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1	H/D Isotope Effects in Brucite at Low-Temperatures
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12	Abstract

13 Joint refinement of powder neutron diffraction data of hydrogenated and 14 deuterated brucite over the temperature range of 10 - 295K clearly shows significant 15 isotope effects in the structural parameters. $Mg(OH)_2$ has a 0.31% larger volume than 16 that of $Mg(OD)_2$ at room temperature, which is mostly due to the *c*-axis expansion of 17 $Mg(OH)_2$ as compared to $Mg(OD)_2$. The isotope affect in the *a*-axis has the opposite, but 18 smaller, behavior as compared to the c-axis. These differences are slightly enhanced with 19 reduction of the cell volume upon cooling. The temperature dependence of the isotropic 20 atomic displacement parameters (ADPs) with the single site model show the ADP of the 21 H atom is approximately larger than that of the D atom by the amount expected from the 22 reduced-mass difference, but this difference is not evident with the split site model. 23 Despite the shorter *c*-axis of the deuterated form, nearest-neighbor D...D distances are 24 longer than the H...H distances, because the O-H distances are longer than the O-D 25 distances, which necessarily places the H atoms closer together by 0.03 Å or more within 26 the interlayer space. This study showcases an example of when a joint Rietveld 27 refinement is ideally suited, by combining data for the deuterated and hydrogenated 28 forms of brucite. The approach reduces the number of least-squares variables, and

reduces the systematic errors. It can be a general method to analyze isotope effects inmaterials studied by neutron diffraction.

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32 Introduction

33 Brucite has been the subject of a number of parametric structural studies owing 34 both to its importance as a building block of hydrous sheet silicates, particularly clay 35 minerals, and as a model system for more complicated rock-forming minerals with both 36 strong and weak chemical bonding. Both high pressure and variable temperature studies 37 for brucite-type compounds have been made using diffraction (Catti et al. 1995; 38 Chakoumakos et al. 1997; Desgranges et al. 1996; dos Santos et al. 2010; Fukui et al. 39 2003; Horita et al. 2010; Jayachandran and Liu 2006; Nagai et al. 2000a, 2000b; Parise et 40 al. 1994; Shim et al. 2006; Xu et al. 2007a, 2007b), Brillouin scattering (Jiang et al. 41 2006), inelastic neutron scattering (Chakoumakos et al. 1997; Kazimirov et al. 2010), 42 vibrational (Raman and infrared) spectroscopy (Shinoda et al. 2002; Shieh and Duffy 43 2002; Shim et al. 2006; Speziale et al. 2005), and a variety of theoretical modeling (Mitev 44 et al. 2009; Azuma et al. 2011; Braterman and Cygan 2006; Hermansson et al. 2008, 45 2009; Civalleri et al. 2009; Jochym et al. 2010; Mookherjee and Stixrude 2006; 46 Ulgliengo et al. 2009; Kirkpatrick et al. 2005; Reynard and Caracas 2009; Sainz-Diaz et 47 al. 2000, Tosoni et al. 2005). We previously undertook a neutron powder diffraction study of the H/D effect in brucite as a function of pressure (Horita et al. 2010). Here we 48 49 report an ambient pressure study of the same effect as a function of temperature using 50 higher resolution data. The present work corroborates the previous findings and gives a 51 more detailed view of the isotope effect and thermal contraction mechanism in brucite.

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53 Experimental Methods

54 After removing fines by elutriation in water, MgO powder (99.95% purity on 55 metal basis, Alfa Aesar) was heated in vacuum at 900° C for 2 days to remove traces of 56 water. The dehydrated MgO (10 g) was loaded into a Teflon cup along with 80 g of either de-ionized water (D/H ~ 0.015%) or heavy water (99.9% purity, Cambridge Isotope 57 58 Laboratory), then sealed inside a 300 mL reaction vessel, and heated to 300° C (ca. 8.5 59 MPa vapor-saturated pressure) for 2 weeks to hydrothermally synthesize the hydrogenated or deuterated brucite. Previous X-ray and neutron powder diffraction and 60 61 Raman spectroscopy confirmed the product of these reactions to be pure brucite with 62 greater >99% of deuteration for the $Mg(OD)_2$ sample (Horita et al. 2010).

63 Neutron powder diffraction measurements were performed at the HB-2A high 64 resolution powder diffractometer located at the High Flux Isotope Reactor at Oak Ridge 65 National Laboratory (Garlea et al. 2010). Powdered samples were loaded in vanadium 66 cans and placed in a top loading low temperature cryostat. The wavelength of 1.1176 Å produced by a Ge [115] monochromator was used and Soller collimation of 12', 21', 6', 67 68 before the monochromator, before the sample, and before the detectors, respectively. For 69 each sample a data set over the 2θ range $10 - 132^{\circ}$ was collected with a step size of ~ 0.067°, counting 60s per step, at the temperatures 20, 30, 40, 50, 60, 70, 80, 90, 100, 120, 70 71 150, 200, and 295 K.

Rietveld refinements were made using the FullProf suite (Rodriguez-Carvajal Scattering lengths used were 5.375, 5.803, 6.671, -3.739 fm for Mg, O, D, and H, respectively. To reduce the number of least squares variables and increase the amount of

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75	data, a combined refinement using both neutron datasets of the hydrogenated and
76	deuterated samples was undertaken. This approach also should lessen the effects of any
77	systematic errors. The resolution of these data was much higher than those from our
78	previous high pressure study on H/D effects in brucite (Horita et al. 2010). Previous
79	studies have employed anisotropic, split-site, or anharmonic to account for the large
80	displacements of the H/D atom (Desgranges et al. 1993, 1996; Parise et al. 1994;
81	Chakoumakos et al. 1997). Here, we have used a split-site model for the H/D atom
82	position.
83	Brucite has space group symmetry $P-3m1$, with atomic positions: Mg 0,0,0; O
84	$\frac{1}{3},\frac{2}{3},z$; and D/H $\frac{1}{3},\frac{2}{3},z$ for the single site model and $\frac{1}{3},y,z$ for the split-site model.
85	Parameters refined in common included the peak shapes and the Mg and O isotropic
86	atomic displacement parameters. Parameters refined for each dataset included lattice
87	parameters; O, D, or H atom y and z coordinates; D or H isotropic atomic displacement
88	parameters; diffractometer zero; preferred orientation, scale factor; and background. The
89	peak shape function used was the Thompson-Cox-Hastings pseudo-Voigt function with
90	axial divergence asymmetry. Asymmetry parameters S/L and D/L were fixed at 0.034
91	and 0.035 respectively, and two additional peak asymmetry parameters were refined.
92	The backgrounds were defined by 3 rd order and 8 th order polynomials with their origin at
93	$60^{\circ} 2\theta$, for the deuterated and hydrogenated data, respectively. The displacements of all
94	atoms were refined isotropically. The deuterium content of the deuterated sample was
95	fixed at unity, based on previous Raman spectroscopy and neutron diffraction analysis.
96	Both single-site and split-site models for the H/D position were made and compared.
97	

98 **Results**

99	Typical neutron powder diffraction patterns are shown for hydrogenated and
100	deuterated brucite in Figure 1. The principal difference between the two patterns, other
101	than the different intensities due to the different scattering factors of the D and H atoms,
102	is the much larger background of the hydrogenated sample due to the nuclear spin
103	incoherent scattering contribution from the H atoms.
104	The temperature dependence of the lattice parameters over the range $20 - 295$ K
105	shows Debye behavior (Fig. 2). The c -axis is smaller for the deuterated brucite as
106	compared to the hydrogenated brucite for all temperatures and this difference increases
107	slightly upon cooling, $\Delta c(295 \text{K}) = 0.015 \text{ Å}$, $\Delta c(20 \text{K}) = 0.02 \text{ Å}$. The <i>a</i> -axis shows the
108	opposite behavior, but the difference is more than an order of magnitude smaller,
109	$\Delta a(295 \text{K}) = -0.001 \text{ Å}, \Delta a(20 \text{K}) = -0.0005 \text{ Å}.$ The net result is that the volume of the
110	deuterated brucite is smaller than that of the hydrogenated brucite at 295K by 0.15 Å ³ , or
111	0.31%, which compares well with the room temperature difference, $0.34%$, determined
112	by X-ray diffraction (Horita et al. 2010).
113	The temperature dependence of the isotropic atomic displacement parameters
114	(ADP) exhibits the expected variation (Fig. 3), except that the values for the oxygen atom
115	are lower than those for the magnesium atom. The ADP of the H atom is larger than that
116	of the D atom by approximately the amount expected from the difference in their reduced
117	masses $\left(\frac{1\times16}{1+16}\right)$ and $\frac{2\times16}{2+16}$ for O-H and O-D, respectively) within a harmonic oscillator model.
118	For the split-site model, the ADP's of H and D coincidentally collapse to common values
119	(Fig. 3b).

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120	The O-D, O-H, and MgO bond lengths of deuterated and hydrogenated brucite
121	show essentially little or no temperature dependence (Fig. 4 and 5). This follows from the
122	lack of any significant temperature dependence of the z-coordinates of the O, D, and H
123	atom positions (Fig. 6). The refined values of the z-coordinates of the O and H atoms for
124	the hydrogenated brucite show a clear model dependence, which gives rise to model
125	dependent Mg-O distance (Fig. 5). The room temperature value of the $z(D)$ is smaller
126	indicating a small increase in $z(D)$ upon cooling from room temperature to ~100K. The
127	z(O) of Mg(OD) ₂ also show a small increase with decreasing temperature. Although, the
128	z(O) was refined independently for each of the hydrogenated and deuterated data sets, the
129	resulting same values indicate that this parameter could have also been held in common
130	for the joint refinement. At 295K the O-D bond length is 0.965 Å as compared to the O-
131	H bond length of 1.025 Å. The O-D distance compares favorable with the average value
132	of 0.95 Å for 131 O-D bonds in the Inorganic Crystal structure Database (ICSD)
133	determined by single-crystal neutron diffraction. The ICSD values are uncorrected for
134	thermal motion, so an additional shortening of \sim 0.01- 0.02 Å is expected. It is known
135	generally that O-D distances are shorter than O-H distances, including liquid water: O-H,
136	0.991 and O-D, 0.970 Å (Ichikawa et al. 1991; Silvestrelli and Parrinello 1999).
137	The nearest-neighbor HH distances are essentially constant within error over
138	the temperature range $20K - 300K$ (Fig. 7). In contrast, the DD distances show a
139	systematic shortening upon cooling and plateau below 100K. Overall the DD
140	distances are longer than the HH distances by 0.03 Å or more. This might seem
141	counter-intuitive given that the volume of the deuterated brucite is smaller than the
142	hydrogenated brucite, but the O-H distances are longer than the O-D distances so the

143	HH distances are necessarily closer within the interlayer space. In addition, a small
144	contribution also comes from the a-axis of the deuterated brucite being slightly longer
145	than that of hydrogenated brucite.

146 The temperature dependence of the splitting of the H/D site basically follows the

same behavior as the atomic displacement parameters (see Fig. 3). The magnitude of the

splitting is larger for the H site as compared to the D site because the H atomic

149 displacement is larger, and the splitting decreases with temperature because the thermal

150 vibration decreases.

151 Temperature dependence of the octahedral layer thickness and the interlayer

152 thickness in deuterated brucite shows that the octahedral layer thickness remains

153 constant, whereas the interlayer thickness contraction accommodates the observed *c*-axis

154 contraction (Fig. 8). This behavior is expected given the weak Van der Waals interaction

between the layers, but is opposite to that reported by Chakoumakos et al. (1997) from an

analysis of t-o-f neutron diffraction data. It is evident that the previous study by

157 Chakoumakos et al. (1997) on the low temperature behavior of deuterated brucite using t-

158 of-f neutron diffraction data has some systematic errors in the refinement of the atomic

159 positions. Even though the data used here does not extend to as high resolution, the

160 temperature dependence of the structural changes is more crystal chemically sensible.

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162 Discussion

163 This study showcases an example of when a joint Rietveld refinement is ideally 164 suited, by combining data for the deuterated and hydrogenated forms of brucite. The 165 approach reduces the number of least-squares variables, and reduces the systematic

166	errors. It can be a general method to analyze isotope effects in materials studied by
167	neutron diffraction. The D/H isotope effect on the lattice parameters of brucite is easily
168	measured by either neutron or X-ray diffraction, but structural details beyond the lattice
169	dimensions require the use of neutron diffraction.
170	The model dependence, albeit small, of the refined z-coordinates for the H and O
171	positions for the hydrogenated brucite is disconcerting (Fig. 6). The origin of this
172	systematic error may stem from the overall lower peak intensities (i.e., higher
173	background) and/or a poorer modeling of the larger ADP of the H atom. Nevertheless,
174	we give more credibility to the split-site model which overall offers slightly better fits to
175	the data for each temperature. The refinements for the deuterated brucite give the same
176	results for both the single-site and split-site models. That the splitting does not reduce to
177	zero at zero temperature, reinforces the view that the H/D atom is positional disordered
178	and does not in fact reside on the 3-fold axis.
179	That the differences in the unit cell dimensions between deuterated and
180	hydrogenated brucite increase with reducing cell volume as show here are consistent with
181	our previous finding that the isotope effect increases with pressure. Our results show that
182	the volume contraction at 20-295 K corresponds to that of pressure increase to ca. 0.4
183	GPa at room temperature (Horita et al. 2010). In essence, the lower zero-point energy of
184	deuterium implies a smaller vibration amplitude, hence smaller molecular or unit-cell
185	size, than that of hydrogen. A quantized harmonic oscillator model suggests that the
186	zero-point motion is inversely proportional to the square-root of the reduced mass, and
187	hence smaller for the deuterated compounds. However, the effect of temperature is
188	opposite, to expand the effective size of deuterium faster than that of hydrogen. Thus, for

189	a system with low-frequency vibrations, deuterated compounds may even become larger
190	than their hydrogenated counterparts at sufficiently high temperatures (Dunitz et al.
191	2008). This study at low temperature and our previous study at elevated pressures (Horita
192	2010) together show that in case of brucite (and likely other hydrous minerals),
193	deuterated minerals have smaller unit-cell volume, probably in the entire P-T space. This
194	measured effect of H–D substitution on the unit-cell volume demonstrates that brucite
195	(and other hydrous minerals) preferentially incorporate deuterium over hydrogen under
196	pressure (Horita et al. 2010), with the implications that the distribution of hydrogen
197	isotopes in deep-earth conditions may differ significantly from that in near-surface
198	environments.
199	Several studies suggested that high-pressure compression accentuates proton
200	disorder in brucite. As the octahedral layers are compressed along the c axis, the
201	interlayer and OH(D) distances become shorter, while the intralayer O-H(D) bond
202	length stays nearly constant or shortens only slightly. Then, hydrogen atoms of two O-
203	H(D) groups from adjacent octahedral layers may start to repulse (H-H repulsion),
204	causing such structural changes. However, two high-pressure neutron diffraction studies,
205	using hydrogenated and deuterated brucite (Catti et al. 1995; Parise et al. 1994) provided
206	somewhat contradicting results. Our study of the two brucites here clearly demonstrate
207	differences in O-H(D) and H(D)H(D) parameters during thermal contraction. If the
208	changes observed in these parameters upon cooling are similar to those under pressure, an
209	implication from the high-pressure neutron diffraction studies cannot be substantiated,
210	namely DD repulsion is stronger than HH repulsion, because the DD distance is
211	greater than HH (Fig. 7). To test this hypothesis, a companion high-resolution neutron

212	diffraction study is needed of $Mg(OH)_2$ and $Mg(OD)_2$ with a joint Rietveld refinement.
213	This and our previous study (Horita et al., 2010) clearly demonstrate that caution must be
214	exercised in using deuterated compounds for neutron scattering experiment to avoid
215	incoherent scattering from hydrogen.
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220	sponsored by the Scientific User Facilities Division, Office of Basic Energy Sciences, U.
221	S. Department of Energy.
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336	

337 Figure Captions

- Figure 1. Typical Rietveld refinement fits for hydrogenated (top) and deuterated (bottom)
 brucite using 1.1176 Å neutrons.
- 340 Figure 2. Temperature dependence of the lattice parameters and the unit cell volume of
- $Mg(OD)_2$ and $Mg(OH)_2$. $Mg(OH)_2$ has a 0.31% larger volume than that of $Mg(OD)_2$ at
- room temperature, which is mostly due to the c-axis expansion of Mg(OH)₂ as
- 343 compared to Mg(OD)₂. The isotope effect in the a-axis is opposite, but smaller, than
- that in the *c*-axis.
- 345 Figure 3. Temperature dependence of the isotropic atomic displacement parameters
- 346 (ADPs) from joint refinement of Mg(OD)₂ and Mg(OH)₂. (left) Single-site model
- 347 shows larger displacements of H as compared to D due to the mass difference. (right)

348 Split-site model results in H/D displacements collapsed to common values.

- Figure 4. Temperature dependence of the O-D and O-H bond lengths of $Mg(OD)_2$ and
- $Mg(OH)_2$. Both the single site and split site models yield same values within the
- 351 standard uncertainties. The distances have been corrected for the large thermal motion
- of the D and H atoms by d(XO) corrected = d(XO) + (Ux Uo)/d(XO), where U's are
- 353 the displacements normal to the bond and X = H or D.
- Figure 5. Temperature dependence of the Mg-O bond length of Mg(OD)₂ and Mg(OH)₂.
- Little change is observed over this temperature range, and larger errors and scatter are
- 356 observed for Mg(OH)₂. For Mg(OH)₂ the Mg-O distance is model dependent.
- Figure 6. Comparison of the temperature dependence of the *z*-coordinates of the O atom
- 358 (top) and the D or H atom (bottom) in $Mg(OD)_2$ and $Mg(OH)_2$, where again the
- $Mg(OH)_2$ values are model dependent.

360	Figure 7. (top) Temperature dependence of the nearest-neighbor DD and HH
361	distances in $Mg(OD)_2$ and $Mg(OH)_2$ for the single-site model. (bottom) Temperature
362	dependence of the H/D site splitting as modeled with isotropic atom displacement
363	parameters.
364	Figure 8. Temperature dependence of the octahedral layer thickness and the interlayer
365	thickness in Mg(OD) ₂ . The octahedral layer thickness remains constant, whereas the
366	interlayer thickness contraction accommodates the observed <i>c</i> -axis contraction. This
367	behavior is expected given the weak interaction between the layers, but is opposite to
368	that reported by Chakoumakos et al. (1997) from an analysis of t-o-f neutron diffraction
369	data.
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375	



Fig 1



Fig 2



Fig 3



Fig 4



Fig 5



Fig 6



Fig 7



Fig 8