

1 Revision 1

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3 **Synthesis of Large and Homogeneous Single Crystals of Water-Bearing**  
4 **Minerals by Slow Cooling at Deep-Mantle Pressures**

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28

29 **Abstract**

30

31 The presence of water in the Earth's deep mantle is an issue of increasing interest in  
32 the field of high-pressure mineralogy. An important task for further advancing research  
33 in the field is to create homogeneous single crystals of candidate deep-mantle

34 water-bearing minerals of 1 mm or larger in size, which is required for applying them  
35 for the Time-of-Flight (TOF) single-crystal Laue diffraction method with a  
36 third-generation neutron instrument. In this study, we perform several experiments to  
37 demonstrate an improved methodology for growing hydrous crystals of such large sizes  
38 at relevant transition zone and lower-mantle conditions via very slow cooling over a  
39 maximum period of 1 day. Successfully synthesized crystals using this methodology  
40 include dense hydrous magnesium silicate (DHMS) phase E, hydrous wadsleyite,  
41 hydrous ringwoodite, and bridgmanite (silicate perovskite). It is also demonstrated that  
42 these hydrous crystals can be grown from deuterium enriched starting materials in  
43 addition to from those having natural hydrogen isotope ratio.

44 Magnitudes of chemical and crystallographic heterogeneities of the product crystals  
45 were characterized by comprehensive analysis of X-ray precession photography,  
46 single-crystal X-ray diffraction (SCXRD), field-emission scanning electron microscope  
47 (FE-SEM), electron probe microanalyser (EPMA), secondary ion mass spectroscopy  
48 (SIMS), powder X-ray diffraction (PXRD), and TOF neutron powder diffraction  
49 (TOF-NPD). The product crystals were confirmed to be inclusion free and  
50 crystallographically homogeneous. Compositional and isotopic homogeneities of major  
51 elements and hydrogen isotope abundances were better than 1 % and 3 %, respectively,

52 among intracrystals and intercrystals within each recovered sample capsule. Phase E  
53 crystals up to 600  $\mu\text{m}$  in the largest dimension were grown at a constant temperature of  
54 1100  $^{\circ}\text{C}$  kept for 3 h. Using a lattice parameter-to-temperature relation of phase E,  
55 thermal gradient in the sample capsules for its synthesis has been evaluated to be  
56 20  $^{\circ}\text{C}/\text{mm}$ . Hydrous wadsleyite crystals with up to 1100  $\mu\text{m}$  in the largest dimension  
57 were grown at 1390  $^{\circ}\text{C}$  with a temperature reduction of 70  $^{\circ}\text{C}$  during heating for 10 h.  
58 Hydrous ringwoodite crystals with up to 1000  $\mu\text{m}$  in the largest dimension were grown  
59 at around 1400  $^{\circ}\text{C}$  with a temperature reduction of 110  $^{\circ}\text{C}$  during heating for 12 h.  
60 Bridgmanite crystals with up to 600  $\mu\text{m}$  in the largest dimension were grown at 1700  $^{\circ}\text{C}$   
61 with a temperature reduction of 30  $^{\circ}\text{C}$  during heating for 12 h. A TOF single-crystal  
62 diffraction instrument has been successfully used for analyzing one of the hydrous  
63 wadsleyite crystals, which demonstrated that appropriate single crystals for their  
64 expected usage are created using the method proposed in the present study.

65

66 **Keywords:** Hydrous minerals, Single crystal growth, Phase E, Wadsleyite,  
67 Ringwoodite, Bridgmanite, Slow-cooling method

68

69

## Introduction

70

71 Water, an essential component of the Earth, exists in surface oceans and in  
72 deep-mantle hydrous minerals (Inoue et al. 1995; Kanzaki 1991; Kohlstedt et al. 1996;  
73 Litasov et al. 2003; Pearson et al. 2014). Recent research has suggested that part of the  
74 oceanic water is being transported into the deep mantle through the subducting  
75 lithosphere to affect the Earth's global dynamic and chemical cycles (Bercovici and  
76 Karato 2003; Ohtani et al. 2004; Schmandt et al. 2014). In the Earth's early evolution,  
77 primordial water was partitioned among deep-mantle minerals, silicate melt, and  
78 metallic Fe alloy (Bolfan-Casanova et al. 2000; Inoue et al. 2010; Okuchi 1997;  
79 Shibasaki et al. 2009). Both transported water and that partitioned into the deep-mantle  
80 minerals affect the relevant mantle redox state (McCammon et al. 2004; Mrosko et al.  
81 2013; Smyth et al. 2014). The existence of such water has been partly confirmed by a  
82 recent discovery of natural hydrous ringwoodite with 1.4 wt% water occurring as an  
83 inclusion in a diamond specimen (Pearson et al. 2014). This discovery indicates that at  
84 least some areas of the transition zone are enriched with a significant amount of water.

85 Hydrous deep-mantle minerals contain water as hydroxyl groups (Keppler and  
86 Smyth 2006). Numerous experimental results indicate that the physical properties of the  
87 host minerals change dramatically with an increase in hydroxyl concentration. These

88 properties include melting behavior (Inoue 1994; Schmandt et al. 2014), phase relations  
89 (Ghosh et al. 2013; Ohtani et al. 2000), transformation kinetics (Kubo et al. 1998),  
90 rheology (Chen et al. 1998), and sound velocities (Inoue et al. 1998; Mao et al. 2011,  
91 2012). An important task for further advancing the research on the effects of these  
92 hydroxyls is to create well-characterized, homogeneous, hydrous crystals of at least 1  
93 mm in size. This is minimum size requirement to use the Time-of-Flight (TOF)  
94 single-crystal Laue diffraction method with a third-generation neutron instrument such  
95 as TOPAZ at the Spallation Neutron Source (SNS) (Schultz et al. 2014). Using this  
96 state-of-the-art instrument will enable for the first time accurate determination of the  
97 position and occupancy of normal hydrogen atoms in the crystal structures of  
98 deep-mantle minerals, and therefore direct observation of chemical bonding distance  
99 and charge substitution mechanism of the hydrogen. Moreover, accurate atomic  
100 displacement parameters of the hydrogen can be measured with this instrument for  
101 crystals 1 mm or larger (Jorgensen et al., 2014). These parameters provide invaluable  
102 information about chemical bonding strength and directionality of the hydrogen, which  
103 is essential to understand and forecast how hydroxyls control the physical properties of  
104 minerals.

105 In the present study, we therefore synthesized such crystals of hydrous wadsleyite  
106 and hydrous ringwoodite, which are highly expected to occur or have been confirmed  
107 within the mantle transition zone. In addition, we synthesized dense hydrous  
108 magnesium silicate (DHMS) phase E with a deuterium-enriched composition in crystals  
109 up to 600  $\mu\text{m}$  in size, which is a representative hydrous mineral that may occur in  
110 hydrated deep-mantle peridotite (Kanzaki et al. 1991; Ohtani et al. 2004). Although the  
111 crystal size is smaller than 1 mm, it may be applicable for TOF single-crystal diffraction  
112 because deuterium scatters neutrons more coherently and induces much less background  
113 than normal hydrogen, in which the required crystal volume is smaller (Wenk 2006).  
114 Moreover, the deuterium-enriched crystals will be useful for investigating the hydrogen  
115 isotope fractionation phenomena occurring within the deep mantle (Chakoumakos et al.  
116 2013; Horita et al. 2010; Okuchi et al. 2014; Yang et al. 2014). In addition to the DHMS  
117 Phase E, therefore, ringwoodite crystals were also synthesized in a deuterium-enriched  
118 growth environment. Finally, lower-mantle bridgmanite (silicate perovskite) crystals  
119 100  $\mu\text{m}$  to 600  $\mu\text{m}$  in size were synthesized in normal water ( $\text{H}_2\text{O}$ )-bearing growth  
120 environments. Although we were unable to grow these bridgmanite crystals into  
121 suitable sizes for TOF single-crystal diffraction, some proved to be ideal for use in

122 advanced high-pressure–high-temperature X-ray and laser experiments (e.g., Goncharov  
123 et al. submitted).

124 It is reasonable to expect that crystals grown into larger sizes may become less  
125 homogeneous. We therefore consider that it is important to evaluate the chemical and  
126 crystallographic heterogeneity of the synthesized crystals. For this purpose, the crystals  
127 were comprehensively analyzed using X-ray precession photography, single-crystal  
128 X-ray diffraction (SCXRD), field-emission scanning electron microscope (FE–SEM),  
129 electron probe microanalyzer (EPMA), secondary ion mass spectroscopy (SIMS),  
130 powder X-ray diffraction (PXRD), and TOF neutron powder diffraction (TOF–NPD). In  
131 particular, we made a special effort to evaluate the hydrogen and deuterium  
132 concentration distributions within each crystal and among the crystals within each  
133 capsule of the selected run charges.

134

### 135 **Experimental strategy**

136

137 Previous studies of high-pressure syntheses of deep-mantle crystals have partly  
138 addressed the aforementioned research task (Ito and Weidner 1986; Jacobsen et al.  
139 2004; Shatskiy et al. 2007, 2009; Smyth et al. 2003). Of particular interest are crystals  
140 of ringwoodite up to 800  $\mu\text{m}$  in size (Jacobsen et al. 2004; Smyth et al. 2003), and those

141 of wadsleyite and bridgmanite up to 1 mm in size (Shatskiy et al. 2007, 2009). Most of  
142 these crystals were grown over a relatively short period within a few hours. For the  
143 latter two products, the chemical gradients were intentionally introduced in addition to  
144 an existing thermal gradient ( $\geq 20$  °C/mm) under the extreme pressure–temperature  
145 environments of the sample capsules in which the crystals were quickly grown over  
146 short periods between 20 min to 3 h, except for one run at 5 h. Although these gradients  
147 played very essential roles in the successful synthesis of large and isometric crystals in  
148 short periods, they may also induce crystallographic heterogeneities due to fast growth,  
149 including lattice defects, crystal inclusions, twinning, and misorientation (Rudolph  
150 2010), which were observed to exist within some of the crystals (Shatskiy et al. 2009).  
151 In order to apply the TOF single-crystal diffraction method, samples having these  
152 crystallographic heterogeneities are not very appropriate (Wenk, 2006).

153 In the present study, we used a significantly slower growth rate over a maximum  
154 period up to 24 h to produce the sample crystals. In addition, we grew the crystals from  
155 a homogeneous silicate melt batch with a volume as large as possible to enable  
156 continuous buffering of the crystal's chemical composition. The temperature of the cell  
157 slowly decreased during the long heating durations so that the crystals were almost kept  
158 in chemical equilibrium with the silicate melt throughout the growth process. This



159 slow-cooling method has been reported as an effective scheme for recovering large and  
160 homogeneous crystals of metal oxides from limited sample volumes confined at high  
161 pressures up to 11 GPa (Loeffert et al. 2002; Shatskiy et al. 2010; Tamura et al. 1974).  
162 In the present study, we extended this method to higher pressures and temperatures of  
163 up to 24 GPa and about 1800 °C, respectively, for the crystal growth of deep-mantle  
164 hydrous mineral phases.

165

### 166 **Experimental methods**

167

168 DHMS phase E, pure magnesian hydrous wadsleyite, and pure magnesian and  
169 iron-bearing hydrous ringwoodite were synthesized at pressures between 14 GPa and 21  
170 GPa and temperatures between 1100 °C and 1400 °C. Bridgmanite with three different  
171 compositions was synthesized at 24 GPa and 1650 °C to 1760 °C. Table 1 shows a  
172 summary of these experimental conditions yielding relatively large crystal products. We  
173 used the scaled-up Kawai-type cell installed at the Institute for Study of the Earth's  
174 Interior at Okayama University (Shatskiy et al. 2011). Two different types of cell  
175 assembly designs suitable for two different pressure regimes were prepared. Fujillo  
176 F08 tungsten carbide anvils with 8 mm or 6 mm in truncation edge length (TEL) were  
177 used to compress these cells, which commonly had MgO–Cr<sub>2</sub>O<sub>3</sub> pressure media of 14

178 mm in the octahedral edge length (OEL) (Figure 1). Anvils 46 mm in size were used for  
179 applied loads of more than 10 MN, and those 32 mm in size were used for loads of less  
180 than 10 MN. The efficiency of the pressure generated by using these applied loads was  
181 previously calibrated for the combination of 14 mm OEL and 8 mm or 6 mm TEL  
182 (Frost et al. 2004; Shatskiy et al. 2009, 2011). Figure 1(a) shows a cross section of the  
183 cell coupled with 8 mm TEL, which was used for the phase E, wadsleyite and  
184 ringwoodite synthesis. This cell design was established in our previous work for mass  
185 production of fully deuterated ringwoodite powder (Purevjav et al. 2014a). In the  
186 present study, the same gold capsule of 4 mm in diameter and 5 mm in length in this  
187 cell, with  $\sim 50 \text{ mm}^3$  in sample volume, was successfully used for the significantly longer  
188 heating durations required for single-crystal growth. Figure 1(b) shows the cross section  
189 of the other cell coupled with 6 mm TEL, which was used for the bridgmanite synthesis.  
190 This capsule was designed and tested to be highly durable at temperatures of at least up  
191 to 1800 °C, at which the required long-duration runs were conducted in a stable manner.  
192 A platinum capsule of 2 mm in diameter and 4 mm in length, with  $\sim 12 \text{ mm}^3$  in sample  
193 volume, was successfully used at even higher temperatures than those used in previous  
194 bridgmanite growth experiments (Ito and Weidner 1986; Shatskiy et al. 2007). The melt

195 fraction and the chemical transport efficiency simultaneously increased at the higher  
196 temperatures at which we expected more homogeneous crystals to grow.

197 The starting materials were weighted mixtures of the components suitable for each  
198 synthesis (Table 1). Magnesium hydroxide [Mg(OH)<sub>2</sub>], aluminum oxide [Al<sub>2</sub>O<sub>3</sub>], iron  
199 [Fe], and its oxide [FeO] were prepared from commercial reagents. Part of the iron and  
200 its oxide were enriched in <sup>57</sup>Fe for the expected applications in Mössbauer spectroscopy  
201 and nuclear resonant inelastic X-ray scattering. Silica [SiO<sub>2</sub>] was prepared from a  
202 high-purity fused silica rod. Magnesium deuterioxide [Mg(OD)<sub>2</sub>] was prepared by using  
203 a hydrothermal reaction of MgO powder and D<sub>2</sub>O in an autoclave (Okuchi et al. 2014).  
204 For all runs except for EP12-6, the starting materials were the mixtures of these  
205 reagents, where H<sub>2</sub>O or D<sub>2</sub>O wt% was controlled by the weight ratio of Mg(OH)<sub>2</sub> or  
206 Mg(OD)<sub>2</sub> and MgO. Then each mixture was cold-sealed into the capsule to be  
207 compressed. For the EP12-6 run for iron-bearing hydrous ringwoodite we adopt a  
208 different procedure; powder of ground San Carlos olivine was first filled into the  
209 capsule, and liquid H<sub>2</sub>O was injected into it. In order to avoid the leakage of H<sub>2</sub>O, the  
210 capsule was spark-welded, weighed, kept in a vacuum for 1 h, and weighed again to  
211 confirm the weight was not changed by evaporation of H<sub>2</sub>O.

212 These capsules were compressed using their relevant cell assembly, anvils, and load  
213 (Table 1). After the load reached the desired pressure, the capsule was heated by  
214 supplying and controlling the applied electric power. After the desired heating duration  
215 passed, the capsule was quenched to room temperature by cutting off the power. We  
216 established the power–temperature relations for each cell assembly by conducting  
217 several calibration runs in which the temperature of the sample capsule was monitored  
218 using a  $W_{97}Re_3/W_{75}Re_{25}$  thermocouple inserted into the cell assembly and attached to  
219 the capsule. Reproducibility of the calibrated power–temperature relations was within  
220  $\pm 5\%$  and  $\pm 6\%$  for the 8 mm and 6 mm TEL cell assembly, respectively. Temperatures  
221 of the actual synthesis runs were controlled using the supplied power derived from the  
222 calibrated power–temperature relations. No correction for the effect of pressure on the  
223 thermocouple electromotive force (emf) was applied.

224 The crystalline portion of the recovered run product from each synthesis experiment  
225 was divided into several samples of crystal aggregate. These pieces were first screened  
226 using a microfocus X-ray diffractometer (Rigaku RINT RAPID II) to find the mineral  
227 phases and to determine whether their Bragg reflections were powder-like rings or  
228 spotty. An incident X-ray beam of  $Cu K\alpha_{1,2}$  radiation operated at 40 kV and 30 mA was  
229 collimated onto a beam size of 100  $\mu m$  for the analyses. The mineral grains were

230 oscillated through movements of the omega and phi axes, whereas X-ray diffraction  
231 patterns were taken in the reflection mode with an exposure time of 5–20 min. By using  
232 this screening process, we quickly evaluated whether the desired crystals with sizes  
233 larger than the collimator aperture were recovered. Then, crystallinities of selected  
234 samples were further evaluated individually by X-ray precession photography or  
235 SCXRD. A precession camera (Rigaku 1533A2) was used for several phase E,  
236 hydrogenated magnesian wadsleyite, and hydrogenated magnesian ringwoodite samples  
237 using Mo  $K\alpha_{1,2}$  radiation operating at 35 kV and 30 mA. Their reciprocal lattice images  
238 were taken on an imaging plate with an exposure time of 10 min. SCXRD  
239 measurements were conducted by using an automated four-circle X-ray diffractometer  
240 (Rigaku AFC-7s) with Mo  $K\alpha$  radiation ( $\lambda = 0.71073 \text{ \AA}$ ) at the power of 50 kV and 30  
241 mA. This instrument was used to measure crystals of deuterated magnesian ringwoodite  
242 and hydrogenated iron-bearing ringwoodite with dimensions of  $80 \times 50 \times 50 \text{ \mu m}^3$  and  
243  $80 \times 70 \times 50 \text{ \mu m}^3$ . A set of  $\{511\}$  X-ray reflections (24 in total) was used to refine the  
244 lattice constants of these ringwoodite crystals.

245 In addition, we evaluated the chemical heterogeneity of the recovered crystals. The  
246 crystals were polished with diamond powder 1  $\mu\text{m}$  in size, coated with carbon, and  
247 analyzed using FE–SEM (JEOL JSM-7001F) and EPMA (JEOL JXA-8800). The

248 electron beam was 5–10  $\mu\text{m}$  in diameter, 15 kV in acceleration voltage, and 8–12 nA in  
249 current, which were adjusted with consideration to the varying fragility of the sample  
250 mineral phases. Analyses of back-scattered electron (BSE) images of the crystals of all  
251 synthesized phases by FE–SEM revealed no evidence of chemical zoning. This result  
252 was confirmed by major element analyses of these crystals by EPMA, which are  
253 detailed subsequently in this paper.

254 Hydrogen and deuterium concentration distributions in the cross sections of phase E,  
255 deuterated magnesian ringwoodite, and hydrated iron-bearing ringwoodite crystals were  
256 measured by SIMS (CAMECA IMS-6F) at Hokkaido University. A primary  $^{133}\text{Cs}^+$   
257 beam operated at 9 nA and 10 kV was focused to a 10  $\mu\text{m}$  spot on the sample, which  
258 was coated by gold 50 nm in thickness. Normal incident electron shower was used for  
259 electrostatic charge neutralization of sputtering area. The secondary negative ions of  
260 hydrogen, deuterium, and  $^{30}\text{Si}$  were collected from the sputtered area sequentially in an  
261 electron multiplier for 10 seconds. Small width of the exit slit was chosen to eliminate  
262 the  $^{28}\text{SiD}$  interference on the  $^{30}\text{Si}$  peak. For quantitative analysis, a natural amphibole  
263 containing 1.66 wt% of  $\text{H}_2\text{O}$  was used as a standard material (Miyagi and Yurimoto  
264 1995). Both  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  contents in the samples were estimated assuming a linear  
265 relationship between their concentrations and the  $\text{H}/^{30}\text{Si}$  or  $\text{D}/^{30}\text{Si}$  ratio (Greenwood et

266 al. 2011). Other measurement conditions were described elsewhere (Greenwood et al.  
267 2011).

268 The space group and lattice parameters of the crystals, including those of all phase E,  
269 wadsleyite, and ringwoodite run products, were determined via analyses of the PXRD  
270 patterns using a few crystalline samples taken from each product (Table 2). All phase E,  
271 wadsleyite, and ringwoodite products were successfully indexed with  $R-3m$ ,  $Imma$  and  
272  $Fd-3m$  space groups, respectively, which is consistent with previous structural analysis  
273 results. Water concentrations in wadsleyite and ringwoodite were estimated from their  
274 refined lattice constants (Jacobsen et al. 2005; Ye et al. 2012). In addition, whole  
275 sample products from the two phase E run charges were analyzed by TOF–NPD. The  
276 observed reflections in TOF–NPD were successfully indexed with the same  $R-3m$  space  
277 group given by the PXRD results, and the linewidth profiles were analyzed in order to  
278 estimate the synthesis environments of these phase E products.

279 For PXRD analysis with the X-ray diffractometer (Rigaku Smartlab), each crystal  
280 aggregate sample was ground into a fine powder in an alumina mortar and was  
281 measured on an Si low-background plate using the para-focusing method with Cu  $K\alpha_1$   
282 radiation operated at 40 kV and 30 mA. The diffraction pattern was obtained by the step  
283 scanning method with a  $2\theta$  range of 10–140°, step size of 0.005°, and scan speed of

284 0.1–0.3°/min. The  $2\theta$  angles were calibrated by the silicon powder standard (NIST640d)  
285 as an external standard, and their lattice constants were refined through least squares  
286 fitting of peak positions by using the UnitCell software (Holland and Redfern 1997).  
287 For TOF–NPD analysis, each sample was coarsely ground and measured using the  
288 TAKUMI Engineering Materials Diffractometer at Materials and Life Science  
289 Experimental Facility (MLF), Japan Proton Accelerator Research Complex (J-PARC)  
290 (Harjo et al. 2006). For the measurements, the optics of the neutron diffractometer were  
291 optimized for analyzing the structure of a hydrous mineral phase with a small sample  
292 mass (Okuchi et al. 2014).

293

294

## Results and Discussion

295

296 Figure 2 shows photographs of the recovered single crystals grown by the  
297 slow-cooling method. Observations of the samples under an optical microscope  
298 confirmed that these crystals had no visible imperfections. In the following sections, the  
299 results of microscopic analysis are given and are discussed in a comprehensive manner,  
300 demonstrating that these crystals were crystallographically and chemically  
301 homogeneous to be suitable for their expected usage. In addition, we show the most



302 suitable experimental conditions for growing the largest crystals for each synthesized  
303 mineral phase.

304

### 305 **DHMS phase E crystals**

306 Single crystals of phase E of up to 600  $\mu\text{m}$  in the largest dimension were synthesized  
307 in the two relevant experimental runs (Table 1, Figure 2(a)). These crystals were  
308 significantly larger than the 100–200  $\mu\text{m}$  size previously reported (Kanzaki 1991;  
309 Kleppe et al. 2001; Kudoh et al. 1993). We found that the growth of phase E crystals  
310 was remarkably fast even at relatively low temperatures around 1100  $^{\circ}\text{C}$ ; therefore, 2–3  
311 h of continuous heating at a constant temperature with a variation of less than  $\pm 10$   $^{\circ}\text{C}$   
312 was sufficient for recovering these large crystals. This fast crystal growth may be due in  
313 part to the very large water concentration ( $\sim 22$  wt%) in the starting material, where  
314 below the silicate solidus, a large fraction of aqueous fluid coexisted with the crystals to  
315 facilitate the growth (Inoue 1994). It was observed that the capsules were completely  
316 filled with phase E crystals that were not sintered to each other, supporting that a large  
317 volume fraction of aqueous fluid had coexisted with the crystals.

318 Figure 3(a) shows a precession photograph of a representative crystal from the PE02  
319 run of  $270 \times 330 \times 530 \mu\text{m}^3$ . The diffraction spots of the phase E crystal are neither

320 diffused nor split, which representatively demonstrates absences of stacking disorder  
321 and mosaicism in the product crystals. The major element compositions in the cross  
322 sections of the other five crystals from the PE02 run showed a slight variation of Mg/Si  
323 =  $1.83 \pm 0.03$  which may include the effect of thermal gradient within the capsule  
324 (Figure 4(a)). The refined lattice constants given in Table 2 were consistent with those  
325 reported in previous works determined by precise X-ray diffraction measurements,  
326 which expanded between  $a = 2.9701(1) \text{ \AA}$ ,  $c = 13.882(1) \text{ \AA}$  and  $a = 2.9853(6) \text{ \AA}$ ,  $c =$   
327  $13.9482(7) \text{ \AA}$  (Kleppe et al. 2001; Kudoh et al. 1993). This variation may have been  
328 caused by the effects of the synthesis temperature in addition to the Mg/Si molar ratio of  
329 the starting materials. Considering the products starting only from Mg/Si = 2, the  
330 variation was reduced to between  $a = 2.9701(1) \text{ \AA}$ ,  $c = 13.882(1) \text{ \AA}$  for the sample  
331 synthesized at 1000 °C (Kudoh et al. 1993), and  $a = 2.975(1) \text{ \AA}$ ,  $c = 13.908(4) \text{ \AA}$  for the  
332 sample synthesized at 1100 °C (Kleppe et al. 2001). Assuming that this variation  
333 reflects only the temperature, we estimated the actual synthesis temperature for PE01 at  
334 ~1060 °C and PE02 at ~960 °C based on their lattice constants. The consistently lower  
335 temperatures by this estimation than the expected values of around 1100 °C may be the  
336 result of an isotope effect because the lattice constants can shrink by deuteration, as in  
337 the case of brucite having a layered structure similar that of phase E (Horita et al. 2010;

338 Chakoumakos et al. 2013). We note that the observed occurrences of minor  
339 clinoenstatite in PXRD and TOF–NPD patterns of the phase E products in the present  
340 study are consistent with a temperature rather close to 1100 °C; otherwise, the  
341 clinoenstatite peaks would have disappeared (Inoue 1994). From these observations it is  
342 inferred that the lattice parameters are slightly different between deuterated and  
343 hydrogenated crystals synthesized at the same temperature, which is desirable to be  
344 more quantitatively evaluated by a future study to update the accuracy of temperature  
345 scale in the following discussion.

346 The thermal gradient within each sample capsule was approximated by using the  
347 abovementioned lattice-to-temperature relation of the hydrogenated phase E. For this  
348 purpose, the full width at half maximum (FWHM) of representative reflections in the  
349 whole-capsule TOF–NPD patterns was determined as a function of *d*-spacing (Figure 5).  
350 Broadening by  $\Delta d/d = 0.07\%$  and  $0.10\%$ , compared with the original resolution of the  
351 diffractometer, was observed for the PE01 and PE02 run charges, which correspond to  
352 the averaged thermal inhomogeneity of 40 °C and 60 °C within the whole capsule,  
353 respectively. The thermal gradient is thus approximated to be 16–24 °C/mm or 20  
354  $\pm 4$  °C/mm for the 2.5 mm length between the center and edge of the capsule shown in  
355 Figure 1(a). We note that the observed FWHM values in TOF–NPD patterns were not

356 increased by particle size and strain effects, because the sample grains are too large to  
357 have the size effect and are loosely packed into the sample holder without strain.

358 We finally discuss hydrogen isotopes of the synthesized phase E crystals. By SIMS  
359 analysis of the cross sections of five different crystals from the PE02 run, we found that  
360 the isotope abundance of hydrogen in these crystals was uniform at  $D/(H + D) = 86 \pm$   
361 1 %, and that total water concentration in these crystals showed little variation ( $7.9 \pm 0.6$   
362 wt% as  $H_2O + D_2O$ ; Figure 6(a)). Considering these SIMS results in addition to the major  
363 element compositions by EPMA (Figure 4(a)), it was confirmed that these deuterated  
364 phase E crystals are suitable for structure refinement or other studies those require  
365 chemical and isotopic homogeneity within their entire crystal volumes. The uniform  
366 hydrogen isotope abundance of 14 % in all measured crystals indicates that hydrogen  
367 was contaminated into the crystals from outside of the capsule such as hydrogen  
368 released from an MgO sleeve or  $MgO-Cr_2O_3$  pressure medium despite calcination at  
369 1000 °C immediately prior to the cell construction and compression, because isotope  
370 abundance of hydrogen was less than 1 % in the deuterated starting material in the  
371 capsule (Purevjav et al., 2014a). We are now conducting a structure refinement study of  
372 phase E from the TOF-NPD pattern of the PE02 sample having known hydrogen  
373 isotope abundance in this study (Tomioka et al. 2014).

374

375 **Hydrous wadsleyite crystals**

376 Single crystals of hydrous magnesian wadsleyite of 400–1100  $\mu\text{m}$  in the largest  
377 dimensions were synthesized in the five relevant experimental runs (Table 1). The  
378 largest crystal size was  $1100 \times 800 \times 800 \mu\text{m}^3$  (Figure 2(b)), which is at least as large as  
379 the crystals previously synthesized using a thermal gradient (Shatskiy et al. 2009). From  
380 the results of these runs, we found that heating at around 1400 °C along with a reduction  
381 of 40 °C to 70 °C was the most effective temperature condition for growing large  
382 hydrous wadsleyite crystals, which is 200 °C to 300 °C higher than the wet solidus of  
383 the MgO–SiO<sub>2</sub> system (Inoue 1994). The observed melt fraction at this temperature  
384 reached about 70% in volume, where efficient buffering of crystal composition is  
385 expected to occur.

386 Figure 3(b) shows a precession photograph of a representative wadsleyite crystal  
387 from the EP13-2 run of  $700 \times 500 \times 500 \mu\text{m}^3$ , which demonstrates the high crystallinity  
388 of the product crystals. The major element compositions in the cross sections of the  
389 other three crystals from the EP13-2 run showed very little variation of Mg/Si = 1.86  
390  $\pm 0.01$  (Figure 4(a)), indicating that slow cooling of the crystals in large melt fraction is  
391 suitable for growing homogeneous hydrous wadsleyite crystals. These analyses also

392 indicate that the thermal gradient of 20 °C/mm, as estimated in the phase E runs, is too  
393 small to induce chemical heterogeneity in these crystals. The averaged H<sub>2</sub>O  
394 concentration ( $C_{\text{H}_2\text{O}}$ ) of wadsleyite in each run product including that from EP13-2 was  
395 separately estimated from each refined lattice constant (Table 2) using the established  
396 relation between the  $b/a$  ratio and  $C_{\text{H}_2\text{O}}$  (Jacobsen et al. 2005). The estimated  $C_{\text{H}_2\text{O}}$  had a  
397 variation between 1.6 % for the EP13-2 run quenched at 1250 °C and 1.1% for the  
398 HR-18 run quenched at 1320 °C. These results are consistent with the reported  
399 systematics of  $C_{\text{H}_2\text{O}} = 637e^{-0.0048T}$  (T in °C), which accurately predicted 1.6 wt% at  
400 1250 °C and 1.1 wt% at 1320 °C (Litasov et al. 2011).

401 A stark contrast of deuterium site positions between hydrous wadsleyite and hydrous  
402 ringwoodite has been reported by neutron diffraction of their deuterated powder  
403 samples (Purevjav et al. 2014; Sano-Furukawa et al. 2011). Whereas both structures  
404 commonly belong to spinel types, it was observed that deuterium in ringwoodite  
405 occupies both octahedral and tetrahedral sites, and that in wadsleyite occupies only an  
406 octahedral site. In other words, ringwoodite structure is uniquely vulnerable to  
407 exchanging silicon cations for hydrogen, which is much more effective for reducing  
408 sound velocity and increasing electrical conductivity than exchanging magnesium or  
409 iron cations for hydrogen (Mao et al. 2012; Panero et al. 2013). However, like the case

410 for the current phase E sample, contamination of hydrogen was observed in the Raman  
411 spectrum of the wadsleyite powder, which might be the reason for smaller deuterium  
412 occupancy around the octahedral site (M3) than the estimation; the observed occupancy  
413 was only ~75 % of the calculated value from D<sub>2</sub>O concentration of the wadsleyite  
414 (Sano-Furukawa et al. 2011).

415 The problem of hydrogen contamination can be completely eliminated by using the  
416 crystals synthesized in the present study. Because single crystal diffraction method is in  
417 principle fully tolerable for incoherent scattering induced by hydrogen, crystals  
418 involving natural hydrogen isotopes are suitably analyzed without taking care of the  
419 contamination. Therefore, the hydrous wadsleyite crystal of sufficient size and  
420 homogeneity in the present study has been measured at TOPAZ instrument at SNS  
421 without any problem. Whereas full structure refinement of this wadsleyite crystal is  
422 currently in progress, it was preliminary shown that hydrogen occupancy in the M3 site  
423 is much closer to the value that estimated from C<sub>H2O</sub> of the crystal (Purevjav et al.  
424 2014b). We will report the refinement result elsewhere in near future.

425

426 **Hydrous ringwoodite crystals**

427       Single crystals of hydrous ringwoodite of up to 1000  $\mu\text{m}$  were synthesized in the four  
428 relevant experimental runs (Table 1). Figure 2(c) shows one of the largest crystals from  
429 the EP13-3 run, which to the best of our knowledge is greater than any size reported in  
430 previous research. Such growth of ringwoodite requires a larger reduction in  
431 temperature (by 110  $^{\circ}\text{C}$ ) compared with that in the case of wadsleyite. The other runs  
432 with smaller changes in temperature (HR14) or constant temperatures (HR12 and  
433 EP12-6) resulted in smaller sizes, between 200  $\mu\text{m}$  and 600  $\mu\text{m}$ . These ringwoodite  
434 crystals were grown between 1260  $^{\circ}\text{C}$  and 1400  $^{\circ}\text{C}$ , where they were observed to  
435 coexist with silicate melt, as expected from the relevant phase relations (Ohtani et al.  
436 2000). For the EP13-3 run, we observed that about half of the sample volume was  
437 melted, whereas the other runs had smaller melt fractions. It was demonstrated from  
438 these observations that the growth rate changed largely with time. The rate was very  
439 rapid during initial crystallization and then became slower and slower during  
440 subsequent recrystallization with the melt-crystal interface kept close to the chemical  
441 equilibrium. The growth rate in the latter stage was moderately enhanced by  
442 precipitation of crystal volume upon reduction of the temperature, and an effective  
443 combination of these two processes made it possible to produce the single crystals of  
444 ringwoodite of 1 mm size in the EP13-3 run.



445 A precession photograph of the crystal shown in Figure 2(c) was captured along the  
446 [001] direction (Figure 3(c)). The photograph shows sharp diffraction spots, and no  
447 extra spot was observed. Moreover, SCXRD results of the two ringwoodite crystals  
448 taken from the HR12 and EP12-6 runs showed no intensity for the forbidden reflections,  
449 and no effect of twinning on their X-ray intensity data. The major element compositions  
450 in the cross sections of five crystals from the EP12-6 run and four crystals from the  
451 HR13-3 run showed little variations of  $Mg/Si = 1.92 \pm 0.01$  and  $(Mg+Fe)/Si = 1.93 \pm$   
452  $0.02$ , respectively (Figure 4(a)). These results indicate that slow cooling of the crystals  
453 in large melt fraction is suitable for growing homogeneous ringwoodite crystals. Using  
454 the refined lattice constants of the ringwoodite crystals by PXRD and SCXRD, the  $C_{H_2O}$   
455 of each run product was estimated along the relation between unit cell volume  $V_0$  and  
456  $C_{H_2O}$  (Ye et al. 2012). The estimated  $C_{H_2O}$  were between 2.0 wt% and 2.6 wt% (Table 2),  
457 which were close to the uppermost values of previously observed  $C_{H_2O}$  in hydrous  
458 ringwoodite (Inoue et al. 1998; Kudoh et al. 2000; Ye et al. 2012). It is suggested that  
459 the  $C_{H_2O}$  in ringwoodite increases with increasing synthesis duration; the highest  $C_{H_2O}$   
460 of 2.6% among the current crystal products was synthesized after 12 h of heating, which  
461 is indeed the most lengthy among all synthesis experiments of the hydrous ringwoodite.  
462 The following SIMS results shows no evidence of zoning of  $C_{H_2O}$ , so that the increase

463 of  $C_{H_2O}$  with time is not because of diffusion of  $H_2O$  into the ringwoodite crystals after  
464 their growth. Another possible process such as  $Mg^{2+}-Si^{4+}$  site disorder involving  
465 hydrogen may be additionally considered (Kudoh et al. 2000).

466 Hydrogen isotope concentrations were analyzed by SIMS for the polished cross  
467 sections of several ringwoodite crystals from the HR12 and EP12-6 runs. Although the  
468 HR12 product crystals were synthesized from fully deuterated starting materials, it was  
469 proved to have been extensively exchanged with hydrogen having natural isotope  
470 abundance. The isotope abundance of hydrogen were uniformly distributed at  $D/(H +$   
471  $D) = 45 \pm 3\%$  among all five crystals from the HR12 run (Figure 6). As in the case of  
472 phase E synthesis, such incorporation of outside hydrogen most likely occurred through  
473 its diffusion from outside of the capsule, which is inevitably more extensive for the  
474 ringwoodite synthesized at a temperature  $300\text{ }^\circ\text{C}$  higher than that for the phase E. It is  
475 therefore concluded that complete deuteration of ringwoodite crystals using the current  
476 cell design and experimental condition is difficult. On the other hand, the EP12-6  
477 product were synthesized from starting materials containing hydrogen with natural  
478 isotope abundance. The five measured crystals consistently showed a very uniform  
479 hydrogen concentration of  $1.6 \pm 0.1\text{ wt}\%$  (Figure 6). By all these analytical results, it  
480 was demonstrated that the slow-cooling method is suitable for growing

481 homogeneously-hydrous ringwoodite crystals which are appropriate for TOF single  
482 crystal diffraction.

483

#### 484 **Bridgmanite crystals**

485 Single crystals of bridgmanite with three different composition types were  
486 synthesized in the five relevant experimental runs (Table 1). The largest were  
487 iron-bearing bridgmanite crystals from the 5k2417 run at up to 600  $\mu\text{m}$  in the largest  
488 dimension (Figure 2(d)). These crystals were grown from the starting materials having  
489 silicate compositions between  $(\text{Mg,Fe})_2\text{SiO}_4$  and  $(\text{Mg,Fe})\text{SiO}_3$ , which are effective for  
490 preventing stishovite crystal inclusions (Shatskiy et al., 2007). Moreover, existence of  
491  $\text{H}_2\text{O}$  significantly facilitates the crystal growth of bridgmanite (Ito and Weidner, 1986;  
492 Shatskiy et al. 2007). The BSE image of these crystals showed their high homogeneity  
493 (Figure 7), which is consistent with the major element compositions measured in the  
494 cross sections of several different crystals from two experimental runs 5k2174 and  
495 5k2179 (Figure 4(b)). These results indicated that the slow cooling in water-bearing  
496 conditions is certainly effective for growing homogeneous bridgmanite crystals of  
497 moderate sizes, while the cell design and the experimental conditions are still necessary  
498 to be improved for growing the crystals 1 mm in size.

499 We confirmed the existence of hydrous silicate melt at temperature conditions of  
500 1650 °C to 1760 °C, from which crystal growth occurred in the present study. This is  
501 consistent with the reported phase relations at hydrous condition (Ito and Weidner.  
502 1986; Ohtani et al. 2000). We found that fine-tuning of the initial temperature and the  
503 extent of its decrease is not essential for increasing the crystal size, whereas a longer  
504 heating duration is definitely effective. From these observations it is inferred that  
505 recrystallization during the lengthy heating is more effective process for increasing  
506 crystal size of bridgmanite, compared with growth through crystal precipitation from  
507 the melt by reduction of temperature. In addition, it was demonstrated that sizes of  
508 aluminous bridgmanite crystals from the 5k2179 run were smaller than those of  
509 non-aluminous bridgmanite from the other runs. It was recently reported that such  
510 aluminous bridgmanite synthesized in hydrous conditions contains a significant amount  
511 of H<sub>2</sub>O up to almost 1 wt% (Inoue et al. 2012), whereas non-aluminous bridgmanite  
512 grown in hydrous conditions contains only a few hundreds of ppm of H<sub>2</sub>O (Litasov et  
513 al. 2003; Shatskiy et al. 2007). Apparently, a considerable amount of future study is  
514 required for growing highly hydrous, alumina-bearing bridgmanite crystals suitable for  
515 analysis of hydrogen by single crystal neutron diffraction.

516 The thermal gradient in the growth environment for the bridgmanite (Figure 1(b)) is  
517 presumably comparable to or smaller than that for the phase E, wadsleyite or  
518 ringwoodite (Figure 1(a)), because the capsule dimensions were smaller and the thermal  
519 insulator was about 70% thicker. In addition, at the significantly higher temperature  
520 conditions used for the bridgmanite synthesis, more efficient radiative heat transfer was  
521 expected along the semi-transparent MgO sleeve body, which had an increased grain  
522 size after the lengthy heating (Hofmeister 2005). Although we cannot evaluate the  
523 actual gradient in a quantitative manner, the expectation was qualitatively consistent  
524 with the observed crystal shape of bridgmanite. The shape was mostly euhedral  
525 throughout the entire capsule volume of all five runs, rather than the fan-shaped,  
526 needle-like aggregate type synthesized in the cells with thermal gradients significantly  
527 larger than 20 °C/mm (Ito and Weidner 1986; Shatskiy et al. 2007).

528 We finally note that a part of the bridgmanite crystal products in the present study  
529 has been successfully applied for in-situ radiative thermal conductivity measurements  
530 via optical absorption spectroscopy in a high-pressure diamond anvil cell (Goncharov et  
531 al. submitted). We are also planning to utilize these crystals for high-pressure studies of  
532 electronic spin and valence states of iron in the lower-mantle bridgmanite using  
533 synchrotron X-ray emission and Mössbauer spectroscopy, in addition to single-crystal

534 elasticity measurements of iron-bearing bridgmanite in the lower mantle using  
535 impulsive stimulated light scattering and Brillouin light scattering.

536

537

### Implications

538

539 It has been demonstrated that slow cooling in a Kawai-type cell is the best suitable  
540 method to synthesize homogeneous silicate crystals of 600  $\mu\text{m}$  to 1100  $\mu\text{m}$  sizes at  
541 pressures to 24 GPa. Degree of temperature reduction to facilitate the growth is  
542 necessary to be evaluated and optimized for each phase. Loeffert et al. (2002) used the  
543 slow-cooling method at 6 GPa in a belt press to grow a  $\text{SrCu}_2\text{O}_3$  crystal 2 mm in size.  
544 The method in the present study should be applicable for the crystal growth of such an  
545 oxide phase at pressures at least up to 20 GPa; it may provide a new application of the  
546 Kawai cell in the field of material sciences by exploiting its high thermal stability and  
547 controllability at such pressure regime.

548 For applications of the product crystals in the present study, we emphasize that  
549 synthetic single crystals at very high pressures have now become measurable by neutron  
550 diffraction. The result is newly allowing us to determine a hydrogen site occupancy of  
551 the contamination-free crystal in terms of isotopes. Moreover, TOF single crystal  
552 diffraction instruments provide about twice or even higher spatial resolution than

553 previously used powder instruments; D20 at Institut Laue-Langevin was used at  $d_{\min} =$   
554  $0.96 \text{ \AA}$  for analyzing the deuterated wadsleyite powder (Sano-Furukawa et al. 2011),  
555 whereas TOPAZ was used at  $d_{\min} = 0.50 \text{ \AA}$  (Schultz et al. 2014) and SXD at ISIS was  
556 used at  $d_{\min} = 0.31 \text{ \AA}$  (Artioli et al. 1995) for analyzing single crystal samples. The  
557 significantly smaller  $d_{\min}$  by these TOF single crystal instruments is because it is the  
558 best sensitive scheme for detecting diffraction at small  $d$ -values which is induced by  
559 neutrons of short wavelengths. Thus, the geometry of chemical bonding around a  
560 hydrogen atom in the deep-mantle hydrous minerals, including the distances of covalent  
561 and hydrogen bonding and the angle between them, are determined with significantly  
562 higher accuracy. Using these we will critically evaluate the proposed relations between  
563 the bonding geometry and Raman and Infrared spectra of hydrogen which includes that  
564 recently proposed by Panero et al. (2013) and their referred previous studies. Bonding  
565 geometries predicted by previous theoretical calculations of deep-mantle hydrogen will  
566 also be evaluated.

567 We finally note that all crystals in each sample capsule in the present study have a  
568 uniform chemical composition. Therefore, lattice parameters, water concentration and  
569 major element compositions of a candidate crystal can be promptly evaluated by PXRD  
570 and EPMA results of the other crystals without destroying the candidate itself, which

571 make it easy and straightforward to select the best sample for single crystal diffraction.  
572 Such a pre-evaluation is essential to avoid wasting invaluable neutron beam time by  
573 measuring an inappropriate crystal.

574

575

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576

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- 791

792 **Figure Captions**

793

794 **Figure 1.** Cell assembly designs for pressure generation (a) to 21 GPa coupled with 8  
795 mm truncation edge length (TEL) and (b) to 24 GPa coupled with 6 mm TEL. Sintered  
796 MgO–Cr<sub>2</sub>O<sub>3</sub> pressure medium with octahedral edge length (OEL) = 14 mm were  
797 commonly used for these cells. Dimensions of the sample capsules are shown at the  
798 central areas of the cells.

799

800 **Figure 2.** Photographs of representative crystals synthesized in and recovered from the  
801 high-pressure cells. Length of the scale bars is 500 μm unless explicitly shown. (a)  
802 Phase E from PE02; (b) hydrous magnesian wadsleyite from HR18; (c) hydrous  
803 magnesian ringwoodite from EP13-3; (d) iron-bearing bridgmanite from 5k2417.

804

805 **Figure 3.** Precession X-ray diffraction photographs of the single crystals. (a) Phase E  
806 from PE02 along the [001] zone axis. (b) Hydrous wadsleyite from EP13-2 along the  
807 [010] zone axis. The weak extra spots come from small wadsleyite crystals associated  
808 with the host crystal. (c) Hydrous ringwoodite from EP13-3 along the [001] zone axis.  
809 Numbers denote Miller indices of diffraction spots indicated by arrows. All of the  
810 patterns show sharp diffraction spots from non-disordered single crystals.



811

812 **Figure 4.** Major-element compositions of recovered single crystals taken by electron  
813 probe microanalyser (EPMA). (a) Results from Phase E (PE02), hydrous wadsleyite  
814 (EP13-2), and hydrous ringwoodite (EP12-6 and EP13-3) run charges plotted in ternary  
815 MgO–FeO–SiO<sub>2</sub> diagrams in molar ratio. (b) Results from MgFe-bearing bridgmanite  
816 (5k2174) and MgFeAl-bearing bridgmanite (5k2179) run charges plotted in ternary  
817 MgO–FeO–(Al<sub>2</sub>O<sub>3</sub>+SiO<sub>2</sub>) and (MgO+FeO)–Al<sub>2</sub>O<sub>3</sub>–SiO<sub>2</sub> diagrams in molar ratio. These  
818 data plots are surrounded by the squares with blue, purple or red color, which indicates  
819 that they are close-up views of the areas indicated by the corresponding colors in the  
820 entire triangle plots shown together. Each different color and symbol within these  
821 close-up views shows the EPMA results taken from each different crystal in the same  
822 run charge. The averaged chemical formula of MgFe- and MgFeAl-bearing bridgmanite  
823 crystals (5k2174 and 5k2179, excluding their H<sub>2</sub>O components which were not  
824 analyzed) are Mg<sub>0.96</sub>Fe<sub>0.07</sub>Si<sub>0.98</sub>O<sub>3</sub> and Mg<sub>0.90</sub>Fe<sub>0.13</sub>Al<sub>0.11</sub>Si<sub>0.90</sub>O<sub>3</sub>, respectively.

825

826 **Figure 5.** Relative neutron diffraction linewidth profiles  $\Delta d/d$  of phase E as a function  
827 of d-spacing.  $\Delta d$  was defined as the full width at half maximum (FWHM) of each  
828 reflection. Thin straight lines show expected  $\Delta d/d$  positions when the described

829 temperature inhomogeneity exists within the entire capsule ( $\Delta Temp$ ). Resolution of the  
830 diffractometer optics is shown by the bold line ( $\Delta Temp = 0$  °C). Filled and open circles  
831 denote the results obtained for the entire PE01 and PE02 run charges, respectively. A  
832 Lorentz-type function was assumed for the lineshape of the observed reflections, where  
833 the linewidth was approximated by the sum of that from the diffractometer optics and  
834 that from  $\Delta Temp$ .

835

836 **Figure 6.** Secondary ion mass spectroscopy (SIMS) results of hydrogen and deuterium  
837 concentration distributions measured on polished cross sections of representative five  
838 crystals selected from each of the three representative run charges (15 crystals in total).  
839 (a, b) H<sub>2</sub>O and D<sub>2</sub>O concentrations of the deuterated phase E crystals recovered from  
840 the PE02 run, respectively. (c, d) H<sub>2</sub>O and D<sub>2</sub>O concentrations of the deuterated  
841 hydrous ringwoodite crystals recovered from the HR12 run, respectively. (e) H<sub>2</sub>O  
842 concentration of hydrous ringwoodite crystals recovered from the EP12-6 run. The  
843 vertical scale shows recalculated H<sub>2</sub>O or D<sub>2</sub>O concentrations from their H/<sup>30</sup>Si or D/<sup>30</sup>Si  
844 ratios, respectively. In order to find compositional zoning or another heterogeneity  
845 within each crystal, the distances between each measurement position and the nearest  
846 rim of each cross section are shown in the horizontal scales. The broken lines show the

847 distances from the center of each cross section to its nearest rim; note that actual  
848 dimension(s) of the measured crystals is larger than twice of these distances, because  
849 these are the observed minimum only within each cross section. The standard deviation  
850 of each water concentration value is estimated to be about 6 % of the value (indicated  
851 by the error bars), which corresponds to the standard deviation of the measurement  
852 results of the hornblende H<sub>2</sub>O standard (15 points in total). D<sub>2</sub>O within the PE02 and the  
853 HR12 crystal products were partly exchanged for H<sub>2</sub>O during the crystal growth.

854

855 **Figure 7.** Back-scattering electron (BSE) image of polished cross section of an  
856 iron-bearing bridgmanite crystal from 5k2417. The scale bar is 100 μm in length.  
857

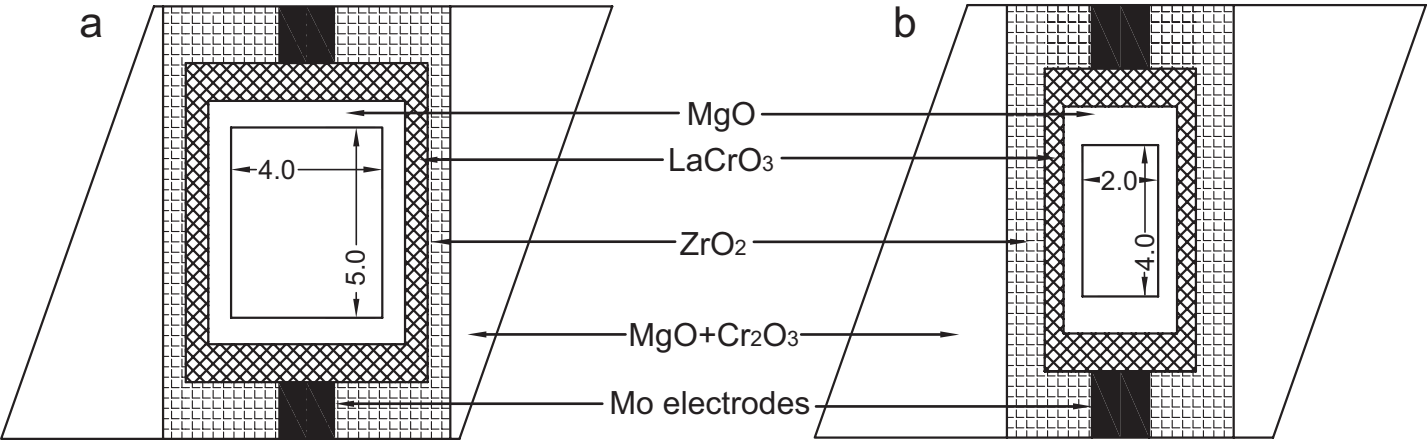
Table 1. Experimental conditions and the maximum size of the recovered crystals

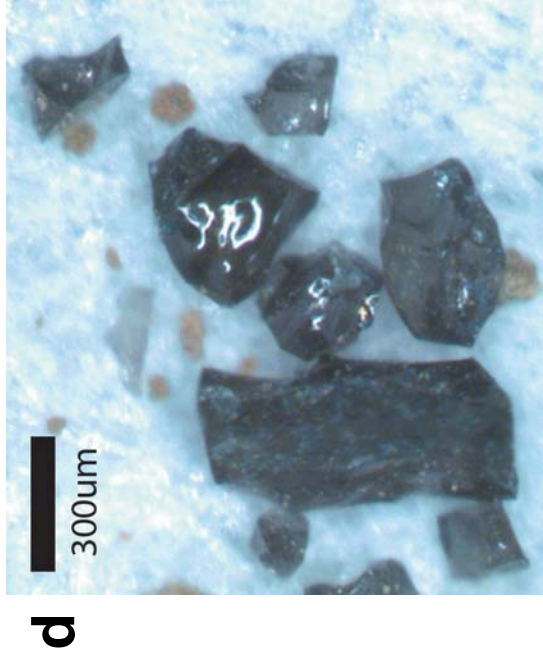
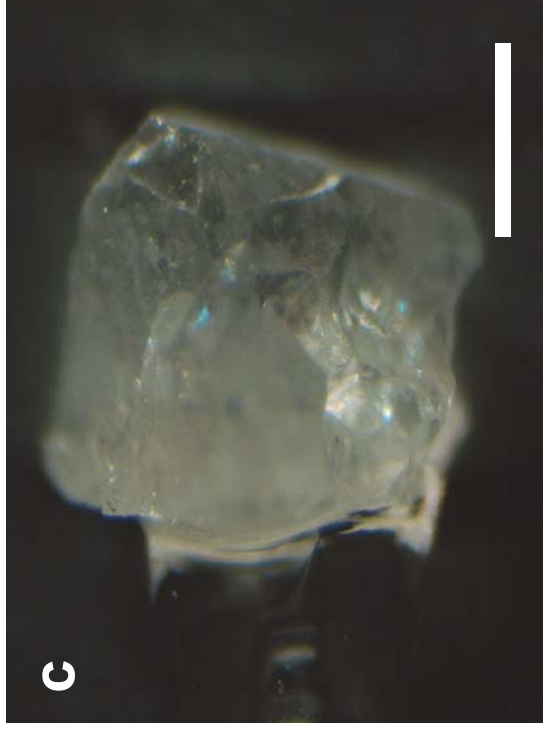
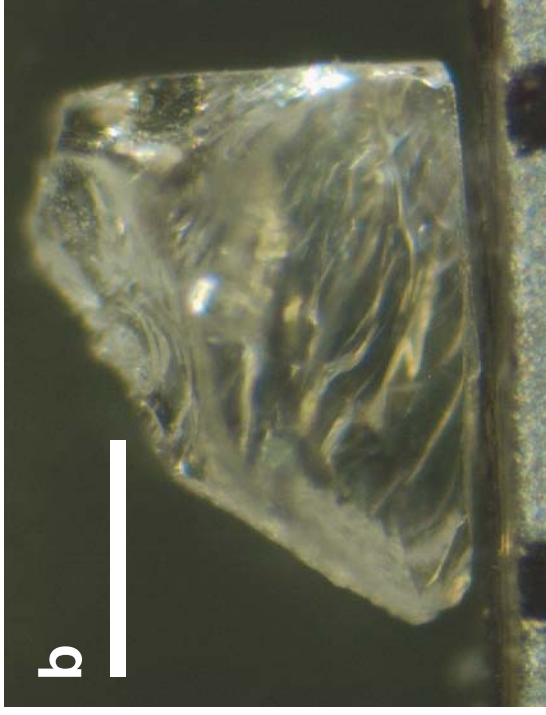
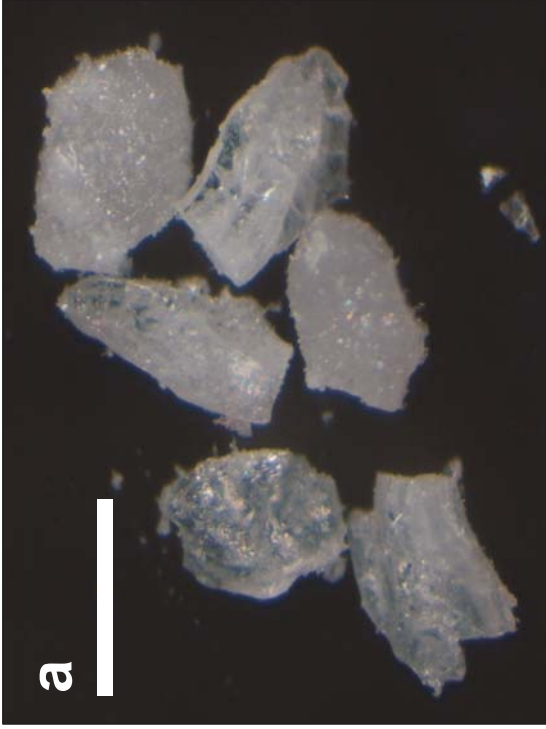
Run No.	Target phase	OEL/TEL [mm]	Load [MN]	Pressure [GPa]	Temperature [°C]	Duration [hours]	Starting materials	Size [ $\mu\text{m}$ ]	Silicate melting
PE01	phase E	14/8	6.5	14	1100 / 1090	2	$\text{Mg}_2\text{SiO}_4$ +22 wt.% $\text{D}_2\text{O}$	500	No
PE02	phase E	14/8	6.5	14	1100 / 1100	3	$\text{Mg}_2\text{SiO}_4$ +22 wt.% $\text{D}_2\text{O}$	600	No
EP13-2	Mg wadsleyite	14/8	8.5	17	1300 / 1250	13	$\text{Mg}_2\text{SiO}_4$ + 10 wt.% $\text{H}_2\text{O}$	600	Yes
EP13-1	Mg wadsleyite	14/8	8.5	17	1400 / 1300	6	$\text{Mg}_2\text{SiO}_4$ + 10 wt.% $\text{H}_2\text{O}$	600	Yes
HW7	Mg wadsleyite	14/8	8.5	17	1400 / 1310	24	$\text{Mg}_2\text{SiO}_4$ + 10 wt.% $\text{H}_2\text{O}$	400	Yes
HW8	Mg wadsleyite	14/8	8.5	17	1400 / 1360	6	$\text{Mg}_2\text{SiO}_4$ + 15 wt.% $\text{H}_2\text{O}$	800	Yes
HW18	Mg wadsleyite	14/8	8.5	17	1390 / 1320	10	$\text{Mg}_2\text{SiO}_4$ + 15 wt.% $\text{H}_2\text{O}$	1100	Yes
HR14	Mg ringwoodite	14/8	19	21	1340 / 1310	1.5	$\text{Mg}_2\text{SiO}_4$ + 5 wt.% $\text{D}_2\text{O}$	200	Yes
HR12	Mg ringwoodite	14/8	19	21	1360 / 1360	2.5	$\text{Mg}_2\text{SiO}_4$ + 11 wt.% $\text{D}_2\text{O}$	600	Yes
EP13-3	Mg ringwoodite	14/8	19	21	1400 / 1290	12	$\text{Mg}_2\text{SiO}_4$ + 10 wt.% $\text{H}_2\text{O}$	1000	Yes
EP12-6	MgFe ringwoodite	14/8	9.6	18	1270 / 1260	1	$(\text{Mg}_{0.91}\text{Fe}_{0.09})_2\text{SiO}_4$ + 10 wt.% $\text{H}_2\text{O}$	300	Yes
5k2247	Mg bridgmanite	14/6	19	24	1710 / 1650	3	$\text{Mg}_{1.4}\text{SiO}_{3.4}$ + 9 wt.% $\text{H}_2\text{O}$	200	Yes
5k2221	Mg bridgmanite	14/6	19	24	1700 / 1700	8	$\text{Mg}_{1.3}\text{SiO}_{3.3}$ + 7 wt.% $\text{H}_2\text{O}$	300	Yes
5k2174	MgFe bridgmanite	14/6	19	24	1650 / 1650	5.5	$(\text{Mg}_{0.92}\text{Fe}_{0.08})_{1.5}\text{SiO}_{3.5}$ + 6 wt.% $\text{H}_2\text{O}$	100	Yes
5k2417	MgFe bridgmanite	14/6	19	24	1700 / 1670	12	$(\text{Mg}_{0.84}\text{Fe}_{0.16})_{1.8}\text{SiO}_{3.8}$ + 7 wt.% $\text{H}_2\text{O}$	600	Yes
5k2179	MgFeAl bridgmanite	14/6	19	24	1760 / 1750	7	$(\text{Mg}_{0.92}\text{Fe}_{0.08})_{1.5}\text{SiO}_{3.5}$ + 0.2 $\text{Al}_2\text{O}_3$ + 5 wt.% $\text{H}_2\text{O}$	100	Yes

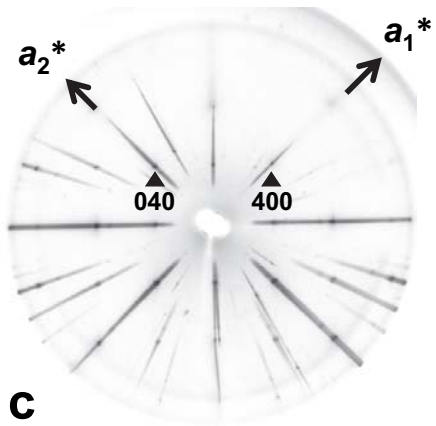
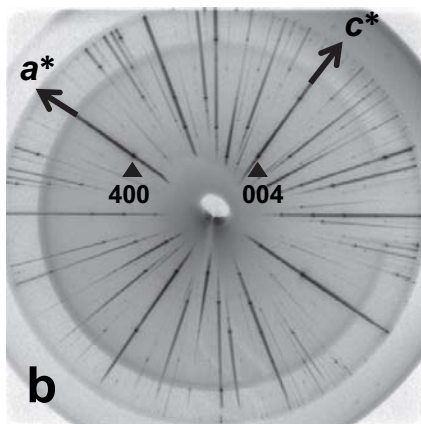
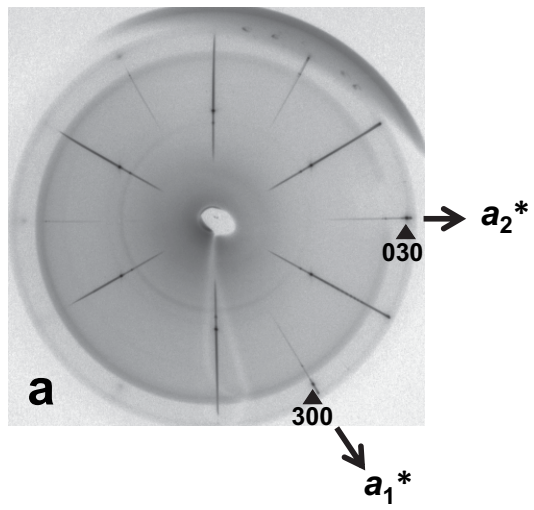
Table 2. Crystallographic parameters of the recovered crystals

Run No.	Phase	Space group	Lattice constants [ $\text{\AA}$ ]			Unit cell vol. $V_0$ [ $\text{\AA}^3$ ]	Hydrogen concentration [wt.%]
			$a$	$b$	$c$		
PE01	phase E	$R\text{-}3m$	2.9731(1)	13.913(1)	106.50(1)	N. A.	
PE02	phase E	$R\text{-}3m$	2.9679(1)	13.888(1)	105.94(1)	1.0(1) (SIMS, H <sub>2</sub> O) 6.9(5) (SIMS, D <sub>2</sub> O)	
EP13-2	Mg wadsleyite	$Imma$	5.6844(8)	11.527(2)	8.2532(7)	540.80(8)	1.6 (from $b/a$ )
EP13-1	Mg wadsleyite	$Imma$	5.6852(4)	11.508(1)	8.2523(8)	539.92(10)	1.3 (from $b/a$ )
HW7	Mg wadsleyite	$Imma$	5.6871(7)	11.512(2)	8.2568(8)	540.61(8)	1.3 (from $b/a$ )
HW8	Mg wadsleyite	$Imma$	5.6879(6)	11.513(1)	8.2542(7)	540.55(8)	1.3 (from $b/a$ )
HW18	Mg wadsleyite	$Imma$	5.6882(6)	11.497(1)	8.2554(8)	539.89(8)	1.1 (from $b/a$ )
HR14	Mg ringwoodite	$Fd\text{-}3m$		8.0795(4)		527.41(8)	2.2(2) (from $V_0$ )
HR12	Mg ringwoodite	$Fd\text{-}3m$		8.0773(3) 8.0771(5)*		526.99(6)	2.1(2) (from $V_0$ ) 0.6(1) (SIMS, H <sub>2</sub> O) 0.5(0) (SIMS, D <sub>2</sub> O)
EP13-3	Mg ringwoodite	$Fd\text{-}3m$		8.0834(2)		528.19(4)	2.6(2) (from $V_0$ )
EP12-6	Mg <sub>92</sub> Fe <sub>08</sub> ringwoodite	$Fd\text{-}3m$		8.0903(6)*		529.53(7)	2.0(2) (from $V_0$ ) 1.6(1) (SIMS, H <sub>2</sub> O)

\*determined by SCXRD

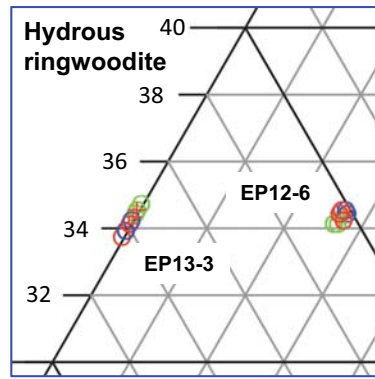
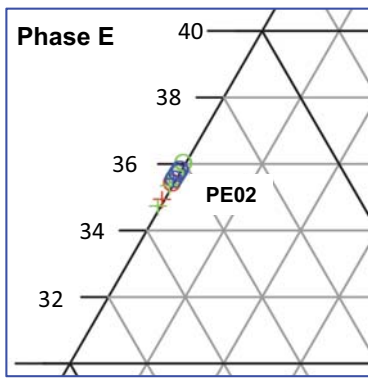
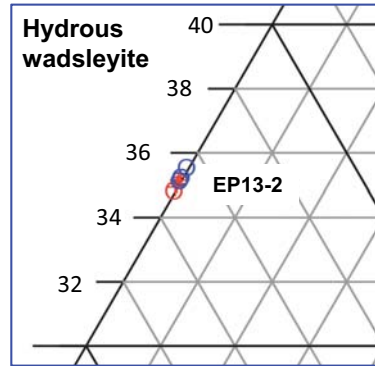
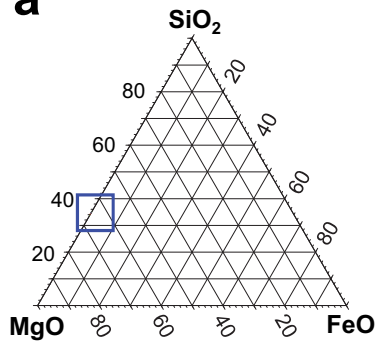








**a**



**b**

