1	REVISION 2
2	Alkali influence on the water speciation and the environment of protons in silicate
3	glasses revealed by ¹ H MAS NMR spectroscopy
4	
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11	Keywords: silicate glasses, proton environment, NMR spectroscopy
12	Abstract
13	Water can form different chemical bonds with the ionic entities composing silicate melts.
14	Because of that, its influence on the physico-chemical properties of magmas can vary with
15	silicate composition and water content, temperature, and pressure. To further our
16	understanding of how silicate chemical composition governs proton distribution in magmas,
17	the environment of protons in hydrous alkali (Li, Na, K) silicate glasses was varied as a
18	function of the type of alkali metal and total water content. From ¹ H MAS NMR
19	spectroscopy, H^+ are distributed among five different structural environments in alkali silicate
20	glasses. One of these environments is in the form of H_2O molecules (H_2O_{mol}). The four others
21	are the proton environments associated with Si-OH bonding, and perhaps also with M-OH
22	bonding (with $M = Li$, Na or K). Those environments differ in their O-O distance and extent
23	of hydrogen bonding. H_2O_{mol} species are located in an environment with an O-O distance of
24	~290 pm. OH ⁻ groups are in environments with O-O distances from 240 to 305 pm. The ionic

25 radius of the alkalis, and hence their ionic field strength, determines the fraction of water 26 dissolved as H_2O_{mol} and OH^2 groups, as well as the distribution of protons in the various OH^2 environments. The mean volume of the H⁺ oxygen coordination sphere was calculated using 27 the ¹H⁺ NMR signal intensity and the mean O-O distance around H⁺. Increasing ionic radius 28 29 of the alkali metal in silicate glasses results in a decrease of this mean volume. The partial 30 molar volume of water in the corresponding melts determined through other technics seems to 31 vary in a comparable way. Therefore, the chemical composition of silicate melts may control 32 the partial molar volume of dissolved water because of its influence on the structural 33 environment of protons. This probably also plays a role in determining water solubility.

34

35 Introduction

Water plays a fundamental role in the dynamics of the past and present Earth because of its major impact on the physicochemical properties of silicate materials. As the main volatile component in natural magmas, water profoundly decreases their viscosity (Richet et al., 1996) so that, for instance, it partly controls the behavior of volcanic eruptions. Liquidus phase relations and volume properties of magmas are also influenced by dissolved water (e.g., Kushiro, 1972; Ochs III and Lange, 1999).

Since the work of Wasserburg (1957), it is commonly assumed that water dissolved in
silicate melts breaks their Si-O-Si linkages, as described by:

44
$$H_2O_{mol (melt)} + Si-O-Si_{(melt)} \Leftrightarrow 2 Si-OH_{(melt)}.$$
 (1)

In this suggested solution mechanism, H_2O_{mol} is molecular water dissolved in the melt. The Si-OH bonds are formed through H_2O_{mol} dissociation into OH⁻ groups via reaction with Si-O-Si bonds. A consequence of reaction (1) is depolymerization of the silicate melt structure, which, in turn, may lead to changes in its physical properties. For example, this simple solution mechanism explains qualitatively the decrease of the viscosity of aluminosilicate
melts observed upon water addition (Richet et al., 1996).

51 Infrared absorption studies of hydrous glasses and melts have demonstrated the 52 presence of both OH⁻ groups and H₂O_{mol} species (Davis and Tomozawa, 1996; Efimov and 53 Pogareva, 2006; Malfait, 2009; Scholtze, 1960; Stolper, 1982). Results of recent NMR 54 spectroscopic studies of depolymerized silicate glasses (Cody et al., 2005; Mysen and Cody, 55 2005; Xue and Kanzaki, 2004, 2006, 2007, 2008) have led to the suggestion that water also 56 can react with alkali and alkaline earth, M, network modifying cations to form M-OH bonds in depolymerized melts. Results from thermodynamic modeling of silicate melts (Moretti, 57 58 2005; Moretti et al., 2014) and Raman studies (Le Losq et al., 2013; Mysen and Virgo, 59 1986a,b) reinforce such hypothesis. Notably, whereas formation of T-OH bonds (T=tetrahedrally coordination cations such as Si^{4+} and Al^{3+}) in melts might cause network 60 61 depolymerisation, formation of M-OH bonds may lead to melt polymerization (e.g., Fraser, 62 1977; Moretti, 2005; Moretti et al., 2014; Xue and Kanzaki, 2004, 2006, 2007, 2008).

63 Temperature controls the ratio of OH⁻ to H₂O_{mol} species in silicate melts (Behrens and 64 Yamashita, 2008; Nowak and Behrens, 1995; Shen and Keppler, 1995). The effect of the melt 65 chemistry is less evident. Most studies focus on glasses, which preserve a record of the 66 structure of melts at their glass transition temperature (Tg). The latter depends on water 67 content and on silicate composition (see for a review Mysen and Richet, 2005). It follows 68 that the OH/H₂O_{mol} ratio in glasses is not only affected by the chemical composition of the 69 precursor melts, but also by their T_g because of the temperature control over the water 70 speciation (Ohlhorst et al., 2001). Therefore, chemical effects on water speciation in melts 71 cannot be accomplished when the glasses have different glass transition temperatures. We 72 note, however, that the *in situ* FTIR study at high temperature of Behrens and Yamashita

(2008) suggested that chemical effects on the water speciation could be quite important. Indeed, the equilibrium constant of the H_2O_{mol} dissociation into OH⁻ species ($H_2O + O^{2-} =$ 2OH⁻) in simple sodium silicate melts is quite different from those previously determined *insitu* for rhyolitic aluminosilicate compositions (Nowak and Behrens, 1995; Shen and Keppler, 1995; Sowerby and Keppler, 1999). However, Behrens and Yamashita (2008) emphasized that the existing dataset is too small to provide sufficient information to infer water speciation and its environment in most silicate melts, and, hence, in natural magmas.

80 In order to provide better insights into the chemical effects on the OH/H₂O_{mol} behavior 81 and the proton environment in amorphous silicate materials, we choose to study simple alkali 82 tetrasilicate glasses ($K_2Si_4O_9$, $Na_2Si_4O_9$, $Li_2Si_4O_9$). By varying the type of the alkali metal in 83 the glasses, their effect on the water speciation can be evaluated. The chosen glasses present a 84 significant advantage in their relatively similar glass transition temperatures (difference in the 85 30 K range) and viscosities (see Bockris et al., 1955; Poole, 1949). Therefore, observations 86 made on the glasses will reflect mainly chemical effects rather than a temperature effect, thus 87 minimizing the difficulties previously discussed. To study the proton environment in the 88 glasses, ¹H MAS NMR spectroscopy was used. This technique brings information about the 89 environment of protons as well as the speciation of water in hydrous glasses (see for instance 90 the studies of Eckert et al., 1988; Kohn et al., 1989; Robert et al., 2001; Schaller and Sebald, 91 1995).

- 93 Experimental methods
- 94
- 95 <u>Sample preparation and water determination</u>
- 96 Starting compositions were anhydrous KS4 (K₂Si₄O₉), NS4 (Na₂Si₄O₉) and LS4

97 $(Li_2Si_4O_9)$ glasses, composed of 20 mol% M₂O (with M = Li, Na, K) and 80 mol% SiO₂ 98 (Table 1). The glasses were synthesized by mixing pure anhydrous SiO₂, K₂CO₃, Na₂CO₃ and 99 Li_2CO_3 powders followed by grinding under ethanol for about 1 hour. The oxide + carbonate 100 mixtures were first decarbonated by placing the samples in a high-temperature, ambient-101 pressure furnace in a Pt crucible. In this step, temperature was increased by about 1.5°/min 102 until reaching 1000°C. This decarbonation step was followed by a further temperature 103 increase for melting near 1200°C for the KS4 glass, 1400 °C for the NS4 glass, and 1630°C 104 for the LS4 glass. These final temperatures were governed by the respective liquidus 105 temperatures of the materials (Eppler, 1963; Schairer and Bowen, 1955, 1956). Melts were 106 quenched to glass by placing the bottom of the crucible in liquid H_2O . The resulting glasses 107 were crushed to a powder under ethanol before the preparation of platinum capsules for high-108 pressure, high-temperature experiments. These crushed powders were heated at 400°C for 1 109 hour to ensure complete removal of ethanol and stored at 110°C before loading in capsules for 110 high-temperature/-pressure experiments to avoid reaction with atmospheric water before an 111 experiment. Such effect is particularly critical for the KS4 glass (Schairer and Bowen, 1955). 112 The NS4 and KS4 starting glasses are transparent, and chemically homogeneous (Table 1). 113 Both observations indicate that NS4 and KS4 glasses are not phase-separated (in accordance 114 with the phase relations of these compositions; Kracek, 1932; Schairer and Bowen, 1955, 115 1956). Phase separation and crystallization may be issues for making the LS4 glass, because 116 of the presence of a miscibility gap near 1000°C and of a tendency to crystalize lithium 117 disilicate (Eppler, 1963; Haller et al., 1974; Kracek, 1930). However, by heating this melt at 118 high temperature (>1600 °C) and quenching it at a rate \geq 500°C/s, we obtained a transparent 119 glass, without "milky" appearance typical of phase-separated glasses. Furthermore, analysis of 120 the glass showed it to be chemically homogeneous chemistry.

121 To prepare hydrous glasses from melts at high temperature and pressure, the starting 122 materials were contained in ~ 10 mm long, 5 mm diameter platinum capsules. The appropriate 123 amount of water was added first (from ~ 1 to ~ 12 µl) with a microsyringe (precision ± 0.1 µl). 124 The exact amount of water was determined by weighing the capsule after addition of water. The crushed, anhydrous silicate glass was then added (typically 150 to 180 x 10^{5} g) and the 125 126 capsules were welded shut. A final weighing step was performed to check if capsules were 127 perfectly sealed by weighing them before and after heating to 300°C for 1 hour. The samples 128 were then placed in $\frac{3}{4}$ diameter furnace assemblies based on the design of Kushiro (1976) 129 and subjected to the desired pressure (1.5 GPa) and temperature (1450°-1650°C) for 90 130 minutes in a solid-medium, high-pressure apparatus (Boyd and England, 1960) (see Table 2). 131 Temperatures were measured with type S thermocouples with no correction for pressure on 132 their *emf*, which may be as much as 10°C (Mao et al., 1971). Pressure was calibrated against 133 the melting point of NaCl and the calcite-aragonite transformation (Bohlen, 1984). Estimated 134 uncertainties are $\sim 10^{\circ}$ C and ~ 0.1 GPa, respectively.

135 To determine water concentrations of glasses, several analytical approaches were used. First, the NS4 glasses were analyzed by FTIR spectroscopy with a Jasco[®] IMV-4000 infrared 136 137 multi-channel infrared spectrometer and using the protocol of Yamashita et al. (2008). 138 Samples were initially double-polished using oil at a thickness of around 100 µm. 139 Transmission spectra were recorded with a 10X objective and a $100 \times 100 \ \mu m$ aperture. A 6-140 order polynomial baseline was employed for subtracting the background in the frequency range of interest (4000 - 5800 cm⁻¹) The extinction coefficients provided in Yamashita et al. 141 142 (2008, see their table 4, protocol A) were used to retrieve the concentrations of OH⁻ and H_2O_{mol} species from the integrated areas beneath the ~4500 and ~5200 cm⁻¹ absorption bands. 143 The water concentration of the LS4 and KS4 glasses was determined with Raman 144

145 spectroscopy because molar absorption coefficients are not available for these glass compositions. The Raman spectrometer was a Jasco[®] NRS 3100 spectrometer equipped with 146 a notch filter, holographic gratings, a single monochromator, and a 1024×128 Andor[®] 147 148 DV401-F1 CCD Peltier-cooled at -69°C. Spectra were collected between 350 and 4000 cm⁻¹ with 1200 lines/mm grating. The samples were excited with a 490 nm solid state Coherent[™] 149 laser, focused $\sim 10 \ \mu m$ below the glass surface through a 50x Olympus[®] lens. Laser power at 150 151 the sample was 44 mW. To determine the water concentration of the glasses, a calibration 152 relating the area of the Raman band attributed to the OH stretching signal to the FTIR-based 153 NS4 glass water concentration was implemented, using a protocol similar to that of Le Losq 154 et al. (2012) (see details in Appendix 1). The area of the Raman OH stretching signal was 155 normalized to that of the Raman silicate signals (located below 1500 cm⁻¹), so that, this 156 calibration does not depend on glass chemistry (Le Losq et al., 2012; Shea et al., 2014). 157 Raman-derived water concentrations for LS4 agree with Loss on Ignition (see the protocol 158 described in the next paragraph and Table 2), reinforcing the previous statement. The water 159 concentrations of glasses are shown in Table 2.

160 A final check of water concentrations was carried out by measuring Loss On Ignition 161 (LOI) by using the fine (< 5 microns) powders also employed for NMR experiments. These 162 powders were first dried one night at 180 °C. Then, weigh loss was measured between 180 °C 163 and 1100 °C. Consistent and reproducible results were obtained for the LS4 glasses (Table 2), 164 but not for the NS4 and KS4 glasses. For those glasses, erratic LOI at 180°C, ranging from 165 1.5 up to 5.0 wt%, indicate that the powders adsorbed atmospheric water. LOI were 166 performed several weeks after the NMR measurements (the powders were obtained right before, see below). Despite the precautions adopted to minimize the time they spent in 167 168 atmosphere, water adsorption by the NS4 and KS4 powders in such a long timescale is not surprising regarding their low granulometry (Newman et al., 1986) and the very hygroscopic
nature of NS4 and KS4 glasses (Schairer and Bowen, 1955, 1956).

171 Because of the pronounced hygroscopic nature of the alkali silicate glasses, particularly of 172 the KS4 composition, the maximum time the starting materials were exposed to ambient 173 atmospheric conditions was minimized. Raman and FTIR spectra of hydrous glasses were 174 acquired, therefore, immediately after their synthesis. In addition, Raman spectra were 175 acquired by focusing the laser beam inside the sample (not at the surface). Moreover, the 176 FTIR spectra analyze a large bulk part of samples. Such analysis conditions ensure that the 177 water IR and Raman signals arise from "structural" water, i.e. water dissolved in the structure 178 of the melt at high pressure and temperature prior to quenching to a hydrous glass. The glass 179 chips from the high-pressure experiments were crushed to powder immediately before 180 performing the NMR experiments. Even so, during the loading of rotors and the installation of 181 the probe, the powder remained in contact with the atmosphere for 30-60 minutes. During the 182 experiments, the powdered sample was spinning at 22 kHz in dry air. This ensures optimal conditions for analyzing powders without hydration. As a result, the ¹H NMR signal is 183 184 representative of water quenched in the glass from the melt.

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186 Solid-state Nuclear Magnetic Resonance (NMR) spectroscopy

187 Hydrogen-1 solid-state NMR experiments were carried out with a Chemagnetics CMX 188 Infinity 300 Solid State NMR. The static field strength of the magnet is \sim 7.05 T. The Larmor 189 frequency of ¹H is \sim 300 MHz. Samples were loaded in 2.5mm diameter zirconia rotors, which 190 were placed in a Chemagnetic double resonance probe. Average sample mass was \sim 12 mg. In 191 order to ensure that only the ¹H signal from the sample was detected (i.e., not from outside the 192 RF coil), a 16-cycle DEPTH sequence was used. This pulse sequence employs four 90 ° pulses at select phases and a cycle of 16 pulses to cause ¹H with a nutation of less than $\pi/4$ (e.g., outside the RF coil) to lose coherence. Only ¹H nucleii inside the coil can experience a nutation angle of $\pi/2$ and will retain coherence and contribute to the acquired signal. The ¹H pulse length was 2.5 µs, with a 100 s recycle delay and 1600 acquisitions. The Magic Angle Spinning (MAS) frequency ($\omega_r/2\pi$) was 22 kHz. All spectra were referenced to tetramethylsilane.

In order to determine the number of components needed to account for the variance of data, the principal component analysis (PCA) algorithm of the Scikit-learn Python package was used (Pedregosa et al., 2011).

202

203 **Results**

The ¹H MAS NMR spectra of hydrous LS4 glasses comprise an asymmetric resonance 204 205 peak centered at 3.7 ppm. The asymmetry reflects a shoulder near 15 ppm and a small peak at 206 1.3 ppm (Figure 1, top). The intensity of the main 3.7 ppm peak increases strongly with 207 increasing water concentration. The 1.3 ppm contribution is sharp and intense in the spectrum 208 of LS4 + 3.3 mol% H_2O , and becomes weaker relative to the main 3.7 ppm peak in the 209 spectrum of LS4 + 9.4 mol% H₂O. With 17.6 mol% H₂O in the LS4 glass, only a small 210 shoulder on the right side of the 3.7 ppm main peak is observed. The intensity of the shoulder at ~15 ppm also increases with increasing water concentration. 211

The intensity patterns of the ¹H MAS NMR spectra of the NS4 hydrous glasses differ significantly from those of hydrous LS4 glasses (Figure 1, middle). The main contribution to the spectra of hydrous NS4 glass is located at 15.9 ppm. This peak is asymmetric, suggesting the presence of another contribution near ~13 ppm. Two other peaks are visible at 4.0 and 1.1 ppm, the latter being the sharpest and most symmetric among the peaks in the ¹H MAS NMR spectra of hydrous NS4 glasses. The 15.9 and 4.0 ppm peaks intensities increase with increasing water content, whereas the 1.1 ppm peak intensity remains constant. This latter behavior is similar to that observed in ¹H MAS NMR spectra of the hydrous LS4 glasses. The \sim 4 ppm maximum splits in two shoulders at 3.5 ppm and 4.6 ppm at 17.6 mol% H₂O (Figure 1, middle).

222 In the case of the hydrous KS4 glasses, most of the ¹H MAS signal is near 16.0 ppm in 223 the form of an asymmetric peak somewhat similar to the 16 ppm peak in the spectra of 224 hydrous NS4 glasses (Figure 1, bottom). However, the intensity at frequencies lower than ~ 10 225 ppm is considerably less pronounced in the spectra of hydrous KS4 compared with those of 226 hydrous NS4 (Figure 1, middle and bottom). At all water contents, there is also a small peak 227 at 1.2 ppm. At 9.4 mol% H₂O, there is an intensity increase of the 16.0 ppm peak, whereas that of the 1.2 ppm peak remains constant, and a small shoulder near \sim 5 ppm appears. A 228 229 further increase of water concentration to 17.6 mol% H₂O results in increased intensity of the 230 16.0 ppm peak and \sim 5 ppm shoulder, which forms a small broad peak at this water content.

In the spectra of hydrous LS4 and NS4 glasses, the intensity of sidebands increases with increasing water concentration (Figure 2). At the same water content, sideband intensities are significantly higher in NMR spectra of hydrous LS4 glasses compared with those of hydrous NS4 glasses. There is little or no sideband intensity in the ¹H MAS NMR spectra of hydrous KS4 glass at any water content.

To identify more clearly the spectral changes as a function of the ionic radius of the alkali metal and of water addition, the portions of signal in the three main visible bands located at ~16.0 ppm, ~4.5 ppm and ~1.3 ppm will be considered. To do so, the spectra were divided in three frequency ranges, the "low frequency" (LF) portion between -30 and 1.8 ppm, the "middle frequency" (MF) one between 1.81 and 7.7 ppm, and the "high frequency" (HF) portion from 7.7 up to 30 ppm (HF). When the ionic radius of the alkali increases, the H⁺ contributions in the HF domain increase of ~200% (Figure 3). In parallel, the H⁺ contributions to the MF domain strongly decrease. The area of the LF domain tends to decrease with increasing the ionic radius of alkali. The water concentration has a much smaller effect (in the 1 - 10% range) than the ionic radius of alkali on the fraction of H⁺ giving signals in the LF, MF and HF domain (Figure 3).

247

248 Discussion

249 Observations summarized in figures 1-3 indicate that the ionic radius of alkalis strongly affect the ¹H MAS NMR signal, whereas total water concentration has a smaller 250 251 effect. The presence of peaks near 1.0, 5.0 and 16.0 ppm, and of shoulders near 3.5 and 12.0 252 ppm, indicate that four and perhaps five spectral components, centered at the corresponding frequency, may be necessary to describe the entire ¹H MAS NMR spectra of the glasses 253 254 (Figure 1). Analysis of the dataset with principal component analysis PCA shows that 4 255 components account for 99.64% of the variance of the data, and 5 for 99.85 (Table 3). Therefore, the visual observation and statistical treatment of ¹H NMR spectra indicate the 256 257 presence of at least five components near 1.0, 3.5, 5.0, 12.0 and 16.0 ppm. The type of alkali 258 metal in the glass affects the relative intensities of those components but not their frequencies 259 (Figure 1). Therefore, the effect of changing the ionic radius of the alkali metal in the silicate glasses is to change the distribution of H⁺ among five structural environments. 260

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262 <u>Origin of the ¹H NMR signals</u>

263 Sideband analysis aids in the structural interpretation of the NMR spectra. To obtain a 264 quantitative isotropic ¹H NMR signal, the ¹H-¹H dipolar coupling was averaged out by using

265 magic angle spinning. Any residual dipolar coupling will be manifested by spinning side 266 bands at $\pm n\omega_r/2\pi$ (n=1, 2, ...). For instance, for H₂O_{mol} entities in a rigid environment (e.g., immobile H₂O molecules), the ¹H-¹H dipolar coupling is on the order of 36 KHz. Therefore, 267 as we spin at 22 kHz, a signal from rigid H₂O_{mol} entities would be expected. However, 268 269 molecular H₂O_{mol} at room temperature in glasses have slight motions (see Eckert and 270 Yesinowski, 1987) that will reduce the intensity of sidebands, because such motions contribute to average out the ¹H-¹H dipolar coupling. Isolated OH⁻ groups will not contribute 271 272 to sideband intensity. In a rigid environment, the sideband signal from OH⁻ groups is governed by the ¹H chemical shift anisotropy together with homonuclear and heteronuclear 273 274 interactions. It is negligible compared to the sideband signal from H_2O_{mol} species (Eckert et al., 1988). Consequently, the intensity of sidebands in the ¹H MAS NMR spectra of hydrous 275 276 alkali silicate glasses mostly reflects the presence and concentration of H₂O_{mol} species (Figure 277 2).

Cody et al. (2005) demonstrated that the sidebands in ¹H MAS NMR spectra of 278 279 hydrous SiO₂ and NS4 glasses primarily are correlated with a central peak near 6.4 ppm, 280 which would suggest H_2O_{mol} . Kohn et al. (1989) also attributed a peak near 5 ppm in ¹H MAS 281 NMR spectra of silica and silicate glasses to H₂O_{mol} species. Additional support for such an assignment is provided in the 2D {¹H}-²⁹Si HETCOR experiment on hydrous NS4 glasses of 282 Robert et al. (2001), who noticed missing intensity near 4 ppm compared to 1D ¹H MAS 283 284 NMR experiments. This means that some protons are not cross-polarized with silicon and, 285 thus, would be in molecules of water (Robert et al., 2001). In the spectra of the NS4 glasses 286 (Figure 1), the shoulder near 5 ppm appears with increasing the water concentration. This 287 development is accompanied by increasing sideband intensities (Figure 2). As a result, and in 288 agreement with the studies of Cody et al. (2005), Kohn et al. (1989), the 5 ppm contribution in the spectra reported here likely is a record of signals of H^+ in H_2O_{mol} species. According to Cody et al. (2005), Kohn et al. (1989), Robert et al. (2001) and Xue and Kanzaki (2004), the 3.5 ppm contribution might arise mainly from signal of H^+ in Si-OH bonds, forming weak

292 hydrogen bonds with their environments.

The sideband intensities in the ¹H MAS NMR spectra of hydrous KS4 glasses are 293 294 negligible, whereas the 12 and 16 ppm contributions are particularly strong (Figures 1, 2). Therefore, these two contributions likely reflect H⁺ in OH⁻ groups bonded to cationic entities 295 in the glass structure. The {¹H}-²⁹Si HETCOR experiments of Robert et al. (2001) performed 296 297 on an NS4 hydrous glass (1.28 wt% H_2O) are consistent with the 16, 12 and 3.5 ppm ¹H NMR lines all correlated with $^{29}\mathrm{Si}$ NMR signals of Q^3 and Q^2 tetrahedral species. Formation of Q^2 298 299 and Q^3 species in sodium silicate glasses is observed upon hydration (Cody et al., 2005; 300 Mysen and Cody, 2005; Zotov and Keppler, 1998). According to those observations, the 12 301 and 16 ppm contributions record signals from OH⁻ bonded to Si, forming the so-called Si-OH 302 groups. However, Cody et al. (2005) also found that the intensity of the 12 ppm peak 303 increases with increasing Na/Si ratio of the glass, which might not be consistent with an 304 assignment of this peak to only Si-OH groups. Furthermore, for the sodium disilicate glass (Na/Si = 1), the H_2O_{mol}/OH^- ratio predicted by using the distribution of Q^n species distribution 305 was too high compared with that inferred from the ¹H NMR spectrum. Combining this 306 observation with considerations about the relaxation times of Q^n species, Cody et al. (2005) 307 308 concluded that at high Na/Si, some Na-OH groups must exist and may result in a signal close 309 to 12 ppm. Therefore, we cannot rule out that some of the signal that contribute to the 12 ppm 310 peak in KS4, NS4, and LS4 hydrous glasses may reflect alkali-OH bonding. This assignment 311 does not disagree with the HETCOR experiment of Robert et al. (2001), because those experiments probe the proximity of Q^n species to protons and not their link. However, the 312

correlation implies that M-OH groups interact with Qⁿ species, through hydrogen bonding for
 example.

315 The ¹H NMR signal near 1.0-1.7 ppm may arise from "free" M-OH groups, i.e. which 316 do not form hydrogen bonds with their environment, according to the work of Xue (2009), 317 Xue and Kanzaki (2004), and Xue and Kanzaki (2009). However, the assignment of the peak 318 near 1 ppm is not totally clear. For example, Si-OH groups without hydrogen bonding can 319 give signals near 1 - 2 ppm, as shown by proton NMR of silica gels (see Bronnimann et al., 320 1988 for instance). Xue and Kanzaki (2004) ruled out such an Si-OH attribution because of 321 the absence of a clear peak near 1 ppm in the Suprasil silica glass, which contains low water concentration (~1200 ppm, and therefore mostly dissolved as OH⁻ group, see the ¹H MAS 322 323 NMR spectra in Kohn et al., 1989). However, without more information, the precise assignment of the ~ 1 ppm peak remains uncertain. But whatever the X cation (X = Si or M) 324 involved in the X-OH bond giving rise to this peak may be, existing data indicate that the OH 325 326 do not form hydrogen bonds with their environment.

327

328 <u> $O^{\cdots}O$ distances of the various H⁺ environments</u>

The ¹H NMR chemical shift in hydrous oxide systems is predominantly dependent on 329 330 the O-H distance, and thus on the O-O distance along the O-H^{...}O bond (with ^{...} an hydrogen 331 bond; Berglund and Vaughan, 1980; Eckert et al., 1988; Xue and Kanzaki, 2001, 2009). Indeed, hydrogen bonding pulls the ${}^{1}\text{H}^{+}$ away from the bonding oxygen, which results in 332 deshielding of the ¹H and in an increase of its resonant frequency. It follows that the NMR 333 334 peaks near 1.0, 3.5, 5.0, 12.0 and 16.0 ppm reflect five environments in the glasses with 335 different mean O-O distances. These distances can be estimated with an uncertainty near ± 10 336 pm by using the equation from Eckert et al. (1988) (Figure 4).

337	The 3	.5 ppm Si-OH signal is related to an environment with an O-O mean distance of
338	~295 pm, a	value similar to that the ${\sim}290~\text{pm}$ O-O distance of the $\mathrm{H_2O_{mol}}$ species
339	environment.	According to the models of Xue and Kanzaki (2004), Si-OH groups in different
340	tetrahedral u	nits and connected between them through hydrogen bonding (Si-OH ··· O(H)Si
341	entities) are g	good candidates for the 295 pm Si-OH environment. The correlation between the
342	¹ H signal nea	ar 3 ppm and that of Q^4 units in the HETCOR data of Robert et al. (2001) also
343	indicates that	t the occurrence of Si-OH BO bonds is plausible (with BO a bridging oxygen
344	connecting ty	wo SiO ₄ tetrahedra). Finally, the formation of Si-OH ^{\dots} OH ₂ (OH ₂ being a water
345	molecule) car	nnot be excluded. Figure 4 summarizes all the proposed attributions.
346	The C	OH ⁻ groups in the 261 and 250 pm environments reflect O-O distances that are
347	shorter or ne	arly equal to that of the typical anhydrous O-O distances in LS4/NS4 glasses
348	(262 pm, Ispa	as et al., 2010). Those two environments can record signals from:
349	(i)	Si-OH groups forming "intratetrahedral" hydrogen bonds. Kohn et al. (1989)
350		proposed this hypothesis to account for the presence of the ¹ H NMR signal
351		between 11 and 17 ppm in hydrous silicate glasses. As the distance between
352		O species on a same tetrahedral edge in a single tetrahedron is close to 264
353		pm, the 261 pm environment may be a good candidate for this attribution.
354		Such "intratetrahedral" hydrogen bonds are observed in crystalline hydrous
355		sodium disilicate (Ai et al., 2002, 2003)
356	(ii)	Si-OH groups connected to other tetrahedral units through strong hydrogen
357		bonds with NBO species (Si-OH O(M)Si bonds), because of the increasing
358		hydrogen bonding strength accompanying introduction of nonbridging
359		oxygen in glasses (Xue and Kanzaki, 2004).
360	(iii)	Si-O[H]-M groups, where the Si-O-M bonds are protonated. Here the protons

361 assist in the electronic balance of NBO when the O-M distance is high (in362 the case of K-bearing silicate glasses for instance).

363 (iv) M-OH groups, as suggested by Cody et al. (2005) for the 261 pm environment.
364 Those M-OH groups must form short hydrogen bonds with surrounding
365 oxygens (M-OH^{...}NBO/BO linkages).

366 It is difficult to discriminate between those hypotheses, and a combination of all of 367 them is possible. For instance, the formation of M-OH^{...}NBO/BO bonds, which could result in 368 a signal near 12 ppm as suggested by the data of Cody et al. (2005), may lead to the Q^{n} -H 369 interactions observed in the HETCOR spectra of Robert et al. (2001). At the same time, the 370 Qⁿ distribution in hydrous NS4 glasses also indicates that OH⁻ groups are mostly involved in 371 Si-OH bonds in those glasses (see also Cody et al., 2005; Mysen and Cody, 2005; Zotov and Keppler, 1998). In addition to that, the ²³Na 1D MAS NMR and ${}^{1}H{}^{-23}Na$ REDOR NMR 372 373 experiments of Robert et al. (2001) indicate that the environment of Na in NS4 hydrous 374 glasses does not seem to vary with water content. This interpretation does not support the 375 presence of significant amount of Si-O[H]-M protonated bonds, or of M-OH bonds. 376 Therefore, in the compositions examined in the present study, the high-frequency 12 and 16 ppm ¹H signals may arise mostly from Si-OH bonds, with the H⁺ forming strong inter- or 377 378 intra-tetrahedral hydrogen bonds with the surrounding oxygen atoms.

379

380 *Origin of the chemical effects on the H⁺ environments*

From the data and discussion above (Figures 1, 2, 3, 4), decreasing ionic field strength of the alkali cation, which acts as network modifiers in silicate glasses, promotes the dissociation of H_2O_{mol} into H^+ and OH^- . Those ionic entities react with the Si-O-Si and Si-O-M bonds to form Si-OH bonds, with the H^+ forming stronger hydrogen bonds the greater the 385 ionic radius of alkali. The formation of M-OH bonds cannot be ruled out and could contribute 386 to the ¹H NMR signals near 12 or 1 ppm, but, according to available data, probably is of minor importance at the silica concentration we studied. In any case, ¹H MAS NMR data 387 388 indicate that the ionic radius of alkalis, and hence their ionic field strength, controls the 389 dissociation constant of H₂O_{mol} species as well as the proton environment. The origin of such 390 control can arise from different effects, but present data do not shed light on it. We can only 391 propose hypothesis about it. For instance, it may arise from electronic/steric effects around alkali and along the M-O-Si bonds that control the H2Omol dissociation constant as well as the 392 strength of hydrogen bonding. It may also arise from the tendency to form H₂O_{mol} hydration 393 394 shells around small alkalis (explaining the high H₂O_{mol} concentration in LS4 compositions for 395 instance; such effect is observed in hydrous solutions, see for instance Mähler and Persson, 396 2012). Simple volume effects can also be proposed: the higher the alkali radius, the lower the 397 volume available for protons, this promoting H₂O_{mol} dissociation in OH⁻ groups that can 398 occupy environments with smaller O-O distances.

399

400 **Implications**

401 The influence of the alkali radius on the mean O-O distance around protons (Figure 4) 402 implies significant modification of the mean volume of water molecules. Assuming that the 403 O-O distances around protons in glasses represent the diameters of the oxygen coordination 404 sphere of protons, the mean volume of this sphere can be calculated using the sum of the O-O distances pondered by their associated relative ¹H NMR intensity (hence proton population). 405 406 For a given glass composition, this mean volume does not vary significantly with water concentrations, because the water concentration does not change the ¹H NMR signal 407 408 significantly (Figure 3). In contrast to that conclusion, increasing the ionic radius of alkali 409 produces a decrease of the mean volume of the proton coordination sphere (Figure 5). There 410 is, for example, a 15.8 % volume difference between the NS4 and KS4 glasses. The partial 411 molar volumes of H₂O in the NS4 and KS4 melts, determined by Mysen and Cody (2004) by 412 using solubility measurements and the Gibbs-free energy of solution of H₂O (Figure 5), differ 413 by 16.9 %. Considering the errors affecting the determinations of the partial molar volume of 414 H_2O and those of the mean volumes of the protons coordination sphere calculated from ${}^{1}H$ 415 NMR spectra, those variations are in very good agreement. Therefore, it is suggested that the 416 change of the distribution of H⁺ in various structural environments explains the observed 417 decrease of the partial molar volume of H₂O that occurs upon increasing the ionic radius of 418 alkalis in silicate melts. As a consequence, it appears that the effects of the ionic radius of 419 alkalis on the water speciation and on the proton environment may govern the variation of the 420 partial molar volume of H₂O in silicate melts. Such relationships will impact the buoyancy of 421 magmas, and also the solubility of water. Indeed, the water solubility is systematically higher 422 in NS4 melts than in KS4 melts (Mysen, 1998; Mysen and Cody, 2004). This leads to the 423 suggestion that, if larger structural sites are available for storing H₂O_{mol} and OH⁻ species, the 424 water solubility will be higher.

In summary, this study demonstrates that the chemical composition of silicate melts and glasses can play an important role in defining the speciation of water and the proton environment. This results in variations in the solubility and partial molar volume of water in magmas. Such a complex relationship will play a critical role in determining the degassing path and the density of a rising magma, and hence will affect in various ways the behavior of volcanic eruptions.

431

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- 440

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614 Appendix 1 – Water analysis with Raman spectroscopy

615 In order to determine the water concentration of glasses and to avoid analyzing 616 powders, which quickly adsorb water from ambient atmosphere, we treated the Raman spectra 617 of glasses with a protocol similar to that described in Le Losq et al. (2012). Slight 618 modifications were made to account for the slightly different OH stretching signal in Raman 619 spectra of silicate glasses compared to that of the aluminosilicate glasses analyzed in Le Losq et al. (2012). The OH stretch signal below 3000 cm⁻¹ in silicate glasses is intense and cannot 620 621 be neglected, as Le Losq et al. (2012) did for aluminosilicate glasses. We modified the 622 baseline definition accordingly, as shown in Figure S1. The corrected spectra were normalized to their total area between 350 and 4000 cm⁻¹. The area A_w under the OH-623 624 stretching signal between 2000 and 4000 cm⁻¹ was measured (Figure S1). Such measurement 625 on area-normalized spectra is equivalent to measuring the ratio between the area of silicate 626 bands and that of the OH stretch band on non-normalized spectra. The water concentrations in 627 NS4 glasses determined through FTIR analysis have been used to standardize the A_w vs. total 628 water concentration relationship. The remaining question was whether Raman cross-sections of the different OH signals observed near 3600, 2800, and even 2300 cm⁻¹, in silicate glasses 629 630 (Figure S1) are equal. The cross-section of the Raman signals from individual species in the 3000-4000 cm⁻¹ range is independent of glass chemistry (Le Losq et al., 2012). The good 631 632 agreement between the Raman-measured water concentrations, the nominal ones, and those 633 determined by Loss On Ignition (Table 2) indicates that the O-H stretching Raman cross-634 section is not significantly dependent on the type of alkali element present in the glass. We 635 conclude, therefore, that the OH stretching Raman cross-sections are not affected by the glass 636 chemical composition. This maybe results from the fact that the OH stretching frequency is 637 mostly controlled by one parameter, the O⁻⁻O distances, and in agreement the OH stretching Raman cross-section at a given frequency must mostly depend on this parameter in silicate materials. The values in Table 2 are affected by a global error estimated from deviations between the NS4 FTIR and Raman values and from numerical errors inherent to the baseline subtraction and the calculation of OH stretching band area. Because we do not have a large quantity of standards for performing the calibration and refining the standard deviation of the technic, we preferred providing an error representing an upper limit, and equal to 0.30 wt% at the 2σ confidence interval.

646	Table 1: Chemical composition of LS4, NS4 and KS4 anhydrous glasses. nom. mol%, nom.
647	wt% and meas. wt% refer respectively to the nominal compositions in mol and wt%, and to
648	the measured composition in wt%. Chemical analysis were conducted with a JEOL FE-SEM
649	equipped with an energy dispersive spectrometer (EDS), and were acquired on 25 x 25
650	microns areas, with a 15 kV and 1.04 nA current. Errors are given at the σ confidence
651	interval. *Lithium has been determined by difference.

	SiO ₂	Li ₂ O*	Na ₂ O	K ₂ O
nom. mol%	80.00	20.00	0.00	0.00
nom. wt%	88.94	11.08	0.00	0.00
meas. wt%	88.77(94)	11.23(75)	n.a.	n.a.
nom. mol%	80.0	20.0	0.00	0.00
nom. wt%	79.50	0.00	20.50	0.00
meas. wt%	80.38(36)	n.a.	19.62(31)	n.a.
nom. mol%	80.0	20.0	0.0	0.0
nom. wt%	71.84	0.00	0.00	28.16
meas. wt%	71.16(49)	n.a.	n.a.	28.84(28)
	nom. mol% nom. wt% meas. wt% nom. mol% nom. wt% nom. mol% nom. mol% meas. wt%	SiO2 nom. mol% 80.00 nom. mol% 88.94 meas. wt% 88.77(94) nom. mol% 80.0 nom. mol% 80.0 nom. wt% 79.50 meas. wt% 80.38(36) nom. mol% 80.0 nom. mol% 80.1 nom. mol% 80.38(36) nom. wt% 71.84 meas. wt% 71.16(49)	SiO2 Li2O* nom. mol% 80.00 20.00 nom. wt% 88.94 11.08 meas. wt% 88.77(94) 11.23(75) nom. mol% 80.0 20.0 nom. mol% 80.0 20.0 nom. mol% 80.0 20.0 nom. mol% 80.38(36) n.a. nom. mol% 80.0 20.0 nom. mol% 80.0 20.0 nom. mol% 80.0 20.0 nom. mol% 80.0 20.0 nom. wt% 71.84 0.00 meas. wt% 71.16(49) n.a.	SiO2 Li2O* Na2O nom. mol% 80.00 20.00 0.00 nom. mol% 88.94 11.08 0.00 meas. wt% 88.77(94) 11.23(75) n.a. nom. mol% 80.0 20.0 0.00 nom. mol% 80.0 20.0 0.00 nom. mol% 80.0 20.0 0.00 nom. mol% 80.38(36) n.a. 19.62(31) nom. mol% 80.0 20.0 0.0 nom. wt% 71.84 0.00 0.00 meas. wt% 71.16(49) n.a. n.a.

Table 2: Water concentration in the glasses in mol% and wt%. Nom.: nominal water concentration. Raman values are given by the calibration determined with the help of the NS4 FTIR water concentrations (Appendix 1). LOI: Loss on ignition (see discussion about them in text). Errors from the Raman-based analyses include the calibration error determined on NS4 glasses and the errors resulting from area determination. Errors are given at the 2σ confidence interval.

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Class	nom.	nom. T° of		Pressure	Wt% H ₂ O		
Glass	mol% H ₂ O	wt% H ₂ O	synthesis (°C)	(GPa)	FTIR	RAMAN	LOI
	3.28	1.12	1650	1.5	-	1.05(30)	1.45(53)
LS4	9.40	3.35	1650	1.5	-	3.61(30)	3.08(22)
	17.64	6.67	1650	1.5	-	7.15(30)	6.89(7)
	3.28	1.00	1450	1.5	1.21(12)	1.05(30)	-
NS4	9.40	3.00	1450	1.5	2.89(18)	3.02(30)	-
	17.64	6.00	1450	1.5	6.02(29)	5.96(30)	-
	3.28	0.90	1550	1.5		1.10(30)	-
KS4	9.40	2.71	1550	1.5	-	2.78(30)	-
	17.64	5.44	1550	1.5	-	4.85(30)	-

- Table 3: Explained fraction of variances as a function of the number of components from the
- 662 PCA analysis of data.

Component n°	Fraction of variance
1	0.59232
2	0.98046
3	0.98984
4	0.99641
5	0.99853
6	0.99967
7	0.99990
8	1.00000

Figure 1: ¹H MAS NMR spectra of LS4 (top), NS4 (middle), and KS4 (bottom) glasses with
3.3, 9.4 and 17.6 mol% water. Intensities are normalized to the total area of spectra and to the
water concentrations.

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Figure 2: The low frequency spinning side bands ¹H MAS NMR spectra of the LS4, NS4 and KS4 glasses. One observes relatively strong side bands for Li, less strong for Na, and very weak for K. ¹H-¹H dipolar coupling mostly controls the side band intensity. In general, such

- 671 dipolar coupling is greatest for H_2O^{mol} and would be non-existent for isolated OH⁻ species.
- 672

Figure 3: Relative proportion of H^+ in the low-frequency LF (dark, signals < 1.8 ppm, see the example ¹H NMR spectrum in the insert), the middle-frequency MF (blue, between 1.8 and 7.7 ppm) and the high-frequency HF (red, signals > 7.7 ppm) spectral regions as a function of the ionic radius of alkali ions (ionic radius values for a six-fold coordination from Whittaker and Muntus, 1970). 3, 9 and 17 values are the mol% water concentrations of glasses truncated of their decimal values.

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Figure 4: Representation of the LS4, NS4, KS4 + 17.6 mol% H₂O ¹H spectra as a function of 680 681 the O-O distances, obtained by converting from the ppm scale with the equation from Eckert 682 et al. (1988). The square, diamond, star, hexagon and circle markers are visual markers of the 683 five environments of H^+ . The legend at the bottom right summarizes the assignments 684 discussed in the text. " represents hydrogen bonds (HB). Intratet. HB refers to intratetrahedral 685 hydrogen bonding. BO is for bridging oxygen (Si-O-Si bonds). O(M)Si and O(H)Si represent 686 non bridging oxygen atoms involved in Si-O-M or Si-O-H bonds respectively; such notation 687 is adopted for clarity purpose.

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Figure 5: Variation of the mean volume of the H^+ oxygen coordination sphere calculated from ¹H MAS NMR spectra (see text), and of the partial molar volume Vm_{H2O} of water at 1100°C determined from solubility measurements (Mysen and Armstrong, 2002; Mysen and Cody, 2004). Error bars represent the standard deviation of the values calculated at different water

693 concentrations.

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696 Figure S1: Raman spectra of the NS4 + 17.6 mol% H₂O glass. A) Raw spectrum (dark) with the spline cubic baseline (blue) constrained at 1400, 2000 and 3900 cm⁻¹; B) Baseline-697 698 subtracted spectrum. Peaks at ~490 and ~1100 results from intertetrahedral bending and intratetrahedral stretching of SiO₄ units respectively, while the \sim 800 cm⁻¹ band results from 699 SiO₄ rocking motions. The sharp band near 1630 cm⁻¹ results from H₂O_{mol} bending, and the 700 2300, 2800 and 3600 cm⁻¹ bands results from OH stretching in OH and H₂O_{mol} units in the 701 702 glass (Zotov and Keppler, 1998). See also Le Losq et al. (2012), Le Losq et al. (2014) and 703 references therein for details on Raman spectra of anhydrous and hydrous glasses. The 704 spectrum presented in B) is used for measuring the area of the signal assigned to OH 705 stretching (in red). Intensities of the spectrum have been previously normalized to the total 706 area of the spectrum, which is now equal to 1 (theoretical value of the sum of the vibrational 707 density of state).









