

1 *Revision 2*

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3 ***P-V-T equation of state and high-pressure behavior of CaCO₃ aragonite***

4
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12
13 **Abstract**

14 The equation of state of aragonite was determined by *in situ* synchrotron X-ray diffraction
15 experiments on a hot-pressed polycrystalline specimen of aragonite within its stability field up to
16 5.8 GPa and 1173 K. As a complement to this experimental study, first principles density
17 functional theory calculations were performed up to 20 GPa at 0 K. Thermoelastic parameters for
18 aragonite (CaCO₃) were derived by a least-squares fit of the experimental P-V-T data to the third-
19 order Birch-Murnaghan equation of state (EOS), yielding the bulk modulus and its pressure and
20 temperature derivatives $K_0 = 65.24 \pm 0.24$ GPa, $K'_0 = 4.95 \pm 0.12$, $(\partial K_T / \partial T)_P = -0.024 \pm 0.002$ GPa
21 K⁻¹ and volume thermal expansion $\alpha_{300} = (6.1 \pm 0.7) \times 10^{-5}$ K⁻¹. The analyses of the axial
22 compressibility at ambient temperature show that the *c*- axis is much more compressible than the
23 *a*- and *b*- axes. Based on first principles calculations, the anisotropic compression behavior of
24 aragonite structure is explained by the heterogeneous shortening of <Ca-O> and <C-O> bond
25 lengths and the rotation of <O-C-O> angles along the *a*-, *b*-, and *c*-axes, whereas the unit cell
26 volume change of aragonite under compression is accommodated by comparable compression rate
27 of the CaO₉ polyhedra and the voids in the crystal lattice. The results attained from this study
28 provide important thermoelastic parameters for understanding the thermodynamic behavior and
29 chemical reactions involving aragonite at subduction zone conditions.

30 **Keywords:** Aragonite, X-ray diffraction, equation of state, axial compressibility, high pressure
31 and high temperature, first principles density functional theory calculation

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33

Introduction

34 Dissolution of carbonates is an important mechanism for transferring carbon into the mantle
35 and ultimately recycling back into the atmosphere to balance the carbon flux (Kerrick and Connolly,
36 1998; Frezzotti et al., 2011). CaCO₃ is a dominant carbon-bearing mineral and exists widely in the
37 crust, mainly in pelagic marine sediments such as clay-rich marls and carbonate ooze (Plank and
38 Langmuir, 1998). When entrained by the subducting process, the polymorphic phases of CaCO₃
39 (e.g., calcite, aragonite and vaterite) can be transported to deep depths in the subduction zone
40 (Dasgupta and Hirschmann, 2010; Perchuk and Korepanova, 2011).

41 In recent years, the high-pressure properties and behaviors of CaCO₃ have attracted considerable
42 interest in order to facilitate our understanding of the global carbon cycle by identifying potential
43 carbon storage pathways in the Earth. Of particular interest is aragonite, a phase stable at the P-T
44 conditions corresponding to the lower crust and the uppermost mantle, which has been identified to
45 be a major carbonate phase that can store oxidized carbon in mantle-related assemblages of high-
46 pressure metamorphic rocks (Litasov and Ohtani, 2010; Grassi and Schmidt, 2011). The
47 thermodynamic calculation results of Kerrick and Connolly (2001) have shown that clay-rich marls
48 along low-temperature geotherms at depths of <180km have undergone virtually no devolatilization
49 in subduction zone; aragonite is also a main carbon storage phase in siliceous limestones and
50 carbonate ooze at these depth ranges.

51 Previously, the behavior of aragonite at high pressures and temperatures has been studied by
52 infrared spectroscopy (Kraft et al., 1991), X-ray diffraction (Martinez et al., 1996; Santillán and

53 Williams, 2004; Ono et al., 2005), Brillouin scattering (Liu et al., 2005) and Raman spectroscopy
54 (Kraft et al., 1991; Williams et al., 1992; Gillet et al., 1993). At ambient pressure, aragonite
55 spontaneously reverts to calcite at temperatures above 673 K (Carlson and Rosenfeld, 1981);
56 whereas at high pressures, metastable phases of CaCO₃, Calcite II and III, have been observed
57 within the stability field of aragonite (Merrill and Bassett, 1975). Calcite III has been considered as
58 a disordered variant of aragonite (Fong and Nicol, 1971). The phase transitions between calcite II,
59 calcite III and aragonite were also identified within the pressure of 0-6 GPa (Suito et al., 2001;
60 Catalli and Williams, 2005; Ono et al., 2005). No phase transitions were observed in aragonite up to
61 30-40 GPa according to Kraft et al. (1991) and Gillet et al. (1993). At pressures above 40 GPa, it
62 transforms to post-aragonite structure and remains stable at the lower mantle pressures (Ono et al.,
63 2005). To date, the bulk modulus and its pressure derivative are still scattered, ranging from 64 to
64 88 GPa and 2.3 to 5.3, respectively, High temperature P-V-T equation of state (EOS) of aragonite
65 has only been studied by Martinez et al. (1996) by in situ X-ray powder diffraction experiment at
66 pressures of 3-8 GPa, and the temperature derivative of the bulk modulus was not conclusively
67 constrained due to the ambiguities in the pressure derivative. In this study, using *in situ* synchrotron
68 X-ray diffraction techniques, the *P-V-T* behavior of aragonite was studied up to 6 GPa and 1173 K.
69 In addition, first principles calculations were performed to gain insight into the crystal chemistry
70 and compression mechanisms of aragonite under high pressure.

71

72 **Experimental procedures and specimen description**

73 A polycrystalline specimen of aragonite was synthesized in a 1000-ton uniaxial split-cylinder
74 apparatus (USCA-1000) at the High Pressure Laboratory of Stony Brook University using reagent-
75 grade calcite powder as starting material. The specimen was hot-pressed at pressure of 8 GPa and

76 temperature of $\sim 700\text{K}$ for 10 hours. Aragonite is quenchable and remains stable after recovery from
77 high pressure and temperature (see also Ono et al., 2005). As evidenced in Fig. 1, the synthesized
78 specimen is a single phase of aragonite. The sample appeared to be well-sintered with a bulk density
79 of $2.92(1)\text{ g/cm}^3$ as measured by Archimedes immersion method (Wilson, 2012), which is within
80 0.2-0.4% of the theoretical density from our X-ray diffraction measurement ($2.932(1)\text{ g/cm}^3$) and
81 those ($2.926\text{-}2.933\text{ g/cm}^3$) attained by Huggins (1922).

82 The high-pressure and high-temperature in situ experiments were conducted using a DIA-type
83 cubic-anvil apparatus at the X17B2 beamline of the National Synchrotron Light Source (NSLS) at
84 Brookhaven National Laboratory (BNL). An energy-dispersive X-ray method was employed to
85 collect the diffracted X-rays at a fixed Bragg angle of $2\theta = 6.44^\circ$. It is worth noting that using a hot-
86 pressed polycrystalline sample instead of a powder decreases the microscopic deviatoric stress
87 experienced by the individual grains so that the diffraction peaks are less likely to exhibit
88 broadening during compression at room temperature (e.g., Li et al., 1996). Details of the
89 experimental assembly setup used in this study can be found elsewhere (Weidner et al., 1992; Li et
90 al., 2004; Zhang et al., 2013). The pressure-transmitting medium was made of pre-compressed
91 boron epoxy (4:1 wt% ratio) and in cubic shape with an edge length about 6.4 mm. The aragonite
92 specimen was embedded in a NaCl-BN powder mixture (10:1 wt% ratio) inside a cylindrical
93 container of boron nitride. The NaCl powder was also used as an internal pressure standard and the
94 cell pressure was calculated from its equation of state (Decker, 1971). The pressure uncertainties are
95 about 0.2 GPa in the P - T range of this study, mainly resulting from the uncertainties of the lattice
96 parameters of NaCl measured at each experimental condition. The temperature was measured by a
97 W/Re25%-W/Re3% thermocouple with the thermal junction placed adjacent to the specimen (e.g.,

98 Li et al., 2004). The vertical and radial temperature gradients over the entire sample chamber were
99 estimated to be less than 20 K/mm (Weidner et al., 1992; Wang et al., 1998).

100 The sample was initially compressed to the highest pressure (7.7 GPa) at ambient temperature
101 and then the temperature was increased to a maximum of 1173 K. Subsequently, the experiment
102 was continued via a series of cooling/heating cycles at various ram loads (Fig. 2). To minimize the
103 effect of deviatoric stress generated during compression/decompression at room temperature, high
104 temperature data were only collected along cooling after reaching the peak temperatures. To avoid
105 the occurrence of phase transition and disordered structure (Suito et al., 2001; Ono et al., 2005), we
106 selected 1173, 973, 773, and 300 K as the temperature conditions for data collection in each cycle;
107 from 773 to 300K, the temperature was quickly quenched by switching off the power supply. The
108 lattice parameters were obtained by Le Bail full spectrum fit (Le Bail, 2005) using the software
109 package GSAS (Larson and Von Dreele, 2000) and the EXPGUI (Toby, 2001).

110 The first principles density-functional theory (DFT) calculations were performed using the
111 Vienna ab initio simulation package (VASP) code (Kresse and Furthmüller, 1996a; 1996b). A plane
112 wave energy cutoff of 700 eV was used for the calculation. The tetrahedron method combined with
113 Blöchl corrections (Blöchl, 1994) has been used for the Brillouin-zone integration over a k mesh of
114 $6 \times 4 \times 5$ for the orthorhombic aragonite. The combined GGA-PAW (generalized gradient
115 approximation and the projector augmented wave) functionals were used for accurate ab initio
116 calculations on energies and structures. The core region cutoff of the potentials is 3.0 a.u. for
117 calcium, 1.5 a.u. for carbon, and 1.52 a.u. for oxygen. The current parameterization assures a
118 convergence of total energy to better than 1meV/atom, forces to better than 0.01eV/Angstrom, and
119 pressure to less than 0.2 GPa.

120

121

Results and discussion

122 Room-temperature equation of state

123 Lattice parameters and unit-cell volumes of CaCO₃ aragonite attained along various isotherms
124 from 300 to 1173 K at pressures up to 7.7 GPa are listed in Table 1. X-ray diffraction spectra
125 obtained at ambient (before and after the experiment) and the peak P-T conditions are compared in
126 Fig. 1. The diffraction peaks were indexed by the orthorhombic unit cell of aragonite with *Pmcn*
127 space group. Within the resolution of the present X-ray diffraction data, no structural changes were
128 observed up to the peak P-T conditions of the current experiment. The unit-cell volume at ambient
129 condition (V_0) was measured to be $226.71 \pm 0.01 \text{ \AA}^3$, which is in good agreement with the values of
130 $226.7 \pm 0.3 \text{ \AA}^3$ and $226.96 \pm 0.01 \text{ \AA}^3$ reported by Santillán and Williams (2004) and Parker et al.
131 (2010), respectively. The current result is $\sim 0.2 \%$ smaller than the values of 227.14 \AA^3 and $227.2 \pm$
132 0.1 \AA^3 reported by Martinez et al. (1996) and Ono et al. (2005), respectively, which could be due to
133 variations of the chemical composition in the samples used in different studies.

134 The P - V data attained at 300 K were fitted by using the third-order Birch-Murnaghan EOS, as
135 shown below,

$$136 \quad R(V, T) = \frac{3}{2} K_0 \left[\left(\frac{V_0}{V} \right)^{\frac{7}{3}} - \left(\frac{V_0}{V} \right)^{\frac{5}{3}} \right] \times \left\{ 1 + \frac{3}{4} (K_0' - 4) \left[\left(\frac{V_0}{V} \right)^{\frac{2}{3}} - 1 \right] \right\} \quad (1)$$

137 where K_0 , K_0' , and V_0 are the isothermal bulk modulus, the pressure derivative of the isothermal
138 bulk modulus, and the unit-cell volume at ambient conditions, respectively. A least squares fit of the
139 present room T data to Eq. 1 with all parameters free yields $V_0 = 226.72 \pm 0.01 \text{ \AA}^3$, $K_0 = 65.24 \pm$
140 0.24 GPa and $K_0' = 4.95 \pm 0.12$, which are in good agreement with those of Ono et al. (2005) ($V_0 =$
141 $227.2 \pm 1.0 \text{ \AA}^3$, $K_0 = 67.1 \pm 6.3 \text{ GPa}$ and $K_0' = 4.7 \pm 0.8$) (see also Fig. 3). The bulk modulus value

142 compares well with the result of 65.4 ± 0.5 GPa from high P - T X-ray study by Martinez et al.
143 (1996), 67.1 ± 6.3 GPa by Ono et al. (2005), and 66.66 GPa reported by Salje and Viswanathan
144 (1976). If we fix K'_0 to 4.0, fitting of Eq. 1 gives a result of $V_0 = 226.7 \pm 2.0 \text{ \AA}^3$, $K_0 = 66.88 \pm 0.05$
145 GPa, which is in agreement within mutual uncertainties with the K_0 value of 64.5 GPa and $64.81 \pm$
146 3.48 GPa reported by Martens et al. (1982) and Martinez et al. (1996), respectively. The K_0 value
147 from this study is also in good agreement with the adiabatic bulk modulus $K_{S0} = 68.9 \pm 1.4$ GPa
148 (corresponding $K_0 \sim 67.4$ GPa) derived by Brillouin spectroscopy (Liu et al., 2005), and in
149 reasonable agreement with 69.6 GPa and 71 GPa attained by previous theoretical calculations
150 (Pavese et al., 1992; Fidler et al., 2000). As shown in Fig. 3, the trend of the compression curve
151 from the current experimental study agrees well with that predicted by our first principles
152 calculations despite there being a $\sim 2.5\%$ difference in the absolute values of volume caused by the
153 well-known underbinding of GGA. By comparison, Santillán and Williams (2004) reported a value
154 of $K_0 = 88$ GPa (with $K'_0 = 4$), which is significantly higher than those from the other studies (see
155 Table 2). A closer examination of Fig. 3 indicates that the high bulk modulus of Santillán and
156 Williams (2004) most likely results from the scatter in the unit cell volumes, especially those at
157 pressure above 15 GPa, presumably caused by deviatoric stress after the solidification of the
158 methanol-ethanol-water pressure medium.

159 **High-temperature equation of state**

160 The P - V - T relations of aragonite were investigated at pressures and temperatures up to ~ 6 GPa
161 and 1173 K by fitting the attained experimental data from present study to the high-temperature
162 Birch-Murnaghan (HTBM) equation of state (Eq. 1). The pressure derivative K'_0 is assumed to be
163 constant in Eq. 1 throughout the whole temperature range. The temperature dependence of the bulk

164 modulus is expressed by a linear function of temperature (Eq. 2), assuming that the temperature
165 derivative is constant throughout the temperature range of present study.

166

$$167 \quad K_T = K_0 + (\partial K_T / \partial T)_P (T - 300) \quad (2)$$

$$168 \quad K'_T = K'_0 \quad (3)$$

169 The unit-cell volume V_T at zero-pressure and a given temperature T , is expressed as

$$170 \quad V_T = V_0 \exp \int_{300}^T \alpha_T dT \quad (4)$$

$$171 \quad \alpha_T = a_0 + a_1 T \quad (5)$$

172 where V_0 is the unit-cell volume at ambient conditions, α_T is the volume thermal expansion at
173 ambient pressure and a given temperature T , a_0 and a_1 are constant parameters. Table 2 lists the
174 values of V_0 , K_0 , K'_0 , $(\partial K_T / \partial T)_P$, a_0 and a_1 derived by fitting our data (Table 1) to Eqs. 1-5. We
175 observe that when none of the parameters is constrained, the uncertainties in the derived parameters
176 are large (Table 2), which might be due to the large number of free parameters relative to the total
177 experimental data points; this observation remains true even though the fitting results ($K_0 = 64.5 \pm$
178 4.2 GPa and $K'_0 = 5.3 \pm 2.3$) are in general agreement with those from room temperature data alone.

179 Alternatively, when K_0 and K'_0 were fixed to 65.24 GPa and 4.95, respectively, as those derived at
180 room-temperature, and the following results were obtained: $V_0 = 226.7 \pm 0.1 \text{ \AA}^3$, $(\partial K_T / \partial T)_P = -0.024$
181 ± 0.002 GPa K^{-1} and $\alpha_{300} = (6.1 \pm 0.7) \times 10^{-5} \text{ K}^{-1}$. Note that the results of $(\partial K_T / \partial T)_P$ and the thermal
182 expansion coefficients from the constrained fit are nearly identical to those from unconstrained fit,
183 but with much smaller uncertainties. It should be pointed out that the apparent difference between
184 $(\partial K_T / \partial T)_P = -0.024 \pm 0.002$ GPa K^{-1} attained in present study and that of -0.018 ± 0.002 GPa K^{-1}
185 reported by Martinez et al. (1996) is believed to be caused by the difference in K'_{T0} used in the

186 respective fits; for instance, if the values of K_0 and K_0' were fixed to those from Martinez et al.
187 (1996), the same value of $(\partial K_T/\partial T)_P = -0.018 \pm 0.002 \text{ GPa K}^{-1}$ was obtained (see Table 2). In
188 addition, the resultant thermal expansion coefficient α_{300} is in excellent agreement with that from
189 Martinez et al. (1996). Figure 4 shows the measured unit cell volumes and the isothermal
190 compression curves derived from this study superimposed with those of Martinez et al. (1996), the
191 good agreement between the two datasets can be clearly seen, except for the data at $\sim 7.5 \text{ GPa}$ and
192 1173 K .

193 **Crystal structure of aragonite under compression**

194 A linear fit to the a/a_0 , b/b_0 and c/c_0 results versus pressure yielded the linear
195 compressibility of aragonite at room temperature (Table 3) of $(2.6 \pm 0.2) \times 10^{-3} \text{ GPa}^{-1}$, $(4.8 \pm 0.4) \times 10^{-3}$
196 GPa^{-1} and $(5.9 \pm 0.3) \times 10^{-3} \text{ GPa}^{-1}$ for a -, b - and c -axis, respectively, indicating a strong anisotropy
197 in compression behavior, with the c -axis much more compressible than the a - and b -axis. These
198 results are in excellent agreement with those reported by Martinez et al. (1996) (2.4 ± 0.2 , 4.2 ± 0.2
199 and $5.8 \pm 0.2 \times 10^{-3} \text{ GPa}^{-1}$ for a -, b - and c -axis, respectively), and comparable to the results of Fisler
200 et al. (2000) and Liu et al. (2005) (see Table 3).

201 The crystal structure of aragonite under pressure was further investigated by first principles
202 calculations for the orthorhombic crystal structure (space group $Pmcn$). As shown in Fig. 5, the
203 aragonite structure is characterized by layers of nine-coordinated Ca^{2+} cations, and slightly aplanar
204 CO_3^{2-} groups stacked perpendicular to the c - axis (Bragg, 1924; Wyckoff, 1925). With increasing
205 pressure, the bond lengths of $\langle \text{Ca-O} \rangle$ and $\langle \text{C-O} \rangle$ decrease continuously (Table 4 and Fig. 6a); at 20
206 GPa, the average $\langle \text{Ca-O} \rangle$ bond length has decreased by 6.24%, in comparison, the $\langle \text{C-O} \rangle$ distance
207 has decreased by only $\sim 1\%$, consistent with previous observations that the CO_3^{2-} groups are rigid
208 units in the aragonite structure (Antao and Hassan et al., 2010; Ye et al., 2012). As a result, in the b -

209 *c* plane (Fig. 5), the $\langle\text{Ca-O}\rangle$ bonds experience significant compression with increasing pressure
210 while the $\langle\text{C-O1}\rangle$ bonds along *b*-axis exhibit nearly no compression with a shortening of only
211 0.85% (Fig. 6a). In the *a-b* plane, the $\langle\text{Ca-O2}\rangle$ bonds in the direction of *b*-axis show slightly more
212 shortening (6.42-6.68%) than that of the $\langle\text{Ca-O1}\rangle$ bonds (6.06%) along the *a*-axis. These
213 anisotropic bond length shortening could account partially for the anisotropic compression behavior
214 of *a*-, *b*- and *c*- axis in aragonite.

215 The unit cell of aragonite can be considered as consisting of the four CaO_9 polyhedra, four
216 aplanar CO_3^{2-} groups, and the voids among polyhedra. Along with the compression of the unit cell
217 of aragonite, the volume of the CaO_9 polyhedra decreases at a slightly lower rate than the voids (Fig.
218 6c); by comparison, the CO_3^{2-} polyhedral undergoes no volume decrease during the compression.
219 Thus, CO_3^{2-} polyhedra have little or no contribution to the volume changes. Moreover, the aplanar
220 characteristics of the CO_3^{2-} groups in the aragonite lattice remain nearly unchanged during
221 compression except for the slight changes in the $\langle\text{O-C-O}\rangle$ angles within the CO_3^{2-} group; this is in
222 good agreement with the small distortion of CO_3^{2-} groups as indicated by the internal vibrational
223 modes under compression (Kraft et al., 1991). As shown in Fig. 6b, from ambient pressure to 20
224 GPa, the $\langle\text{O1-C-O2}\rangle$ angle increases from 120.260° to 120.497° , while the $\langle\text{O2-C-O2}\rangle$ angles
225 decrease from 119.338° to 118.816° , resulting in an increase in the difference between $\langle\text{O1-C-O2}\rangle$
226 and $\langle\text{O2-C-O2}\rangle$ from 0.9° to 1.7° . This trend is consistent with the increase from 0.7° to $4.5(\pm 1.2)^\circ$
227 within 0 - 40 GPa inferred from vibrational spectroscopy data by Kraft et al. (1991). Previously,
228 Holl et al. (2000) carried out a single-crystal X-ray diffraction study on witherite (BaCO_3), an
229 isostructural phase of aragonite, up to 8 GPa at ambient temperature. Compared to aragonite,
230 witherite has a slightly different anisotropic compression behavior with $\beta_c > \beta_a = \beta_b$ in the axial
231 compressibility; the pressure dependence of the bond angles $\langle\text{O1-C-O2}\rangle$ and $\langle\text{O2-C-O2}\rangle$, however,

232 exhibit the same trend as that of CaCO₃ aragonite during compression, suggesting that the same
233 compression mechanisms might be operative in the isostructural carbonate mineral groups.

234 **Implications**

235 In this study, high pressure and temperature data for the equation of state and compressibility of
236 aragonite were obtained by in situ X-ray diffraction at pressures up to 5.8 GPa and temperatures up
237 to 1173 K. Thermoelastic parameters of aragonite (CaCO₃) were derived by a least-squares fit of the
238 experimental P-V-T data to the third-order Birch-Murnaghan equation of state, yielding $K_{T0} = 65.24$
239 ± 0.24 GPa, $K'_{T0} = 4.95 \pm 0.12$, $(\partial K_T / \partial T)_P = -0.024 \pm 0.002$ GPa K⁻¹ and $\alpha_{300} = (6.1 \pm 0.7) \times 10^{-5}$ K⁻¹.

240 The *c*-axis of aragonite is much more compressible than *a*-axis and *b*-axis, which is related to the
241 anisotropic shortening of <Ca-O> and <C-O> bonds along the *a*-, *b*-, and *c*- axes in the unit cell.

242 The compression of the unit cell of aragonite is mostly the result of the compression of the CaO₉
243 polyhedra and the voids in the structure, while the CO₃²⁻ groups in the aragonite structure remains
244 applanar with the <O1-C-O2> angle increasing and the <O2-C-O2> angle decreasing respectively

245 with elevated pressure. The results attained from this study are useful to attain and evaluate the
246 thermoelastic and thermodynamic parameters of aragonite such as the temperature derivative of
247 bulk modulus, density, Grüneisen parameter and its volume dependence at high pressure and high

248 temperature. By comparing the density data of aragonite with other volatile-containing phases in the
249 subduction zone, such as superhydrous phase B or phase D (Litasov et al., 2007; 2008), the effect of

250 the presence of aragonite on the buoyancy of the subducting slabs in the deep mantle can be well
251 estimated. For instance, according to the current P-V relation, the density of aragonite at 12 GPa

252 and 1173 K is about 0.2 g/cm³ lower than that of superhydrous phase B (Litasov et al., 2007),
253 leading the existence of aragonite more effective for buoyance of the subduction slabs in the deep

254 mantle, compared to H₂O containing phases. In addition, using the method of minimizing the Gibbs

255 free energy, the EOS parameters attained presently are useful for the thermodynamic calculations
256 on the stability field and decarbonation reactions associated with CO₂ degassing and CaCO₃
257 dissolution at great depths in the mantle: CaCO₃ = CaO (lime) + CO₂ (fluid), CaCO₃ + SiO₂ =
258 CaSiO₃ (wollastonite) + CO₂ (fluid) and 4 CaCO₃ + 4 SiO₂ + (Mg, Fe)₃Al₂Si₃O₁₂ (garnet) = 3
259 Ca(Mg, Fe)Si₂O₆ (clinopyroxene) + CaAl₂SiO₆ (Tschermack's molecule)+ 4 CO₂ (fluid)
260 (Hammouda, 2003). The result will be helpful for the estimation of carbonates amount and the CO₂
261 release in the subduction zone conditions.

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401 **Figure captions**

402 **Figure 1.** Selected X-Ray diffraction patterns of aragonite.

403 **Figure 2.** Experimental P-T conditions. Phase diagram of aragonite after Suito et al. (2001). Grey
404 long dashed lines represent the P-T path to approach the peak conditions. Bold black lines represent
405 the cycles under different loading pressures. Black short dashed lines represent the phase
406 boundaries between calcite I, calcite II and III. Solid lines represent the phase boundaries between
407 calcite III, aragonite and disordered aragonite. Br = After Bridgman (1938), H.G.&E = After Hess
408 et al. (1991), S.N.&H = After Suito et al. (2001).

409 **Figure 3.** P-V relations of aragonite at 300 K obtained in this study, in comparison with previous
410 studies. Solid black curve represents the result of fitting the data of this study yielding $V_0 = 226.72 \pm$
411 0.01 \AA^3 , $K_0 = 65.24 \pm 0.24 \text{ GPa}$ and $K'_0 = 4.95 \pm 0.12$. Cold compression data were included for
412 comparison but not used in the fit.

413 **Figure 4.** P-V-T relations of aragonite obtained in this study. Isothermal compression curves at
414 various temperatures are calculated by using the thermoelastic parameters attained in the present
415 study ($V_0 = 226.7 \pm 0.1 \text{ \AA}^3$, $K_0 = 65.24 \text{ GPa}$, $K'_0 = 4.95$, $\partial(K_T/\partial T)_P = -0.024 \pm 0.002 \text{ GPa K}^{-1}$ and
416 $\alpha_{300} = (6.1 \pm 0.7) \times 10^{-5} \text{ K}^{-1}$). The errors of unit-cell volumes are less than the size of symbols.

417 **Figure 5.** Crystal structure of aragonite CaCO_3 . The numbers label the <Ca-O> bonds listed out in
418 Table 4.

419 **Figure 6.** Variations of the bond lengths (a), angles (b), and volumes of CaO9 polyhedra and voids
420 (c) of aragonite with pressure.

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435 **Table 1** Experimental P - T conditions, lattice parameters and unit-cell volumes of aragonite obtained by *in-*
 436 *situ* X-ray diffraction. Numbers in the parenthesis represent 1σ errors.

| T (K) | P (GPa) | a (Å) | b (Å) | c (Å) | V (Å ³) |
|----------------------------|-----------|-----------|------------|-----------|-----------------------|
| Compression before heating | | | | | |
| 300 | 0.00 | 4.9592(2) | 7.9653(2) | 5.7392(2) | 226.71(1) |
| 300 | 1.16 | 4.9485(3) | 7.9293(4) | 5.6991(3) | 223.63(2) |
| 300 | 3.35 | 4.9153(4) | 7.8563(6) | 5.6314(4) | 217.46(2) |
| 300 | 5.13 | 4.8928(5) | 7.8040(7) | 5.5749(4) | 212.87(2) |
| 300 | 6.56 | 4.8797(5) | 7.7682(9) | 5.5442(5) | 210.16(2) |
| 300 | 7.70 | 4.8643(6) | 7.7418(10) | 5.5152(5) | 207.70(3) |
| Compression after heating | | | | | |
| 1173 | 4.55 | 4.9380(5) | 7.9060(6) | 5.7638(5) | 225.09(2) |
| 1173 | 5.37 | 4.9308(4) | 7.8818(6) | 5.7326(5) | 222.79(2) |
| 1173 | 5.78 | 4.9211(5) | 7.8426(8) | 5.7335(3) | 221.28(2) |
| 973 | 2.55 | 4.9552(5) | 7.9772(6) | 5.7890(4) | 228.83(2) |
| 973 | 3.38 | 4.9546(3) | 7.9219(8) | 5.7552(6) | 225.89(2) |
| 973 | 3.76 | 4.9422(4) | 7.9152(6) | 5.7358(6) | 224.38(2) |
| 973 | 4.78 | 4.9394(4) | 7.8722(6) | 5.6888(3) | 221.21(2) |
| 973 | 5.29 | 4.9286(2) | 7.8506(3) | 5.6777(2) | 219.69(1) |
| 773 | 2.00 | 4.9532(4) | 7.9547(5) | 5.7634(3) | 227.08(2) |
| 773 | 2.82 | 4.9445(3) | 7.9222(7) | 5.7332(5) | 224.58(2) |
| 773 | 3.26 | 4.9415(3) | 7.8976(6) | 5.7090(5) | 222.80(2) |
| 773 | 4.23 | 4.9272(3) | 7.8644(3) | 5.6689(3) | 219.67(1) |
| 773 | 4.82 | 4.9170(3) | 7.8467(4) | 5.6564(3) | 218.24(1) |
| 300 | 0.00 | 4.9468(3) | 7.9537(4) | 5.7393(3) | 225.82(2) |
| 300 | 0.90 | 4.9454(4) | 7.9107(7) | 5.7172(4) | 223.67(2) |
| 300 | 1.65 | 4.9305(4) | 7.8930(5) | 5.6949(4) | 221.62(2) |
| 300 | 2.18 | 4.9249(7) | 7.8601(8) | 5.6694(6) | 219.47(2) |
| 300 | 3.13 | 4.9194(4) | 7.8393(4) | 5.6369(5) | 217.39(2) |
| 300 | 3.63 | 4.9098(4) | 7.8204(5) | 5.6189(4) | 215.75(2) |

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440 **Table 2** Thermoelastic parameters for aragonite derived from the Birch–Murnaghan (BM) and High

441 Temperature Birch–Murnaghan (HTBM) EOS. Numbers in bracket are errors of the least-squares fitting.

| V_0 (Å ³) | K_0 (GPa) | K'_0 | $(\partial K_T/\partial T)_P$ (GPa K ⁻¹) | α_{300} (10 ⁻⁵ K ⁻¹) | a_0 (10 ⁻⁵ K ⁻¹) | a_1 (10 ⁻⁸ K ⁻²) | References |
|----------------------------|----------------|----------|---------------------------------------------------------|-------------------------------------------------------|----------------------------------------------|----------------------------------------------|------------------------------|
| 226.8 (2) | 64.5 (42) | 5.3 (23) | -0.025 (7) | 6.1(14) | 4.2 (7) | 6.6 (22) | This study |
| 226.7 (1) | 65.24* | 4.95* | -0.024 (2) | 6.1(7) | 4.2 (4) | 6.4 (10) | |
| 226.9 (1) | 64.8* | 4* | -0.018 (2) | 6.2(7) | 4.7 (4) | 4.9 (10) | |
| | 64.8 (43) | 4* | -0.018 (2) | 6.5(1) | 6.5 (5) | 0.3 (3) | Martinez et al., 1996 |
| | 65.4(5) | 2.7 (7) | -0.013 (2) | 6.7(5) | | | Salje and Viswanathan, 1976 |
| | 66.66 | | | 6.1 | | | Ono et al., 2005 |
| | 73.1(22) | 4* | | | | | |
| | 67.1(63) | 4.7 (8) | | | | | |
| | 88 | 4 | | | | | Santillán and Williams, 2004 |

442 *Values were fixed during data fitting

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445 **Table 3** The linear compressibility of *a*-, *b*- and *c*-axis for aragonite

| Linear compressibility | This study | Martinez et al., 1996 | Fisler et al., 2000 | Liu et al., 2005 |
|------------------------|-----------------------------------|-----------------------------------|-----------------------------------|-----------------------------------|
| | $\times 10^{-3} \text{ GPa}^{-1}$ | $\times 10^{-3} \text{ GPa}^{-1}$ | $\times 10^{-3} \text{ GPa}^{-1}$ | $\times 10^{-3} \text{ GPa}^{-1}$ |
| β_a | 2.6 ± 0.2 | 2.4 ± 0.2 | 2.1 | 3.0 ± 0.2 |
| β_b | 4.8 ± 0.4 | 4.2 ± 0.2 | 4.2 | 4.6 ± 0.2 |
| β_c | 5.9 ± 0.3 | 5.8 ± 0.2 | 5.8 | 7.3 ± 0.6 |

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452 **Table 4** Variation of bond lengths (Å), angles (°), polyhedra and voids (Å³) of aragonite at

453 high pressures calculated by DFT (Density Functional Theory) method

| Bond & Angle | | (A-B)/A $\times 100\%$ | A-0 GPa | 3 GPa | 5 GPa | 8 GPa | 12 GPa | 16 GPa | B-20 GPa |
|----------------------|------------|------------------------|---------|---------|---------|---------|---------|---------|----------|
| Ca-O1 ⁽⁵⁾ | $\times 1$ | 5.07 | 2.425 | 2.398 | 2.387 | 2.363 | 2.340 | 2.319 | 2.302 |
| Ca-O1 ⁽¹⁾ | $\times 2$ | 6.06 | 2.691 | 2.656 | 2.635 | 2.608 | 2.581 | 2.553 | 2.528 |
| Ca-O2 ⁽²⁾ | $\times 2$ | 6.42 | 2.571 | 2.538 | 2.516 | 2.488 | 2.456 | 2.429 | 2.406 |
| Ca-O2 ⁽⁴⁾ | $\times 2$ | 6.38 | 2.475 | 2.437 | 2.417 | 2.391 | 2.362 | 2.338 | 2.317 |
| Ca-O2 ⁽³⁾ | $\times 2$ | 6.68 | 2.546 | 2.507 | 2.482 | 2.456 | 2.423 | 2.398 | 2.376 |
| <Ca-O>[9] | | 6.24 | 2.555 | 2.519 | 2.499 | 2.472 | 2.443 | 2.417 | 2.395 |
| C-O1 | $\times 1$ | 0.85 | 1.292 | 1.290 | 1.289 | 1.287 | 1.285 | 1.282 | 1.281 |
| C-O2 | $\times 2$ | 1.08 | 1.302 | 1.300 | 1.298 | 1.296 | 1.293 | 1.290 | 1.288 |
| <C-O>[3] | | 1.00 | 1.299 | 1.297 | 1.295 | 1.293 | 1.290 | 1.287 | 1.286 |
| O1-C-O2 | $\times 2$ | -0.20 | 120.260 | 120.296 | 120.349 | 120.379 | 120.458 | 120.490 | 120.497 |
| O2-C-O2 | $\times 1$ | 0.44 | 119.338 | 119.256 | 119.147 | 119.078 | 118.910 | 118.836 | 118.816 |
| CaO9 dodecahedron | $\times 4$ | 16.82 | 30.6977 | 29.5057 | 28.8481 | 27.9724 | 27.0207 | 26.2203 | 25.5338 |
| CO3 tetrahedron | $\times 4$ | | 0.0207 | 0.0213 | 0.0214 | 0.0219 | 0.0225 | 0.0229 | 0.0232 |
| Voids | | 18.98 | 111.512 | 106.818 | 103.913 | 100.499 | 96.502 | 93.222 | 90.342 |

454 ⁽¹⁾⁻⁽⁵⁾ Number labels the different Ca-O bond by their lengths. ^{(1),(5)} represent the longest and shortest

455 Ca-O bond, respectively. [] Number of bond lengths/angles used to compute the average.

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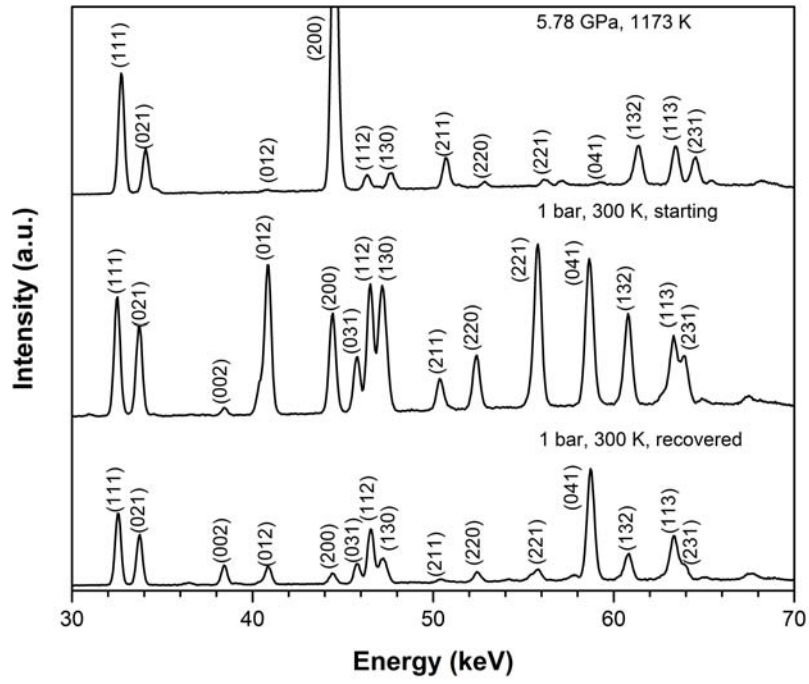


Fig. 1.

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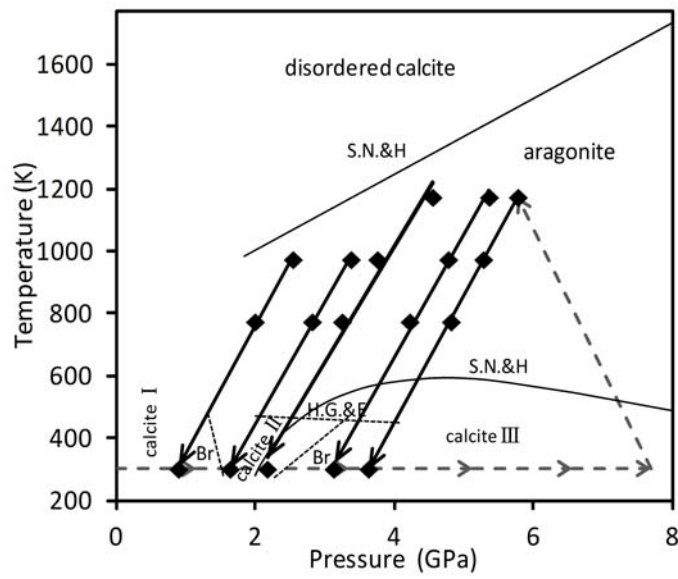
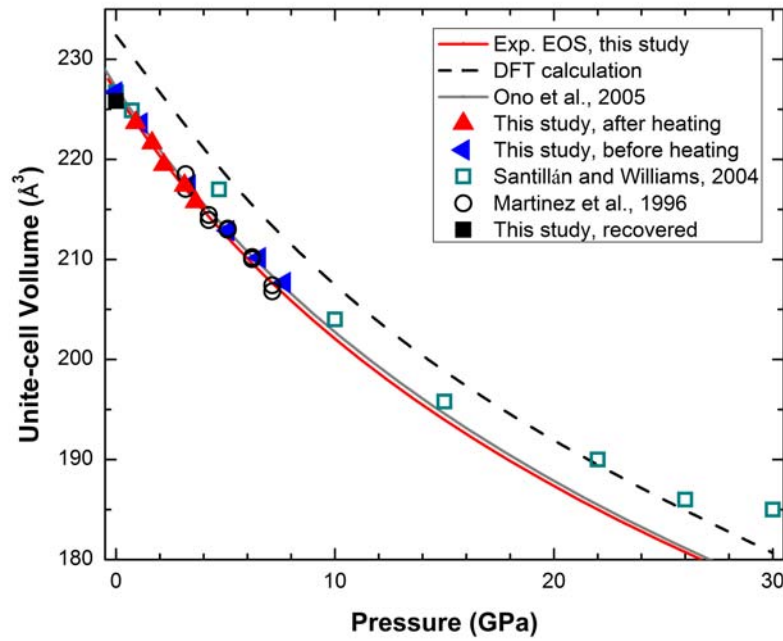


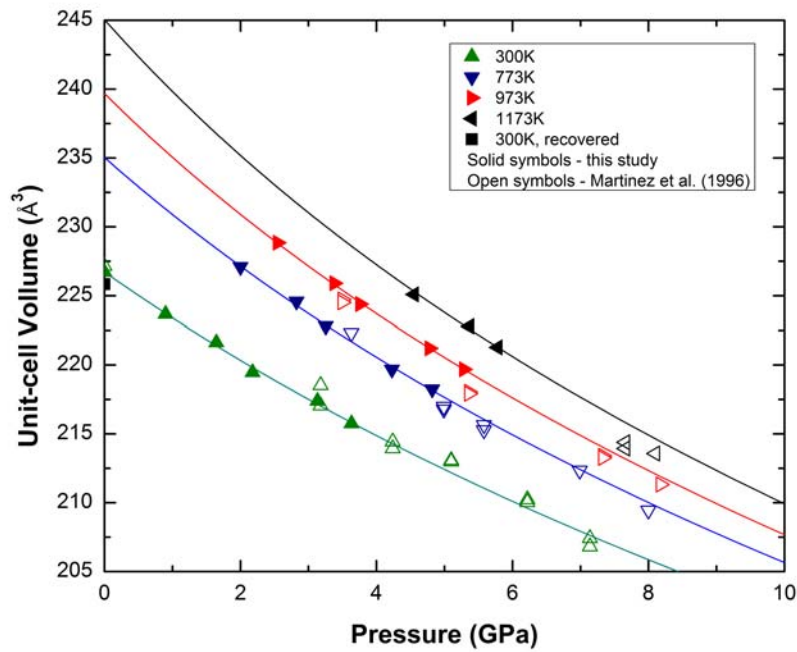
Fig. 2.

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Fig. 3.



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Fig. 4.

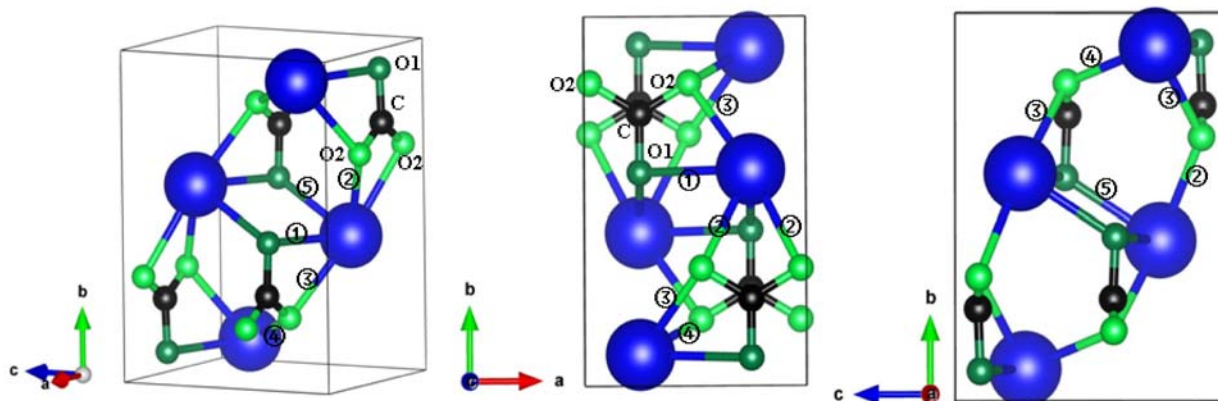


Fig. 5.

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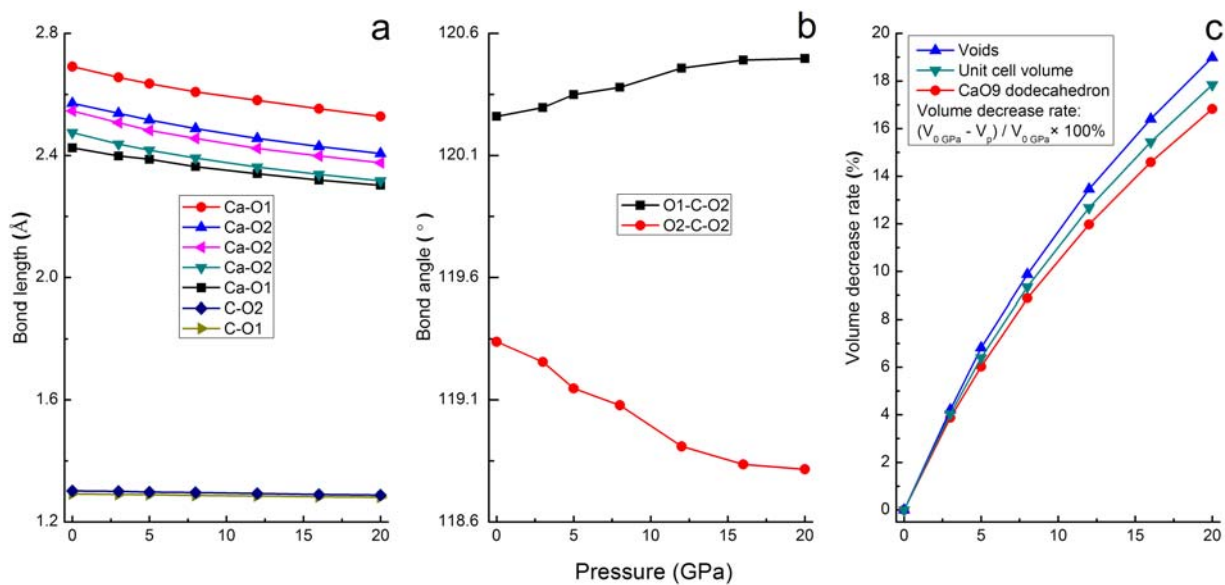


Fig. 6.

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