

1 **REVISION 1**

2 **Complex IR spectra of OH⁻ groups in silicate glasses: implications for the use of the 4500**
3 **cm⁻¹ IR peak as a marker of OH⁻ groups concentration**

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9
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11
12 **Abstract**

13 Previous studies of hydrous glasses and melts with infrared spectroscopy have led to the
14 conclusion that the IR combination peaks near 4500 and 5200 cm⁻¹ reflect the existence of
15 OH⁻ (hydroxyl) groups and H₂O_{mol} water molecules in those materials. Here, we show that the
16 glass chemical composition can impact profoundly the intensities and frequencies of the
17 fundamental O-H stretching signal and, therefore, potentially those of the 4500 and 5200 cm⁻¹
18 combination peaks. In alkali silicate glasses, compositional effects can give rise to peaks
19 assigned to fundamental O-H stretching at frequencies as low as 2300 cm⁻¹. This expanded
20 range of Raman intensity assigned to O-H stretch is increasingly important as the ionic radius
21 of the alkali metal increases. As a result, the combination of the fundamental O-H stretch in
22 OH⁻ groups with the Si-O-H stretch located near 910 cm⁻¹ gives rise to a complex
23 combination signal that can extend to frequencies much lower than 4200 cm⁻¹. This
24 combination signal then becomes unresolvable from the high-frequency limb of the band

25 assigned to fundamental O-H stretch vibration in the infrared spectra. It follows that, when O-
26 H stretch signals from OH⁻ groups extend to below 3000 cm⁻¹, the 4500 cm⁻¹ peak does not
27 represent the total OH⁻ signal. Under such circumstances, this infrared peak may not be a
28 good proxy for determining the concentration of OH⁻ hydroxyl groups for glassy silicate
29 materials.

30

31 **Introduction**

32 Water dramatically affects the physical properties of silicate melts and glasses (see the
33 review by [Mysen & Richet, 2005](#)), and, as a result, is of key importance in both industrial and
34 geologic processes. An important feature of water dissolved in glassy and molten silicate
35 materials is that it exists both as molecules (H₂O_{mol}) and as hydroxyl groups (OH⁻) bonded to
36 the silicate network. It has often been stated that those two different species give rise to IR
37 combination peaks at ~5200 and ~4500 cm⁻¹ respectively ([Davis & Tomozawa, 1996](#); [Efimov](#)
38 [& Pogareva, 2006](#); [Malfait, 2009](#); [Scholtze, 1960](#); [Stolper, 1982](#)). The OH⁻ groups are bonded
39 to the tetrahedral network of silicate glasses and melts, forming Si-OH and Al-OH bonds by
40 breakage of bridging oxygen bonds (Si-O-Si, Al-O-Al, and Si-O-Al) (see for instance
41 [Moulson & Roberts, 1961](#); [Scholtze, 1960](#); [Stolen & Walrafen, 1976](#); [Stolper, 1982](#)). This
42 solution mechanism results in depolymerization of the silicate network (see for instance the
43 study of [Mysen & Cody, 2005](#)). In contrast, water dissolved in the form of H₂O_{mol} species
44 does not change melt polymerization. As a result, the dissolution of water as OH⁻ groups or as
45 H₂O_{mol} species has different effects on the polymerization and hence properties of hydrous
46 silicate amorphous materials. These differences are reflected in transport properties of
47 hydrous melts, glass transition temperature (see for instance [Deubener et al., 2003](#)), solidus

48 temperatures, melt/mineral phase equilibria, and element partitioning (see for a review [Mysen](#)
49 [and Richet, 2005](#)).

50 The water speciation, in terms of $\text{H}_2\text{O}_{\text{mol}}$ species and OH^- groups, varies with total water
51 concentration and temperature ([Behrens & Yamashita, 2008](#); [Nowak & Behrens, 1995](#);
52 [Stolper, 1982](#)). It also depends on composition ([Deubener et al., 2003](#); [Moretti et al., 2014](#)).
53 However, it is not known how temperature effects vary with bulk composition, and how bulk
54 compositional variables govern solution mechanisms and solubility.

55 FTIR spectroscopy may be used to examine such effects by using the combination peaks
56 often located near 4500 and 5200 cm^{-1} in the spectra. However, as a first step, it is necessary
57 to ascertain exactly what governs the frequency and integrated intensity of the 4500 and 5200
58 cm^{-1} combination peaks in the IR spectra of hydrous silicate glasses (and melts). An example
59 of this is the recent work of [Malfait \(2009\)](#) who concluded that the 4500 cm^{-1} IR peak is
60 formed by the combination of the fundamental O-H stretching mode at $\sim 3600 \text{ cm}^{-1}$ with the
61 Si-O-H, and presumed Al-O-H, stretching vibrational mode. The Al-O-H vibrational mode
62 can occur near 800 cm^{-1} , whereas the Si-O-H is near 920 cm^{-1} . Combined with the
63 fundamental O-H stretching mode at 3600 cm^{-1} in simple aluminosilicate glasses, the exact
64 frequency of the 4500 cm^{-1} would be somewhat correlated, therefore, with the Al/Si of the
65 glass. Concerning the 5200 cm^{-1} peak, it is attributed to arise from the combination of the
66 3600 cm^{-1} stretching and 1630 cm^{-1} bending modes of $\text{H}_2\text{O}_{\text{mol}}$ species ([Scholtze, 1960](#);
67 [Stolper, 1982](#)),

68 The infrared signal assigned to O-H stretch in OH^- and $\text{H}_2\text{O}_{\text{mol}}$ species, usually
69 characterized by an asymmetric band centered near 3600 cm^{-1} , can have a complex shape in
70 some silicate glasses, where broad bands extend into the 2000-3000 cm^{-1} frequency range (see
71 for instance data on sodium silicate glasses of [Uchino et al., 1991](#); [Zotov & Keppler, 1998](#)). If

72 those bands are from OH⁻ groups, then their combination with the Si-O-H stretch mode
73 located near 910 cm⁻¹ (Zotov & Keppler, 1998) will give OH⁻ combination signals at
74 frequencies significantly lower than 4500 cm⁻¹. However, if the low-frequency signals are
75 from O-H stretch in H₂O_{mol} species, their combination with the H₂O_{mol} bending mode near
76 1630 cm⁻¹ will result in signals in the range of the OH⁻ combination peak. In both cases, using
77 the intensities of the 4500 and 5200 cm⁻¹ peaks would be not appropriate to determine the
78 [OH⁻] and [H₂O_{mol}] concentration in the glass. To evaluate impacts of such effect and their
79 consequence for the determination of [OH⁻] and [H₂O_{mol}], we studied alkali (Li, Na, K)
80 silicate glasses containing different amount of dissolved water by using FTIR and Raman
81 spectroscopies.

82

83 **Experimental methods**

84 Anhydrous KS4 (K₂Si₄O₉), NS4 (Na₂Si₄O₉) and LS4 (Li₂Si₄O₉) glasses were synthesized
85 by mixing pure anhydrous SiO₂, K₂CO₃, Na₂CO₃ and Li₂CO₃ powders and grinding them
86 under ethanol for about 1 hour. The powders were placed in a Pt crucible and heated at about
87 1.5°/min until reaching 1000°C before bringing the samples to complete melting. Melting was
88 accomplished near 1200°C for the KS4 composition, 1400 °C for the NS4 composition, and
89 1600°C for the LS4 composition, in agreement with their respective liquidus temperature
90 (Epler, 1963; Schairer & Bowen, 1955; Schairer & Bowen, 1956). Melts were quenched to
91 glass by placing the bottom of the crucible in liquid H₂O. The resulting starting materials
92 were stored at 110°C to avoid reaction with atmospheric water, an effect particularly critical
93 for the KS4 glass (Schairer & Bowen, 1955). Chemical analysis of the starting glasses, shown
94 in Table 1, were obtained by using a JEOL FE-SEM equipped with an energy dispersive

95 spectrometer. Measurements were performed on 25 x 25 micron areas with a 15 kV and 1.04
96 nA current. Lithium was quantified by difference.

97 Hydrous glasses were synthesized by melting the starting glasses together with known
98 amounts of H₂O in a piston-cylinder apparatus (Boyd & England, 1960). To this end, glasses
99 were first crushed to powder in ethanol, and then heat-treated at 400°C for 1 hour to ensure
100 complete removal of the ethanol. Then, glass powder and liquid H₂O were added in ~10 mm
101 long, 5 mm diameter platinum capsules, which were welded shut and placed in ¾-inch
102 diameter furnace assemblies, based on the design of Kushiro (1976). The products were then
103 subjected to the desired pressure (1.5 GPa) and temperature (1450 °C for NS4, 1550 °C for
104 KS4 and 1650 °C for LS4) for 90 minutes. Temperatures were measured with type S
105 thermocouples with no correction for pressure on their *emf*. Such pressure effects may be as
106 much as 10°C (Mao et al., 1971). Pressure was calibrated against the melting point of NaCl
107 and the calcite-aragonite transformation (Bohlen, 1984). Estimated uncertainties are ~10°C
108 and ~0.1 GPa, respectively.

109 Transmission infrared spectra were recorded using a Jasco® IMV-4000 multi-channel
110 infrared spectrometer with a 10X objective, an MCT detector and the standard light source, a
111 100µm aperture, and 1000 acquisitions. The hydrous glasses were double-polished to a
112 thickness in the 20-110 µm range by using oil in order to avoid any reaction of their surface
113 with water. For KS4 composition glasses, two sets of samples were used. One was ~20 µm
114 thick and was used for recording the 3600 cm⁻¹ O-H stretching band. However, those ~20 µm
115 samples were too thin for acquisition of combination signals. Consequently, another thicker set
116 of samples (~100 µm) were prepared for this purpose. The glass samples were placed on BaF₂
117 windows during measurements. After each spectrum, the sample was moved aside and a
118 background spectrum was recorded at the same spot. Division of sample spectra by

119 background spectra gave the final absorbance spectra, which were normalized to samples
120 thicknesses.

121 Raman spectra were recorded using a Jasco[®] NRS3100 spectrometer, equipped with a
122 holographic notch filter, a single monochromator, and a 1024x128 Andor[®] DV401-F1 CCD
123 Peltier-cooled at -71°C. A Coherent[®] solid laser line of 490 nm was used to excite samples,
124 with a power of ~44 mW on the sample. The laser beam was focused through a 50X
125 Olympus[®] lens. Measurements were done with a 1200 grooves/mm grating. Frequency
126 accuracy of the Raman spectrometer was checked by using the 520.7 cm⁻¹ Raman peak of
127 pure silicon. All spectra are unpolarized. After acquisition, a baseline constrained in the
128 spectral portions devoid of signal near 2000 and 3900 cm⁻¹ was subtracted from the spectra.
129 The intensities of baseline-subtracted spectra were normalized to their total spectral area
130 (calculated between 250 and 3900 cm⁻¹). Further details about spectra processing are available
131 in Le Losq et al. (2014a).

132 Total water concentration of glasses was estimated by using the FTIR and Raman spectra.
133 Absorption coefficients for the 4500 and 5200 cm⁻¹ IR peaks in NS4 glasses are available in
134 the literature. We used the protocol of Yamashita et al. (2008) together with their absorption
135 coefficients, which were derived by using total water concentrations determined
136 independently with Karl-Fisher titration. Because of that and despite the fact that the present
137 study will discuss the use of the 4500 cm⁻¹ IR peak for inferring [OH⁻] abundance, the FTIR-
138 based [OH⁻] + [H₂O_{mol}] sum is not affected by possible error in the assignments and
139 attributions of the IR absorption peaks. As IR absorption coefficients are not available for
140 KS4 and LS4 glasses, Raman spectroscopy was used to estimate the water concentration of
141 these glasses. A protocol similar to that described in Le Losq et al. (2012) was used. The
142 variations of the areas of the Raman signals assigned to O-H stretching, located between 2000

143 and 4000 cm^{-1} , was calibrated with the FTIR-based total water concentrations of the same
144 NS4 glasses. This calibration was then used with the other LS4 and KS4 glass series to
145 determine their water concentration. More details are available in Le Losq et al. (2014a).
146 Table 2 reports the measured water concentrations of the glasses.

147

148 **Results**

149 In the FTIR spectra of hydrous alkali silicate glasses, strong peaks at ~ 1630 (ν_B H-O-
150 H), ~ 2350 (A in Fig. 1), ~ 2800 (B in Fig. 1) and $\sim 3570\text{ cm}^{-1}$ (C in Fig. 1) are observed. The
151 IR spectra of hydrous NS4 glasses are similar to those previously published for NS4 hydrous
152 compositions (e. g., Uchino et al., 1991; Zotov & Keppler, 1998). Spectra of hydrous KS4 and
153 LS4 glasses have not been reported in the literature to our knowledge.

154 The three peaks (A, B, and C) in the $2000\text{-}4000\text{ cm}^{-1}$ portion of FTIR spectra are
155 attributed to stretching of O-H bonds in $\text{H}_2\text{O}_{\text{mol}}$ and OH^- groups. The 1630 cm^{-1} peak is
156 assigned to the bending vibrational mode of $\text{H}_2\text{O}_{\text{mol}}$ species (Bartholomew et al., 1980;
157 Scholtze, 1960; Uchino et al., 1991; Wu, 1980; Zotov & Keppler, 1998). The other small
158 peaks located at frequencies below 2000 cm^{-1} probably arise from overtones of the vibrational
159 modes of the glass silicate network (see the study of Efimov & Pogareva, 2006 for instance)
160 and are neglected because of their minor importance regarding the topic of this study. Above
161 3800 cm^{-1} , peaks from the combination modes are present but not visible at the scale of
162 Figure 1. They will be discussed in detail latter.

163 In the IR spectra of all glasses, the peaks A and B are broad, whereas the peak C is
164 considerably sharper. Absorbance values of those peaks are reported in Table 2. At given
165 water concentration, the absorbance intensities vary strongly as a function of the ionic radius
166 of the alkali metal (Fig. 1). The absorbance intensities of peaks A and B increase strongly

167 with increasing ionic radius of alkali, whereas that of peak C decreases. In a given
168 composition, increasing water content results in a general absorbance increase of all
169 contributions, but the absorbance of the $\sim 3600\text{ cm}^{-1}$ contribution increases slightly more
170 rapidly than the others (Table 2 and Fig. 1).

171 The three peaks in the infrared spectra, A, B and C, are also visible in Raman spectra
172 of the same glasses (Fig. 2). As already emphasized by FTIR data (Fig. 1), the comparison of
173 the spectra of the silicate glasses containing 17.6 mol% H₂O highlights that with increasing
174 the ionic radius of the alkali, peaks A and B grow at the expense of the peak C. For
175 comparison purposes, the Raman spectrum of a pure SiO₂ glass with 20 mol% H₂O is also
176 reported in Figure 2. In this latter spectrum, the signal intensity below 3000 cm^{-1} is negligible
177 compared with the intensity between 3000 and 4000 cm^{-1} . The two peaks located near 3595
178 and 3650 cm^{-1} can be attributed to O-H stretching from OH⁻ groups in different environments
179 (*the 3595 cm^{-1} peak is only observed in H₂O-rich silica glasses*), but their attribution remains
180 unclear and is subjected to speculation (see for instance the studies of [Davis & Tomozawa,](#)
181 [1996](#); [Efimov & Pogareva, 2006](#); [Holtz et al., 1996](#); [Mysen & Virgo, 1986a](#)).

182 The shape of the peaks near 4500 and 5200 cm^{-1} in the IR spectra also varies
183 significantly with the glass composition at given total water content (Fig. 3). In IR spectra of
184 the LS4 glasses, two asymmetric peaks are observed at ~ 4520 and $\sim 5240\text{ cm}^{-1}$. Their
185 intensities increase with increasing water concentration, but in different ways. The $\sim 4520\text{ cm}^{-1}$
186 peak is more intense than the $\sim 5240\text{ cm}^{-1}$ peak in the IR spectrum of the LS4 + 3.3 mol%
187 H₂O glass (Figs. 3, 4). In contrast, at 17.6 mol% H₂O, the $\sim 5240\text{ cm}^{-1}$ peak is more intense
188 than the $\sim 4520\text{ cm}^{-1}$. Therefore, the intensity of the $\sim 4520\text{ cm}^{-1}$ peak increases less and less
189 with water addition, whereas that of the $\sim 5240\text{ cm}^{-1}$ peak shows an opposite trend (Fig. 4). If
190 these two IR peaks are attributed to OH⁻ (4520 cm^{-1}) and H₂O_{mol} (5240 cm^{-1}) species, as has

191 commonly been the case (see Introduction), their intensity variations are consistent with a
192 decrease of the $\text{OH}^-/\text{H}_2\text{O}_{\text{mol}}$ ratio with increasing water concentration in LS4 glass. In addition
193 to those two peaks, a small and broad peak near 4000 cm^{-1} is visible. This peak is well
194 separated from the $\sim 4520\text{ cm}^{-1}$ peak (Fig. 3, LS4 spectra). Its intensity increases with water
195 concentration, but is also affected by the high-frequency tail of the 3600 cm^{-1} peak. Indeed,
196 this tail shifts toward high frequency and probably varies in intensity with increasing water
197 concentration (Fig. 3). This topological evolution probably affects the shape of the peak
198 located near 4000 cm^{-1} . The origin of the latter is unresolved but may involve unidentified
199 combination modes of both $\text{H}_2\text{O}_{\text{mol}}$ and OH^- species (see the discussion of [Stolper, 1982](#) and
200 also [Davis & Tomozawa, 1996](#); [Efimov & Pogareva, 2006](#)).

201 The three peaks near 4000 , 4520 and 5236 cm^{-1} also are observed in the IR spectra of
202 NS4 hydrous glasses. Their intensities vary with water concentration in a way similar to that
203 observed in the LS4 spectra (Figs. 3, 4). However, the separation of the ~ 4000 and $\sim 4520\text{ cm}^{-1}$
204 1 peaks is less pronounced in NS4 spectra than in LS4 spectra (Fig. 3). It becomes even less
205 pronounced in the KS4 spectra, in which the peaks in the $3900\text{-}4700\text{ cm}^{-1}$ range are poorly
206 defined (Fig. 3). In the KS4 + 3.3 mol% H_2O spectrum, a small peak at $\sim 4510\text{ cm}^{-1}$ is
207 observed, and there is no peak near 5200 cm^{-1} . No peak is observed near 4000 cm^{-1} , but the
208 high-frequency end of the 3600 cm^{-1} O-H stretch band affects the intensity of the spectrum in
209 the $3800\text{-}4200$ portion of spectra (Fig. 3). In the KS4 + 9.4 mol% H_2O glass spectrum, the
210 $\sim 4510\text{ cm}^{-1}$ peak is barely visible. The IR absorbance between 3800 and 4800 cm^{-1} shows a
211 global, nearly linear, decay. There is no intensity near 5200 cm^{-1} in this spectrum. At 17.6
212 mol% H_2O , the $\sim 4510\text{ cm}^{-1}$ contribution is very broad and preceded by another broad small
213 peak centered near 4050 cm^{-1} . Finally, a small symmetric peak is present at $\sim 5230\text{ cm}^{-1}$ in the
214 spectrum of this water-rich composition.

215 The absence of the $\sim 5200\text{ cm}^{-1}$ peak in the IR spectra of the KS4 with 3.3 and 9.4
216 mol% H₂O and its low intensity in the spectrum of KS4 with 17.6 mol% H₂O (Figs. 3, 4)
217 indicate that essentially all the water is present as OH⁻ groups in KS4 compositions, in
218 agreement with results from ¹H NMR spectroscopy (Le Losq et al., 2014a). In addition, the IR
219 results and previous ¹H NMR data (Le Losq et al., 2014a) both indicate that, at fixed water
220 concentration, the OH⁻/H₂O_{mol} ratio increases strongly with increasing ionic radius of alkali in
221 the order Li < Na < K (Fig. 3; Le Losq et al., 2014a). However, despite their very high OH⁻
222 /H₂O_{mol} ratios, the KS4 hydrous glasses show the less resolved OH⁻ combination peaks in
223 their IR spectra in comparison to the NS4 and LS4 compositions (Fig. 3). Interestingly, they
224 also show the highest intensities in the 4000-4700 cm⁻¹ portion of spectra (see comparison in
225 Fig. 5).

226

227 Discussion

228 Increasing the ionic radius of the alkali cation in hydrous silicate glasses, formed by
229 quenching melts equilibrated at high temperature and pressure, produces in both the FTIR and
230 Raman spectra a strong intensity increase of two peaks at frequencies below 3000 cm⁻¹ and
231 attributed to O-H stretch in H₂O_{mol} and OH⁻ groups (Figs. 1, 2 and 5). The exact frequency of
232 the contribution B also decreases by $\sim 100\text{ cm}^{-1}$ with increasing the alkali ionic radius (Figs. 2,
233 5). Such effect probably arises from the modification of the inter-oxygen O-O distance around
234 protons (e.g. the distance of O-H \cdots O bonds, with \cdots hydrogen bond) as a function of the ionic
235 radius of the alkali. Indeed, from ¹H NMR spectroscopy of the same alkali silicate glasses, a
236 decrease of the mean O-O distance around protons does occur with increasing the ionic radius
237 of the alkali (Le Losq et al., 2014a). This mean O-O distance is $\sim 287\text{ pm}$ in LS4 glasses and
238 $\sim 255\text{ pm}$ in KS4 glasses (Le Losq et al., 2014a). The frequency of the IR and Raman bands

239 assigned to O-H stretching is related to the inter-oxygen O-O distance around protons
240 (Nakamoto et al., 1955; Novak, 1974; Wall & Hornig, 1965). From the data compilation of
241 Novak (1974), the mean O-O distances of ~287 and ~255 pm correspond to mean O-H stretch
242 frequency of ~3500 and ~2300 cm^{-1} , respectively. These frequencies correspond to the main
243 peaks C and A in the FTIR and Raman spectra of LS4 and KS4 glasses (Figs. 2, 3). Therefore,
244 FTIR and Raman data confirm that increasing the alkali ionic radius in hydrous binary silicate
245 glasses promotes a decrease of O-O distances around protons.

246 An increase of the alkali ionic radius also results in a significant increase of the OH^-
247 $/\text{H}_2\text{O}_{\text{mol}}$ ratio, so that the hydrous KS4 glasses, for example, contain mostly OH^- groups at
248 least up to the maximum total water content of this study (17.6 mol % H_2O). Given that peaks
249 A and B are well expressed in the KS4 glasses (Figs. 1, 2, 3 and 5), it appears that those two
250 peaks arise from O-H stretch in OH^- groups. Therefore, their combination with the Si-O-H
251 stretch vibrational mode located near 910 cm^{-1} (Fig. 2) must give a complex OH^- combination
252 signal. This is indeed observed when superimposing the spectra of the 17.6 mol% glasses
253 (Fig. 5, bottom). As the alkali ionic radius increases, the intensity of peaks A and B increase.
254 As a result, the intensity between 4100 and 4400 cm^{-1} also increases (arrow in Fig. 5, bottom),
255 and the signal between 4000 and 4700 cm^{-1} becomes less and less resolved. Consequently, the
256 combination of the O-H fundamental stretch signals from OH^- groups with the signal assigned
257 to Si-O-H stretching can result in a broad combination signal that extends from 4700 cm^{-1}
258 down to at least 3800 cm^{-1} , and thus becomes mixed with the unattributed ~4000 cm^{-1} peak
259 and with the high-frequency end of the O-H stretching band in the IR spectra (Fig. 5).

260

261 **Implication**

262 The IR combination signal of OH⁻ groups is often reported to be near 4500 cm⁻¹
263 because most of the signals assigned to O-H stretch are located above 3000 cm⁻¹ in many
264 aluminosilicate glasses (Stolper, 1982). However, the data reported here show that IR
265 combination signals used to study the speciation of water in glasses and melts can be complex
266 because of the strong dependence of the O-H stretch frequency on the O-O distances around
267 protons, and hence, on the glass structure. In any study of hydrous silicate glasses, the strong
268 dependence of the proton environment as a function of the type of metallic cation within the
269 silicate network of glasses must be kept in mind, because this dependence affects the IR
270 combination signals and therefore the accuracy of the determination of water speciation with
271 IR spectroscopy. The combination peaks in IR spectra of alkali-rich or alkaline-earth rich
272 natural compositions, such as alkali phonolite and phono-trachyte for instance, may,
273 therefore, be broadened as illustrated in the simple compositions used in this study. However,
274 this is purely speculative at this time and has to be investigated. In all cases, we advise users
275 to check the occurrence and importance of low-frequency (<3000 cm⁻¹) O-H stretching
276 signals in glasses before using the 4500 and 5200 cm⁻¹ combination peaks as proxies for [OH⁻
277] or [H₂O_{mol}]. If those low-frequency O-H stretching signals are significant, any use of the
278 combination peaks to determine [OH⁻] or [H₂O_{mol}] is problematic.

279

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286

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- 379

380 Table 1: Chemical composition of LS4, NS4 and KS4 anhydrous glasses. nom. mol%, nom.
381 wt% and meas. wt% refer respectively to the nominal compositions in mol and wt%, and to
382 the measured composition in wt%. *Lithium has been determined by difference. Errors are
383 given at the 1σ confidence interval.

		SiO₂	Li₂O*	Na₂O	K₂O
LS4	nom. mol%	80.00	20.00	0.00	0.00
	nom. wt%	88.94	11.08	0.00	0.00
	meas. wt%	88.77(94)	11.23(75)	<i>n.a.</i>	<i>n.a.</i>
NS4	nom. mol%	80.0	20.0	0.00	0.00
	nom. wt%	79.50	0.00	20.50	0.00
	meas. wt%	80.38(36)	<i>n.a.</i>	19.62(31)	<i>n.a.</i>
KS4	nom. mol%	80.0	20.0	0.0	0.0
	nom. wt%	71.84	0.00	0.00	28.16
	meas. wt%	71.16(49)	<i>n.a.</i>	<i>n.a.</i>	28.84(28)

384
385

386 Table 2: Water concentration in mol and wt% of glasses. Nom.: nominal water concentration.
 387 I_A , I_B and I_C FTIR absorbance intensities correspond to those observed on Fig. 1 for peaks A,
 388 B and C (as this is for indicative purpose, no error are indicated). Errors are given at the 2σ
 389 confidence interval.

390

Glass	nom.	nom.	Wt% H ₂ O		FTIR absorption, cm ⁻¹		
	mol% H ₂ O	wt% H ₂ O	FTIR	RAMAN	I_A	I_B	I_C
	3.28	1.12	-	0.97(30)	38.6	49.5	76.1
LS4	9.40	3.35	-	3.61(30)	98.3	152.8	277.2
	17.64	6.67	-	6.75(30)	190.0	316.3	603.7
	3.28	1.00	1.21(12)	0.98(30)	96.4	100.0	37.1
NS4	9.40	3.00	2.89(18)	3.25(30)	238.6	264.3	124.8
	17.64	6.00	6.00(29)	5.96(30)	500.5	604.2	383.7
	3.28	0.90		1.04(30)	172.4	166.0	31.8
KS4	9.40	2.71	-	2.76(30)	344.6	345.4	65.1
	17.64	5.44	-	4.84(30)	702.1	729.9	138.2

391

392

393

394 Figure 1: FTIR spectra of LS4 (top), NS4 (middle) and KS4 (bottom) glasses with 0.0, 3.3,
395 9.4 and 17.6 nom. mol% water (from bottom to top in the figures respectively). A, B and C
396 are visual markers for the three observed peaks attributed to the O-H stretch (ν_s O-H) in
397 H_2O_{mol} and OH^- groups. The peak at 1630 cm^{-1} arises from the bending of H_2O_{mol} (ν_B H-O-
398 H). Signal below 1300 cm^{-1} is at saturation and therefore not shown here. Spectra are offset of
399 5 absorbance units for clarity.

400

401 Figure 2: Raman O-H stretching signal in the KS4, NS4 and LS4 glasses with 17.6 nom.
402 mol% H_2O . For comparison, the peak assigned to O-H stretch in a SiO_2 glass with 20 mol%
403 H_2O is also shown (made at 1.5 GPa - 1650°C). A, B and C are visual markers for the three
404 observed peaks (see also Figure 1). Small sharp peaks near 2330 cm^{-1} arise from the N_2
405 atmospheric signal. Insert: signals from the Si-O-Si bending (the 500 cm^{-1} band) and the Si-O
406 stretching modes (the 1100 cm^{-1} band) of the KS4 + 17.6 mol% H_2O are shown (see [Le Losq](#)
407 [et al., 2014b](#); [Mysen, 1990](#); [Mysen & Cody, 2005](#); [Mysen et al., 1982](#) for more details and
408 attributions for silicate species); the signal from the Si-O-H stretch may give rise to the peak
409 near 910 cm^{-1} that is perfectly visible on this spectrum ([Malfait, 2009](#); [Spiekermann et al.,](#)
410 [2012](#); [Zotov & Keppler, 1998](#)).

411

412 Figure 3: Details of the $3800 - 5600\text{ cm}^{-1}$ region of the FTIR spectra of KS4, NS4 and LS4
413 glasses. From bottom to top in the figures: 0.0, 3.3, 9.4, 17.6 nom. mol% H_2O . Spectra are
414 offset of 1 absorbance unit for clarity. The dotted lines indicate the presence of peaks in the
415 spectra (see text).

416

417 Figure 4: Absorbance of the peaks located at ~ 4500 and ~ 5200 cm^{-1} , observed on the spectra
418 of hydrous glasses presented in Figure 3, reported as a function of the nominal water content
419 of glasses. Absorbance was measured at the maximum intensity of peaks. Please note that
420 those values were measured on spectra for which no background subtraction was applied, so
421 that those values cannot be used to retrieve the water concentration of the glasses. Lines are
422 guides for the eyes.

423

424 Figure 5: Top: $2000\text{-}4000$ cm^{-1} portion of spectra of 17.6 mol% H_2O KS4, NS4 and LS4
425 glasses, the A, B and C marker correspond to those on figure 1; Bottom: $3800 - 5600$ cm^{-1}
426 region of same spectra. The arrows are used as visual marker for highlighting the effects
427 produced by an increase of the ionic radius of the alkali in the glasses (see text).

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