1 **"Revision 1"**

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Phase diagram and P-V-T equation of state of Al-bearing seifertite at lowermost mantle conditions D. Andrault ¹, R.G. Trønnes ², Z. Konôpková ³, W. Morgenroth ⁴, H.-P. Liermann ³, G. Morard ⁵ and M. Mezouar ⁶ ¹ Laboratoire Magmas et Volcans (LMV-OPGC), Université B. Pascal, Clermont-Ferrand, France ² Natural History Museum and Centre for Earth Evolution and Dynamics, University of Oslo, Norway

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17 Abstract

We investigated the properties of Al-bearing SiO₂ (with 4 or 6 wt% Al₂O₃) at pressures and temperatures corresponding to the lowermost mantle, using laser-heated diamond anvil cell coupled with synchrotron-based in-situ X-ray diffraction. The phase transition from CaCl₂-structured to α -PbO₂-structured (seifertite) polymorphs occurs between 113 and 119 GPa at 2500 K. The range of pressure where the two phases coexist is small. There is a slight decrease of the transition pressure with increasing Al-content. We propose a tentative phase diagram reporting the minerals composition as a function of pressure in the SiO₂-Al₂O₃ system.

25 We also refine the P-V-T equation of state of Al-bearing seifertite based on volume

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measurements up to more than 160 GPa and 4000 K (V_0 = 92.73(10) Å³, K_0 = 304.2(3.0) GPa, K'₀=4.59 (fixed), Θ_{D0} =1130 K (fixed), γ_0 =1.61 (3)). At 300 K, the volume decrease at the CaCl₂ to α -PbO₂ transition is 0.5(1) %, a value slightly lower than the 0.6% reported previously for Al-free samples. At high temperature, the Grüneisen parameter of seifertite is found to be similar to that of stishovite. Nevertheless, the $\Delta V/V$ across the CaCl₂-form to seifertite transition is found to increase slightly with increasing temperature.

Across the phase transition, volume changes can be translated into density changes only when the Al substitution mechanisms in both CaCl₂-form and seifertite are defined. The analysis of all available data sets suggests different substitution mechanisms for the two SiO₂ polymorphs. Alsubstitution could occur via O-vacancies in the CaCl₂-form and via extra interstitial Al in seifertite. This would result in a density increase of 2.2(3) % at 300 K for SiO₂ in basaltic lithologies. Alternatively, the same Al-substitution mechanism in both of the SiO₂-dominated phases would yield a density increase of 0.5(1) %.

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41 **Keywords:** Seifertite, phase transition in SiO₂, P-V-T equation of state, lowermost Earth mantle.

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43 Introduction

Seismic tomography suggests that the slab behavior depends primarily on the original plate 44 velocity, the penetration angle of the slab into the mantle and the age of the subducted crust. In 45 several regions the slabs appear to flatten and be delayed near the 660 km discontinuity. In most 46 47 regions, however, slabs eventually descend through the lower mantle (e.g. (van der Hilst et al., 2007; van der Meer et al., 2010)). The sinking of subducting slabs in the deep mantle is controlled 48 primarily by the density difference between the slab materials and the ambient mantle, while the 49 50 mantle viscosity governs the speed of the descent. Detailed tomographic mapping of high-velocity regions in the mantle and comparison with surface paleogeography indicate that most slab remnants 51

sink through the entire mantle in about 250 Ma, i.e. with an average speed of about 1.2 cm/year
(van der Meer et al., 2010).

The slab density is partly related to the slab temperature but also to the intrinsic density of the 54 55 various minerals and their relative proportions. Older and relatively colder slabs are expected to penetrate the deep mantle more easily. Still, a major complication is the layered structure of the 56 57 subducted slab, including its sedimentary, basaltic and harzburgitic components. The density inversion between peridotite and basaltic material in the uppermost 50-100 km of the lower mantle 58 59 caused by gradual transformation of majoritic garnet (e.g. (Ringwood, 1991)) led to the suggestion of partial segregation (flotation) of garnetite near the 660 km discontinuity. The low temperature 60 and high viscosity of slab material, however, make such a lithological separation unlikely. Another 61 possible site of basalt-peridotite separation is the core-mantle boundary (CMB) region. Sinking slab 62 material deflected laterally along the CMB thermal boundary layer will be heated considerably, 63 resulting in reduced viscosity (e.g. (Steinberger and Calderwood, 2006)). Even moderate basalt-64 peridotite density contrasts in the D" zone could therefore lead to such separation. 65

Difference in chemical composition between basaltic and peridotitic materials yields largely 66 different mineral assemblages and mineral compositions at deep mantle P-T conditions. Peridotites 67 varying from fertile lherzolite to depleted harzburgite are dominated by Mg-perovskite (pv, 70-68 69 80%) with minor amounts of ferropericlase (15-25%) and Ca-perovskite (Ca-pv, 0-10%). Basalts have 3-10 times higher concentrations of Al, Ti, Ca and Na and more than 5 times lower 70 71 concentration of Mg than peridotite. The resulting lower mantle basaltic mineralogy has no ferropericlase and a low proportion of Mg-perovskite (or post-perovskite) with high Fe/Mg-ratio. 72 73 Over-abundance of silica and alumina produces the separate silica-dominated (Irifune et al., 1994) and Al-rich phases (NAL and Ca-ferrite phases, (Akaogi et al., 1999; Miyajima et al., 1999)). At 74 intermediate lower mantle depths, common basalts crystallize 15-20% CaCl₂-structured silica, 15-75 25% Ca-ferrite, 35-40% Mg-perovskite and 20-30% Ca-perovskite (Hirose et al., 1999). It is 76 77 generally estimated that the Fe-rich Mg-perovskite makes basaltic material denser than peridotite throughout the lower mantle below 720 km depth (e.g. (Guignot and Andrault, 2004; Hirose et al., 2005; Ricolleau et al., 2010)). Partly due to its elevated bulk modulus, the CaCl₂-structured silica phase becomes the least dense phase of the subducted slabs at 70-120 GPa, with a density difference of more than 3 % relative to the most abundant Mg-rich silicate perovskite.

At ~2500 K, pure SiO₂ was reported to undergo a first-order phase transformation from the 82 83 CaCl₂-form to seifertite (adopting the α -PbO₂-structure (El Goresy et al., 2008)) at ~120 GPa, with a dp/dT-slope of about 10 MPa/K (Murakami et al., 2003). In parallel, (Hirose et al., 2005) showed 84 85 that the silica-dominated phases in a basaltic composition display considerable Al_2O_3 solubility. The presence of Al₂O₃ produces a considerable pressure reduction of the stishovite to CaCl₂-phase 86 87 transition, from ~50 GPa for pure SiO₂ to ~25 GPa for Al-bearing silica (Bolfan-Casanova et al., 2009; Lakshtanov et al., 2007b). On the other hand, the effect of Al_2O_3 on the CaCl₂-form to 88 seifertite transformation and on the density change associated with this transition remains poorly 89 90 constrained. Our objective here is to further evaluate the phase and density relations in the binary silica-alumina system at the P-T conditions of the lowermost mantle. 91

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93 Experimental methods

94 Silica glasses with 4 and 6 wt% Al_2O_3 were prepared in Pt-crucibles at 1650-1670 °C in air, using a vertical furnace at Centre for Material Science, University of Oslo and a Nabertherm 95 furnace at Bayerisches Geoinstitut, Bayreuth, Germany. Tiny air bubbles persisted even after 3-4 96 97 repetitions of re-grinding and re-melting. Electron probe microanalyses yield Al-contents of 4.1(4) 98 and 5.8(3) for glasses Al-4% and Al-6%, respectively. Relatively high values of standard deviations 99 denote persistence of some chemical heterogeneity in the glass. Our analytical precision yielded element totals of 99.7(2) and 99.5(2) for more than 20 measurements on glasses Al-4% and Al-6%, 100 respectively. The glasses were finely powdered and mixed with 10-15 wt% Pt-powder which served 101 as a YAG laser absorber and a pressure indicator at high temperatures. Overall, our samples should 102 103 be homogeneous on a 2-3 micrometer scale, which corresponds to the size of our X-ray probe.

104 Pressures up to 120 GPa were generated by membrane-type diamond anvil cells (DAC) 105 equipped with beveled diamonds of 150/300 or 100/300 µm diameters. Rhenium gaskets were preindented to $\sim 25 \,\mu m$ thickness before laser drilling or spark erosion of 40-60 μm diameters holes. 106 107 Sample pellets were loaded between two pieces of dry NaCl or KCl, which served as pressure 108 transmitting medium and thermal insulator from the diamonds. Temperatures up to more than 4000 K were generated by two NIR Ytterbium fiber lasers. The size of the hot spot on the sample was 109 kept between 15 and 25 µm by adjusting the distance between the focusing-lenses and the sample. 110 Using the spectral radiometry method, precision in temperature determination is better than 50 K. 111 This value is also a maximum for temperature fluctuations at high temperature. Because phase 112 transformations between dense silica polymorphs are known to be sluggish and depend on the 113 starting material and pressure-temperature history (Prakapenka et al., 2004), we heated for 20-40 114 115 min at each temperature point.

Pressures at 300 K were determined based on the equation of state (EoS) of platinum (Holmes et 116 al., 1989). It enables quantitative comparisons with previous reports on the formation of seifertite 117 (mostly (Hirose et al., 2005; Murakami et al., 2003)), in which Pt (Holmes et al., 1989) or Au 118 (Tsuchiya, 2003) EoS were used. Pressures derived from Pt and Au were reported to be fully 119 compatible between each other (Hirose et al., 2005). Alternatively, we used the NaCl-B2 EoS which 120 121 has been calibrated against Pt (Sata et al., 2002). NaCl provides accurate pressure measurements after annealing at low temperature, when the Pt diffraction peaks remain broad due to poor release 122 of the deviatoric stresses in the silica glass. At high temperature, the Pt that is intimately mixed with 123 the starting material was used as a reliable sample-pressure indicator. We observe that the pressures 124 measured at high temperatures are consistent with a pressure correction $\Delta P = 2.5 \ 10^{-3} \ GPa/K$ 125 (Andrault et al., 1998) compared to the nominal pressure measured at 300 K. The pressure error is 126 estimated to be about 3%. 127

Angle dispersive X-ray diffraction was performed *in situ* at the P02.2 (PETRA III, DESY) and ID27 (ESRF) beamlines. We used a monochromatic X-ray beam tuned at 42.7 keV (P02.2) or 33

keV (ID27) focused to less than $2x2 \mu m^2$ FWHM by two mirrors in the Kirkpatrick-Baez 130 configuration. Diffraction images were acquired on a Perkin Elmer flat panel (P02.2) or Mar-CCD 131 (ID27) detector with exposure times of 10-30 seconds, with or without sample rotation. Distance 132 between sample and detector of 451 mm (P02.2) or 223 mm (ID27) was calibrated using a CeO₂ 133 134 (P02.2) or LaB₆ (ID27) standard. Diffraction patterns were integrated and treated using the Fit2D 135 (Hammersley, 1996) and the GSAS (Larson and Von Dreele, 1988) codes (Figures 1 and 2). The orthorhombic lattice (Pbcn) of seifertite contains 4 independent parameters for atomic positions plus 136 thermal factors for Si and O atoms. Due to the limited number of crystallites present in the X-ray 137 spot, we preferred to avoid refinement of all those parameters. We first performed full X-ray profile 138 refinement, using the Rietveld mode and keeping all atomic positions fixed to the values previously 139 reported in the literature (upper frames in Figure 2; (El Goresy et al., 2008)). Although the atomic 140 parameters may evolve slightly with P and T, the changes will be minor because these parameters 141 are closely linked to the atomic arrangement within the lattice. Consequently, the evolution of the 142 diffraction peaks intensity with pressure should also be reduced. This technique allows to model the 143 experimental peaks with d_{hkl} lines of comparable intensities and prevents evolution of the 144 orthorhombic lattice toward solutions that are favorable from a mathematically rather than physical 145 point of view. In a second step, we verified the adequateness between all experimental and 146 theoretical diffraction features, using the Le Bail mode in which the peak intensities are adjustable 147 parameters (lower frames in Figure 2; (Le Bail, 1992)). For the NaCl or KCl pressure medium and 148 149 platinum, we adjusted each diffraction line intensity independently using the Le Bail method, in order to obtain a best refinement of the diffraction profile. More than 20 diffraction lines were used 150 to identify the CaCl₂-structure and seifertite in two-theta ranges between 5 and 18 (P02.2) and 6 and 151 22 (ID27) degrees. The reduced chi-squared value are typically 0.15 or 0.07 for siefertite (or the 152 CaCl₂-form) refined in Rietveld or Le Bail modes, respectively. The error in volume determination 153 is less than 0.3%. The XRD data were acquired repeatedly during heating and after quenching to 154 room temperature. 155

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157 Phase transformation from CaCl₂-form to seifertite

We conducted 5 different runs (4 at P02.2 and 1 at ID27) of stepwise compression at 300 K 158 followed by laser heating at successively higher temperatures. When the same sample was used for 159 160 repeated experiments at increasing pressure, previously unheated sample areas were investigated at 161 each new pressure. Crystallization of the CaCl₂-form from the silica glass was obtained after laser 162 heating to 2500 K at 91, 98 and 106 GPa and at 90 and 107 GPa for the samples containing 4 and 6 wt% Al₂O₃, respectively. After further steps of incremental pressurization, heating led to the 163 formation of seifertite at nominal pressures of 119 and 121 GPa in the samples with 4 and 6 wt% 164 Al₂O₃, respectively. In the 5th experiment performed on the Al-6% glass at the ID27 beamline, we 165 observed pure CaCl₂-phase up to 108 GPa and 2500 K, a mixture of the two silica polymorphs 166 between 108 and 118 GPa, and pure seifertite starting from this pressure. Those results pinpoint at 167 ~2500 K the transition pressures at 108-119 GPa and 108-118 GPa for Al-4% and Al-6% 168 compositions, respectively. Those values are compatible with previous studies: At similar 169 temperatures, the transition pressure was reported at 122 GPa for pure SiO₂ (Murakami et al., 2003) 170 and between 100 and 113 GPa for MORB compositions (Hirose et al., 2005). Because our starting 171 material is a glass, we did not observe any metastable formation of seifertite at about 80 GPa as 172 reported previously from cristobalite starting material (Dubrovinsky et al., 2001; Shieh et al., 2005). 173 174 The glass crystallization is completed above ~3000K (Figures 1-3). In all cases, no further phase transformation was observed upon heating after the glass crystallization. Although additional Al-175 bearing phases were not detected, we cannot rule out Al₂O₃ saturation of the CaCl₂-structured and 176 seifertite phases. Such saturation would imply that the Al₂O₃-content in SiO₂-dominated phases are 177 lower than the 4 or 6 wt% of our starting materials. However, Al₂O₃ saturation in our seifertite is 178 unlikely, because seifertite crystallized from a natural basalt was previously reported with 12.8 wt% 179 Al₂O₃ (Hirose et al., 2005). 180

181 Based on the available data, we draw a tentative phase diagram for the Al-bearing SiO_2 182 compositions as a function of pressure (Figure 4). Due to the presence of Al-rich Ca-ferrite (near the MgAl₂O₄-NaAlSiO₄ join), the Al₂O₃ activity in a MORB composition is expected to be high but 183 lower than unity. Therefore, the Al₂O₃-contents of 3.4 and 12.6 wt% of the CaCl₂-phase and 184 seifertite, respectively, reported by Hirose et al. (2005), provide lower bounds for the Al₂O₃ 185 186 solubility limit in silica. These values can hardly be used to calculate an Al partition coefficient between the two phases (K_{Al}^{CaCl2/Seifertite}), because the changes in mineralogy in their samples, 187 between 60 GPa and 113 GPa, is likely to affect the Al-activity significantly. On the other hand, the 188 pressure of 113 GPa for their run #6 corresponds to a higher bound for the CaCl₂-form to seifertite 189 phase transition in case of saturation in Al_2O_3 . The results for our two compositions provide 190 constraints on the position of the phase loop between the two minerals. The binary phase loop tends 191 towards lower pressure with increasing Al-content. We draw straight lines for phase boundaries 192 compatible with all data sets in the phase diagram (Figure 4). By comparing the Al-contents in 193 CaCl₂-phase and seifertite at a given pressure, we calculate a K_{Al}^{CaCl2/Seifertite} value of ~0.55(5). For 194 this value, the phase transformation for e.g. an Al-bearing silica composition with 5 wt% Al_2O_3 195 would extend from ~112 GPa to ~116 GPa, a pressure gap corresponding to a ~70 km thick 196 transition interval at a depth of ~2500 km. 197

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199 **Results for the CaCl₂-form of silica**

We report our SiO₂ volumes measured at 300 K for the CaCl₂-form below the transition pressure [Figure 5A]. We also report the data sets available for starting materials made of pure SiO₂ (Andrault et al., 2003; Murakami et al., 2003), MORB (Hirose et al., 2005), and an Al-bearing SiO₂ glass with 4 wt% Al₂O₃ (Bolfan-Casanova et al., 2009). In the latest study, stishovite was synthesized with an estimated Al₂O₃-content of 2.2-3.6 wt%, for a starting material containing 4 wt%. This range of values is comparable to 3.4 wt% Al₂O₃ measured in the CaCl₂-form by Hirose et al. (2005) and slightly less than the 4 and 6 wt% Al₂O₃ present in our starting materials. At this

point, it is interesting to note significant differences in the CaCl₂-form volumes between previous 207 208 studies. The lowest compression curve, by ~1% of volume or 6 GPa compared to our previous studies, is that of pure SiO_2 reported by Murakami et al. (2003). At mid distance between those two 209 210 curves plot the compression curve of MORB (Hirose et al., 2005). These differences could be due to a difference in pressure determination or in volume measurement of the silicate phase. Because 211 212 no pressure medium was used in the previous studies, the most likely cause of experimental difference originates from pressure determination. The volume relaxation from high-temperature to 213 214 300 K for metallic blobs of Pt or Au pressure markers enclosed in a hard silica matrix can be hampered when shutting down the IR-laser. However, the volume change at the transition from 215 CaCl₂-form to seifertite should be precise and significant for all studies, because the same 216 experimental procedure was used throughout. 217

Our data set plots in perfect agreement with our prior studies performed on Al-free (Andrault et 218 al., 2003) and 4 wt% Al₂O₃ (Bolfan-Casanova et al., 2009), irrespective of Al-composition. Using 219 220 the same experimental methods, compression curves of the 4 starting materials (2 in previous studies and 2 in the present one) are found within the experimental error. It shows that Al presents a 221 222 negligible effect on the volume of CaCl₂-structured silica at very high pressures. A study specifically dedicated to the effect of Al on the stishovite bulk modulus, using Brillouin scattering 223 at ambient conditions, showed that addition of 6 wt% Al₂O₃ induces a ~8% decrease in bulk 224 modulus K₀ (Lakshtanov et al., 2007a). However, the well-known tradeoff between K₀, its first 225 226 pressure derivative K'₀, and maybe also second derivative K''₀, makes the compression curve very similar at high pressures. In our previous studies (dotted black line in Figure 5A), the Birch-227 Murnagham formalism was used to refine an isothermal (300 K) compression curve with the 228 following parameters: $V_0 = 46.77 \text{ Å}^3$, $K_0 = 292 \text{ GPa}$ and $K'_0 = 4.59$ (see sample PV-20 in Table 1 of 229 (Bolfan-Casanova et al., 2009)), where V_0 is the unit cell volume at standard P-T conditions. In this 230 study, we preferred the use of the Vinet EoS formalism for reference isotherms, in order to facilitate 231 232 the comparison with a recent study on stishovite (Wang et al., 2012). We therefore adjusted the previous compression curve using the Vinet EoS and obtained $V_0 = 46.77$ Å³ (fixed), $K_0 = 297(7)$ GPa and $K'_0 = 4.59$ (fixed).

To our knowledge, there is no study reporting the EoS of the CaCl₂-form of Al-bearing SiO₂ 235 phases at high pressures and temperatures. A previous work performed up to 54 GPa and 1700K 236 reports the following parameters for the high temperature properties of stishovite: $\Theta_{D0}=1130$ K, 237 $\gamma_0=1.67$, a=1, b=3, where Θ_{D0} and γ_0 are the Debye temperature and the Grüneisen parameter at 238 standard P-T conditions, respectively (Wang et al., 2012). The EoS formalism was based on Mie-239 Grüneisen-Debye to model the high temperature properties. In absence of a more reliable data set, 240 we assume the following parameters to estimate volumes of the CaCl₂-form at all P-T conditions: 241 $V_0 = 46.77 \text{ Å}^3$, $K_0 = 297 \text{ GPa}$, $K'_0 = 4.59$, $\Theta_{D0} = 1130 \text{ K}$, $\gamma_0 = 1.67$, a = 1, and b = 3. Based on this PVT 242 EoS, we calculate volumes of the CaCl₂-form at 300, 2000 and 4000 K in the pressure range 243 investigated in this study (Figure 5B). 244

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246 **P-V-T equation of state of seifertite**

247 We also report our SiO_2 volumes measured at 300 K [Figure 5A] and at temperatures up to 4300 K [Figure 5B] for seifertite between the transition pressure and 162 GPa. Because seifertite 248 and the CaCl₂-structured phase have 4 and 2 formula units per unit cell, respectively, we divide the 249 seifertite volume by 2 to simplify the comparison. As for the CaCl₂-form, we also report data sets 250 for pure SiO₂ (Murakami et al., 2003) and MORB (Hirose et al., 2005). We first refine the P-V-251 300K EoS of seifertite of Al-6 composition to $V_0 = 92.73(10) \text{ Å}^3$, $K_0 = 304.2(3.0)$ GPa and $K'_0 = 4.59$ 252 (fixed), using the 14 data points recorded at ambient temperature in this study [Table 1]. We fix K'₀ 253 254 because our experimental data set does not allow to adjust precisely the 3 parameters independently, due to the dependency between K_0 and K'_0 values. The seifertite bulk modulus appears very similar 255 to that of the CaCl₂-form, only ~2% higher than the value of 297 GPa recalculated for $K'_0 = 4.59$, 256 based on a study using precisely the same experimental methods (Bolfan-Casanova et al., 2009). If 257 we fix K'₀=5.0, we obtain V₀= 93.75(10) Å³ and K₀= 270.3(3.0) GPa. The bulk modulus is found 258

~7% lower than the 292 GPa reported for stishovite in a previous study using a large volume press
apparatus (Wang et al., 2012).

261 At high temperature, we refine $\Theta_{D0}=1130$ K (fixed), $\gamma_0=1.61(3)$, a=1 (fixed) and b=3 (fixed) [Table 2], using the 16 data points recorded in this study [Table 1]. The choice of fixing Θ_{D0} , a, and 262 b is to facilitate the quantitative comparison with a previous study on stishovite (Wang et al., 2012). 263 With these parameters, the maximum difference between the experimental and theoretical pressures 264 is 2.7 GPa, which is within the experimental uncertainties. The refined room pressure Grüneisen 265 parameter of seifertite appears very similar to that proposed recently for stishovite, which was 266 reported as 1.67 (Wang et al., 2012). In the P-T range investigated, we calculate Debye temperature 267 and Grüneisen parameter mean values of Θ_D =1488(20) K and γ =0.78(4), respectively (the Θ_{D0} and 268 γ_0 reported above are ambient values). Based on this PVT EoS, we calculate volumes of seifertite at 269 300, 2000 and 4000 K for the pressure range investigated in this study (Figure 5B). 270

We finally analyzed the evolution of the seifertite cell parameters with pressure (Figure 6). We treated the fictive volumes a^3 , b^3 and c^3 in a similar manner like seifertite unit cell volumes. It results in a set of isothermal and high temperature elastic parameters [Table 2]. The Seifertite unit cell is quite anisotropic at high P and T, with the a parameter being 30% less incompressible than b and c, and with the b parameter showing a Grüneisen parameter about twice of a and c.

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277 Density change across the phase transition

A comparison between the 300 K isothermal compression curves of the CaCl₂-form and seifertite of Al-6 composition evidences a volume change $\Delta V/V = -0.5(1)\%$ across the phase transition [Figure 5A]. This value is very comparable to the volume changes of -0.6% refined for pure SiO₂ (Murakami et al., 2003). For temperatures of 2000 and 4000 K, we refine volume changes of -0.52(10) % and -0.62(10) %, respectively, based on the PVT EoS [Figure 5B]. The fact that $\Delta V/V$ increases slightly with temperature is compatible with a Grüneisen parameter lower in seifertite compared to the CaCl₂-form.

The volume reduction of -0.6% observed for pure SiO₂ (Murakami et al., 2003) translates 285 directly in a +0.6% density change at the transition. However, the relationship between volume and 286 density is more complex for Al-bearing SiO₂ compounds. Three different mechanisms are possible 287 for the Al₂O₃ insertion in SiO₂ (Bromiley et al., 2006; Hirose et al., 1999): (a) Formation of one 288 oxygen vacancy per two Si^{4+} substituted by Al^{3+} ; (b) Addition of one extra-interstitial Al^{3+} (i.e. a 289 total of four Al^{3+}) when three Si⁴⁺ are replaced; (c) A coupled (Al^{3+}, H^{+}) substitution in the presence 290 291 of hydrogen. Also, one should keep in mind that different substitution mechanisms could still coexist in the same silica phase. Deriving density changes from volume changes at the phase 292 transition requires assumptions on the predominant Al substitution mechanism in both the low and 293 high pressure forms of SiO₂. The reason is that the number of atoms per unit cell, and thus the 294 corresponding atomic mass, is not the same for the three different Al-substitution mechanisms. 295 Based on volume variations available for various Al-bearing compositions, we calculate density 296 changes that should be expected at the transition for all possible types of substitution mechanisms in 297

298 CaCl₂-form and seifertite (Table 3).

Thermodynamic principles require a density increase (a positive $\Delta \rho$) for the pressure-induced 299 phase transformation irrespective of composition. Therefore, some Al-substitution mechanisms can 300 be readily eliminated. First of all, having an identical substitution mechanism (ISM) for both 301 polymorphs would yield a 2.3% density reduction at the transition pressure for Al-bearing silica in 302 the MORB material (Hirose et al., 2005). This situation is not impossible, because this negative $\Delta \rho$ 303 could be counterbalanced by complementary mineralogical changes and element partitioning in the 304 305 multi-component material. However, a largely negative $\Delta \rho$ would be surprising and we therefore 306 search for a more likely situation. The O-vacancy substitution mechanism in seifertite is unlikely, because this would cause negative $\Delta \rho$ in our samples whatever is the Al₂O₃-substitution 307 mechanisms in the CaCl₂-form (once the ISM case is discarded). In the same way, the Al-interstitial 308 mechanism is unlikely in the CaCl₂-form. In a study at pressures up to 25 GPa, it was shown that 309 the mechanism involving coupled Al^{3+} and H^+ substitution can only explain 40% of the Al-310

substitution in stishovite (Litasov et al., 2007). In our experiments, particular care was taken to prevent the presence of water. Also, any water present in one polymorph is also likely to be present in the other, which would yield a similar Al-substitution mechanism in both polymorphs (the ISM case).

The most likely mechanisms are therefore the O-vacancy substitution in the CaCl₂-form and the 315 Al^{3+} -interstitial substitution in seifertite (the O->Al mechanism in Table 3). This is the only 316 scenario in which the phase transition obeys the thermodynamic rule of density increase for all the 317 data sets available. It is compatible with a previous study showing that formation of O-vacancies is 318 319 the predominant mechanism for the Al-substitution in stishovite (Bromiley et al., 2006). It remains unknown how it evolves with the increasing pressure in the rutile-structure, as well as across the 320 321 structural distortion that yields the CaCl₂-type structure. For our samples with 4 or 6 wt% Al₂O₃, 322 the O->Al mechanism would imply a $\sim 2.19\%$ density increase at the phase transition. This density increase is significantly larger than for the pure SiO₂ system ($\Delta \rho / \rho \sim 0.6\%$), in agreement with a 323 lowering of the transition pressure with increasing the Al-content. In MORB, the density change is 324 smaller, with a value of $\sim 0.23\%$, because of the strong Al-partitioning into seifertite. 325

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327 **Buoyancy of the subducted slabs**

Recent studies agree that basaltic materials of subducted slabs have a slight, but significant, 328 329 density excess compared to a peridotitic lower mantle, even if the slabs are thermally equilibrated. 330 Depending on composition and pressure, a density difference of 0.4 to 4.0 wt% was reported (Guignot and Andrault, 2004; Hirose et al., 2005; Ono et al., 2005; Ricolleau et al., 2010). MORBs 331 with elevated Fe-content are clearly denser. Despite the density excess, basaltic material under 332 lower mantle conditions contains free silica phases, which are the lowest density phases at depths 333 334 exceeding 1600-1800 km. Altogether, some studies suggest a decrease of the MORB density excess with increasing mantle depth (Guignot and Andrault, 2004; Ricolleau et al., 2010) and others 335 suggest a constant $\Delta \rho$ (Hirose et al., 2005; Ono et al., 2005). For sedimentary material, the density 336

is significantly lower than for MORB, due to low Fe and high Si-contents (Guignot and Andrault,2004).

339 Our experimental results on compositions with 4-6 wt% Al_2O_3 suggest a ~2.19 % density increase for the transition from the CaCl₂-structured form to seifertite at 113-119 GPa and 300 K. 340 At temperatures above 2000 K, this value should reach ~3%, due to higher volume change found 341 for this transition at high temperature. The strong Al-partitioning from the CaCl₂-structured phase 342 343 into seifertite will lead to compositional and density adjustments of the other Al-bearing phases in basaltic lithologies, in addition to the density effects caused by the silica-dominated phase 344 transition. In subducted slab material the silica-dominated phases can represent up to about 20 and 345 40 vol% of MORB and sedimentary components, respectively (Hirose et al., 1999; Irifune et al., 346 1994). The CaCl₂-structured phase to seifertite transition could have a minor effect on the buoyancy 347 in the peridotitic mantle of the basaltic ocean crust of subducted slabs, because free-SiO₂ is not 348 highly abundant and the MORB density is already significantly higher. For the sedimentary 349 350 component, however, the CaCl₂ to seifertite transition may have a larger impact and facilitate the descent of the sedimentary material to the lowermost part of the mantle. 351

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441	Table 1 : Unit cell parameters of seifertite of 6 wt% composition recorded from 100 to 162 GPa and
442	300 to 4300 K. Pressure is derived from the PVT-EoS of platinum (Holmes et al., 1989) and
443	temperature was determined by radiometric measurements. Uncertainties are 3% for pressures, 50 K
444	for high temperatures, ~0.005 Å for (a, b, c) cell parameters and ~0.1 Å ³ for unit cell volumes.
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		Seifertite				
T [K]	P [K]	A [Å]	B [Å]	C [Å]	V [Å ³]	V/2 [Å ³]
300	102.4	3.785	4.722	4.201	75.10	37.55
3410	117.8	3.778	4.736	4.197	75.10	37.55
300	105.7	3.767	4.723	4.198	74.69	37.35
3400	121.2	3.779	4.717	4.190	74.69	37.35
3200	121.0	3.768	4.718	4.191	74.50	37.25
300	118.9	3.755	4.676	4.194	73.64	36.82
3020	130.0	3.761	4.688	4.188	73.83	36.91
3260	127.4	3.764	4.693	4.193	74.07	37.03
300	116.5	3.753	4.681	4.185	73.52	36.76
300	120.9	3.747	4.688	4.172	73.30	36.65
300	121.5	3.743	4.686	4.179	73.30	36.65
3225	130.5	3.745	4.694	4.193	73.71	36.85
3240	130.3	3.746	4.713	4.183	73.86	36.93
300	123.5	3.738	4.690	4.177	73.22	36.61
3250	131.7	3.741	4.696	4.182	73.46	36.73
3475	138.8	3.730	4.689	4.170	72.93	36.46
3330	135.9	3.740	4.692	4.171	73.18	36.59
300	133.3	3.717	4.664	4.168	72.26	36.13
3850	145.7	3.736	4.667	4.158	72.50	36.25
300	131.5	3.721	4.660	4.160	72.14	36.07
4300	150.8	3.731	4.661	4.161	72.35	36.18
300	136.2	3.715	4.651	4.158	71.84	35.92
4170	156.6	3.723	4.663	4.134	71.77	35.89
3350	152.9	3.719	4.660	4.151	71.95	35.97
300	140.8	3.716	4.659	4.136	71.60	35.80
3835	159.0	3.712	4.659	4.139	71.58	35.79
300	144.1	3.709	4.653	4.130	71.27	35.64
3600	162.3	3.703	4.652	4.134	71.21	35.61
300	147.4	3.695	4.634	4.136	70.82	35.41
300	153.4	3.692	4.628	4.135	70.65	35.33

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Table 2: Elastic parameters refined for seifertite, as well as for its 3 fictive a^3 , b^3 , c^3 volumes. We used the Vinet isothermal compression curves and the Mie-Grüneisen approach for modeling the effect of temperature. Θ_{D0} and γ_0 are the Debye temperature and Grüneisen parameter at ambient conditions. Some parameters were fixed following the reports of (Bolfan-Casanova et al., 2009)¹ and (Wang et al., 2012)².

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Seifertite	Volume	a ³	b ³	c ³
V_0 [Å ³]	92.73(10)	71.0(5)	129.8(3)	92.9(5)
K ₀ [GPa]	304.2(3.0)	205(9)	310(4)	291(10)
K' ₀	4.59^{1}	4.59 ¹	4.59 ¹	4 .59 ¹
$\Theta_{ m D0}$ [K]	1130 ²	1130 ²	1130 ²	1130 ²
γ_0	1.61(3)	3.01(7)	5.96(20)	2.83(10)
А	1^2	1 ²	1 ²	1 ²
В	3^{2}	3 ²	3 ²	3 ²

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Table 3: Measured volume changes ($\Delta V/V$) and calculated density changes ($\Delta \rho/\rho$) across the 459 CaCl₂-form to seifertite transition (Figure 5A). M-03 and H-05 stand for (Murakami et al., 2003) 460 and (Hirose et al., 2005), respectively. The $\Delta \rho / \rho$ are calculated using the measured $\Delta V / V$ taking 461 into account all possible Al substitution mechanisms in SiO₂ polymorphs. The Al³⁺ substitution can 462 imply formation of an oxygen vacancy (O), insertion of extra-interstitial Al³⁺(Al), or coupled Al³⁺-463 H⁺ substitution (H). The term "O->Al", for example, means a transition from an oxygen-deficient 464 CaCl₂-form to a seifertite containing interstitial Al³⁺. ISM: Identical substitution mechanism in both 465 polymorphs. 466

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	ΔV/V (%)	Δρ/ρ (%)						
		ISM	O->Al	Al->0	0->H	H->0	Al->H	H->Al
SiO ₂ (M-03)	-0.60	0.60	0.60	0.60	0.60	0.60	0.60	0.60
Al-6 (This study)	-0.50	0.50	2.19	-1.16	1.56	-0.55	-0.12	1.12
MORB (H-05)	2.30	-2.30	0.23	-5.39	-1.72	-4.49	-3.23	-0.38

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470 **Figure Captions**

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Figure 1: Selected regions of the 2D diffraction image (upper frames) and its spectral integration (lower frame) recorded *in situ* at 120 GPa and 2500 K. Regions where major seiferte diffraction peaks occur are colored in red. The number of seifertite crystallites is insufficient for a full structural refinement using the Rietveld method. However, the diffraction peak intensities are in good agreement with the structural model of seifertite (see Figure 2).

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Figure 2: The diffraction spectra are refined in two steps. We first perform a full X-ray profile 478 refinement, using the Rietveld mode and the structural model of seifertite (upper frames). 479 Differences between experimental and calculated peak intensities are due to (i) change with 480 pressure and temperature of the seifertite atomic parameters; the structural model that we used 481 corresponds to ambient P-T conditions (El Goresy et al., 2008)), and (ii) an insufficient number of 482 crystallites in the X-ray spot. In a second step, we perform a profile refinement using the Le Bail 483 mode (lower frames; see text). In each frame, green, upper red and lower red spectra correspond to 484 the experimental data, fully refined profile, and seifertite contribution, respectively. We also present 485 the major d_{hkl} indexation of the seifertite diffraction peaks (we report "several" when more than 2 486 487 seifertite diffraction peaks plot at the same 2-theta position). Left and right frames present the same experimental spectrum in 2-theta ranges of 5-10 and 10-20 degrees, respectively. 488

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Figure 3: Examples of integrated X-ray diffraction patterns for the CaCl₂-structured silica phase and seifertite. These patterns were recorded (A) on ID-27 beamline at 3400 K and 121 GPa and (B) on P02.2 beamline after heating at 2500 K at 103 and 114 GPa. Small bumps on the diffraction patterns up to 2500 K are due to still incomplete transformation of the starting material. Platinum and the NaCl or KCl pressure medium are also visible. The reported pressure is based on the equation of state of Pt for the diffraction pattern recorded at high temperature before quenching. Black crosses, green and red lines are experimental measurements, fitted background and
 theoretical diffraction spectra, respectively.

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Figure 4: Approximate phase diagram for the SiO₂-Al₂O₃ system at 2500 K. For pure SiO₂, the phase transition is well defined at ~122 GPa (M-03; (Murakami et al., 2003)). For the mid-ocean ridge basalts used in H-05 (Hirose et al., 2005), the CaCl₂-form and seifertite phases were observed at 100 GPa, and 123 GPa, respectively. Concerning our samples, the rectangles represent the pressure range between our samples where CaCl₂-form or seifertite was synthesized. The pressure uncertainty of 3 GPa is included in the rectangle sizes. The loop diagram represented here is the widest model remaining compatible with the different data sets.

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Figure 5: (A) Volumes of the two silica polymorphs measured around the transition pressure at 300 507 K. We also report the data sets reported for pure SiO₂ (red: (Murakami et al., 2003); black dotted 508 line: (Andrault et al., 2003)), a MORB-type material (green: (Hirose et al., 2005)), and an Al-509 bearing SiO₂ glass with 4 wt% Al₂O₃ (black dotted line: (Bolfan-Casanova et al., 2009)). Closed 510 and open symbols represent the $CaCl_2$ -phase and seifertite, respectively. (B) High temperature 511 512 volumes (red circles) recorded for seifertite from 100 to 162 GPa and up to 4300 K. We report the isothermal compression curves calculated at 300 K, 2000 K and 3000 K for the CaCl₂-form (blue 513 514 dotted lines) and seifertite (dashed black lines) based on the PVT EoS refined in this study.

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Figure 6: Seifertite (a, b, c) unit cell parameters recorded at 300 K (blue diamonds) and up to 4300
K (purple rectangles) for pressures from 100 to 162 GPa.

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