

1     **In-situ high-pressure transmission electron microscopy for earth and**  
2                                   **materials sciences**

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8                                   **REVISION 1**

9  
10                                  **ABSTRACT**

11     Transmission electron microscopy in combination with in-situ high-pressure and  
12     high-temperature measurements is uniquely able to provide high-resolution data about  
13     materials under conditions resembling those in Earth's interior. By using  
14     nanocontainers made of graphitized carbon, it is possible to achieve pressures and  
15     temperatures up to at least 40 GPa and 1500 °C, respectively. A wide range of relatively  
16     simple minerals have been studied using this approach. Results to date show the  
17     influence of crystallographic defects in concentrating and storing carbon within analogs  
18     to minerals occurring deep inside Earth.

20       **Keywords:** in-situ transmission electron microscopy, high-pressure measurements, carbon  
21 nanocontainers, carbon nanotubes (CNTs), carbon nanofibers (CNFs), carbon nano-  
22 onions (CNOs)

23

24

## INTRODUCTION

25       Transmission electron microscopy (TEM) has long been used to study the products  
26 of high-pressure experiments at the near-atomic scale. However, in all cases it has been  
27 necessary to quench the samples before they could be imaged at high resolution (Mao  
28 and Hemley 1998). Diamond-anvil cells (DACs) and multi-anvil presses (MAPs), the  
29 instruments currently used for pressure generation, prevent the *in-situ* use of TEM  
30 because their substantial sizes preclude the necessary electron transparency. As a  
31 consequence, in-situ TEM applications for experiments at gigapascal pressure ranges,  
32 particularly meaningful to the earth sciences, have been impossible up to now.

33       X-ray diffraction and other spectroscopic techniques available for in-situ high-  
34 pressure research acquire statistical information averaged over the relatively large  
35 sample volumes interacting with the source radiation. However, in many cases, studies  
36 of crystal defects and mineral reactions at unit-cell dimensions are central to  
37 understanding geophysics and geochemistry in Earth's interior (Cordier 2002; Karato  
38 2010; Stixrude and Lithgow-Bertelloni 2012). TEM is one of the most useful techniques,

39 and commonly the only one, for observing defect features and analyzing chemical  
40 compositions at down to atomic resolutions (Buseck 1992; Veblen 1985). Therefore, in-  
41 situ TEM capabilities at high pressure have long been desired within the earth and  
42 materials science communities.

43 The goal of this paper is to provide an overview of recent efforts to develop and  
44 refine an in-situ, high-pressure TEM method for the earth and materials sciences. With  
45 successful applications to geophysically significant minerals and mineral analogs, we  
46 demonstrate the feasibility and potential of this new technique.

#### 47 **GRAPHITIC NANOCONTAINERS AND NANOPRESSES**

48 Graphitic networks can lose carbon atoms through displacement damage and  
49 vacancy formation when exposed to electrons with acceleration voltages over ~86 kV in  
50 an electron microscope (Smith and Luzzi 2001). If the graphitic networks are curved on  
51 the nanometer scale and the temperature is raised to above ~300 °C, structural  
52 reorganization occurs around the relatively immobile vacancies in the networks,  
53 causing their shrinkage (Fig. 1) (Banhart 1999; Banhart 2004; Krashennikov et al. 2005).  
54 If they are in the form of closed containers that enclose condensed materials,  
55 compression of the enclosed materials occurs (Banhart and Ajayan 1996). Calculations  
56 indicate that if the containers are sufficiently small, what we call nanocontainers, then

57 the internal pressures can reach 40 GPa in, for example, multi-walled carbon nanotubes  
58 (CNTs) (Sun et al. 2006a).

59 Pressure generation in carbon containers can be understood in terms of Laplace's  
60 law, which relates internal pressure (P) of a fluid-filled hollow vessel to wall tension (T)  
61 and its hollow radius (R). For a cylindrical vessel,  $T = P \cdot R$ , whereas  $T = P \cdot R/2$  for a  
62 spherical vessel. The wall tension of a 19-shelled CNT is at least 140 N/m (Sun et al.  
63 2006b). Atomistic calculations suggest that internal pressures in multi-walled CNTs  
64 converge to a maximum with only ~6 graphitic shells, such that further increases in the  
65 number of walls do not produce proportional pressure increases (Sun et al. 2006a).  
66 Therefore, for an inner sample diameter of 100 nm, the electron-transparent thickness  
67 limit for most materials, maximum internal pressures of greater than 2.8 and 5.6 GPa  
68 would be expected in tubular and spherical graphitic containers, respectively.

69 The workable wall thicknesses are limited by half of the mean absorption distance ( $\lambda$ )  
70 for graphite since the container walls both below and above an enclosed sample interact  
71 with the incident electrons. For a typical TEM acceleration voltage and collection angle,  
72 e.g., 300 kV and 3 mrad, respectively,  $\lambda$  is ~225 nm (Akhtar et al. 2012).

73 Carbon nanocontainers enclosing samples of interest can be prepared through either  
74 insertion of samples into pre-existing containers or growth around the minerals of  
75 interest. If sufficiently thin, the walls of these carbon nanocontainers permit the use of  
76 standard TEM techniques that include selected-area electron diffraction (SAED),



77 convergent-beam electron diffraction (CBED), high-resolution TEM (HRTEM), energy-  
78 dispersive X-ray spectroscopy (EDS), electron energy-loss spectroscopy (EELS), and  
79 energy-filtered TEM (EFTEM).

## 80 **METHOD DESCRIPTION**

### 81 **Growing suitable containers**

82 Carbon nanocontainers are essential for TEM measurements at elevated pressures.  
83 The containers can be: (a) CNTs consisting of rolls of graphene layers that form hollow  
84 tubes (Fig. 2a), (b) carbon nanofibers (CNFs) that consist of graphene layers arranged as  
85 stacked cones or cups (Fig. 2b), (c) carbon nano-onions (CNOs) that are spherical and  
86 made of multiple graphene layers surrounding an empty core (Fig. 2c), or (d) poorly  
87 graphitic cages or hydrocarbon coatings that can be graphitized under controlled  
88 conditions within the column of the electron microscope. In all cases, the containers  
89 should be as perfectly graphitized as possible prior to starting the high-pressure  
90 experiments. CNTs, CNFs, and CNOs are available commercially but can also be  
91 readily synthesized in laboratories by using arc discharge, chemical vapor deposition  
92 (CVD), or laser ablation (Fig. 3).

### 93 **Loading minerals into carbon nanocontainers**

94 Placing minerals of interest into tiny carbon containers is a major challenge. We used  
95 two approaches: (a) filling existing containers, and (b) growing the containers around

96 nanoparticles of interest. For convenience, we call them type-A and type-B containers,  
97 respectively.

98 Type-A containers are tubular CNTs or CNFs and can be loaded through capillary  
99 introduction of solutions or melts from which the solids of interest crystallize (Dujardin  
100 et al. 1994). The fluids can be of mineral samples or their precursors, either in the  
101 molten state or in solution, commonly followed by thermal decomposition (Ajayan and  
102 Iijima 1993; Tsang et al. 1994). This method has proven particularly useful for loading  
103 oxide minerals such as magnetite ( $\text{Fe}_3\text{O}_4$ ), shcherbinaite ( $\text{V}_2\text{O}_5$ ), and bunsenite ( $\text{NiO}$ )  
104 (Fig. 4). A typical loading yield is at least 1% of all containers in the product. Improving  
105 the loading efficiency remains an area of study.

106 A challenge with using type-A containers is that they typically form with closed  
107 ends and therefore must be opened prior to sample insertion. The containers can be  
108 opened mechanically through milling or chemically by oxidizing agents such as hot  
109 solutions of  $\text{HNO}_3$  or other acids or by air at high temperature. If desired, the open ends  
110 of a container containing a sample can be closed by intentionally irradiating the ends  
111 (Ugarte, 1992).

112 Silicates are important for high-pressure geoscience studies but are difficult to put  
113 into nanocontainers. Specific procedures are needed for different minerals. Here we  
114 illustrate the use of olivine, for which capillary wetting works well for the type-A  
115 containers (Fig. 5). To fill CNFs with single-crystal olivine, we immersed open CNFs

116 into a sol created by mixing tetraethyl orthosilicate ( $\text{Si}(\text{OC}_2\text{H}_5)_4$ ), magnesian nitrate  
117 hexahydrate ( $\text{Mg}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ ), and 1-M nitric acid ( $\text{HNO}_3$ ) in a mass ratio of 1:2.4:36  
118 (Sanosh et al. 2010). The mixture was stirred for 12 hours, followed by programmed  
119 annealing of the CNFs sifted from the solution.

120 An alternate approach to filling CNFs (or CNTs) is to deposit graphitic layers  
121 directly onto nanosized mineral grains, thereby creating nanocontainers. The resulting  
122 type-B containers tightly encapsulate nanoparticles of interest when the latter are  
123 suspended freely in or passed through an atmosphere containing carbon vapor. It also  
124 seems that coating can occur when the particles move on the substrate, presumably  
125 because of thermal vibration. An advantage of the coating approach is that it avoids  
126 formation of defects during opening the closed type-A containers.

127 Type-B containers can be produced by the methods illustrated in Figure 3, of which  
128 CVD and laser ablation have been the most useful. To coat nanoparticles with graphitic  
129 carbon we either dispersed them onto TEM grids, which were then placed into the hot  
130 zone of a tube furnace, or we suspended the nanoparticles in a flowing inert gas, which  
131 transported them through the hot zone of the furnace.

132 When employing the arc-discharge method, a graphite rod impregnated with a  
133 target material or its components is used as an electrode. Carbon nanocontainers filled  
134 with the target material form in the discharge product. For example, we used a silicon-  
135 impregnated graphite rod in an arc and obtained moissanite ( $\text{SiC}$ )-filled CNOs. A

136 limitation of this procedure is that the target materials have to be stable under the high  
137 temperature (up to 3500 °C) and reducing carbon atmosphere in an arc.

138 Another way to create type-B containers is to take advantage of the fact that  
139 amorphous or poorly graphitic carbon is graphitized under electron radiation (Ugarte  
140 1992). Also, CVD, laser ablation, and arc discharge can all grow amorphous and poorly  
141 graphitic carbon within their reaction regions if the temperature is lower than ~600 °C.  
142 Yet one more convenient way to deposit amorphous carbon is to dip nanocrystals into  
143 an organic solvent such as acetone and then decompose this coating under electron  
144 radiation (Wu and Buseck 2013).

145 The preferable choice of method for loading minerals into carbon nanocontainers  
146 depends on the sample. Table 1 summarizes our sample-loading methods used to date  
147 and the minerals that have been loaded successfully.

#### 148 **Achieving high temperatures plus thermal and mechanical stability within an** 149 **electron microscope**

150 Sample heating to above ~300 °C is necessary for shrinking carbon nanocontainers  
151 during electron irradiation. At elevated temperatures, displaced carbon interstitials are  
152 sufficiently mobile to prevent their clustering and thus losing their ability to recombine  
153 with vacancies. As a result, self-rearrangement of the atomic structure around vacancies  
154 leading to shrinkage of these containers can occur. On the other hand, because of  
155 limited direct exposure to electrons of samples plus their good contact with conductive

156 carbon container walls, sample heating through electron radiation results in only small  
157 temperature changes. Using data from Williams and Carter (1996), the temperature  
158 changes are estimated to be <10 °C in our experiments.

159 Controlling sample temperature is important for high-pressure research. It is far  
160 easier to precisely control heating (up to 1500 °C) within an electron microscope than in  
161 high-pressure instruments such as MAPs and DACs. However, maintaining positional  
162 stability during rapid sample heating and cooling is both necessary and difficult for  
163 reliable TEM measurements. This concern arises because sample drift and stage  
164 vibration when changing the temperature can degrade high-resolution imaging and  
165 effectively inhibit TEM examination.

166 Controlling mechanical stability is especially problematic when using standard  
167 furnace-type TEM heating holders. The recent development of heating holders that  
168 utilize a MEMS (micro-electro-mechanical systems) design provides a good solution.  
169 Significantly improved sample stability and heating rate relative to standard stages can  
170 be achieved in this way (Fig. 6). More remarkably, even a substantial temperature  
171 increase of 300 °C caused no significant sample drift so that no sample translation was  
172 needed to record the image at 400 kX. Therefore, it is possible to capture features at  
173 atomic resolution in the sample without concern about instability-induced resolution  
174 loss. A MEMS holder also allows in-situ EDS analyses with a windowed detector at

175 temperatures up to ~800 °C, which is difficult or impossible with standard furnace-type  
176 holders.

### 177 **Compressing minerals through electron irradiation**

178 The rate of pressure increase during electron irradiation depends on the current  
179 density of the electron beam. Because electrons have a small mass, their displacement  
180 ability is relatively weak, with a cross section of 180 barn for interaction at 100 kV with  
181 carbon atoms (Cosslett 1978). As a result, electron-induced vacancy creation in the  
182 graphitic networks is slow, as is the pressure buildup. We used electron beams with  
183 current densities on the order of 10-100 A/cm<sup>2</sup> and irradiation durations of tens to  
184 hundreds of minutes.

185 Using such energetic electrons, radiation damage to the enclosed mineral samples  
186 can be a problem, especially for beam-sensitive silicates. One way to mitigate this  
187 situation is to minimize sample exposure to the electron beam. We achieved this goal by  
188 using the condenser-lens stigmators of the microscope to distort the electron beam into  
189 an elongated shape that was positioned along the walls of the carbon nanocontainer  
190 adjacent to the enclosed sample. An alternate way to achieve the same goal is to use a  
191 nanoprobe in STEM mode and then raster the beam along the walls along a pre-defined  
192 path under software control.

193 Figure 7 illustrates a forsterite nanocrystal that was compressed in a shrinking CNF.  
194 By comparing diffraction patterns recorded before and after compression, a ~3.8%

195 radial strain was measured in the nanocrystal. In light of its unknown crystallographic  
196 orientation, we approximate the Young's modulus and Poisson's ratio to those (~200  
197 GPa and ~0.24, respectively) for polycrystalline forsterite. The generalized Hooke's law  
198 then yields a radial pressure of ~10 GPa in the forsterite.

199 Whether the compression is hydrostatic or not depends on the morphology of the  
200 carbon nanocontainers. Shrinking tubular nanocontainers (CNTs and CNFs) provides  
201 non-uniform compression, whereas spherical containers (CNOs) generate relatively  
202 hydrostatic pressures. Both hydrostatic and directional (deviatoric) stresses can be  
203 useful experimentally, depending on the problem being addressed.

#### 204 **Measuring pressures**

205 As with many high-pressure studies, we used the equation of state (EoS) of the  
206 crystalline samples to estimate the pressure they experienced. Pressure-induced  
207 decreases in lattice spacings appear as increases in the distances of the relevant  
208 reciprocal spots from the central spot in either diffraction patterns or diffractograms  
209 derived from fast Fourier transforms (FFT) of HRTEM images (Fig. 8). The volume  
210 change estimated from the change in sample lattice spacing was then used in the Birch-  
211 Murnaghan EoS to calculate the pressure. Plots of pressure versus time or incident  
212 electron current density can also be attained by measuring temporal changes in lattice  
213 spacing at different beam current densities.

214       Uncertainties in pressure estimates using SAED patterns or FFT-derived  
215       diffractograms are dominated by systematic errors generated when comparing lattice  
216       spacings before and after compression. When the comparison is done for identical  
217       experimental conditions except pressure, the errors arise primarily from the finite sizes  
218       of the nanocrystal and CCD pixels. For the zincite case (Fig. 8), we estimate a ~16%  
219       uncertainty in the measurement of the lattice-spacing decrease, which corresponds to a  
220       ~20% uncertainty in the pressure estimate of ~18 GPa. An improved way of measuring  
221       internal pressure that we plan to explore is to take advantage of pressure-dependent  
222       features in electron-generated spectra such as cathodoluminescence from the sample  
223       under compression.

#### 224       **Observing minerals in situ under pressure**

225       Observation of pressurized mineral samples is the same as for TEM examination of  
226       materials at ambient conditions, i.e., high-resolution imaging and chemical analyses can  
227       be readily done while the pressure and temperature are maintained. More important,  
228       direct observation of defects and real-time TEM of phase transitions at elevated  
229       pressures become possible for the earth sciences for the first time. For example, Figure 9  
230       shows the real-time monitoring of anatase to  $\alpha$ -PbO<sub>2</sub>-type structure and concordant  
231       development of an associated stacking fault in titanium dioxide at >8 GPa and 770 °C.

232       High-pressure study of samples at low temperatures is also possible by cooling the



233 samples after a desired internal pressure is reached. A subsequent temperature change  
234 would have only a slight effect on the internal pressure maintained inside the  
235 containers. However, inevitable agglomeration of radiation-induced defects at low  
236 temperatures will lead to structural disorder in the container walls at a rate that  
237 depends on the incident beam current density, and this disorder causes a gradual  
238 relaxation in the internal pressure. Therefore, a relatively low electron dose is suggested  
239 for such TEM observations at low temperatures.

240 Despite its power, this technique also has limitations. The small sizes of the carbon  
241 containers, as well as the sample-thickness limit imposed by TEM, make it difficult to  
242 apply the technique to studies of coarse-grained samples. Also, high-pressure  
243 properties of nano-sized samples measured using this method may not represent those  
244 of their bulk counterparts, and possible modifications of the samples might occur  
245 because of the affinity of displaced carbon atoms to certain materials. Fortunately,  
246 carbon nanocontainers likely lose only a small fraction of their total carbon during  
247 electron irradiation before their walls become unstable and break up (Sun et al. 2006a),  
248 and carbon is incompatible with most geophysically important minerals.

## 249 APPLICATION TO EARTH SCIENCES

250 We used pressurization during TEM to address a long-standing geochemical  
251 question regarding the storage mode of carbon within mantle minerals (Wu and Buseck

252 2013). Using EFTEM, carbon concentrations of several atomic percent were detected  
253 along a stacking fault in titanium dioxide *in situ* at over 8 GPa and 770 °C (Fig. 10). Such  
254 measurement is impossible using other currently available high-pressure techniques.  
255 Thus, this procedure helped in the discovery of a new mechanism for hosting mantle  
256 carbon by segregation to crystal defects in an analog of a mantle mineral. This result has  
257 potential applications for geoscientists and others concerned with mantle geochemistry  
258 and geophysics, as well as with the deep carbon cycle.

## 259 IMPLICATIONS

260 High-pressure measurements in the earth and materials sciences have been  
261 expanded from DACs, MAPs, and shock experiments to include transmission electron  
262 microscopes using carbon nanocontainers as sample containers and presses. The  
263 advantage of the TEM approach is that dynamic changes can be observed in progress  
264 and, perhaps more importantly, materials can be examined at almost the atomic scale  
265 while at elevated pressure and temperature. The technique provides a new way of  
266 directly observing mineral structures and reactions at elevated conditions within an  
267 electron microscope, so that changes that occur deep within Earth or other planets can  
268 be studied in detail at high spatial resolution. Also, being able to observe materials *in*  
269 *situ* removes problems induced upon quenching if one wants to observe details of non-  
270 quenchable phases. Finally, determining the characteristics and properties of individual

271 nanograins using this technique can be extremely useful, whereas it is difficult with  
272 other methods currently available. A result is that the versatile power of TEM under  
273 ambient conditions that has long been familiar to earth scientists can now be extended  
274 to similar studies at elevated pressures and temperatures.

275

276

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283

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332 analogues. *Nature Geoscience*, 6, 875-878.
- 333

334 **TABLE 1.** Minerals placed into carbon nanocontainers and the methods used

335

Loading Methods		Loaded Minerals
Type-A containers	Capillary wetting by molten salts	molybdate ( $\text{MoO}_3$ ), shcherbinaite ( $\text{V}_2\text{O}_5$ )
	Capillary introduction of soluble materials in solution, followed by annealing	magnetite ( $\text{Fe}_3\text{O}_4$ ), bunsenite ( $\text{NiO}$ ), zincite ( $\text{ZnO}$ ), yttriaite ( $\text{Y}_2\text{O}_3$ ), olivine ( $\text{Mg}_2\text{SiO}_4$ )
Type-B containers	Chemical vapor deposition with dispersed nanoparticles	olivine, anatase ( $\text{TiO}_2$ ), rutile ( $\text{TiO}_2$ ), barioperovskite ( $\text{BaTiO}_3$ ), perovskite ( $\text{CaTiO}_3$ ), tausonite ( $\text{SrTiO}_3$ )
	Laser ablation with suspended nanoparticles	olivine
	Arc discharge of target or precursor materials	cohenite ( $\text{Fe}_3\text{C}$ ), iron ( $\text{Fe}$ ), moissanite ( $\text{SiC}$ ), naquite ( $\text{FeSi}$ )

336

337 **FIGURE 1.** Shrinkage of a carbon nanotube (CNT) that is caused by structural reorganization

338 around vacancies in its graphitic networks. (Adapted from Sun et al. 2006a)

339

340 **FIGURE 2.** Sketch of structures of some carbon nanocontainers. (a) CNT. (b) CNF. (c) CNO.

341

342 **FIGURE 3.** High-temperature synthesis apparatuses for producing carbon nanocontainers. (a)

343 Arc-discharge system and (b) tube-furnace assembly used for CVD and laser ablation. The

344 green light is from the laser beam. Both units were built at Arizona State University for growing

345 CNTs and CNOs.

346

347 **FIGURE 4.** Methods for loading carbon nanocontainers. (a) Magnetite ( $\text{Fe}_3\text{O}_4$ ) nanocrystals

348 within a CNF. The mineral was placed into the CNFs by aqueous wetting in a  $\text{Fe}(\text{CH}_3\text{COO})_2$ -

349  $\text{CH}_3\text{COOH}$  solution followed by annealing in argon at 600 °C for 2 hours. The composition was

350 confirmed by the EELS spectrum, in which the oxygen and iron edges are indicated. (b)

351 Shcherbinaite ( $\text{V}_2\text{O}_5$ ) nanocrystal within a CNT. The mineral was introduced into pre-opened

352 CNTs by capillarity from the molten state. The composition was confirmed by its EDS spectrum.

353 The C peak in this and other panels is from the CNT, and the Cu peak is from the TEM grid. (c)

354 Bunsenite ( $\text{NiO}$ ) nanocrystal within a CNT. The mineral was placed into CNTs by aqueous

355 wetting in a  $\text{Ni}(\text{NO}_3)_2$ - $\text{HNO}_3$  solution followed by annealing in air at 500 °C for 2 hours. The

356 composition was confirmed by the EDS spectrum; the indicated fringe spacings match those of

357 bunsenite (111) and (020).

358

359 **FIGURE 5.** Forsterite ( $\text{Mg}_2\text{SiO}_4$ ) nanocrystal within a CNF. Aqueous wetting of the CNF in a

360  $\text{Si}(\text{OC}_2\text{H}_5)_4$ - $\text{Mg}(\text{NO}_3)_2$ - $\text{HNO}_3$  solution was followed by annealing in air at 900 °C for 2 hours. The

361 composition was confirmed by the EDS spectrum. The C peak is from the CNT and the Cu peak

362 is from the TEM grid.

363

364 **FIGURE 6.** Time sequence of a cohenite ( $\text{Fe}_3\text{C}$ ) nanocrystal in a contracting CNT. The panels

365 track changes occurring to the crystal during heating for 190 s at 500 °C, followed by an



366 essentially instantaneous increase to 800 °C. The morphology change of the cohenite crystal,  
367 which resulted from surface-energy minimization of the carbide, occurred in less than 1 s.  
368 Extended heating at that temperature resulted in additional CNT contraction and consequent  
369 pressurizing of the encapsulated nanocrystal. There is a magnification increase between the two  
370 panels at the bottom right. The images were recorded with a MEMS-based (Protochips Aduro)  
371 heating holder within a Philips CM200 field-emission-gun electron microscope operated at 200  
372 kV. The electron current densities were ~100-250 A/cm<sup>2</sup>.

373

374 **FIGURE 7.** Compression series of an olivine nanocrystal subjected to non-hydrostatic stress in a  
375 CNF at 410 °C. The compression directions are indicated by the increase in the vertical  
376 dimension of the nanocrystal and decrease in its horizontal dimension. To minimize radiation  
377 damage to the olivine, the electron beam was distorted into a thin band and then positioned  
378 alternately along the outer edge of the CNF wall on opposite sides of the nanocrystal, as  
379 illustrated by the blue dashed ellipses.

380

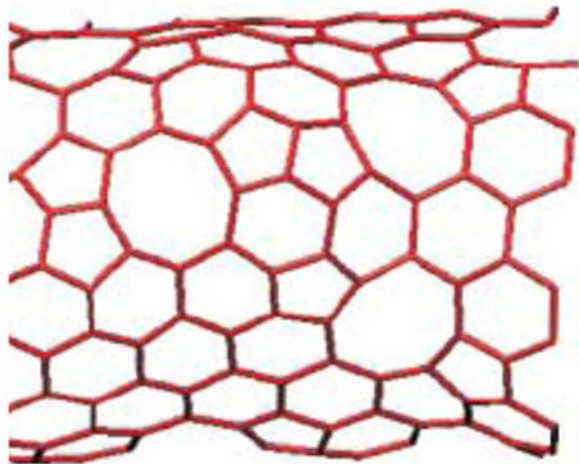
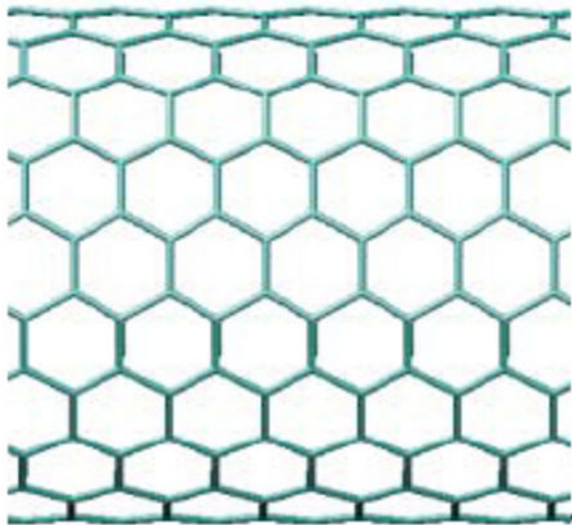
381 **FIGURE 8.** Radiation-induced compression of a zincite (ZnO) nanocrystal in a CNT at 500 °C. **(a)**  
382 The crystal prior to irradiation. **(b)** Same crystal after irradiation for 5 minutes. **(c)**  
383 Superimposition of FFTs from **(a)** in yellow and **(b)** in green. Compression is indicated by an  
384 approximately 6% decrease in the (100) lattice spacing.

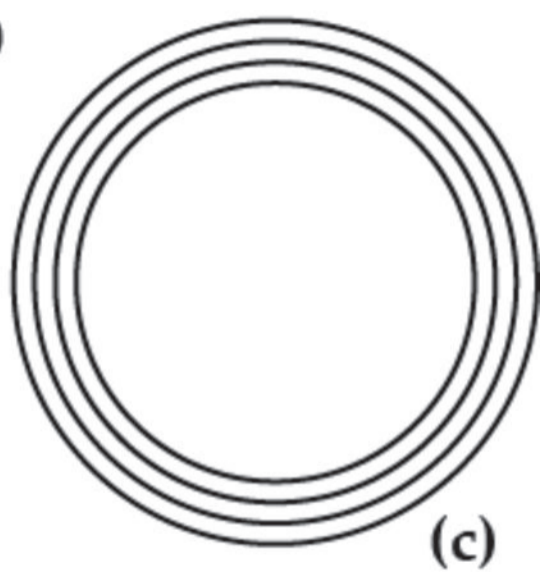
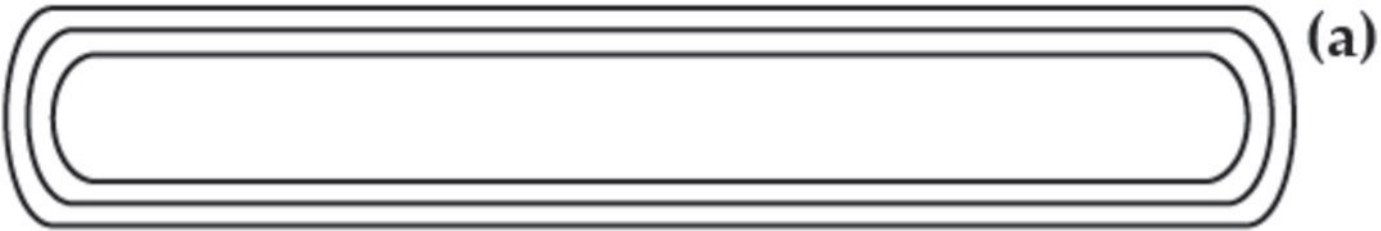
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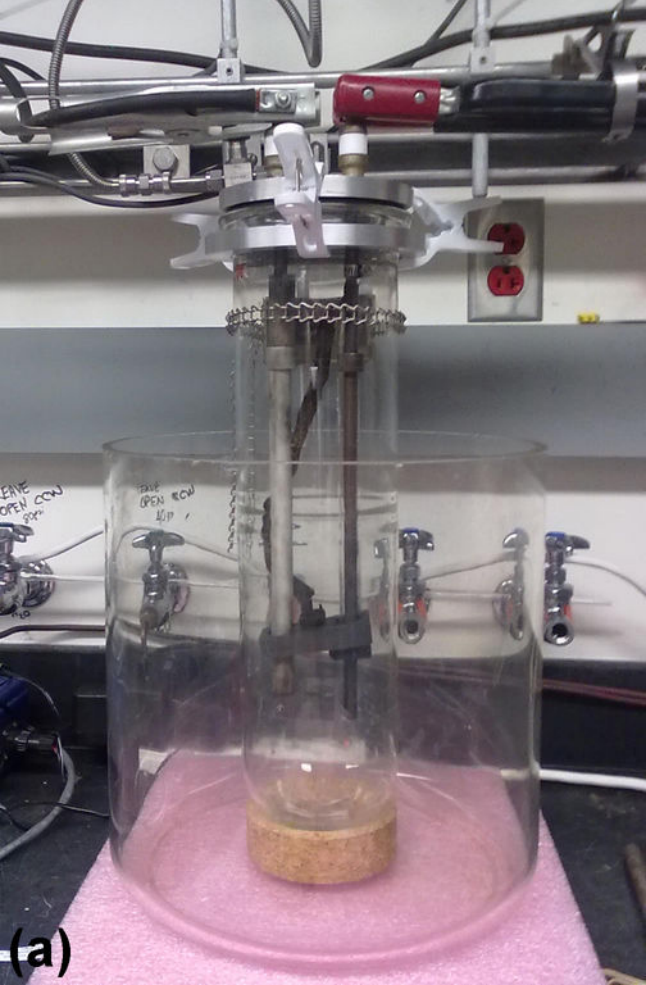
386 **FIGURE 9.** Compression sequence of an anatase nanocrystal enclosed within a CNO at (a) 0, (b)  
387 10, (c) 13, (d) 26, (e) 40, (f) 43, (g) 54, and (h) 58 minutes after the start of electron irradiation. (i)  
388 The Fourier-filtered image from the dashed region in (h) showing a stacking fault in the TiO<sub>2</sub>.  
389 Pressure was generated by shrinkage of the CNO at 770 °C (using ~30 A/cm<sup>2</sup> electrons).  
390 Disappearance of anatase (101) planes and emergence of α-PbO<sub>2</sub>-type TiO<sub>2</sub> (110) planes indicate  
391 a phase transition between (e) and (f) (cf. insets). The stacking fault is indicated by a line  
392 segment in (h) and an arrow in (i). (Adapted from Wu and Buseck 2013)

393

394 **FIGURE 10.** EFTEM images showing appreciable carbon signals at the location of the stacking  
395 fault in the pressurized TiO<sub>2</sub> nanocrystal. The fault is illustrated in the upper inset. (a) Carbon  
396 map. (b) Jump-ratio image. The yellow boxes highlight the faint carbon intensities. The line  
397 profile in the lower inset shows the intensities integrated along the arrowed direction within the  
398 box in (a), with the central peak representing the carbon signals along the stacking fault. (From  
399 Wu and Buseck 2013)





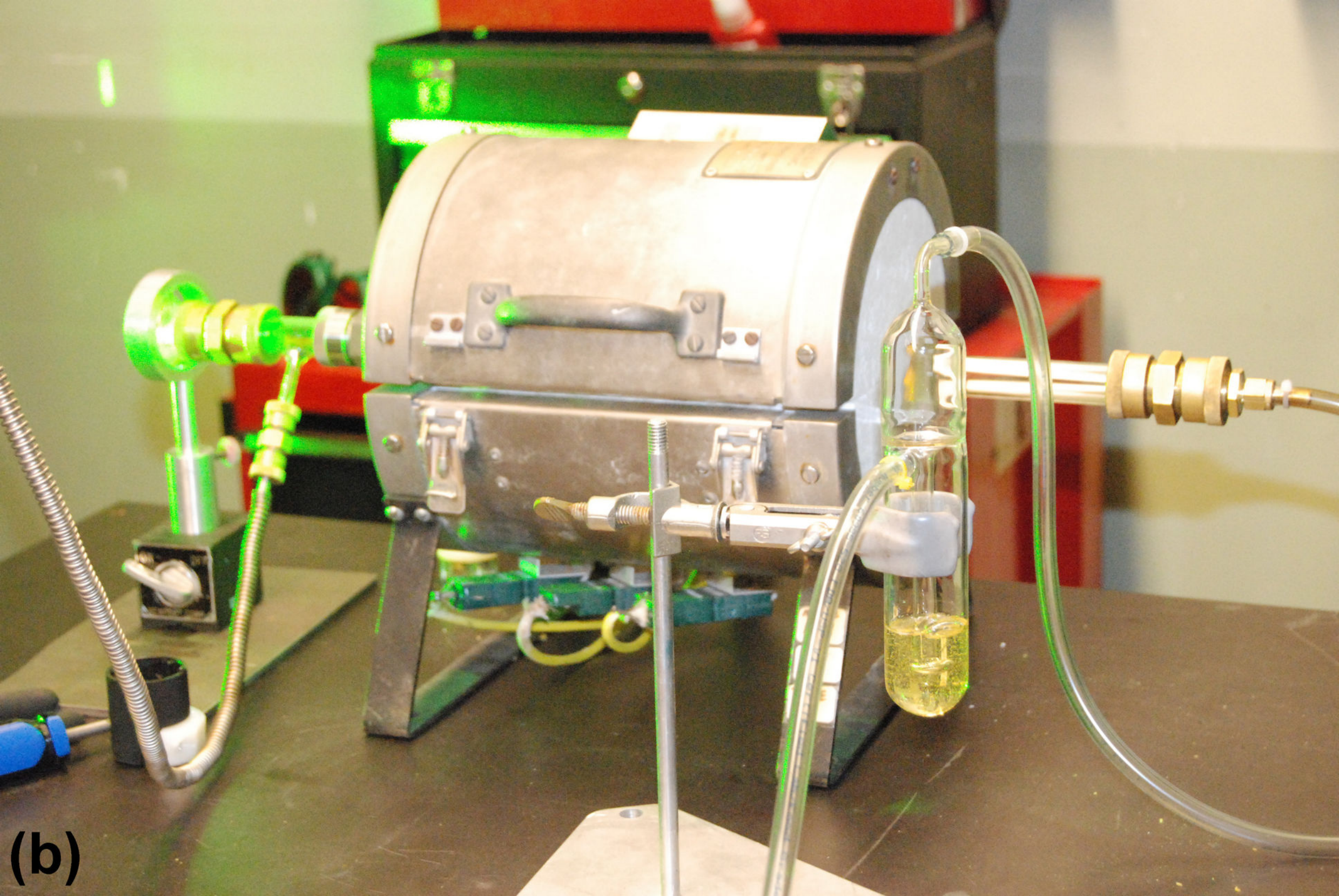


LEAVE  
OPEN CW  
8/2/10

FAVE  
OPEN CW  
10/17

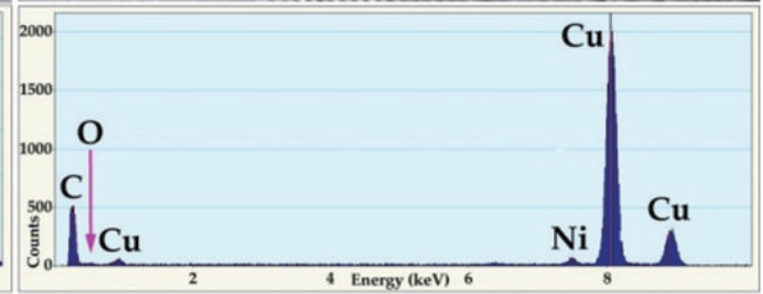
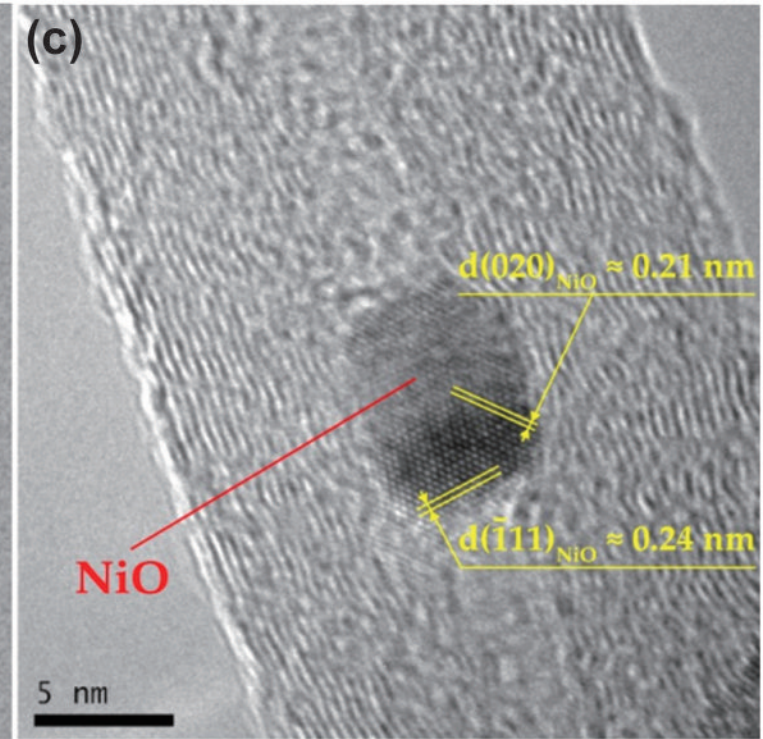
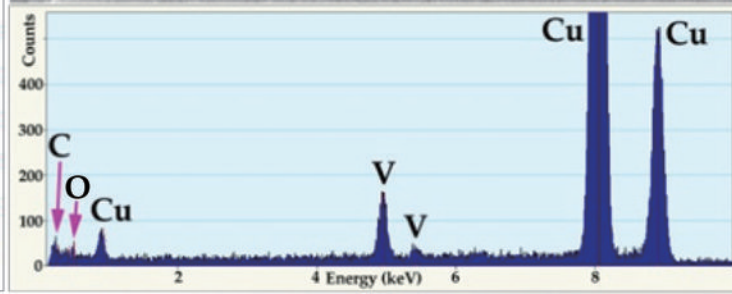
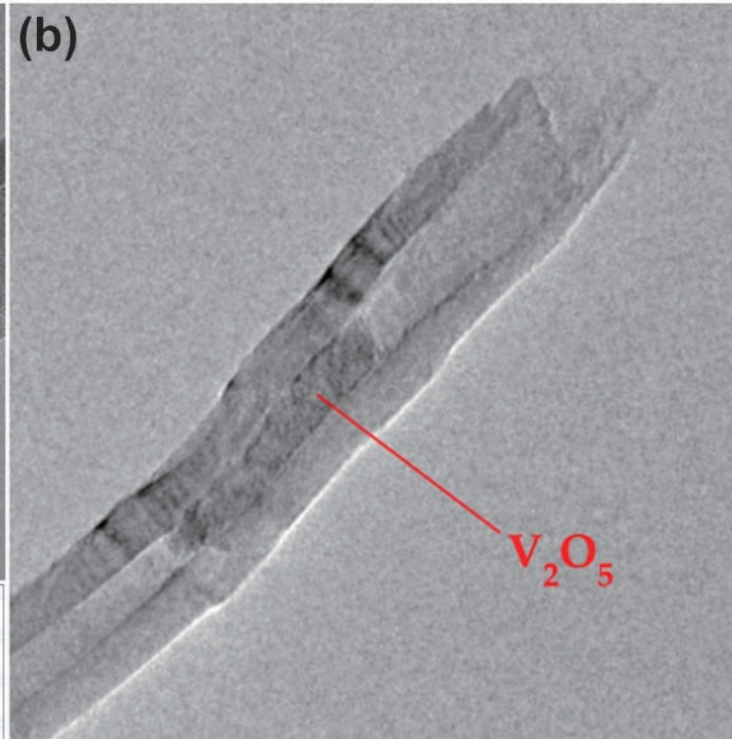
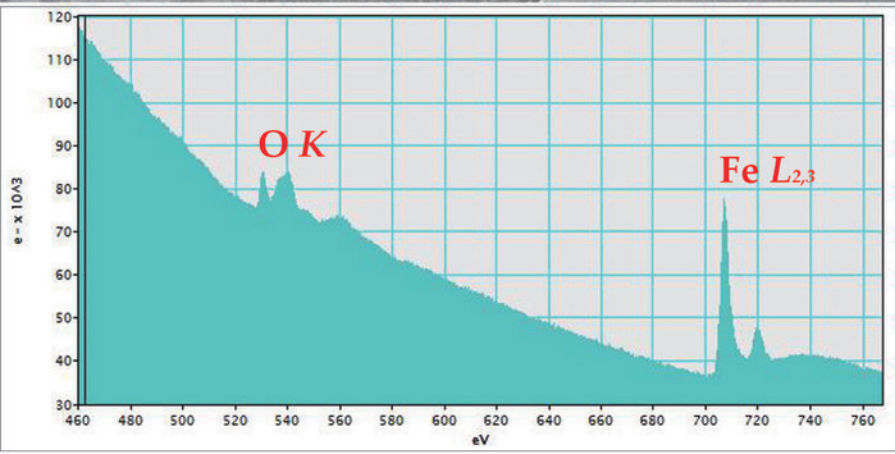
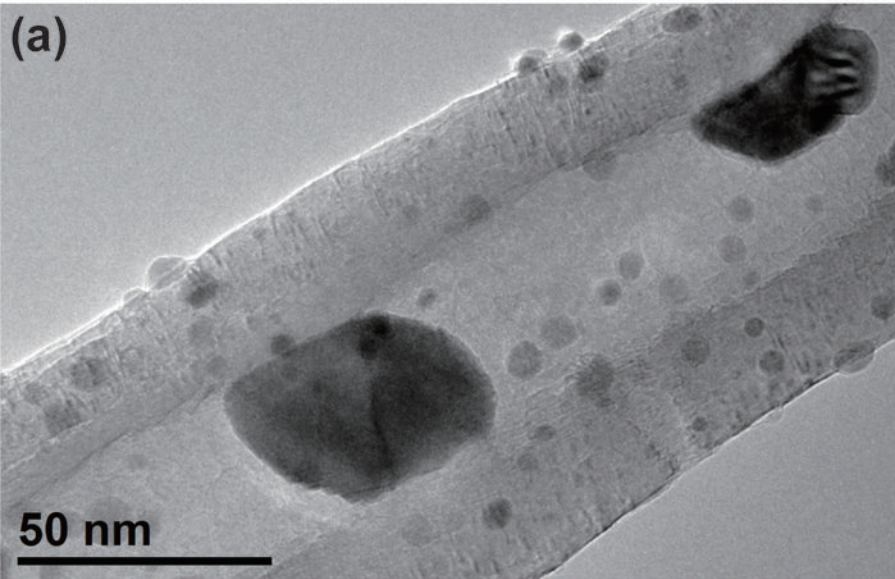
(a)



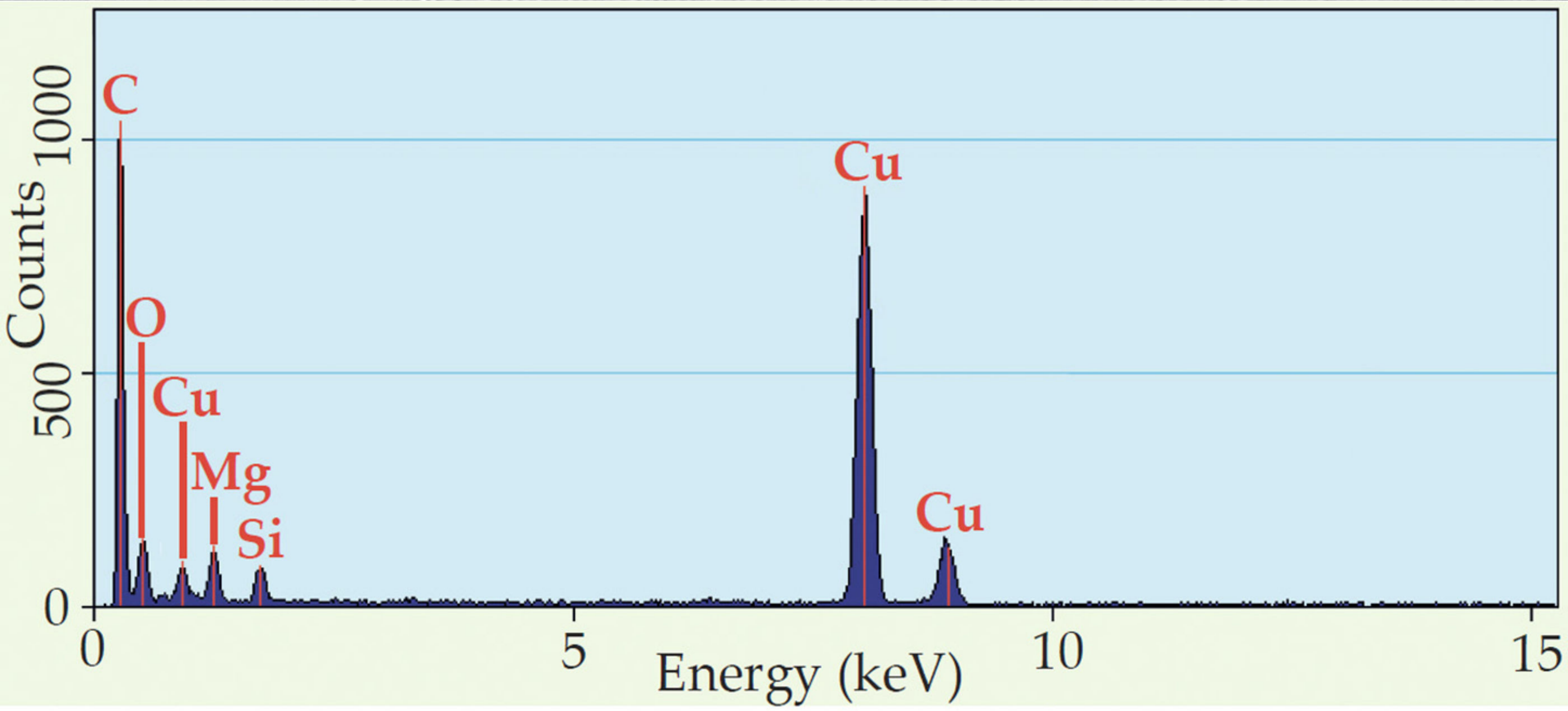
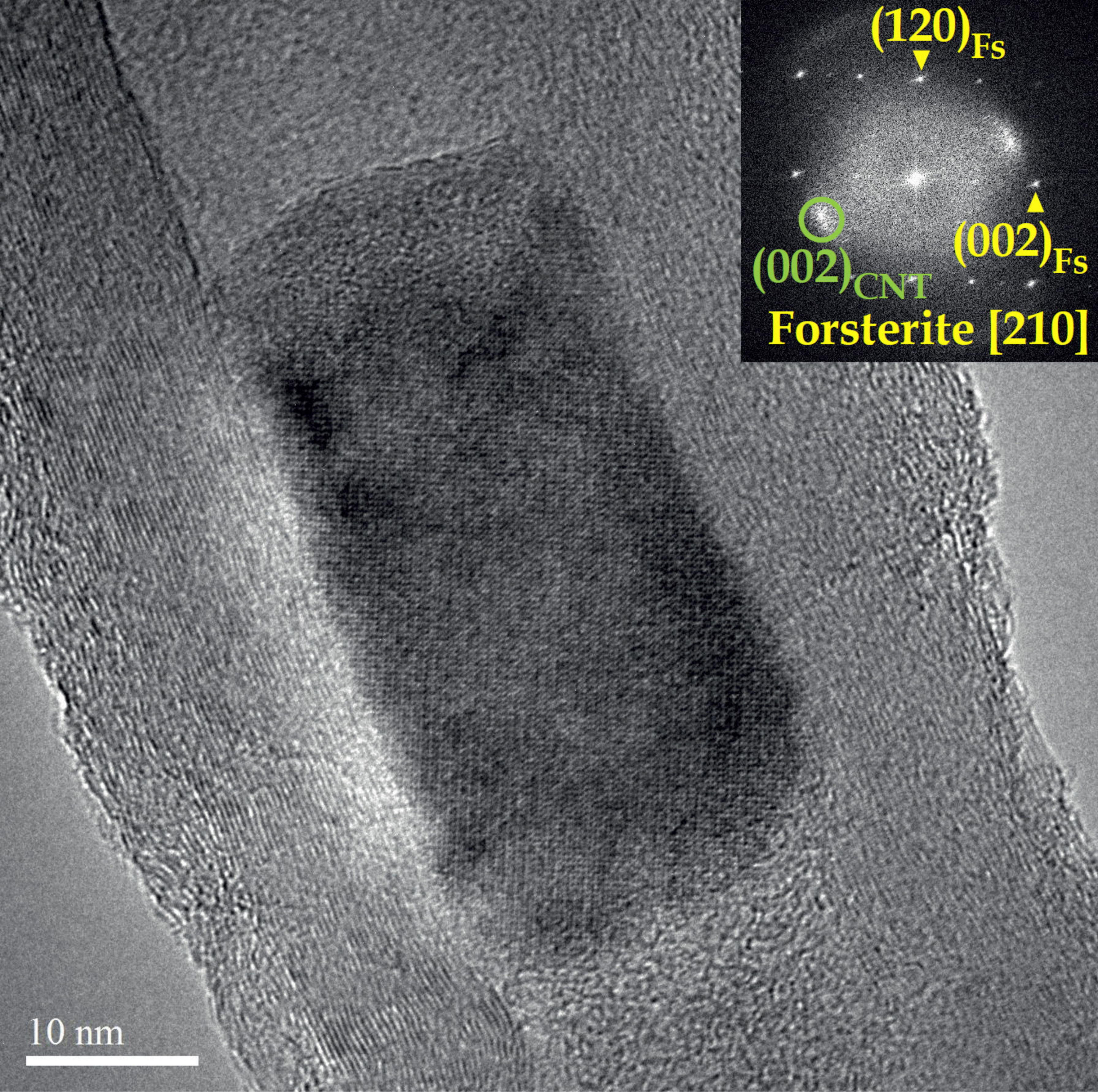


(b)

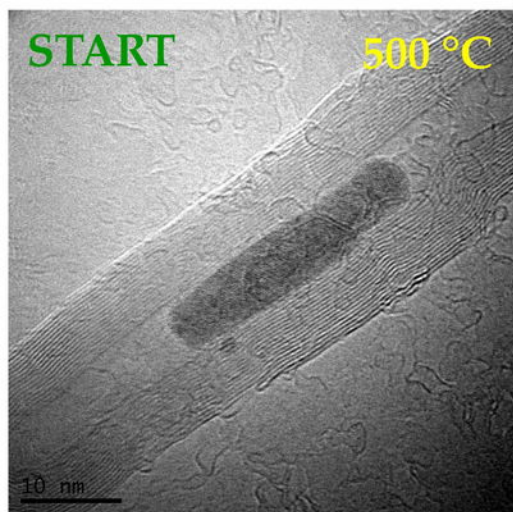




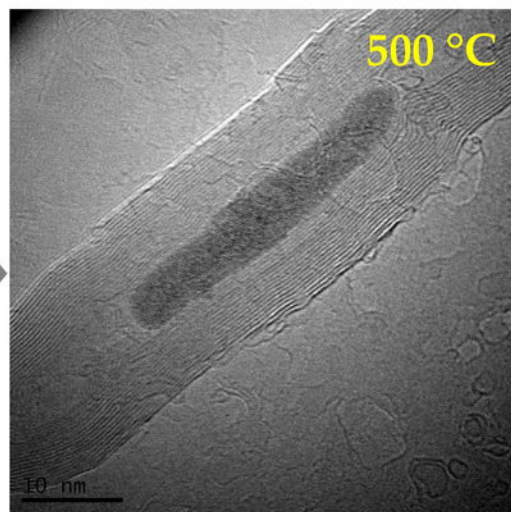




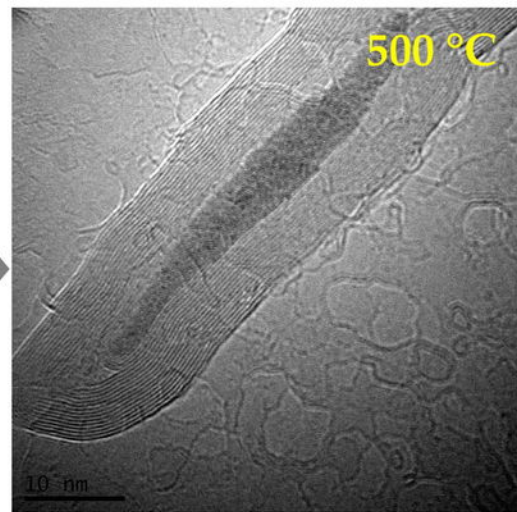




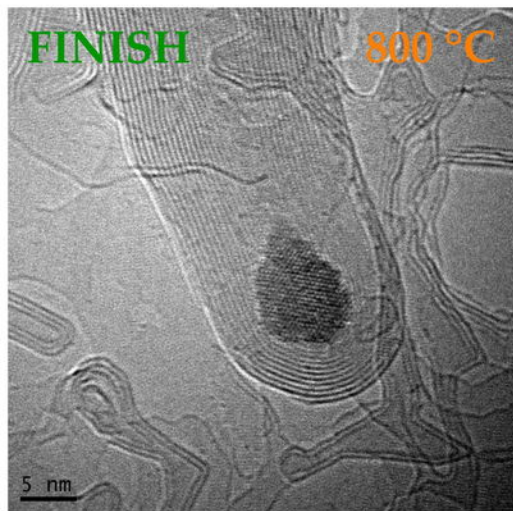
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seconds



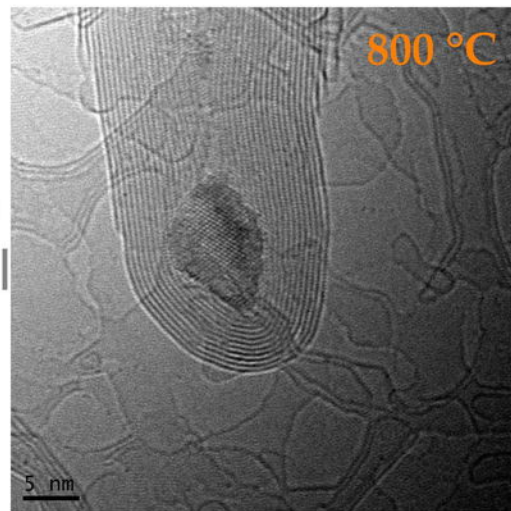
160  
seconds



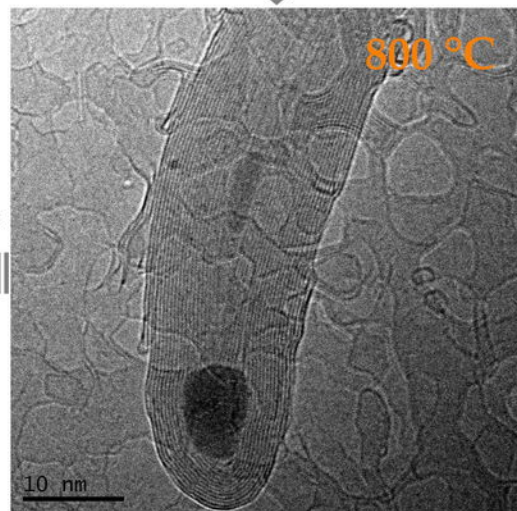
<1 SECOND!



60  
seconds

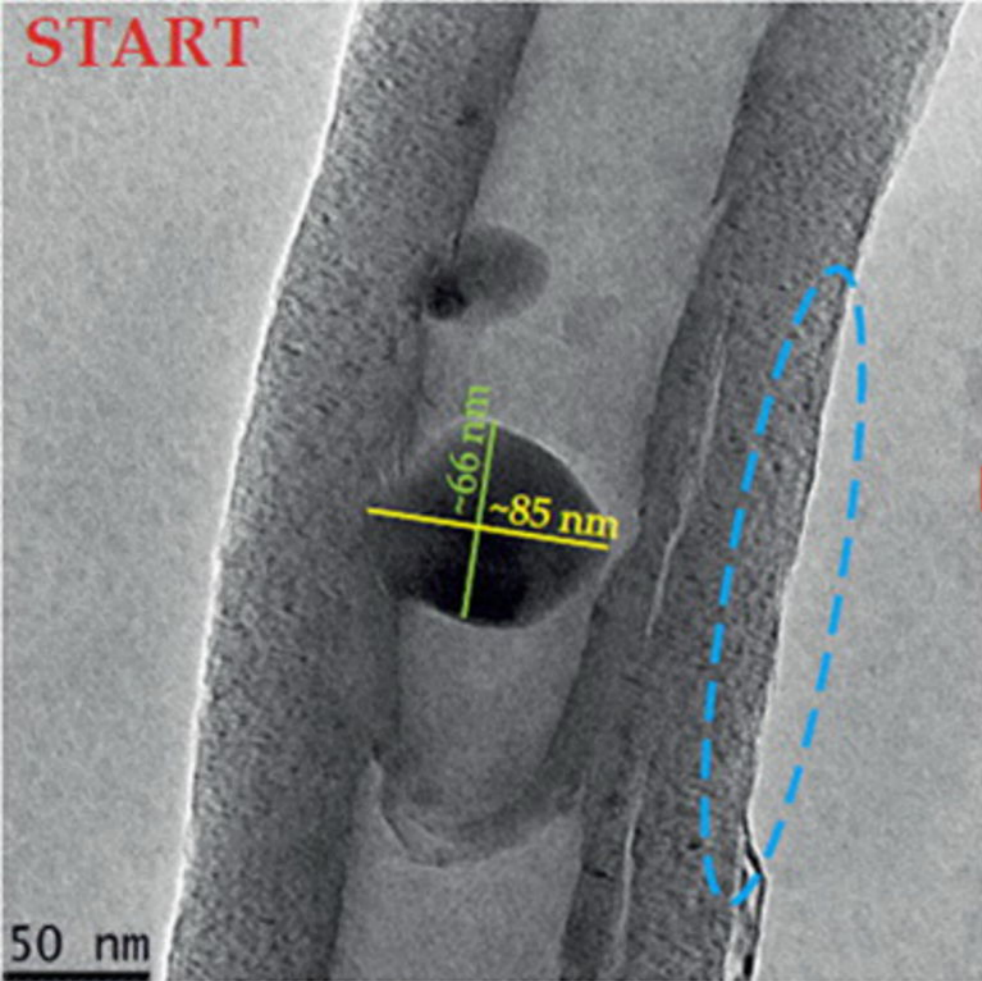


600  
seconds

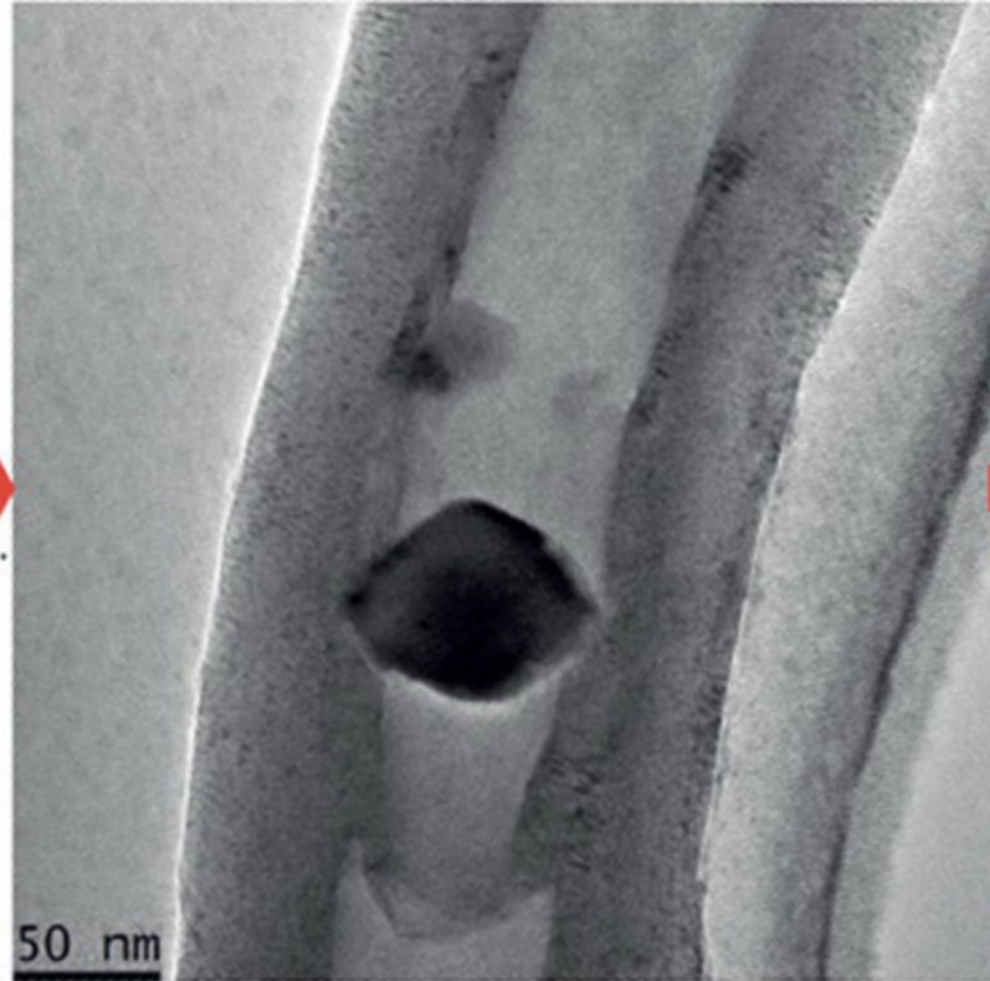




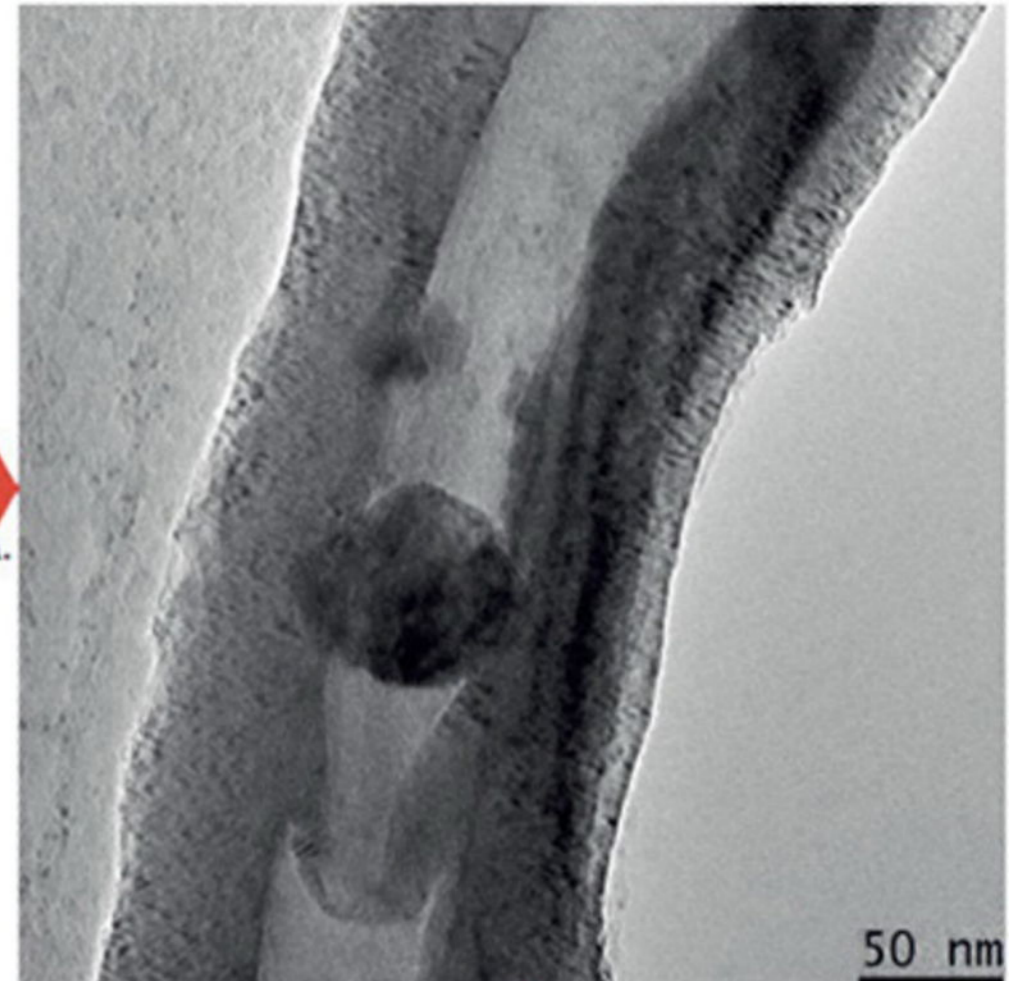
START



17  
min.

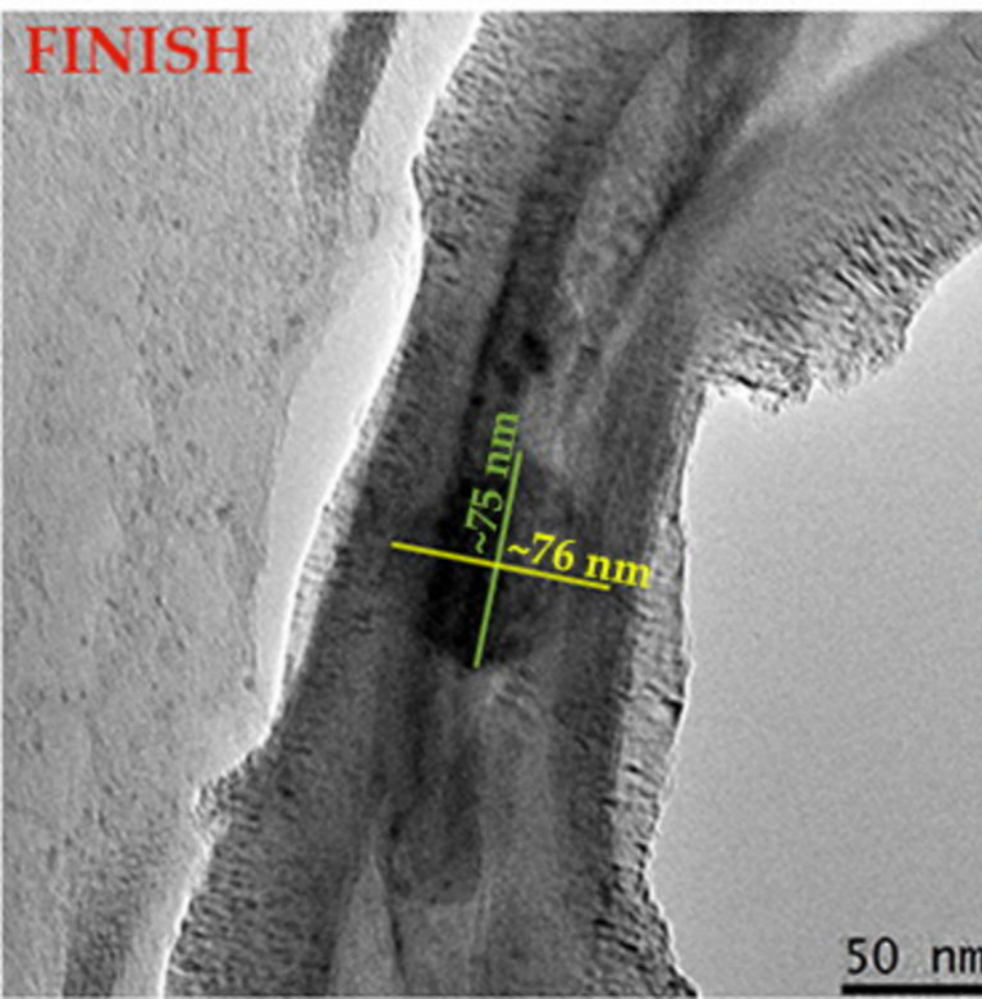


59  
min.

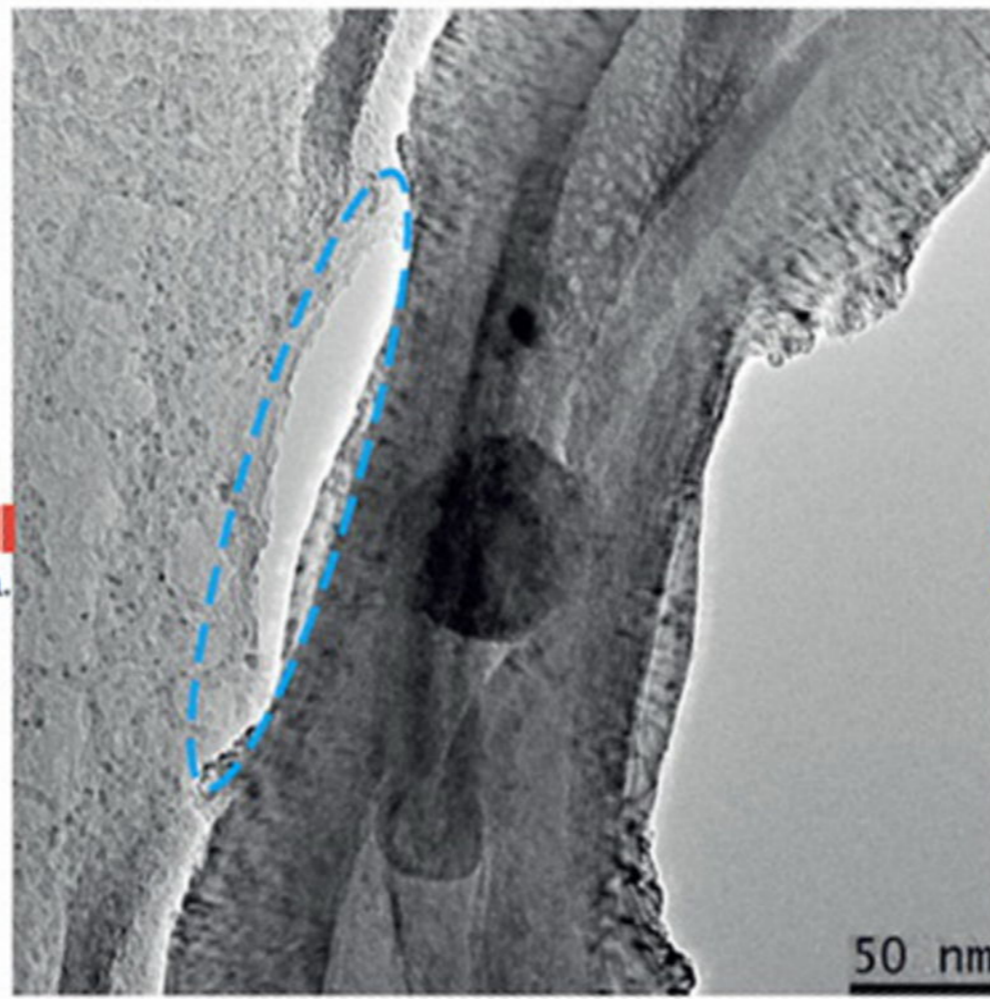


100 min.

FINISH



34  
min.



37  
min.

