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1	REVISION 3
2	Stability and spectroscopy of Mg sulfate minerals.
3	Role of hydration on sulfur isotope partitioning
4 5	
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13	ABSTRACT
14 15	We study hydrated magnesium sulfate minerals from first-principles calculations
16	based on density-functional theory. We determine the heat of hydration for
17	MgSO ₄ .nH ₂ O, compute the Raman and infrared spectra for several phases and
18	calculate the S isotope partitioning as a function of hydration. We find that
19	epsomite and meridianiite, with respectively n=7 and n=11 water molecules per
20	MgSO ₄ unit, are particularly stable with respect to other individual or
21	combinations of hydration states. The Raman spectra of all phases present clear
22	SO ₄ features that are easily identifiable. We use this to show one can use the
23	vibrational spectroscopic information as an identification tool in a remote
24	environment, like the martian surface. We discuss the character and atomic
25	displacement pattern of all vibration modes and compute the ³⁴ S/ ³² S partitioning;
26	this work shows that hydration favors enrichment in the lighter S isotope ³² S with
27	respect to the heavier ³⁴ S, which is accumulated in the less hydrous structures.
28	We show for the first time that the signature of ${}^{34}S/{}^{32}S$ partitioning could be
29	observed by in situ spectroscopy on the surface of Mars. Moreover this can be
30	related to the diurnal cycle of hydration and dehydration and hence it can
31	improve the modeling of the water circulation on Mars.
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INTRODUCTION

The presence of monohydrate and multiply-hydrated sulfate minerals at the surface of Mars as major water-bearing minerals is now a well-established fact [*e.g.* Gendrin et al., 2005, Wray et al., 2009, Flahaut et al., 2010, Noe Dobrea et al., 2012]. Magnesium sulfates play a major role amongst these minerals, but sulfates of other cations have been also proven to exist, like gypsum, jarosite, etc. [Blaney and McCord, 1995, Aubrey et al., 2006, Roach et al., 2009, Horgan et al., 2009, Murchie et al., 2009].

Magnesium sulfates are compounds exhibiting an amazingly large number 43 44 of hydration states. Apart from two anhydrous phases it can contain almost 45 anything between 1 and 7 water molecules per formula unit, with even up to 11 46 water molecules for each MgSO₄ unit. The anhydrous phases do not exist as 47 minerals. Their affinity for water is so large that they are currently used as a 48 desiccant. The family of related minerals are, in order of increasing hydration: the 49 monohydrate, MgSO₄.H₂O, also known as the mineral kiesierite [Hawthorne et 50 al., 1987], sanderite, MgSO₄.2H₂O [Ma et al., 2009a], MgSO₄.2.5H₂O [Ma et al., 51 2009b], starkeyite [Baur, 1964a] and cranswickite [Peterson, 2011], both with 52 MgSO₄.4H₂O formula, pentahydrite, MgSO₄.5H₂O, hexahydrite, MgSO₄.6H₂O 53 [Zalkin et al., 1964], epsomite, MgSO₄.7 H_2O [Baur, 1964b], and finally, 54 meridianiite, MgSO₄.11H₂O [Fortes et al., 2008]. Additional minerals are formed 55 in binary and ternary diagrams with other halogenide and/or sulfate systems 56 [Steiger et al., 2011]. The minerals of this group have been traditionally studied in 57 the context of evaporitic deposits, or of alteration products of sulfide minerals. 58 More recently, with the exploration of the Martian surface and of the icy moons 59 from the outer solar system, the planetological importance of these minerals has 60 been fully appreciated.

61 The relative stability of these phases is conditioned not only by pressure 62 and temperature but also by the partial pressure of water vapor [Hogenboom et al., 1995, Feldman, et al., 2004, Chou et al., 2013]. Some of these phases can
be easily transformed from one to another, but not all such transitions are
reversible [Steiger et al., 2011]. The presence of other specific minerals like
hydrous silicates [Wilson and Bish, 2012] or organic matter [Aubrey et al., 2006],
may affect the stability field of the hydrated MgSO₄ phases by changing and/or
preserving humidity conditions locally.

Experiments on the phase diagram of the $MgSO_4 - H_2O$ system reveal that the surface conditions on the planet Mars, and more precisely the day/night cycle, span the static stability field of the mono-, di-, tetra-, hepta- and hendecahydrates. However, as mentioned above, the precise mineralogy may depend on kinetics of reactions and on the thermodynamical path [Vaniman et al., 2004, Wang et al., 2011, Chou et al., 2013].

Here we use first-principles calculations to compute the energies of hydration to help interpret the water cycle on Mars, and to determine the effect of hydration on the Raman spectra and on the isotope partitioning.

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METHODOLOGY

We perform first-principles calculations using standard density-functional theory (Kohn and Sham 1965; Payne et al. 1992; Martin 2003) in the VASP (ref) and ABINIT (Gonze et al. 2002, 2009) implementations, both based on planewaves and pseudopotentials.

The crystal structures, the heat of formation, and the hydration energies are obtained using VASP simulations, with planar-augmented wavefunction (PAW) (Blochl, 1991) pseudopotentials and the Perdew-Burke-Erzernhof formulation of the Generalized Gradient Approximation (Perdew et al., 1996). The energy cutoffs for the planewaves on respectively the standard grid and the augmented grid are 800 eV and 850 eV kinetic, ensuring an accuracy of the energy better than \pm 0.2 meV. The sampling of the electron density in the reciprocal space is 93 performed on special k points (Mokhorst and Pack, 1976) with varying density 94 from one structure to the other. With these values we achieve energy 95 convergence on the order of the meV/formula unit. We start from the 96 experimental structures from the references mentioned above, and we fully relax 97 them, i.e. we minimize both the theoretical stresses on the unit cell and the 98 forces on the atoms.

99 Then we use ABINIT to compute the lattice dynamical properties (Baroni et 100 al., 2001, Gonze et al., 2005), including the Raman spectra with both peak 101 position and intensities (Veithen et al., 2005) on four selected phases: the two 102 anhydrous compounds, kiesierite, and starkeyite. These calculations are 103 performed using the Local Density Approximation for the exchange-correlation, 104 and Troullier-Martins norm-conserving pseudopotentials. We store all the Raman 105 spectra on the WURM website (http://www.wurm.info). For more details of the 106 Raman calculations, choice of parameters, pseudopotentials, etc. see Caracas 107 and Bobocioiu (2011).

Finally the reduced partition function ratios (RPFR) for sulfur isotopes are obtained from comparing the phonon frequencies obtained for various S masses in the zone-center of the Brillouin zone. The procedure has been successfully applied to a variety of minerals (Balan et al., 2009, Reynard and Caracas, 2009).

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RESULTS AND DISCUSSION

114Energies of hydration115The basic thermodynamic parameters of the hydrated sulfate minerals have116been recently measured (Grevel and Majzlan, 2009). The heat of formation for117the hydrous phases was reported both from elements and from ionic groups and118water molecules, and it was shown that the relative humidity, i.e. the partial water119vapor pressure, plays a critical role in the phase diagrams and the relative120stability of various phases.

Here, we study and compute the heat of hydration for all the hydrated phases of MgSO₄, based on our calculations at 0K. The results are summarized in Table 1, where the enthalpy difference between the mechanical mixture of less
hydrated phase plus water molecule(s) and the more hydrated phase is
computed as follows:

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 $\Delta H_{\text{hydration}} = H_{\text{MgSO4.mH2O}} - [H_{\text{MgSO4.nH2O}} + (\text{m-n}) H_{\text{H2O}}], \text{ with } \text{m>n} (1)$

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129 Negative values favor hydration and positive values favor dehydration. In 130 the limit of thermodynamical equilibrium, which applies to our calculations, the 131 hexahydrite turns out to be metastable with respect to any other type of mixture. 132 This is coherent with the experimental phase diagram of the MgSO₄ – H_2O 133 system, where hexahydrite is metastable at temperatures below freezing 134 [Vaniman et al., 2004, Wang et al., 2006, 2011]. Except for hexahydrite all the 135 other phases are more stable with respect to water vapor and the 136 correspondingly less-hydrated phase.

Heats of reactions between two hydrated phases resulting in a third hydrated phase are listed in Table 2. As expected from the results shown in Table 1, hexahydrite is metastable with respect to the other hydrated sulfates, and epsomite is one of the most stable phases. Reactions between the other sulfates, not involving hexahydrite are easily reversible.

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Raman spectra

We calculate the Raman spectra of the two anhydrous phases, as well as of kieserite, MgSO₄.H₂O, starkeyite, MgSO₄.4H₂O, and epsomite, MgSO₄.7H₂O, with both peak positions and intensities. The spectra are represented in Figure 1. The results are fully reported on the WURM website, where the reader can find as well the crystal structures, the dielectric properties, and animations with the vibration pattern for each phonon mode in the zone center.

The spectra present similarities between the various phases. Strong signatures of the SO₄ groups are clearly distinguishable in all phases, corresponding to the asymmetric stretching and breathing modes. They can be used as identification criteria for this class of minerals. The vibration modes ofthe water molecules lay at high frequency.

Below 700 cm⁻¹ the peaks are dominated by lattice modes, where SO_4 anionic groups and, for the two hydrated minerals, H_2O molecules, vibrate against each other or against the Mg sublattice as rigid units. These modes have relatively low intensity.

The modes in the 900-1200 cm⁻¹ range correspond to vibrations of the SO₄ 159 160 tetrahedra. In the anhydrous phases the highest-intensity peaks of the spectra are the breathing modes of the SO₄ tetrahedra, at 998 and 1006 cm⁻¹ in the 161 respectively alpha and beta phases. The second most intense mode is the 162 asymmetric stretching of the tetrahedra, respectively at 1146 and 1128 cm⁻¹. In 163 164 the hydrated phases, the corresponding breathing modes are at 1004, 985, and 952 cm⁻¹ in respectively kieserite, starkeyite, and epsomite, and the strongest 165 asymmetric stretching at 1120 cm⁻¹ in kieserite. The asymmetric stretching 166 modes are degenerate into several modes with similar intensity over the 1047-167 1185 cm⁻¹ range in starkeyite and over the 1001-1142 cm⁻¹ in epsomite. The 168 169 water molecules show a small participation to these modes, probably due to 170 some dipole-dipole or higher-order dielectric couplings. These modes appear as 171 small peaks in the same frequency interval in experiments.

In the hydrated phases, the bending water modes form a distinct group, at 173 1527 and 1543 cm⁻¹ in kiesierite, in the 1515-1660 cm⁻¹ in starkeyite, and in 174 1560-1655 cm⁻¹ range in epsomite.

Above 2900 cm⁻¹ the peaks correspond to pure water vibrations and are found only in the hydrated phases. One can observe symmetric and asymmetric stretching, signatures that would appear in Raman or infrared. Experimentally the water vibration peaks are convoluted into large asymmetric wide bands around 3300- 3500 cm⁻¹.

Our theoretical results are in relatively good agreement with the experimental measurements [Wang et al., 2006]. The characteristic breathing modes of the SO_4 tetrahedra in the anhydrous phases are at about 20-30 cm⁻¹ lower than the experimental values, typical for vibrational calculations (Baroni et

al., 2001, Caracas and Bobocioiu, 2011). For the hydrated phases, this departure 184 is increased to 40 cm⁻¹ for kieserite, but only 15 cm⁻¹ for starkeyite. Globally the 185 inverse correlation between the position of the breathing mode and the hydration 186 187 state is less pronounced in calculations (about 5 cm⁻¹/water molecule) than in experiments (about 10 cm⁻¹/water molecule). Differences can also be seen in the 188 189 water peaks. In experiments the water peaks are less resolved than in 190 calculations, due to the inherent H disorder present in the samples at finite 191 temperature. The measured vibrational frequencies of water molecules are also 192 higher than in calculations.

These various discrepancies can be attributed to inherent DFT approximations, temperature and anharmonic effects neglected in the calculations, and differences in specific volume between experiment and calculations.

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Isotope partitioning

199 The difference in mass associated with various isotopes induces changes in 200 the phonon frequencies, which in turn have an effect on the free energy. The 201 structure with lighter isotopes will see the frequency of the phonon modes 202 increased with respect to the structure with heavier isotopes. The higher the 203 participation of that particular exchanged isotope to one particular mode, the higher is the effect on its frequency. Thus the calculations of the Raman and 204 205 infrared spectra can be used to determine the effect of hydration on the 206 fractionation of sulfur isotopes.

The reduced partition function ratios, $\beta^{34}S_{/}^{32}S$ for anhydrous β -MgSO₄, 207 208 kieserite, starkeyite, and epsomite are listed in Table 3 as a function of 209 temperature for a range relevant for the surface of Mars. We obtain the ${}^{34}S/{}^{32}S$ 210 partitioning functions between any two sulfate phases as the ratios of the 211 corresponding β functions. The results are plotted in Figure 2. The largest values 212 are for the β -MgSO₄ and epsomite pair, considerably larger than for the β -MgSO₄ 213 and any other hydrated sulfate. Hydrating kieserite and starkeyite yields an 214 intermediate value, as expected. Even if temperature reduces the partitioning

effect, this still occurs at the maximum temperatures characteristic for Mars'surface.

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IMPLICATIONS

220 We compute the hydration energies of hydrous and anhydrous phases of 221 magnesium sulfate minerals. We predict that hydration of anhydrous phases, 222 mono-, tri- and tetra- hydrates is particularly easy to achieve. Reactions between 223 sulfates, not involving hexahydrite, are close to equilibrium and thus easily 224 reversible. Hence both meridianiite and epsomite should be found in large 225 amounts on the surface of Mars, as they are remarkably stable with respect to 226 dehydration; we anticipate to see much less hexahydrite, as this phase is 227 metastable with respect to the other hydrated sulfates.

The simulated Raman spectra show clear and distinctive peaks of the SO_4^{2-} groups and of the H₂O molecules, which can help in identification of this group of minerals. The theoretical results are in good agreement with the existent experimental measurements and provide a nice complement thereof. The <u>http://www.wurm.info</u> website provides all details of the spectra.

Finally, our calculations show that the ${}^{34}S/{}^{32}S$ partitioning increases with increasing difference in hydration. This is consistent with having the less hydrated species enriched in the heavier isotope, S³⁴. Or alternatively, this can be viewed as hydration increasing the amount of light isotope, S³².

This suggests that the signature of the diurnal hydration/dehydration cycles on Mars should become visible over geological time in the isotopic composition of the various hydrous magnesium sulfates. The light isotopes would accumulate in the more hydrated phases, and the heavy isotopes in the more anhydrous phases. This trend could be captured by in situ spectroscopic measurements. If kinetic data are available, the amount of partitioning could be even correlated to the duration of that particular deposit.

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436	FIGURE CAPTIONS
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440	FIGURE 1. Theoretical Raman spectra of anhydrous α -MgSO ₄ , β -MgSO ₄ ,
441	kieserite MgSO ₄ .H ₂ O, starkeyite MgSO ₄ .4H ₂ O, and epsomite MgSO ₄ .7H ₂ O.
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445	FIGURE 2. Theoretical S^{34}/S^{32} partitioning as a function of temperature between
446	β -MgSO ₄ , kieserite MgSO ₄ .H ₂ O, starkeyite MgSO ₄ .4H ₂ O, and epsomite
447	MgSO ₄ .7H ₂ O. The results suggest that hydration favors enrichment in lighter
448	isotopes. This feature should be visible by in situ measurements on the surface
449	of Mars and could be related to interpret the diurnal cycle of hydration/hydration.
450	\boldsymbol{m} and \boldsymbol{n} denote the number of molecules per MgSO4 unit in the two phases
451	between which the partitioning is plotted.
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TABLES

459 460 Table 1.

Heat of hydration for the various magnesium sulfate phases considered in this study. The hydrate with *m* water molecules needs the amount of energy listed in the table to add (*n*-*m*) water molecules/MgSO₄ group unit to form a hydrate with *n* water molecules. *m*=*n*=0 is for β -MgSO₄. Units are kJ/mol. Negative numbers favor the hydrated phase, positive numbers favor dehydration. In general higher hydrated states are preferred, especially epsomite. Hexahydrite is the only metastable phase at 0K.

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n	0	1	3	4	5	6	7	11
m								
0		-86.92	-185.98	-245.75	-307.86	130.11	-447.32	-685.31
1			-99.07	-158.83	-220.95	217.03	-360.41	-598.40
3				-59.77	-121.88	316.09	-261.34	-499.33
4					-62.11	375.86	-201.57	-439.56
5						437.97	-139.46	-377.45
6							-577.43	-815.42
7								-237.99
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473 Table 2.

474 Heat of reaction between sulfates with various hydration states: $\Delta H_{hydration} =$ 475 $H_{MgSO4.mH2O} + H_{MgSO4.nH2O} - 2 H_{MgSO4.(m-n)/2H2O}$]. Except for the formation of 476 hexahydrite, the other transformations are easily reversible. Epsomite is 477 particularly stable with respect to other hydrated sulfates.

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m	n	(m+n)/2	ΔH_{min} (kJ/mol)
1	5	3	-0.24
1	7	4	-0.44
1	11	6	-10.70
3	5	4	-0.02
3	7	5	-0.18
3	11	7	0.24
5	7	6	-10.52

482 Table 3.

- 483 Reduced partition function ratios (the $\beta_{s}^{34}{}_{/s}^{32}$) for β -MgSO₄, kieserite, starkeyite,
- and epsomite in the temperature range relevant to Mars's surface.

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		Kieserite,	Starkeyite,	Epsomite,
Temperature (K)	β-MgSO ₄	$MgSO_4.H_2O$	$MgSO_4.4H_2O$	MgSO ₄ .7H ₂ O
100	1.4402	1.4338	1.4265	1.4161
120	1.3396	1.3351	1.3295	1.3219
140	1.2726	1.2692	1.2647	1.2588
160	1.2249	1.2223	1.2186	1.2138
180	1.1894	1.1873	1.1842	1.1802
200	1.1621	1.1603	1.1576	1.1542
220	1.1404	1.1389	1.1366	1.1336
240	1.1229	1.1216	1.1196	1.1169
260	1.1085	1.1073	1.1056	1.1032
280	1.0965	1.0955	1.0939	1.0917
300	1.0863	1.0855	1.0840	1.0821
320	1.0777	1.0769	1.0756	1.0738
340	1.0703	1.0696	1.0684	1.0668
360	1.0639	1.0633	1.0622	1.0606
380	1.0583	1.0577	1.0567	1.0553
400	1.0534	1.0529	1.0519	1.0506



