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3 **Compressibility and structural stability of two variably hydrated olivine**
4 **samples (F_{0.97}Fe₃) to 34 GPa by X-Ray diffraction and Raman spectroscopy**

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ABSTRACT

27 Two hydrous olivines of composition $\text{Fo}_{97}\text{Fa}_3$ with water content of 4883 parts per million by
28 weight (ppmw) (SZ0407A) and 8000 ppmw (SZ0407B) were synthesized at 1250 °C and 12
29 GPa. Single-crystal X-ray diffraction was used to determine unit cell parameters of SZ0407A
30 and SZ0407B at pressures up to 7.1 GPa at room temperature. Synchrotron powder X-ray
31 diffraction and Raman scattering experiments were performed on sample SZ0407A in a
32 diamond-anvil cell to 34 GPa at room temperature. For both samples, the compressibility is
33 the largest along the *b*-axis and smallest along the *a*-axis. Using the compression (V/V_0) vs.
34 pressure data for sample SZ0407A to 29.08 GPa, in conjunction with the third-order Birch-
35 Murnaghan equation of state, we calculate the isothermal bulk modulus and its pressure
36 derivative as $K_0 = 119.4(15)$ GPa and $K'_0 = 6.6(5)$. Single-crystal compression data for sample
37 SZ0407A to 7 GPa give $K_0 = 121.5(6)$ GPa and $K'_0 = 5.7(2)$; and for sample SZ0407B $K_0 =$
38 $122.2(12)$ GPa and $K'_0 = 6.2(4)$. High-pressure Raman spectra for SZ0407A up to 34 GPa
39 show a continuous shift of all the observed bands to higher frequency with increasing pressure;
40 there is no indication of any first-order phase transition. However, the Raman spectra indicate
41 subtle discontinuous changes around 22 GPa, unobserved in previously reported studies on
42 anhydrous olivines.

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44 **Keywords:** Hydrous olivine, X-ray diffraction, Raman spectroscopy, hydration mechanism,
45 high pressure

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57 INTRODUCTION

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59 The nominally anhydrous mineral phases in the Earth's deep upper mantle and
60 transition zone (410-660 km depth) may serve as a large internal reservoir of water that has
61 profound implications for Earth's evolution as a water planet. Water plays important roles in
62 the phase transformation kinetics in the Earth's interior, which affects the mantle dynamics
63 (e.g. Kubo et al. 1998). Trace or minor hydroxyl in these phases affects strength and rheology
64 (Chen et al. 1998; Kavner 2003; Mei and Kohlstedt 2000), and electrical conductivity (Poe et
65 al. 2010). Therefore, it is important to understand the state of water and its role in the
66 structure of minerals in the Earth's interior. Much effort has been put into estimating the OH⁻
67 storage capacity of nominally anhydrous mineral phases to determine factors controlling OH⁻
68 incorporation into the structure (Bai and Kohlstedt 1993; Bell et al. 2003; Hushur et al. 2009;
69 Khisina et al. 2001; Kohlstedt et al. 1996; Mackwell and Kohlstedt 1990; Matveev et al. 2001;
70 Miller et al. 1987; Mosenfelder et al. 2006; Rossman and Smyth 1990; Skogby et al. 1990;
71 Smyth 1994; Walker et al. 2007). These studies have revealed that hydroxyl occurrence in
72 nominally anhydrous mantle phases is closely related to cation vacancy defects. Although
73 point defects have been linked to hydrogen incorporation, there remain several questions as to
74 which point defect predominates under various *P,T* conditions.

75 Olivine (Mg,Fe)₂SiO₄ is the most abundant mineral phase in the upper mantle to a
76 depth of 410 km, and it is nominally anhydrous. Although natural olivines typically contain
77 less than 200 ppmw H₂O by weight, samples synthesized at pressures of 12-14 GPa have been
78 obtained with up to 8900 ppmw, enough to significantly affect elastic and other physical
79 properties (Litasov et al. 2007; Smyth and Jacobsen 2006) as well as the overall hydrogen
80 budget of the planet. This has generated broad interest in systematic investigation of the
81 various structural and elastic (seismic) properties, as well as the stability of these phases as a
82 function of water content. X-ray diffraction (XRD) can give the full image of a crystal
83 structure averaged over a sample volume; however information on local structural disorder,

84 dynamics and defects with small concentrations may be limited. In contrast, spectroscopy is
85 very sensitive to the local structural disorder and defects. A combination of both XRD and
86 spectroscopic studies may give important additional information on crystal structures such as
87 disorder and dynamics, and may even help avoiding pitfalls in structure solution (Libowitzky
88 2006).

89 The XRD technique has been applied in various previous studies to understand the
90 structure and compressibility of both hydrous and anhydrous olivines. Isothermal compression
91 methods, deploying powder or single-crystal X-ray diffraction techniques, yield unit cell
92 volume as a function of pressure (P - V relations) which enables calculation of isothermal bulk
93 modulus (K_T) and its pressure derivative ($\partial K_T/\partial P$). The isothermal bulk modulus K_T is related
94 to the adiabatic bulk modulus K_S by

$$95 \quad K_S = K_T (1 + \alpha\gamma T),$$

96 where α is the thermal expansion coefficient ($\sim 2.6 \times 10^{-5} \text{ K}^{-1}$) (Fei 1995; Kroll et al. 2012; Ye
97 et al. 2009), γ is the Grüneisen parameter (≈ 1) (Anderson 1989), and T is temperature (K). In
98 general, K_S is roughly one to two percent larger than K_T . Table 1 summarizes elasticity data
99 for anhydrous and hydrous olivines. One of the major findings is that the effect of hydration is
100 to decrease bulk modulus (K_0) and increase K_0' by increasing the vacancy or unoccupied
101 volume in the structure.

102 Anhydrous olivine has been widely studied by vibrational spectroscopies. The
103 pressure dependencies of the Raman and IR spectra at room temperature have been measured
104 in several studies for the end-members forsterite and fayalite, as well as for some intermediate
105 compositions. Both Raman (Besson et al. 1982; Chopelas 1990; Durben et al. 1993; Gillet et
106 al. 1991; Wang et al. 1993) and IR data (Hofmeister 1997; Hofmeister et al. 1989; Williams et
107 al. 1990) show that the high frequency modes ($> 500 \text{ cm}^{-1}$), related to the internal vibrations
108 of the isolated SiO_4 tetrahedra of the olivine structure have smaller mode Grüneisen
109 parameters than the lattice modes, which involve vibrations of the MgO_6 octahedra. This
110 effect is consistent with the larger compressibility and relative weakness of the Mg-O bond
111 compared to the Si-O bond as shown by high pressure x-ray diffraction studies (Hazen 1976;
112 Kudoh and Takéuchi 1985), and discussed in detail by Smyth et al. (2000). Vibrational studies
113 to date have not been carried out for Fe-bearing olivines with significant hydration. We report

114 here single-crystal XRD data up to 7.1 GPa, powder synchrotron XRD data up to ~29 GPa,
115 and Raman spectroscopic data of synthesized hydrous $\text{Fo}_{97}\text{Fa}_3$ olivine samples in the diamond
116 anvil cell up to 34 GPa. Raman spectra of high-quality single crystals of hydrous olivine were
117 measured in the wide spectral range of 190 – 3700 cm^{-1} . Our aim was to investigate the effect
118 of water on the compressibility and vibrational properties, the local structure of hydrous Fe-
119 bearing olivine, and to track changes in H bonding and proton environments under pressure.

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121 **EXPERIMENTAL METHODS**

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123 Synthesis of these hydrous olivines was carried out in double-capsule experiments in
124 the 5000-ton multi-anvil press at Bayerisches Geoinstitut at 12 GPa and 1250°C (Smyth et al.
125 2006). In these experiments, pyroxene-normative and periclase-normative compositions were
126 formulated with about 3 weight percent H_2O . The two compositions were welded into
127 separate inner Pt capsules and packed with brucite into a welded outer capsule. When the
128 inner capsules were opened free liquid water emerged. The silica-excess capsule contained
129 olivine, enstatite, and a small amount of apparent quench melt phase. The silica deficient
130 capsule contained olivine, clinohumite and a quenched melt phase. The olivine, clinohumite
131 and clinoenstatite were identified by Raman spectroscopy and by single crystal X-ray
132 diffraction. The grain size of the silica excess sample was about 100 μm , whereas that of the
133 silica deficient capsule was about 250 μm . H_2O contents were measured by polarized FTIR
134 spectroscopy on X-ray-oriented, faceted single crystals. The contents were computed from
135 integrated absorbances using the calibration of Bell et al. (2003) for olivine. Sample SZ0407A
136 contained 4883 (± 500) ppmw H_2O , whereas sample SZ0407B contained 8000 (± 500) ppmw.

137 Unit-cell parameters were measured as a function of pressure up to 7.1 GPa by single-
138 crystal X-ray diffraction. For the single crystal X-ray diffraction we used a steel gasket T-301
139 and methanol:ethanol 4:1 mixture was used as transmitting pressure,
140 which is hydrostatic at least up to 9.4 GPa as demonstrated by Angel et
141 al. (2007). The X-ray source was the MoK α with working conditions at
142 50 kV and 40 mA. Two crystals (SZ0407A and B) were loaded into the diamond cell along
143 with a quartz crystal for an internal pressure standard. Cell parameters of the two crystals and
144 quartz were measured at 13 different pressures up to 7.1 GPa using a Huber four-circle

145 diffractometer operating under the program SINGLE (Angel and Finger 2011). Pressures
146 were determined using the quartz calibration curve of Angel et al. (1997).

147 High-pressure synchrotron powder X-ray diffraction studies in angle-dispersive mode
148 on the hydrous olivine (SZ0407A) were carried out up to 29 GPa at the GeoSoilEnviroCARS
149 beam line (13-BM-D) at the Advanced Photon Source (APS) using a monochromatic X-ray
150 beam of energy of 37.07 keV ($\lambda = 0.3344 \text{ \AA}$). Finely powdered samples of grain size 1-3 μm
151 were loaded in a piston-cylinder type diamond-anvil cell (DAC) of culet size 250 μm . We
152 used a steel gasket T-301 and silicone oil was used as the pressure transmitting medium. The
153 pressure was estimated using the shift in the R1 line of the ruby loaded along with the sample
154 (Mao et al. 1986). The average time for measurement at each pressure was 10 minutes. Fit2D
155 (Hammersley et al. 1996) was used to reduce the data to 2-theta vs. intensity profiles. The
156 unit-cell parameters were estimated from the d -spacings at each pressure using the program
157 Unit Cell (Holland and Redfern 1997). We have used the peaks (1 1 2), (1 3 0), (1 0 2), (0 2 1)
158 and (0 2 0) for the lattice parameter calculation.

159 Raman spectra were recorded using a triple monochromator Dilor XY spectrometer
160 equipped with a liquid nitrogen-cooled charge-coupled-device (CCD) detector. The 514.5 nm
161 green line of an Ar ion laser from Spectra Physics was used to excite the sample. The laser
162 light was focused using an OLYMPUS microscope with a long distance 50 \times objective to a
163 spot of about 3 μm diameter in the sample with the laser power being 25 mW on the sample.
164 No laser heating effects should occur under such conditions. The spectrometer was calibrated
165 using single-crystal silicon as a reference. For high-pressure measurements in the diamond
166 anvil cell, a rhenium gasket was pre-indented to 42 μm thickness and an 85 μm diameter hole
167 was drilled at the center of the indentation. A randomly oriented single crystal fragment 35 \times
168 40 \times 20 μm^3 in size was mounted in a diamond-anvil cell where diamonds had 350 μm
169 diameter culets. Two pieces of ruby \sim 5 μm in size were placed at the two sides of the sample
170 for pressure calibration. Liquid argon was loaded as a pressure-transmitting medium using a
171 liquid nitrogen cooling method. Pressure was estimated using the shift in the R1 line of the
172 ruby loaded along with the sample (Mao et al. 1986). All spectra were recorded in the
173 backscattering geometry with no polarization used for the collected signal.

174

175 **RESULTS AND DISCUSSION**

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177 Unit cell parameters of hydrous Fo₉₇ containing 4883 ppmw H₂O (SZ0407A) and
178 hydrous Fo₉₇ containing 8000 ppmw H₂O (SZ0407B) measured at ambient conditions and 12
179 pressures by single-crystal XRD are reported in Table 2. The measured pressure dependencies
180 of the unit-cell volumes are plotted in Figures 1 and 2. The program EOSFIT5.2 (Angel 2002)
181 was used to fit equation of state parameters to the observed compression data. The fitting is
182 weighted according to the estimated error of pressure determination and fitting error of
183 volume determination. The isothermal bulk modulus and its derivative were estimated using
184 the third order Birch-Murnaghan equation of state.

185
$$P(V) = \frac{3K_o}{2} \left[\left(\frac{V_o}{V} \right)^{\frac{7}{3}} - \left(\frac{V_o}{V} \right)^{\frac{5}{3}} \right] \left\{ 1 + \frac{3}{4} (K'_o - 4) \left[\left(\frac{V_o}{V} \right)^{\frac{2}{3}} - 1 \right] \right\}$$

186 where K_o is the bulk modulus at ambient pressure, K'_o is its first pressure derivative,
187 and V_o is the reference volume. Single crystal data to 7.1 GPa give $K_o = 121.5(6)$ GPa and K'_o
188 $= 5.7(2)$ for sample SZ0407A, and $K_o = 122.2(12)$ GPa and $K'_o = 6.2(4)$ for sample SZ0407B.
189 The K_o and K'_o of SZ0407A and SZ0407B show that slightly different water content (~ 3200
190 ppmw) has minimal effect on the bulk modulus within experimental uncertainty. The values
191 are $K_o = 119.2(12)$ GPa and $K'_o = 6.6(0.4)$ for combined single-crystal and powder diffraction
192 data for sample SZ0407A, and no distinctive changes were observed in the isothermal volume
193 compressibility. The compression data for all samples are thus in very good agreement and
194 indicate a decrease of about 6 percent in bulk modulus relative to anhydrous forsterite data
195 cited above. Comparison of bulk modulus of hydrous Fo₉₇ containing 8000 ppmw H₂O with
196 hydrous Fo₉₅ containing 8000 ppmw H₂O (Smyth et al. 2005; Smyth and Jacobsen 2006)
197 which is 120(2) GPa shows that the 2 mol% Fe has no effect on K_o within experimental
198 uncertainty. Although we observe a decrease in bulk modulus with hydration, the effects of
199 hydration on elasticity of forsterite may not be simple. Mao et al. (2010) report a crossover of
200 the velocities of hydrous and anhydrous forsterite measured by Brillouin scattering at high
201 pressure.

202 The axial compressibility is the largest along the b axis and smallest along the a axis
203 as shown in Figure 3. The axial compression values were $\beta_a = 0.00166(2)$, $\beta_b = 0.00309(4)$,
204 and $\beta_c = 0.00243(3)$ GPa⁻¹ for SZ0407A, respectively. For SZ0407B, the axial compression
205 values were $\beta_a = 0.00158(3)$, $\beta_b = 0.00308(4)$, and $\beta_c = 0.00240(3)$ GPa⁻¹. The $a:b:c$ axial
206 compressibility ratio was 1:1.9:1.5 for both SZ0407A and SZ0407B. The anisotropic nature of
207 the compression is a common feature for all anhydrous olivines. The compression along the b
208 axis depends on the compression of the most compressible structural unit, the M2 octahedron.
209 Our result is consistent with studies on anhydrous olivines, with the b axis being the most
210 compressible axis. However, our results also indicate an increase of the compressibility along
211 the a and c axis as compared to the anhydrous forsterite and anhydrous olivine Fo₈₃Fa₁₇
212 (Andraut et al. 1995; Smyth et al. 2000; and references therein), which may be related to the
213 hydration mechanism in hydrous olivine (SZ0407A, SZ0407B).

214 Raman spectra in the lattice and silicate vibrational region of hydrous olivine Fo₉₇Fa₃
215 containing 4883 ppmw water (SZ0407A) measured as a function of pressure to 34 GPa are
216 shown in Figure 4. The ambient spectra are similar to those reported in previous Raman
217 studies (Besson et al. 1982; Chopelas 1990, 1991; Durben et al. 1993; Gillet et al. 1991; Iishi
218 1978; Piriou and McMillan 1983; Wang et al. 1993). With increasing pressure, all the twelve
219 observed modes show a continuous shift to higher frequency. Of these, four modes exhibit a
220 linear pressure dependence of frequency; however the remaining eight modes are
221 characterized by a nonlinear response. For these 8 modes, the frequency vs. pressure data
222 were fitted with quadratic equations as shown in Figures 5 and 6. According to the original
223 definition for a vibrational mode of frequency ν_i , a mode Grüneisen parameter is defined by

$$224 \quad \gamma_i = -\frac{d \ln(\nu_i)}{d \ln(V)}$$

225 where V is volume. γ_i were determined by fitting $\ln(\nu_i)$ versus $\ln(V)$ values linearly, and are
226 shown in Table 3. γ_i of the mode related to Mg(2)O₆ translation (Table 3) does not show
227 marked changes in hydrous olivine Fo₉₇Fa₃ (SZ0407A) relative to anhydrous forsterite. γ_i of
228 all the observed modes related to SiO₄ tetrahedra show an increase as compared to those for
229 anhydrous forsterite, indicating Si-O bonds in hydrous olivine Fo₉₇Fa₃ (SZ0407A) are more
230 compressible than Si-O bonds in anhydrous olivine. This may be responsible for the increase
231 of the compressibility along the a and c axis as compared to the anhydrous forsterite and

232 anhydrous olivine $\text{Fo}_{83}\text{Fa}_{17}$ (Andrault et al. 1995; Smyth et al. 2000; and references therein).
233 Our Raman study on Fe-free hydrous forsterite does not show such an increase in γ_i of the
234 modes related to SiO_4 tetrahedra (Hushur et al. 2009), which may indicate that the hydration
235 mechanism is different in hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ (SZ0407A) as compared to the hydrous
236 forsterite. May be, a significant protonation of Si-vacancies also exist in this hydrous olivine
237 $\text{Fo}_{97}\text{Fa}_3$.

238 In the entire pressure range of this study, the sharp Raman bands assigned to
239 vibrations of the olivine crystal structure are retained, indicating no evidence of any phase
240 transition. However, the high-pressure spectra do indicate some subtle structural modification.
241 Specifically, in the low frequency spectra range between 200 cm^{-1} and 800 cm^{-1} , the M2
242 translation mode at 302.5 cm^{-1} , SiO_4 rotation mode at 325.2 cm^{-1} and ν_4 mode at 542.3 cm^{-1}
243 disappears, whereas the ν_4 mode at 606.9 cm^{-1} intensifies in the high pressure Raman spectra
244 as shown in Figure 5. Both 542.3 cm^{-1} and 606.9 cm^{-1} modes are initially A_g modes and
245 represent anti-symmetric bending of the oxygens in SiO_4 tetrahedra. The disappearance of the
246 542.3 cm^{-1} mode above 22 GPa, while the mode at 606.9 cm^{-1} is intensified, is not an effect of
247 sample orientation. This may indicate that the two bands represent vibrations of different
248 structural units that appear or disappear above 22 GPa. Interestingly, the similarity of the
249 high-pressure spectra with those of γ -spinel becomes more pronounced at high pressures: an
250 intense band with a zero pressure location near 600 cm^{-1} , and a doublet at higher frequencies
251 (Chopelas et al. 1994; Kleppe et al. 2002). The decrease in number of the Raman active
252 modes at high pressure may indicate that the olivine structure is modified locally above 22
253 GPa, with the partial appearance of Si-O-Si linkages accompanied by a possible metastable
254 increase in the silicon coordination. In our recent Raman study on Fe-free hydrous forsterite
255 we observed a change in the pressure derivatives of many of the observed Raman modes at
256 ~ 20 GPa (Hushur et al. 2009). However, in the current study on Fe-bearing hydrous forsterite
257 we do not observe such a slope change at ~ 20 GPa. This may be related to different local
258 environments in Fe-bearing hydrous forsterite as compared to the Fe-free hydrous forsterite.

259 The olivine structure is orthorhombic containing isolated SiO_4 units linked by
260 octahedrally coordinated divalent cations. The divalent cations are in two crystallographically
261 distinct sites, the smaller M1 (C_i symmetry) site and larger M2 (C_s symmetry) site. The M1
262 octahedron shares six edges with neighboring polyhedra, two with SiO_4 tetrahedra. The M2

263 octahedra share three edges with neighboring polyhedra but only one with a SiO₄ tetrahedron.
264 The unit cell parameters increase in Fe-bearing hydrous forsterite as compared to the Fe-free
265 hydrous forsterite that has the same amount of hydrogen present. Kudoh et al. (2006) and
266 Kudoh et al. (2007) reported on a single-crystal X-ray structural analysis of Fe-bearing
267 (Mg_{1.85}Fe_{0.14}Si_{0.99}H_{0.06}O₄) and Fe-Free (Mg_{1.985}Si_{0.993}H_{0.06}O₄) hydrous forsterite, in which
268 increases in the *a*, *b*, and *c* axes due to the effect of Fe were 0.13%, 0.08%, and 0.03%,
269 respectively, and showing a 0.2 % increase in the unit cell volume due to the presence of Fe.
270 In Fe-bearing hydrous forsterite, cation vacancies predominantly occurred at the M2 site
271 (Kudoh 2008; Kudoh et al. 2007), in contrast to the case of Fe-free hydrous forsterite in which
272 the cation vacancies predominantly occur at the M1 site (Hushur et al. 2009; Kudoh 2008;
273 Kudoh et al. 2006; Smyth et al. 2006). A vacancy at the M1 site leads to an increase in the
274 M2-O bond length in Fe-free hydrous forsterite. In case of the Fe-bearing hydrous forsterite,
275 Fe atoms lead to an increase in both the M1-O and M2-O bond lengths. The difference in
276 proton occupancy between Fe-bearing hydrous forsterite and the Fe-free hydrous forsterite
277 may be responsible for the different pressure dependence of the Raman modes observed in the
278 Fe-bearing hydrous forsterite as compared to the Fe-free hydrous forsterite.

279 The Raman modes of the hydroxyl groups in hydrous olivine (SZ0407A) are shown in
280 Figure 7. We fit the observed spectra at ambient pressure by four Gaussian functions using
281 ORIGIN Pro 7.5 software. The spectra above 9 GPa are fitted with three Gaussian functions.
282 The mode frequency as a function of pressure obtained by fitting is shown in Figure 8. At
283 ambient pressure, four Raman bands were observed at 3555, 3578, 3613, and 3656 cm⁻¹. It is
284 similar to the FTIR spectra measured on the sample SZ0407B and SZ0408A (Smyth et al.
285 2006), and Raman spectra of the sample SZ408A (Hushur et al. 2009). However, an
286 additional Raman mode at 3656 cm⁻¹ is observed only on this hydrous olivine (SZ407A). As
287 pressure increases, the Raman active OH stretching modes of hydrous olivine (SZ0407A)
288 with frequencies < 3600 cm⁻¹ have strong, negative pressure dependence whereas OH modes
289 at higher frequencies are less sensitive to pressure. According to the empirical relation
290 between OH-stretching frequencies and O···O distances (Libowitzky 1999), a negative
291 frequency response of OH modes under pressure requires a reduction of the respective O···O
292 distances and indicates a strengthening of the hydrogen bond.

293 The protonation of hydrous olivine has been studied extensively by FTIR spectroscopy.
294 Previous studies show that there are two different mechanisms related to silicon and
295 magnesium vacancies by which hydrogen is incorporated into olivine crystal structure. Bai
296 and Kohlstedt (1993) and Kohlstedt et al. (1996) suggested protonation of Mg-vacancies in
297 the olivine structure, whereas Litasov et al. (2009), Matveev et al. (2001) and Lemaire et al.
298 (2004) preferred significant protonation of Si-vacancies. Assuming linear O-H...O topologies,
299 we can correlate the observed OH stretching frequencies with O...O bond distances
300 (Libowitzky 1999). Octahedral edges in the olivine structure range from 2.85 to 3.33 Å except
301 for edges shared with the tetrahedron, which are about 2.55 Å. Tetrahedral edges range from
302 2.55 to 2.76 Å. The largest and broad peak in the pattern at 3555 cm⁻¹ (in the *b*-direction) is
303 consistent with protonation of the shared O3-O3 edges (2.76 Å) of the M2 octahedron. The
304 mode at 3578 cm⁻¹ could correspond to the unshared O1-O3 edge of M1 (2.84 Å). The peak at
305 3613 cm⁻¹ corresponds to protonation of an O-O distance of > 2.8 Å (Libowitzky 1999), and
306 is consistent with the O1-O2 edge shared between M1 octahedra (2.85Å). There is also the
307 possibility of the O3-O3 edge of the M2 octahedron (2.99 Å), but a proton on this edge would
308 require a strongly bent O-H...O angle to be consistent with polarizations. The mode at 3656
309 cm⁻¹ corresponds to protonation of an O-O distance of > 2.9 Å (Libowitzky 1999) and is
310 consistent with the O3-O3 edge of the M2 octahedron (2.99 Å) in this Fe-bearing olivine.
311 There might be significant deviations from the linear O-H...O topologies, or the stretching
312 mode at 3656 cm⁻¹ is not from the hydrogen bonded O-H group. Khisina et al. (2001) reported
313 an OH band at 3656 cm⁻¹ which is related to the 10 Angstrom phase. However the strong
314 intensity of this band in our spectra indicates that it is not from the minor phases in olivine.
315 And also the single crystal data of our synthetic sample did not show any monoclinic phase in
316 our studied sample. Therefore, it is not likely from the 10 Angstrom phase. May be the band
317 at 3656 cm⁻¹ is related to the Si vacancies. The band at 3656 cm⁻¹ shows splitting into two
318 modes above 15 GPa and this may be related to the observed local structural modification in
319 this study. It would be enlightening to carry out complementary FTIR measurements on such
320 a hydrous olivine sample. The above initial findings may aid in refining the models for
321 structural stability and compressibility (seismic velocities) of hydrous phases in the upper
322 mantle.
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Table and Figure Captions

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511 Table 1. Elastic properties of olivines.

512 Table 2. Unit-cell parameters of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water

513 (SZ0407A) and hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 8000 ppmw water (SZ0407B) as

514 a function of pressure as determined by single-crystal X-ray diffraction.

515 Table 3. Mode frequencies (in cm^{-1}), mode Grüneisen parameters, and mode assignments for

516 observed Raman modes of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water

517 (SZ0407A), and of anhydrous forsterite.

518 Figure 1. Unit-cell volumes of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water

519 (SZ0407A, open squares) and hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 8000 ppmw water

520 (SZ0407B, open circles) as a function of pressure obtained from single-crystal XRD

521 studies. Error bars are also shown in the Figure. Solid and dashed lines represent the

522 two best-fit equations of state.

523 Figure 2. Unit-cell volumes of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water

524 (SZ0407A) as a function of pressure. Open symbols represent powder XRD data.

525 Solid symbols are single-crystal XRD data. Solid and dashed lines represent the two

526 best-fit equations of state, based on powder and single-crystal XRD, and single-crystal

527 XRD, respectively. Note that the symbols are larger than the uncertainties.

528 Figure 3. Axial compressions a/a_0 , b/b_0 and c/c_0 as a function of pressure for a hydrous

529 olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) and for hydrous olivine $\text{Fo}_{97}\text{Fa}_3$

530 containing 8000 ppmw water (SZ0407B). Note that the symbols are larger than the

531 uncertainties.

532 Figure 4. Raman spectra of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A)

533 to 33.9 GPa.

534 Figure 5. Pressure dependences of the observed low frequency Raman modes for hydrous

535 olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) from 200 ~ 750 cm^{-1} .

536 Figure 6. Pressure dependences of the four high frequency Raman modes for hydrous olivine

537 $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) from 800 ~ 1200 cm^{-1} .

538 Figure 7. Pressure dependence of Raman spectra of hydrous olivine $Fo_{97}Fa_3$ containing 4883
539 ppmw water (SZ0407A) in the OH-stretching region. The spectra are offset for clarity.
540 The solid red line in the 1 bar spectrum is a fit using four Gaussian functions. Labels
541 at right side indicate pressure in gigapascals and numbers on the 1 bar spectrum
542 indicate O-H stretching mode frequencies obtained from the fit.

543 Figure 8. Pressure dependence of the Raman modes of hydroxyl group in hydrous olivine
544 $Fo_{97}Fa_3$ containing 4883 ppmw water (SZ0407A).

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Tables and Figures

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551 Table 1. Elastic properties of olivines

Olivine composition	Method	Pressure range (GPa)	K_s (GPa)	K_s'	K_T (GPa)	K_T'	References
Anhydrous							
Forsterite (Fo ₁₀₀)	Brillouin spectroscopy	32	129	4.2			Zha et al. (1998)
Forsterite (Fo ₁₀₀)		Ambient pressure	129.5				Bass (1995)
Fo ₉₁ Fa ₉	Ultrasonic	3 GPa	129.5	4.65			Bass (1995)
Fo ₉₃ Fa ₇		Ambient pressure	129.4				Bass (1995)
Fo _{91.3} Fa _{8.7}		Ambient pressure	128.1				Bass (1995)
Fo ₉₂ Fa ₈		Ambient pressure	126.7				Bass (1995)
San Carlos olivine Fo ₉₀	Impulsive-stimulated scattering	12	129.4(5)	4.29	126.3	4.28	Abramson et al. (1997)
Forsterite (Fo ₁₀₀)	Single-crystal XRD	17.2			125(2)	4.0(4)	Downs et al. (1996)
Forsterite (Fo ₁₀₀)	Single-crystal XRD	15			122.6	4.3	Kudoh and Takéuchi (1985)
Forsterite (Fo ₁₀₀)	Powder Synchrotron XRD	30			135.7(10)	3.9(10)	Will et al. (1986)
Hydrous							
Fo ₉₅ (8000 ppmw H ₂ O)	Single-crystal XRD	8			120(2)		Smyth et al. (2005)
Fo ₁₀₀ (0.8-0.9 wt% H ₂ O)	Brillouin scattering		125.7(2)				Jacobsen et al. (2008, 2009)
Fo ₉₇ (0.8-0.9 wt% H ₂ O)	Brillouin scattering		125.2(8)				Jacobsen et al. (2008, 2009)
Fo ₉₇ Fa ₃ (4883 ppmw H ₂ O)	Single-crystal XRD	7			121.5(6)	5.7(2)	This study
Fo ₉₇ Fa ₃ (8000 ppmw H ₂ O)	Single-crystal XRD	7			122.2(12)	6.3(10)	This study

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Table 2. Unit cell parameters of hydrous olivines SZ0407A and SZ0407B as a function of pressure as determined by single-crystal X-ray diffraction.

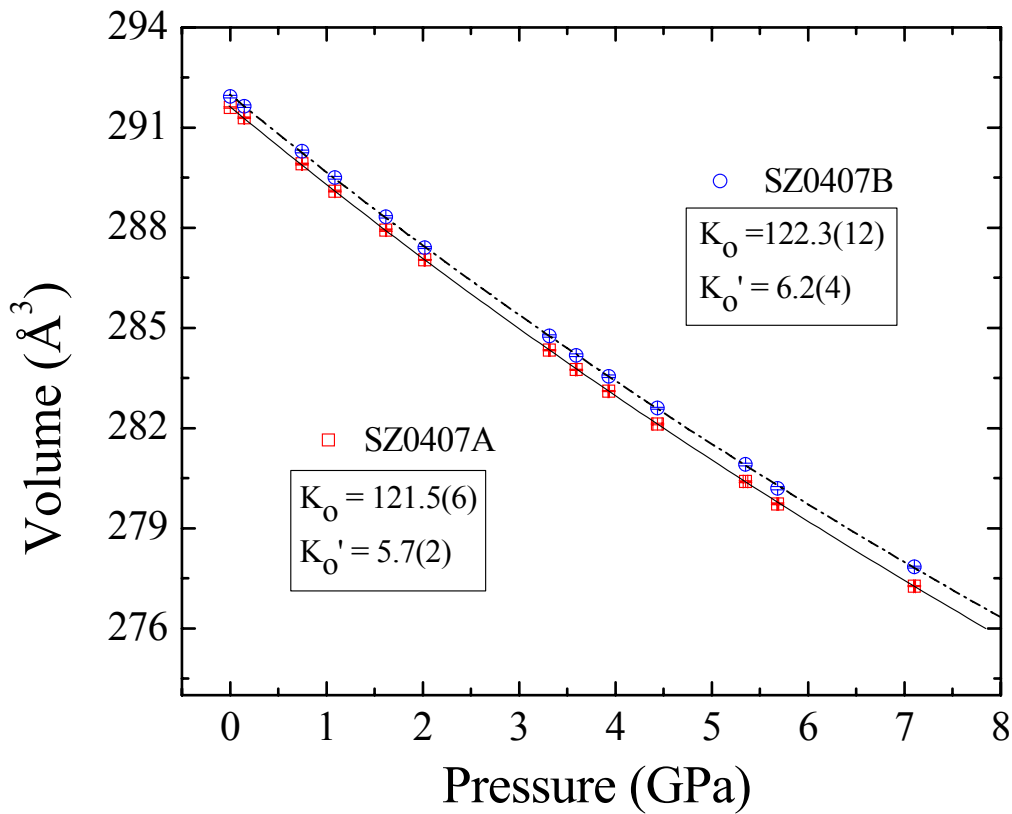
Pressure(GPa)	SZ0407A					SZ0407B				
P(GPa)	<i>a</i> (Å)	<i>b</i> (Å)	<i>c</i> (Å)	V(Å ³)	V/V ₀	<i>a</i> (Å)	<i>b</i> (Å)	<i>c</i> (Å)	V(Å ³)	V/V ₀
0.0001	4.7590(3)	10.2238(2)	5.9932(1)	291.60(2)	1.00000	4.7624(6)	10.2260(3)	5.9936(2)	291.93(4)	1.00000
0.147(4)	4.7583(4)	10.2187(3)	5.9909(2)	291.30(3)	0.9990(1)	4.7624(6)	10.2209(3)	5.9914(2)	291.64(4)	0.9990(1)
0.748(4)	4.7532(5)	10.1974(3)	5.9813(2)	289.91(3)	0.9942(1)	4.7572(6)	10.2003(3)	5.9821(2)	290.29(4)	0.9944(1)
1.089(4)	4.7493(5)	10.1861(3)	5.9760(2)	289.10(3)	0.9914(1)	4.7544(7)	10.1883(3)	5.9767(2)	289.50(4)	0.9917(1)
1.619(5)	4.7455(5)	10.1674(3)	5.9675(2)	287.93(3)	0.9874(1)	4.7503(6)	10.1698(3)	5.9684(2)	288.33(3)	0.9877(1)
2.020(6)	4.7418(2)	10.1543(2)	5.9615(1)	287.04(1)	0.98436(5)	4.7461(6)	10.1561(3)	5.9625(3)	287.41(3)	0.9845(1)
3.314(8)	4.7314(3)	10.1130(2)	5.9424(2)	284.33(2)	0.97508(6)	4.7365(6)	10.1152(3)	5.9436(2)	284.76(3)	0.9754(1)
3.593(7)	4.7292(3)	10.1040(2)	5.9383(1)	283.75(2)	0.97309(6)	4.7341(6)	10.1066(2)	5.9396(2)	284.18(4)	0.9735(1)
3.929(7)	4.7267(3)	10.0941(2)	5.9337(1)	283.11(2)	0.97088(6)	4.7322(7)	10.0963(3)	5.9347(2)	283.55(4)	0.9713(1)
4.44(5)	4.7231(4)	10.0790(3)	5.9266(2)	282.13(2)	0.96750(8)	4.7288(6)	10.0813(2)	5.9279(2)	282.60(3)	0.9680(1)
5.35(2)	4.7164(4)	10.0524(3)	5.9142(2)	280.40(3)	0.96157(9)	4.7222(6)	10.0553(2)	5.9161(2)	280.91(4)	0.9623(1)
5.684(9)	4.7138(4)	10.0421(3)	5.9094(2)	279.73(3)	0.95929(9)	4.7190(8)	10.0448(3)	5.9111(3)	280.20(5)	0.9598(2)
7.104(8)	4.7041(1)	10.0050(3)	5.8916(1)	277.27(2)	0.95084(6)	4.7109(5)	10.0073(2)	5.8936(2)	277.84(3)	0.9517(1)

582 Table 3. Mode frequencies (in cm^{-1}), mode Grüneisen parameters, and mode assignments for observed Raman modes of hydrous
 583 olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A), with hydrous and anhydrous forsterite.

Hydrous $\text{Fo}_{97}\text{Fa}_3$ (This study)			Hydrous Fo_{100} (Hushur et al. 2009)		Anhydrous Fo_{100} (Chopelas 1991)		
Symmetry	Mode frequency	Mode Grüneisen Parameters	Mode frequency	Mode Grüneisen Parameters	Mode frequency	Mode Grüneisen Parameters	Mode type
A_g	227.0	0.88	228.4	0.689	226	0.674	SiO_4 translation
A_g	302.5	1.62	305.6	1.31	304	1.63	M2 translation
B_{2g}	325.2	1.65			323		Mix (M2 translation)
B_{2g}	368.9	1.37			365		Mix (SiO_4 rotation)
B_{3g}	432.9	1.82	435.7	1.51	435	1.4	Mix (SiO_4 rotation)
A_g	542.3	0.69	545	0.467	545	0.528	ν_4
B_{2g}	588.4	0.85	588.7	0.481	592		ν_4
A_g	606.9	1.01	608.2	0.732	608	0.704	ν_4
A_g	823.1	0.77	824.4	0.582	824	0.483	$\nu_1+\nu_3$
A_g	854.5	0.655	856.2	0.426	856	0.489	$\nu_1+\nu_3$
B_{3g}	920.0	0.58	921		920	0.382	ν_3
A_g	960.0	0.878	967	0.586	965	0.661	ν_3

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587 Figure 1. Unit-cell volumes of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water

588 (SZ0407A, open squares) and hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 8000 ppmw

589 water (SZ0407B, open circles) as a function of pressure obtained from single-

590 crystal XRD studies. Error bars are also shown in figure. Solid and dashed lines

591 represent the two best-fit equations of state

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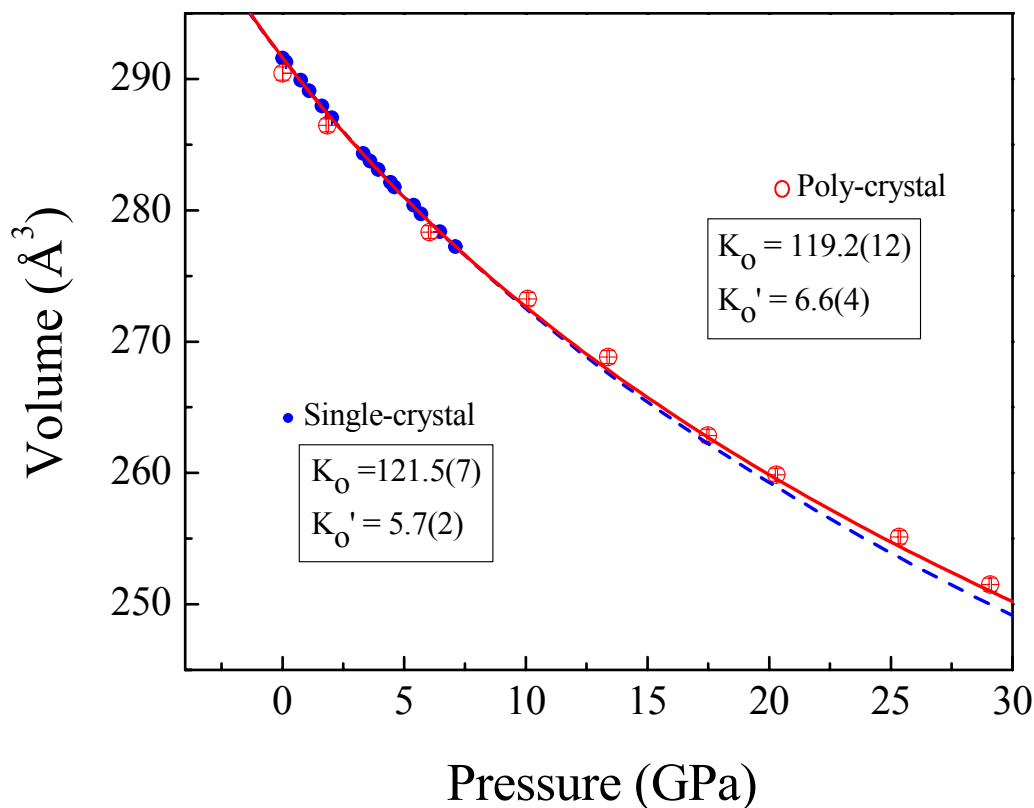
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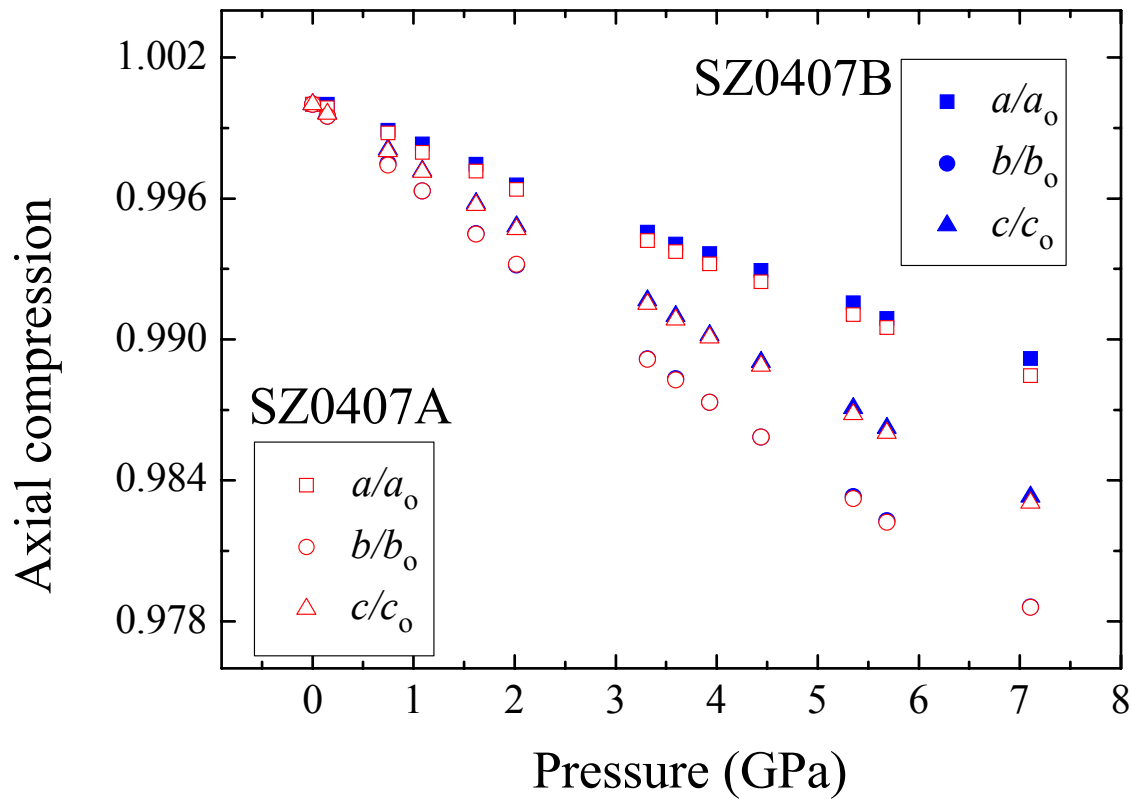
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Figure 2. Unit-cell volumes of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) as a function of pressure. Open symbols represent powder XRD data. Solid symbols are single-crystal XRD data. Solid and dashed lines represent the two best-fit equations of state, based on powder and single-crystal XRD, and single-crystal XRD respectively. Note that the symbols are larger than the uncertainties



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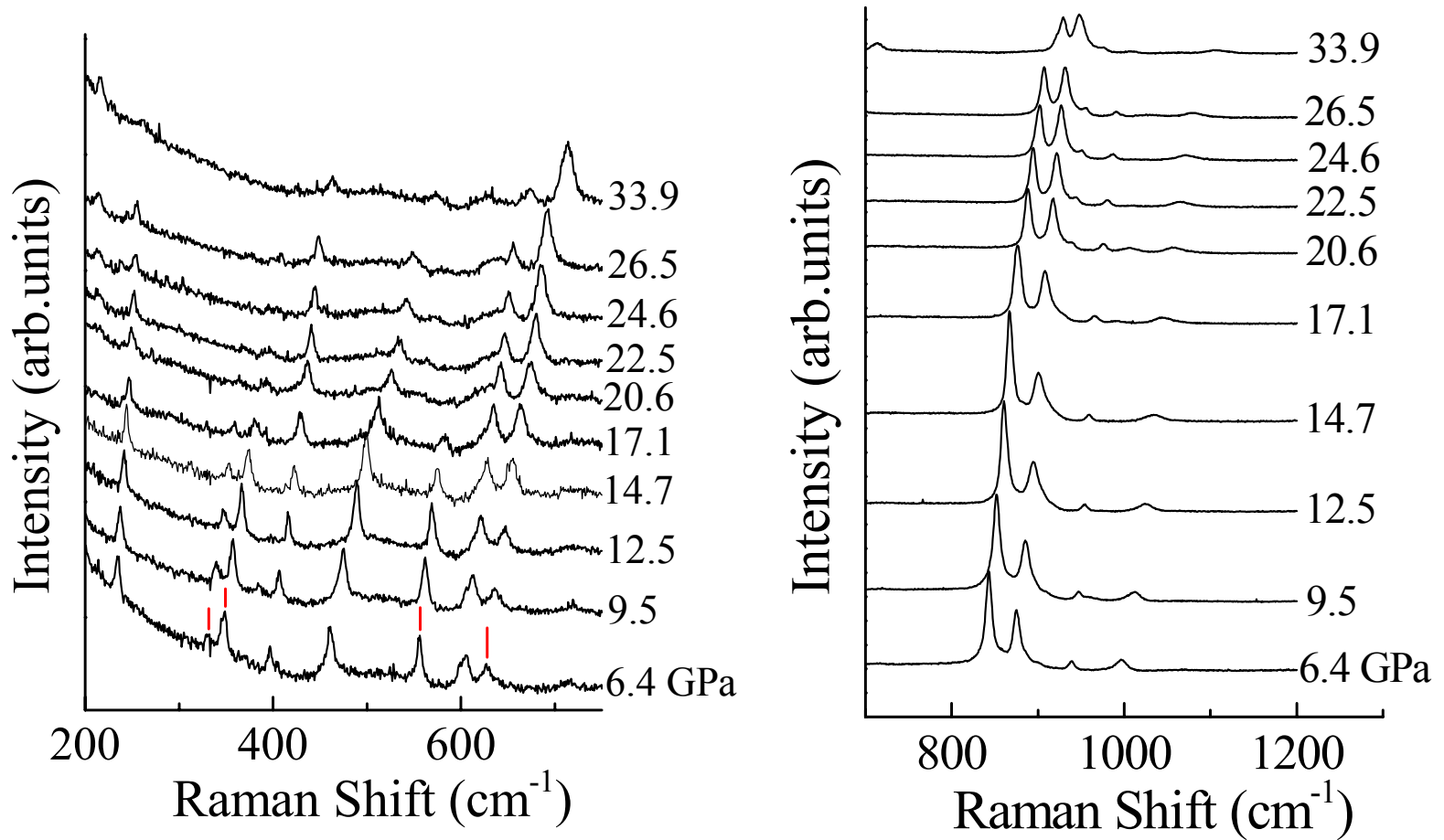
622 Figure 3. Axial compressions a/a_0 , b/b_0 and c/c_0 as a function of pressure for a hydrous

623 olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) and for hydrous

624 olivine $\text{Fo}_{97}\text{Fa}_3$ containing 8000 ppmw water (SZ0407B). Note that the symbols

625 are larger than the uncertainties

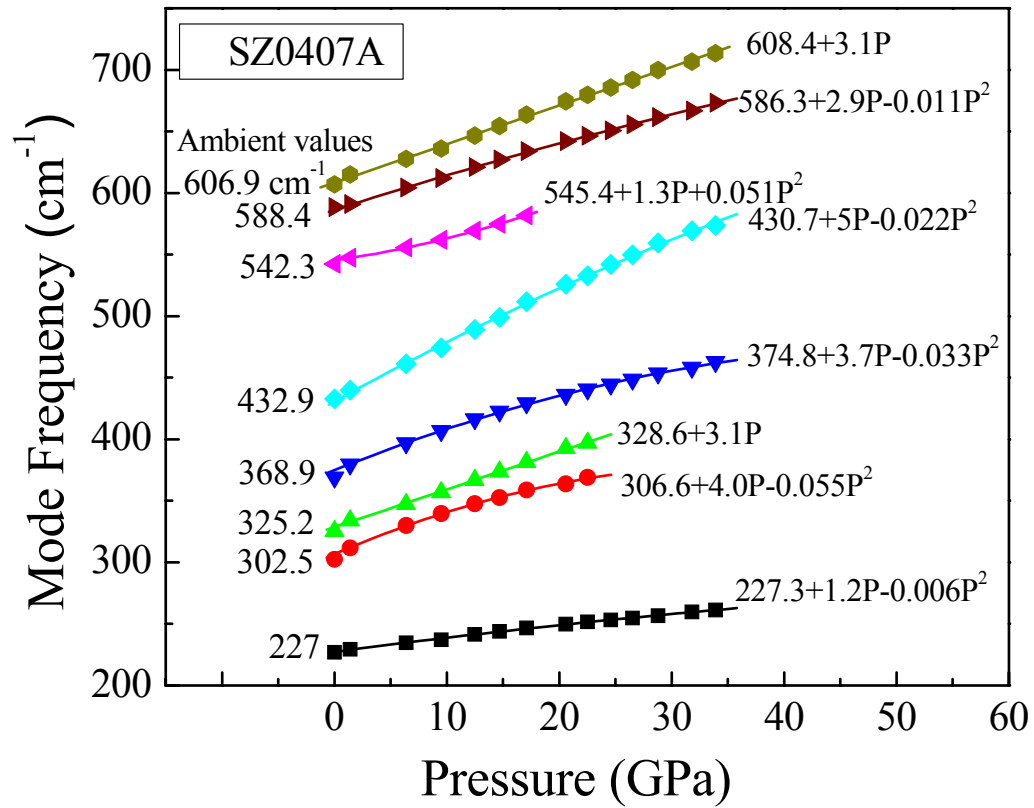
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Figure 4. Raman spectra of hydrous olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) to 33.9 GPa.

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639 Figure 5. Pressure dependences of the observed low frequency Raman modes for hydrous
640 olivine $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A) from 200 ~ 750 cm^{-1} .

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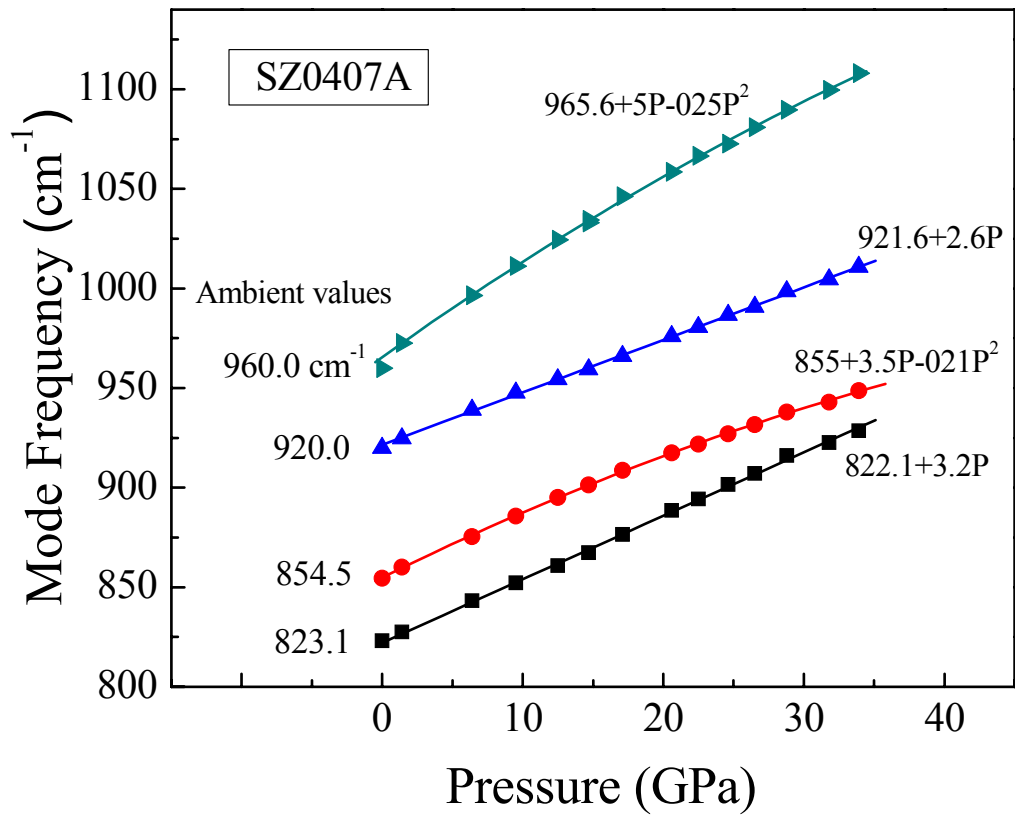
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650 Figure 6. Pressure dependences of the four high frequency Raman modes for hydrous
651 olivine Fo₉₇Fa₃ containing 4883 ppmw water (SZ0407A) from 800 ~ 1200 cm⁻¹.

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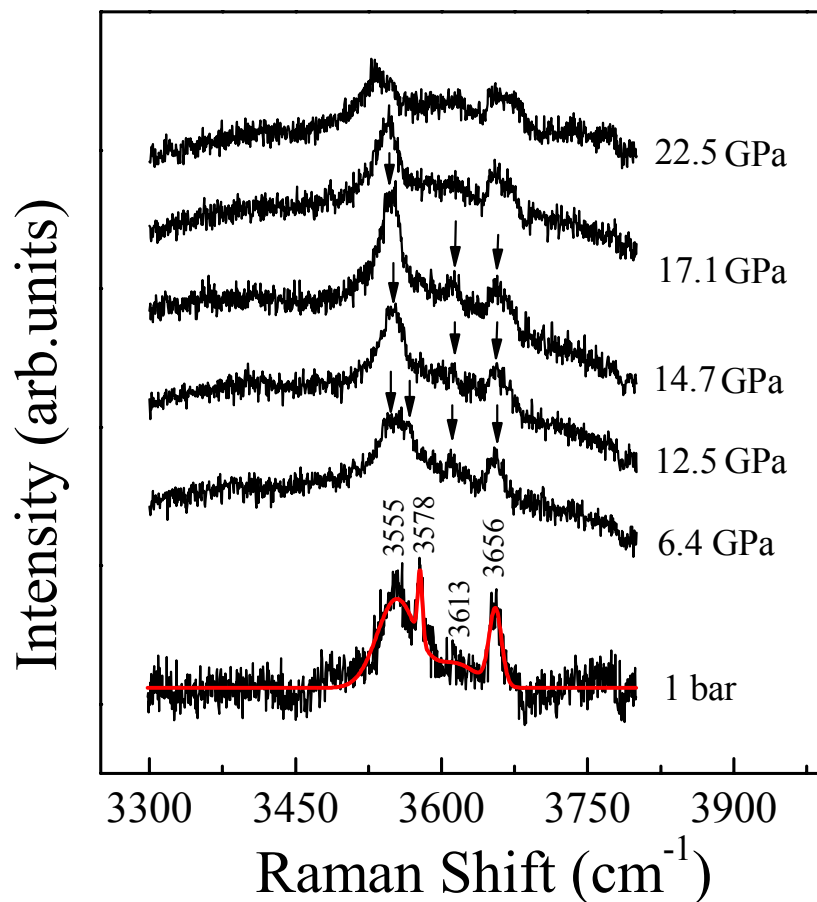
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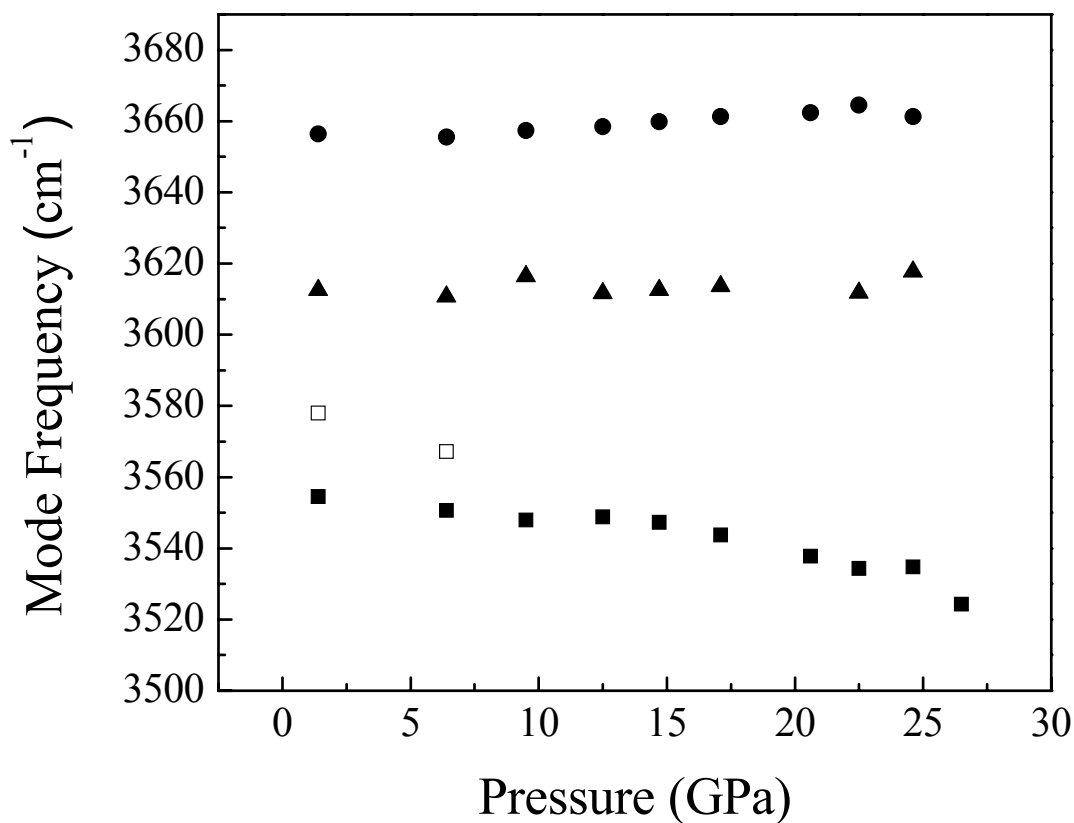
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662 Figure 7. Pressure dependence of Raman spectra of hydrous olivine Fo₉₇Fa₃ containing
663 4883 ppmw water (SZ0407A) in the OH-stretching region. The spectra are
664 offset for clarity. The solid red line in the 1 bar spectrum is a fit using four
665 Gaussian. Labels at right side indicate pressure in gigapascals and numbers on
666 the 1 bar spectrum indicate O-H stretching mode frequencies obtained from the
667 fit.

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670 Figure 8. Pressure dependence of the Raman modes of hydroxyl group in hydrous olivine

671 $\text{Fo}_{97}\text{Fa}_3$ containing 4883 ppmw water (SZ0407A).

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