nanophase fayalite samples.

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EXPERIMENTAL PROCEDURES

Materials

Several samples of fayalite were analyzed in this experiment. Powdered nanophase samples, indicated as fayalite FS21, 22, 23, 24, and 26, were used as received from Michael DeAngelis. A sample of bulk fayalite was prepared by Donald Lindsley (5-15-17 fayalite). The synthesis of the nanophase fayalite was described in detail by DeAngelis et al. (2012). The bulk Fe_2SiO_4 (5-15-17 fayalite) was prepared by mixing silica and Fe_2O_3 and grinding them together for 2 hours under ethanol in an agate mortar. "Fe-sponge" was then added and ground for an additional 40 minutes. The mix was wrapped in silver foil and inserted into a silica-glass tube, which was evacuated and sealed, then annealed in a furnace at 930–940 °C for 10 days. X-ray diffraction (XRD) showed that the product was ~97 % fayalite, ~1.4 % quartz, ~1.3 % wüstite, and ~0.5% metallic Fe. This material was then re-ground and annealed for 15 days to react the small amount of silica, wüstite, and Fe to fayalite. A third heating was performed with additional Fe-sponge due to a small leak that caused a partial oxidation of one end of the sample during the second run. This was run for about 45 days, XRD of the final product showed ~1 % free quartz, ~99 % fayalite, with no unreacted wüstite. Optical examination in refractive index oil confirmed the presence of quartz; other than a few globules of silver, there were no opaque particles, strongly suggesting the absence of hematite, wüstite, magnetite, and metallic Fe. Absence of the latter two was further confirmed by testing with a magnet, to which no sample adhered.

Characterization

The amount of surface water on each of the nano-sized samples was determined by simultaneous thermogravimetry and differential scanning calorimetry (TG-DSC). Each fayalite sample was heated in argon from room temperature to 1073 K at 10 K/min in a Setaram Labsys Evo instrument.

X-ray diffraction (XRD) of the bulk sample was performed on a Rigaku Miniflex 600, using $\text{Cu K}\alpha$ radiation. Data were processed with Match! 3TM software, and the unit-cell parameters were determined by Rietveld refinement using FullProfTM. The characterization of the nano fayalite samples was described in detail by DeAngelis et al. (2012).

Calorimetry

High temperature oxide melt solution calorimetry was performed using the Setaram AlexSYS - the commercial version of the Tian Calvet twin calorimeter described previously by Navrotsky (1977, 1997, 2014). In the drop solution calorimetry experiment, samples in the form of pellets (between 4 and 6 mg) were dropped from room temperature (298 K) into molten $2\text{PbO}\cdot\text{B}_2\text{O}_3$ solvent at 1073 K in a platinum crucible. Pure oxygen was flushed through the system and bubbled through the solvent at 110 and 5 ml/min, respectively, to remove the evolved water vapor and assist in the complete oxidation of Fe^{2+} to Fe^{3+} . More details about the reproducible final oxidation state of Fe-containing compounds dissolved in lead borate can be found in Lilova et al. (2012).

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95 RESULTS

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97 Characterization

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99 The XRD of the bulk sample showed fayalite ($a = 4.815$, $b = 10.468$, $c = 6.084$ Å) with
100 about 1 % free silica as quartz. According to DeAngelis et al. (2012), the nano fayalite samples
101 contain 3 wt. % or less metallic Fe (0.6 wt. % for FS22). The weight loss from the
102 thermogravimetric analysis corresponds to the total water content (physisorbed and chemisorbed)
103 on the surface of the sample. While it is possible that simultaneous weight gain may have
104 occurred due to oxidation, the product was checked by XRD after the TGA measurement and
105 found to be a coarsened fayalite.

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107 Enthalpy of formation and surface enthalpy of fayalite

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109 The calorimetric data for the six samples analyzed in this study are given in Table 1. The
110 enthalpy of formation of bulk fayalite from oxides at 298 K is calculated as -25.91 ± 2.42 kJ/mol,
111 using the thermodynamic cycle shown in Table 2. Calculation of the enthalpy of formation of all
112 samples, however, requires that the measured drop solution enthalpies be corrected for the water
113 and the small amount of Fe reported in the nano fayalite, and the silica in the bulk one (Table 1).
114 The 3 wt.% of Fe reported in the nanophase Fe_2SiO_4 sample would result in a change of
115 approximately -36 kJ/mol. This is a significant change relative to the values measured for each
116 sample. For example, the drop solution enthalpy of sample FS23 is -38.63 kJ/mol. If FS23
117 contains 3 wt % Fe, ΔH_{ds} becomes -2.66 kJ/mol. After the water correction, the final drop
118 solution enthalpy would be -47 kJ/mol, which is practically the drop solution enthalpy of bulk
119 fayalite (Table 3), i.e. the nanophase compound would appear to be as thermodynamically stable
120 as its bulk counterpart. It should be noted, however, that FS23 has the smallest particle size.
121 Similar calculations performed for FS21 would result in a drop solution enthalpy of -38 kJ/mol
122 after all corrections, which would indicate higher energetic stability for the nano fayalite than the
123 bulk one. In fact, any Fe content above or equal to 1 wt % would result in drop solution
124 enthalpies more endothermic than -48.15 kJ/mol for at least one of the nano fayalite. This seems
125 unreasonable and we, therefore, conclude that the Fe content in these nano samples is likely to be
126 significantly less than 3 % and may be closer to the value found for FS22 (0.6 wt %) (DeAngelis
127 et al., 2012). Hence all drop solution enthalpies for the nano samples have been corrected for 0.6
128 wt% Fe. Similarly, the drop solution enthalpy of bulk fayalite was corrected for 1 wt % quartz.
129 We realize that the nanophase fayalite samples have compositional uncertainties but realistically,
130 better samples cannot be readily prepared at present. Though the cumulative errors arising from
131 iron content, water content, and possible sample heterogeneity cannot be rigorously quantified,
132 the data, as the first experimental measurement of fayalite surface energy, are new and useful.

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134 To compare the results from this study with reported values for fayalite the values must
135 also be corrected for the composition of wüstite, as the enthalpy of formation from the oxides
136 should be to end-member FeO. The enthalpy of the reaction $0.106\text{Fe} + 2\text{Fe}_{0.947}\text{O} = 2\text{FeO}$ is -7.11
137 ± 5.39 kJ/mol (Robie and Hemingway 1995). The drop solution enthalpy of pure Fe in lead
138 borate solvent at 1073 K was calculated using the thermodynamic cycle in Table 4. Correcting
the measured values to an enthalpy of formation of fayalite from stoichiometric FeO and SiO₂

139 then yields -18.80 ± 5.91 kJ/mol. This is reasonably consistent with values from the literature (-
140 17.2 ± 6.3 kJ/mol, Thierry et al. 1981; -20.5 kJ/mol, Hewitt 1978; -24.6 ± 2.1 kJ/mol, Robie et al.
141 1978).

142 The corrected drop solution enthalpies of all fayalite samples as a function of surface area
143 are shown on Figure 1. The data have been fitted to a straight line, the absolute value of the slope
144 of which, 2.47 ± 0.25 J/m², is equal to the surface energy. The error is two standard deviations of
145 the slope. Because the correction for water content was made using the heat content of pure H₂O,
146 neglecting any interactions of water with the surface, this energy refers to that of the hydrated
147 surface, as discussed previously (Ushakov et al. 2005). Water adsorption calorimetry on the nano
148 fayalite samples would be needed to make appropriate corrections to obtain the enthalpy of the
149 anhydrous surface. Unfortunately, insufficient material was available for such studies.

150 For all the calculations below, we have used the reported surface energies of the hydrated
151 surfaces for consistency. This is reasonable, as nanoparticles hold on strongly to their surface
152 water and its total removal often coincides with coarsening (Navrotsky, 2011). Thus, calculation
153 of phase equilibria under most laboratory and natural conditions where nanoparticles occur
154 generally should be done using the energetics of the hydrated surfaces.

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157 DISCUSSION

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159 Oxygen buffers at the nanoscale

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161 Oxygen buffers are used to control or estimate the oxygen fugacity of synthetic and
162 natural systems. Several of these, such as magnetite/hematite, quartz/fayalite/magnetite,
163 magnetite/wüstite, and wüstite/iron are based on oxidation - reduction reactions involving
164 metallic, divalent and trivalent Fe. It is possible that the phases involved may have small particle
165 sizes and thus the differences in surface energies between the bulk and nano-phases may shift the
166 redox equilibria (Navrotsky et al. 2010), especially at low to moderate temperatures, where
167 recrystallization/crystal growth are less likely, or in cases of shear or fault zones (Krot et al.
168 2000; Sammis and Ben-Zion 2008),

169 In order to evaluate the importance of nanoparticle size on fayalite-bearing redox buffers,
170 we have calculated the positions of the quartz/fayalite/magnetite and quartz/iron/fayalite buffers

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175 for 100, 50, and 30 nm diameter particles. The fayalite/wüstite buffers were not assessed,
176 because Fe_{1-x}O nanoparticles are not thermodynamically stable below around 1000 K, and at
177 such temperatures it is likely wüstite will coarsen (Navrotsky et al. 2010). This limitation,
178 combined with the non-stoichiometry of wüstite complicates the calculations and significantly
179 increases the uncertainty.

180 The oxygen fugacities of the bulk and nano QFM buffers were calculated for reaction 1
181 using eqs. 4 – 6.

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$$183 \quad \log p_{\text{O}_2} = \frac{\Delta G_r}{2.3RT} \quad (4)$$

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$$\log p_{O_2} = \frac{\Delta G_r + \Delta G_r^{ex}}{2.3RT} \quad (5)$$

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$$\Delta G_r^{ex} = 2(\Delta H_{surface, Fe_3O_4}) + 3(\Delta H_{surface, SiO_2}) - 3(\Delta H_{surface, Fe_2SiO_4}) \quad (6)$$

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189 where ΔG_r is the Gibbs free energy of the bulk QFM reaction, ΔG_r^{ex} is the excess Gibbs free
190 energy associated with the nano QFM reaction, and $\Delta H_{surface}$ is the total enthalpy of the surface
191 of each phase, i.e the surface enthalpy times the surface area per mole of phase corresponding to
192 the given particle size. Analogous formulae were used for the QIF buffer with the respective
193 coefficients. All values are shown in Tables 5 and 6, and the results of the calculations are shown
194 in Figures 2-3. The results for the bulk redox buffers between 300 and 900 K are consistent with
195 the data of Misra (2012). The results suggest that decreases in fayalite particle size significantly
196 destabilize it, decreasing the oxygen fugacity of the QFM buffer and favoring the stability of
197 magnetite (Fig. 2). The calculated shift in QFM is 3-7 log units for the smallest particle size (30
198 nm), depending on temperature. The effect on the oxygen fugacity of the QIF buffer (Fig. 3) is to
199 increase its oxygen fugacity, again destabilizing fayalite to a small extent although, since the
200 amount of oxygen in the QIF reaction is smaller than in QFM (equations 1 and 2) the effect is
201 smaller. Thus, we conclude that the QFM buffer is more sensitive to particle size effects than
202 QIF.

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204 **High pressure olivine - spinel phase transition**

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206 Mantle olivines are predominantly composed of forsterite, but also contain a significant
207 fayalite component (Ringwood 1975; Irfune and Isshiki 1998). The olivine/spinel transition is a
208 well-known feature of the mantle (Hart 1969; Ringwood 1970). It includes the two transitions at
209 the Mg-rich end (forsterite to wadsleyite and wadsleyite to ringwoodite) and the one transition at
210 the Fe-rich end (fayalite to ahrensite).

211 The surface enthalpy (hydrated surface) of fayalite (2.47 J/m^2) is higher than that of
212 typical Fe spinels (0.80 J/m^2 for magnetite and ulvospinel, Lillova et al. 2012a, 2012b), and
213 smaller than that of the hydrous forsterite surface (3.37 J/m^2 , Chen et al. 2009). If the surface
214 enthalpy of ahrensite $\gamma\text{-Fe}_2\text{SiO}_4$ is similar to that of other Fe-bearing spinels, this difference
215 between fayalite and spinel will result in thermodynamic stabilization of the spinel relative to the
216 olivine, which will, therefore, reduce the pressure of the Fe-rich olivine-spinel transition for
217 small particle sizes. To estimate the effect of this difference on the pressure of the phase
218 transition, the following approximate calculation was performed.

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$$P = \frac{T\Delta S - \Delta H}{\Delta V_m} \quad (8)$$

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224 where P, ΔS , ΔH , ΔV_m are the pressure, entropy, enthalpy, and molar volume of the transition.
225 The effects of thermal expansion and compressibility were neglected, and the enthalpy, entropy,
226 and volume of the transition were assumed constant for illustrative purposes, as our goal is to
227 estimate the magnitude of the pressure change resulting from the decrease in particle size, not to

228 make an accurate calculation of the equilibrium pressure. Assuming 2.47 J/m^2 surface energy
229 (hydrated surface) for fayalite and 0.8 J/m^2 (the same as for magnetite) for ahrensite (γ -
230 Fe_2SiO_4), for a particle size of 100 nm, the enthalpy of the transition from α - Fe_2SiO_4 (fayalite) to
231 γ - Fe_2SiO_4 (ahrensite) decreases by about 5 kJ/mol, yielding a pressure drop of approximately 1
232 GPa. Since the equilibrium pressure for this transition near 1273 K is around 6 GPa (Akaogi et
233 al. 1989) this represents a drop of about 17 %, which is outside the error of pressure calibrations
234 and may be significant. While the expected temperature of this transition is large enough that
235 nanoparticles are likely to recrystallize, diminution of particle size during shearing or as a
236 consequence of phase transition (Riedel and Karato 1997; Jackson 2000; Michibayashi et al.
237 2006) could generate such effects. This pressure drop is comparable, on a percentage basis, to
238 that calculated for the forsterite – wadsleyite transition in Mg_2SiO_4 calculated by Chen and
239 Navrotsky (2010), who measured the surface enthalpy of forsterite for both hydrated and
240 anhydrous surfaces by a combination of oxide melt solution calorimetry and water adsorption
241 calorimetry. They then calculated the change in the transition pressure for the forsterite -
242 wadsleyite transition assuming that β - Mg_2SiO_4 has a surface enthalpy similar to that of MgAl_2O_4
243 spinel. For an average grain size of 100 nm, 4.4 J/m^2 surface energy for the olivine and 1.8 J/m^2
244 for the β - Mg_2SiO_4 , the calculated transition pressure was lowered by about 2 GPa (above 19 %)
245 and the transition enthalpy by around 7 kJ/mol. Again, particle size reduction favors the high
246 pressure phase, which has the lower surface enthalpy. If ringwoodite and wadsleyite have similar
247 surface energies, both being spinel-related structures, then the β - γ transition should be much less
248 affected by particle size reduction than the α - β transition.

249 There are no experimental data available for the surface energy of olivine, wadsleyite, or
250 ringwoodite solid solutions with $\text{Mg}/(\text{Mg}+\text{Fe})$ near 0.9, which is characteristic of mantle
251 olivines. Near that composition, the high pressure transitions are solid solution loops in the $\alpha+\beta$
252 and $\beta+\gamma$ fields. The assumptions made by Chen and Navrotsky (2010) in calculating the
253 forsterite - wadsleyite transition in Mg_2SiO_4 can, however, be generalized to the solid solutions.
254 One can reasonably assume that in the olivine phase the surface energy varies linearly between
255 those of forsterite and fayalite. Unfortunately, there are few constraints on the surface energies of
256 wadsleyite and ringwoodite solid solutions. Nonetheless, within the framework of the
257 assumptions made by Chen and Navrotsky (2010), one can perform very approximate
258 calculations suggesting that, within rather large uncertainties, particle size diminution to 100 nm
259 would lower the pressure of first appearance of wadsleyite for an olivine of $\text{Mg}/(\text{Mg}+\text{Fe}) = 0.9$
260 by about 2 GPa, similar to that for pure Mg_2SiO_4 .

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264 **IMPLICATIONS**

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266 The surface energy of fayalite (α - Fe_2SiO_4) is significantly higher than that of magnetite
267 (FeFe_2O_4) and other spinels, including, presumably, ahrensite γ - Fe_2SiO_4 . Thus, Fe-bearing spinel
268 phases are thermodynamically stabilized relative to fayalite at the nanoscale. This difference in
269 surface energy causes a significant shift in the quartz/fayalite/magnetite (QFM) buffer to lower
270 oxygen fugacity and the Fe_2SiO_4 olivine-spinel transition to lower pressure. The
271 quartz/iron/fayalite (QIF) buffer is less affected. Thus, under petrologic regimes where fine
272 particles are likely to form and persist (e.g. sedimentary and low grade metamorphic conditions,
273 high strain rates, phase transition induced grain size diminution), one must be cautious in

274 applying established oxygen buffers and geobarometers/geothermometers to estimate conditions
275 of phase assemblage formation. For particle sizes below 100 nm, corrections of several orders of
276 magnitude in oxygen fugacity and 1-2 GPa in pressure may be required. In addition, particle-
277 size-driven local shifts in the oxygen fugacity of the QFM buffer could be relevant to redox-
278 related processes in the mantle transition zone such as mid-ocean ridge basalt glass melting and
279 crystallization, element partitioning, and volcanic degassing. As the oxidized to total Fe ratio is
280 expected to be a function of particle size this can affect the bulk composition and geophysical
281 properties and of the mineral assemblage, and thus related geologic processes.

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414 **Fig. 1.** Corrected drop solution enthalpy of five nanophase and one bulk fayalite samples fitted
415 as a linear function of the surface area with $R^2 = 0.997$. The fit yields a surface energy of the
416 hydrated surface of $2.47 \pm 0.25 \text{ J/m}^2$

417 **Fig. 2.** QFM buffer curve change with particle size. The nano curves are calculated for 100, 50,
418 and 30 nm particle diameter with data for hydrated surfaces.

419 **Fig. 3.** QIF buffer curve change with particle size. The nano curves are calculated for 100, 50,
420 and 30 nm particle diameter with data for hydrated surfaces.

421

Table 1. Calorimetric data for all fayalite samples. The drop solution enthalpies are corrected for impurities and total water content as discussed in the text. Reported errors are two standard deviations of the mean, (n) is number of experiments performed.

Sample	BET area m ² /mol	Particle size, calculated from BET, nm	ΔH_{ds} kJ/mol	ΔH_{ds} kJ/mol, corrected for water and impurities
bulk	0	>1000	$-47.19 \pm 0.92(8)$	-48.15 ± 0.95
FS22	2568	108	$-54.71 \pm 1.72(8)$	-51.79 ± 1.73
FS21	4320	64	$-50.85 \pm 0.82(8)$	-63.51 ± 0.84
FS24	5456	50	$-45.56 \pm 2.03(9)$	-66.27 ± 2.04
FS26	5616	49	$-45.28 \pm 1.36(12)$	-67.24 ± 1.37
FS23	9577	29	$-38.63 \pm 1.30(12)$	-72.91 ± 1.32

Table 2. Thermodynamic cycle for calculating enthalpy of formation of fayalite

Fe_2SiO_4 (s, 298 K) + 0.5O ₂ (g, 1073 K) → Fe ₂ O ₃ (dis, 1073K) + SiO ₂ (dis, 1073K)	ΔH_{ds}
Fe (s, 298 K) + 0.75O ₂ (g, 1073 K) → 0.5Fe ₂ O ₃ (dis, 1073 K)	ΔH_2
O ₂ (g, 298 K) → O ₂ (g, 1073 K)	ΔH_3
Fe _{0.947} O (s, 298 K) + 0.21025 O ₂ (g, 1073 K) → 0.4735 Fe ₂ O ₃ (dis, 1073 K)	ΔH_5
SiO ₂ (s, 298 K) → SiO ₂ (dis, 1073 K)	ΔH_6
0.106Fe (s, 298 K) + 2Fe_{0.947}O (s, 298 K) + SiO₂ (s, 298 K) → Fe₂SiO₄ (s, 298 K) $\Delta H_{f,ox} = -\Gamma H_{ds} + 0.106\Gamma H_2 + 2\Gamma H_5 + \Gamma H_6$	

Table 3 Iron, wüstite, magnetite and hematite drop solution enthalpies (ΔH_{ds}) in lead borate, 1073 K.

Sample	ΔH_{ds} [kJ/mol]	ΔH_{ox} [kJ/mol]	$\Delta H_{f,el}$ [kJ/mol]
Fe	-344.21 ± 4.44 (calculated)	-830.51 ± 4.53	0
Fe _{0.947} O	-42.68 ± 1.08^1	-124.17 ± 1.25	-268.49 ± 1.70
Fe ₃ O ₄	147.58 ± 1.12^1	-120.37 ± 2.30	-1124.20 ± 4.30
α -Fe ₂ O ₃	182.29 ± 1.34^1	0	-829.70 ± 2.42
quartz SiO ₂	47.79 ± 0.32^2	0	-910.7 ± 1

¹Lilova et al., 2012

²Navrotsky, 2014

Table 4. Thermodynamic cycle for calculating enthalpy of drop solution of iron metal in lead borate at 1073 K.

$\text{Fe}_3\text{O}_4 (\text{s}, 298 \text{ K}) + 0.25\text{O}_2 (\text{g}, 1073 \text{ K}) \rightarrow 1.5 \text{Fe}_2\text{O}_3 (\text{dis}, 1073 \text{ K})$	ΔH_1
$\text{Fe} (\text{s}, 298 \text{ K}) + 0.75\text{O}_2 (\text{g}, 1073 \text{ K}) \rightarrow 0.5\text{Fe}_2\text{O}_3 (\text{dis}, 1073 \text{ K})$	ΔH_2
$\text{O}_2 (\text{g}, 298 \text{ K}) \rightarrow \text{O}_2 (\text{g}, 1073 \text{ K})$	ΔH_3
$3\text{Fe} (\text{s}, 298 \text{ K}) + 2\text{O}_2 (\text{g}, 298 \text{ K}) \rightarrow \text{Fe}_3\text{O}_4 (\text{s}, 298 \text{ K})$	ΔH_4
$\Delta H_4 = -\Delta H_1 + 2*\Delta H_3 + 3*\Delta H_2$	

Table 5. Gibbs free energies of the bulk and nano buffer reactions. The Gibbs energies of the bulk reaction are calculated using the Gibbs energies of formation from Robie et al. (1978), Hemingway (1990), and Jacob et al. (1989). The excess Gibbs energies for the nano reaction are calculated using the surface energies from Lilova et al. (2012), Navrotsky et al. (2010) and this work. The excess Gibbs energy is assumed equal to the excess enthalpy as obtained from the surface energy differences, as done previously (Navrotsky et al., 2010).

Reaction	ΔG_r , kJ/mol	ΓG_r^{ex} 100 nm, kJ/mol	ΓG_r^{ex} 50 nm, kJ/mol	ΓG_r^{ex} , 30 nm kJ/mol
QFM buffer	$-530.82 + 0.22T \pm 3.52$	-12.24 ± 2.13	-24.47 ± 4.24	-40.79 ± 7.07
QIF buffer	$-565.36 + 0.14T \pm 1.62$	4.16 ± 0.74	8.33 ± 1.47	13.88 ± 2.44

Table 6. The surface enthalpies and surface areas of all phases (hydrated surface) for 100, 50, and 30 nm particle diameter

Phase	Surface energy, J/m ²	Surface area for 100 nm, m ² /mol	Surface area for 50 nm, m ² /mol	Surface area for 30 nm, m ² /mol	$\Delta H_{\text{surface}}$, 100 nm kJ/mol	$\Delta H_{\text{surface}}$, 50 nm kJ/mol	$\Delta H_{\text{surface}}$, 30 nm kJ/mol
Fe	1.59 ± 0.22^1	426	851	1418	0.68 ± 0.09	1.35 ± 0.19	2.25 ± 0.31
Fe ₃ O ₄	0.80 ± 0.05^2	2697	5395	8992	2.16 ± 0.13	4.32 ± 0.27	7.19 ± 0.45
Fe ₂ O ₃	0.75 ± 0.16^1	1808	3616	6026	1.36 ± 0.29	2.71 ± 0.58	4.52 ± 0.96
SiO ₂	1.0 ± 0.1^3	1361	2723	4538	1.36 ± 0.14	2.72 ± 0.27	4.54 ± 0.45
Fe ₂ SiO ₄	2.47 ± 0.25	27850	5570	9283	6.88 ± 0.7	13.76 ± 1.39	22.93 ± 2.32

¹Navrotsky et al., 2010

²Lilova et al., 2012

³Parks 1990

Figure 1

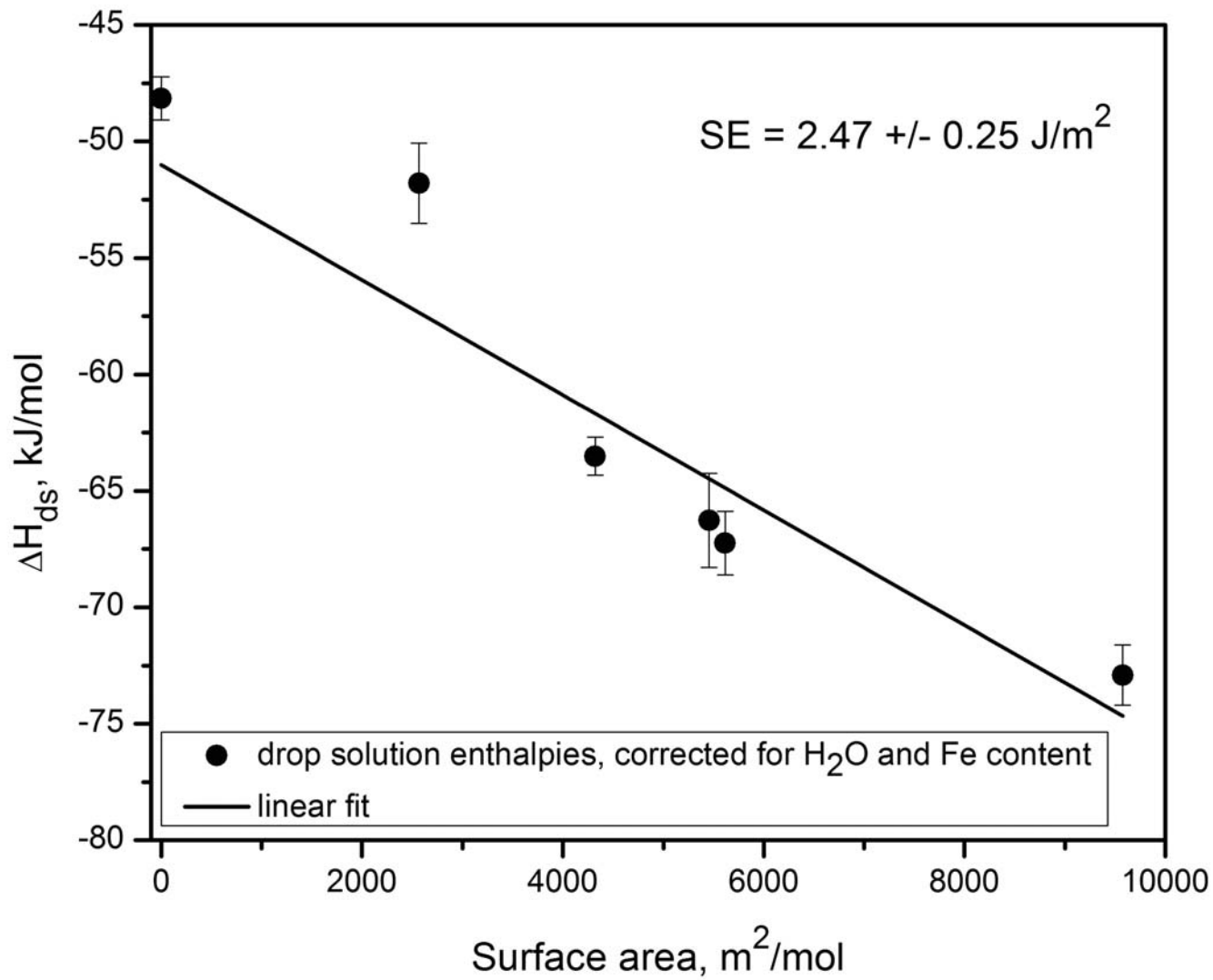


Figure 2

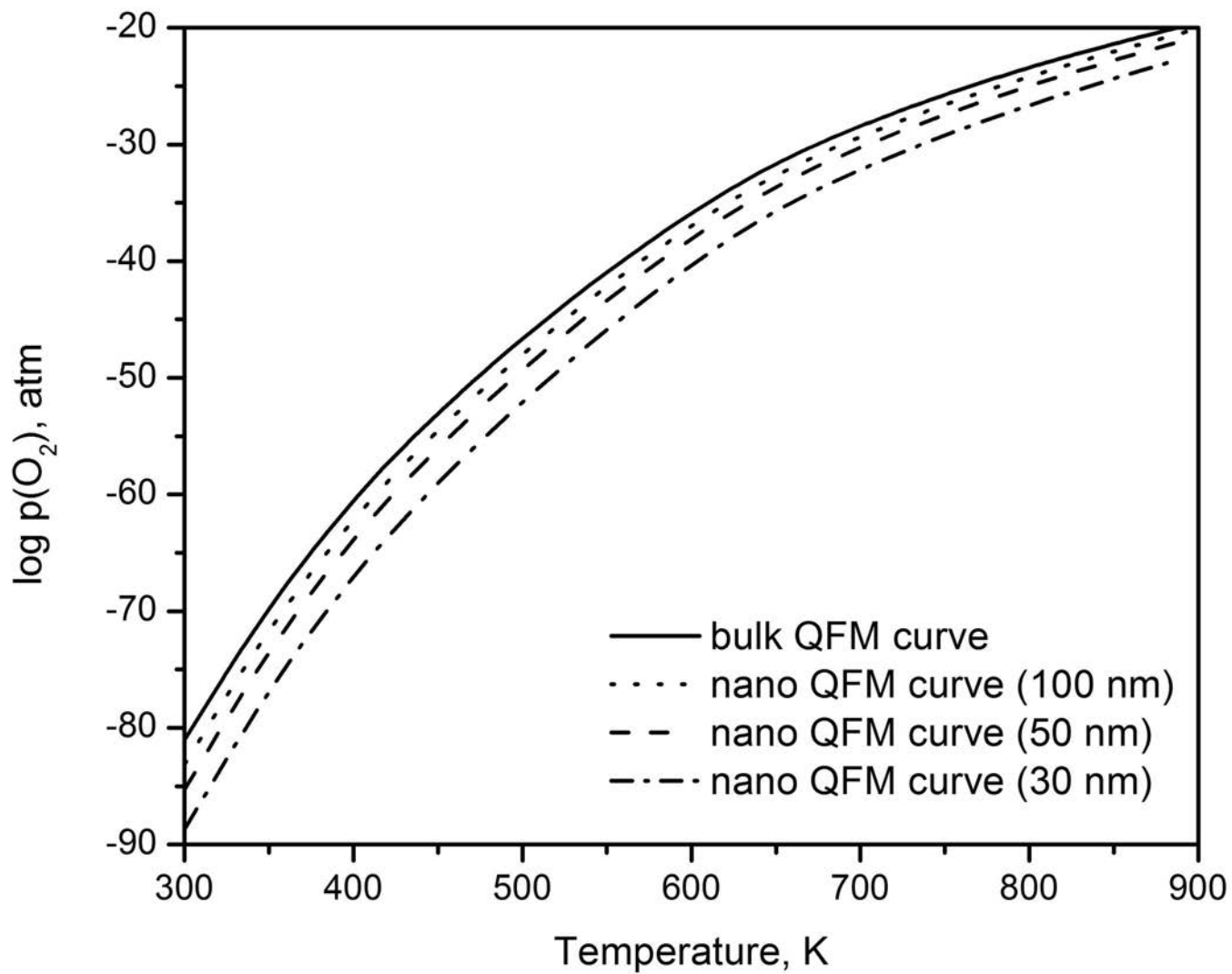


Figure 3

